

Time-like gravitational formfactors and shear viscosity



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Dedicated to the memory of
Anatoli Vasil'evich Efremov (1933-2021)

and

Maxim Polyakov (1966-2021)

Gravitational Formfactors (spin 1/2)

$$\langle p' | T_{q,g}^{\mu\nu} | p \rangle = \bar{u}(p') \left[A_{q,g}(\Delta^2) \gamma^{(\mu} p^{\nu)} + B_{q,g}(\Delta^2) P^{(\mu} i \sigma^{\nu)\alpha} \Delta_\alpha / 2M \right] u(p)$$

- Conservation laws - zero Anomalous Gravitomagnetic Moment : $\mu_G = J$ (g=2)

$$P_{q,g} = A_{q,g}(0) \quad A_q(0) + A_g(0) = 1$$

$$J_{q,g} = \frac{1}{2} [A_{q,g}(0) + B_{q,g}(0)] \quad A_q(0) + B_q(0) + A_g(0) + B_g(0) = 1$$

- May be extracted from high-energy experiments/NPQCD calculations
- Describe the partition of angular momentum between quarks and gluons
- Describe interaction with both classical and TeV gravity

Gravity and hadron structure: (OT'99)

- Interaction – field vs metric deviation

$$M = \langle P' | J_q^\mu | P \rangle A_\mu(q)$$

$$M = \frac{1}{2} \sum_{q,G} \langle P' | T_{q,G}^{\mu\nu} | P \rangle h_{\mu\nu}(q)$$

- Static limit

$$\langle P | J_q^\mu | P \rangle = 2e_q P^\mu$$

$$\sum_{q,G} \langle P | T_i^{\mu\nu} | P \rangle = 2P^\mu P^\nu$$
$$h_{00} = 2\phi(x)$$

$$M_0 = \langle P | J_q^\mu | P \rangle A_\mu = 2e_q M \phi(q)$$

$$M_0 = \frac{1}{2} \sum_{q,G} \langle P | T_i^{\mu\nu} | P \rangle h_{\mu\nu} = 2M \cdot M \phi(q)$$

- Mass as charge – equivalence principle



EP and hadron structure

“Microscopic” EP (coupling of gravity to EMT)

+

Conservation law

(Momentum SR to get local from LC: $\int dx \ x (\Sigma q(x) + G(x)) = 1$)

=

“Macroscopic” EP (universal falling) :

Tested VERY precisely



Gravitomagnetism

- Gravitomagnetic field (weak, except in gravity waves) – action on spin from $M = \frac{1}{2} \sum_{q,G} \langle P' | T_{q,G}^{\mu\nu} | P \rangle h_{\mu\nu}(q)$

$$\vec{H}_J = \frac{1}{2} \text{rot} \vec{g}; \quad \vec{g}_i \equiv g_{0i}$$

spin dragging twice
smaller than EM

- Lorentz force – similar to EM case: factor 1/2 cancelled with 2 from same as EM

$$h_{00} = 2\phi(x) \quad \text{armor frequency}$$

$$\omega_J = \frac{\mu_G}{J} H_J = \frac{H_L}{2} = \omega_L \quad \vec{H}_L = \text{rot} \vec{g}$$

- Orbital and Spin momenta dragging – the same - Equivalence principle



Experimental test of PNEP

Reinterpretation of the data on G(EDM) search

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Search for a Coupling of the Earth's Gravitational Field to Nuclear Spins in Atomic Mercury

B. J. Venema, P. K. Majumder, S. K. Lamoreaux, B. R. Heckel, and E. N. Fortson

Physics Department, FM-15, University of Washington, Seattle, Washington 98195
(Received 25 September 1991)

If (CP-odd!) $G_{EDM}=0 \rightarrow$ constraint for AGM (Silenko, OT'07) from Earth rotation – was considered as obvious (but it is just EP!) background

$$\mathcal{H} = -g\mu_N \mathbf{B} \cdot \mathbf{S} - \zeta \hbar \boldsymbol{\omega} \cdot \mathbf{S}, \quad \zeta = 1 + \chi$$

$$|\chi(^{201}\text{Hg}) + 0.369\chi(^{199}\text{Hg})| < 0.042 \quad (95\% \text{C.L.})$$



Quantum measurement and EP

If spin is just a (pseudo) vector : EP due to Earth rotation is trivial

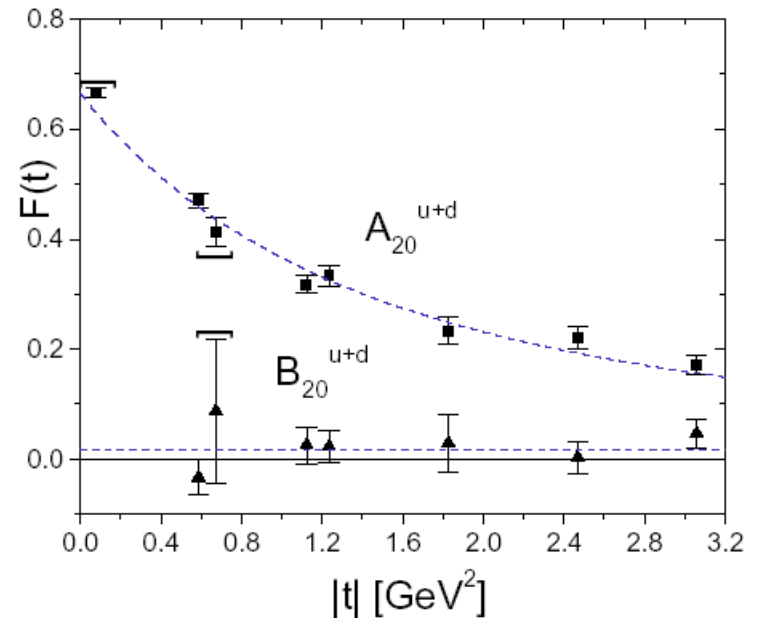
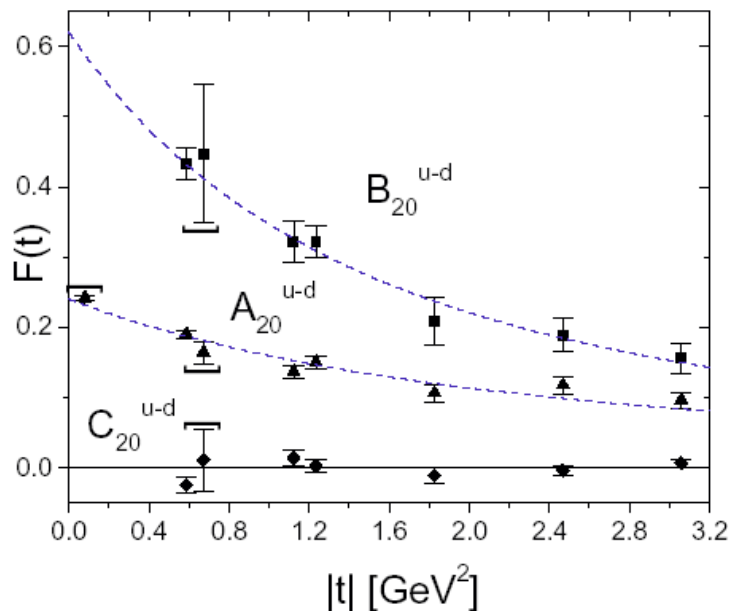
Crucial if measured by a device in rotating frame

Quantum measurement problem becomes practical

Cf Unruh effect in HIC (Prokhorov, OT, Zakharov'19) – talk of G. Prokhorov on Wednesday

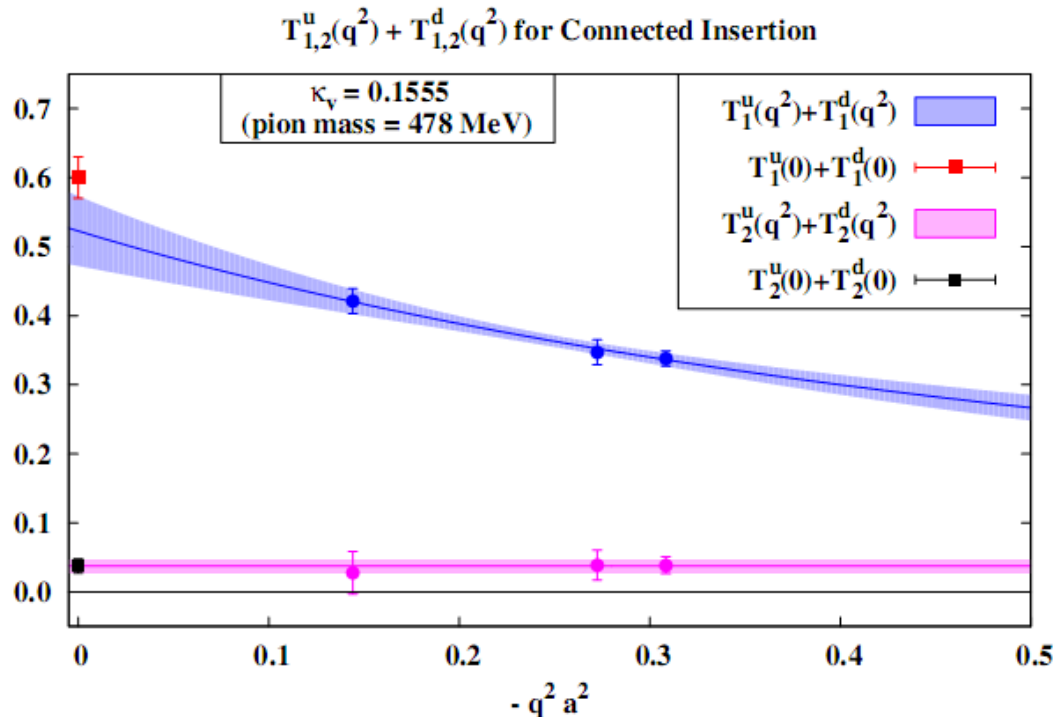
Generalization of Equivalence principle

- Various arguments: $AGM \approx 0$ separately for quarks and gluons – most clear from the lattice (LHPC/SESAM)



Recent lattice study (M. Deka et al. Phys.Rev. D91 (2015) no.1, 014505)

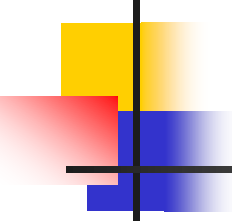
- Sum of u and d for Dirac (T1) and Pauli (T2) FFs



Extended Equivalence

Principle=Exact EquiPartition

- In QED, pQCD – violated (Brodsky et al)
- Reason – in the case of ExEP- no smooth transition for zero fermion mass limit (Milton, 73)
- Conjecture (O.T., 2001 – prior to lattice data) – valid in NP QCD – zero quark mass limit is safe due to chiral symmetry breaking
- Gravityproof confinement? Nucleons do not break even by black holes?
- Support by recent observation of smallness of (nucleon “cosmological constant”) C_{bar}



*One more gravitational formfactor
(related to "D-term" of Maxim
Polyakov and Christian Weiss)*

- *Quadrupole*

$$\langle P + q/2 | T^{\mu\nu} | P - q/2 \rangle = C(q^2)(g^{\mu\nu} q^2 - q^\mu q^\nu) + \dots$$

- *Cf vacuum matrix element –
cosmological constant*

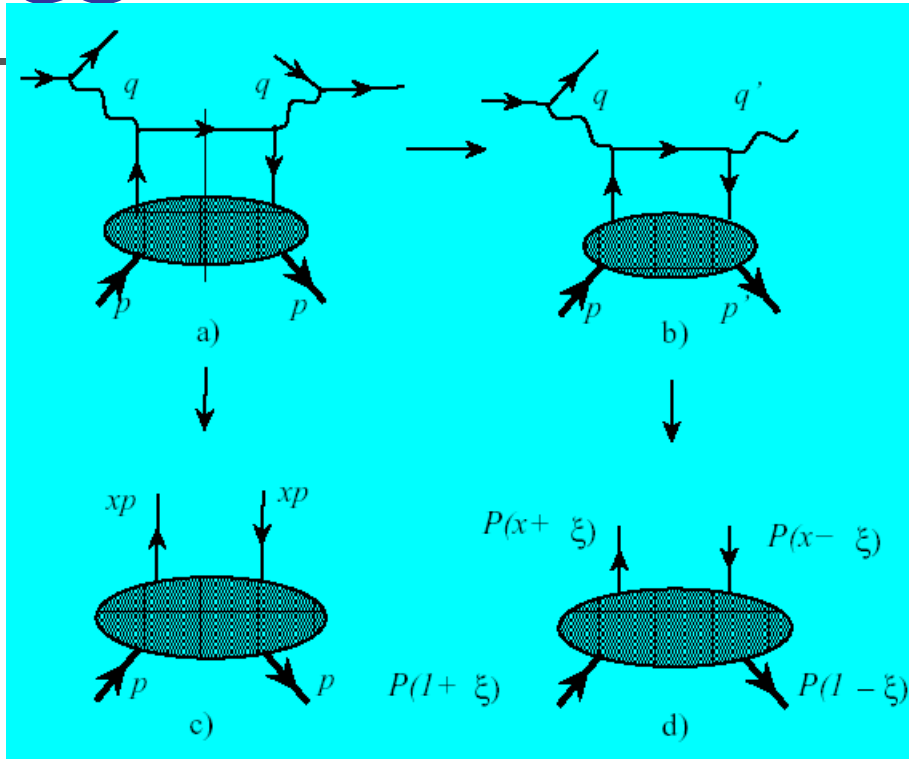
$$\langle 0 | T^{\mu\nu} | 0 \rangle = \Lambda g^{\mu\nu}$$

$$\Lambda = C(q^2) q^2$$

- *NO "vacuum-like" term – EP, Smallness
-ExEP*

- *How to measure experimentally – DVCS
(and DVMP?)*

QCD Factorization for DIS and DVCS



- *Manifestly spectral*

$$\mathcal{H}(x_B) = \int_{-1}^1 dx \frac{H(x)}{x - x_B + i\epsilon}$$

- *Extra dependence on ξ*

$$\mathcal{H}(\xi) = \int_{-1}^1 dx \frac{H(x, \xi)}{x - \xi + i\epsilon}$$



Unphysical regions

- DIS : Analytical function – polynomial in $1/x_B$ if $1 \leq |X_B|$

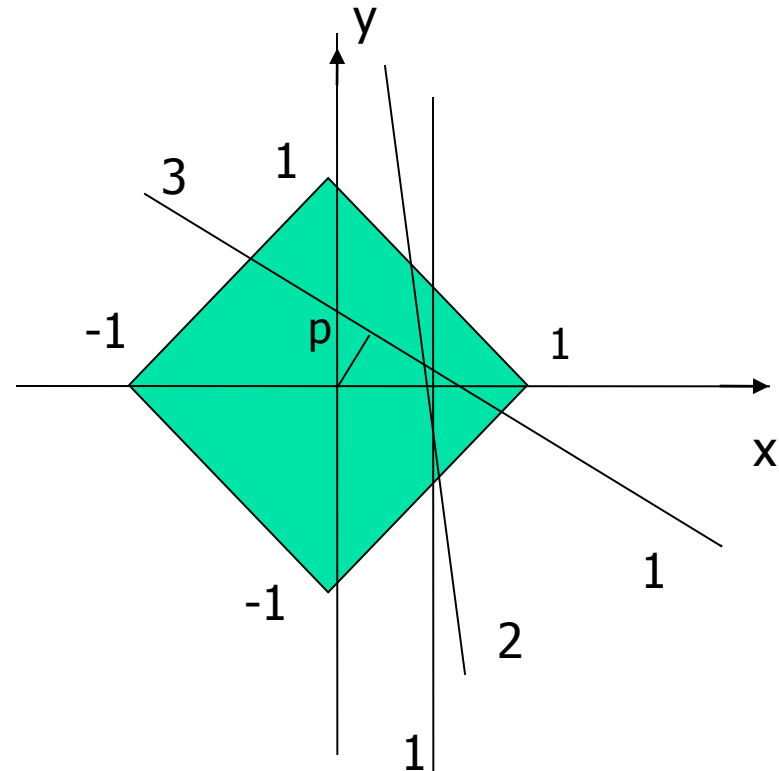
$$H(x_B) = - \int_{-1}^1 dx \sum_{n=0}^{\infty} H(x) \frac{x^n}{x_B^{n+1}}$$

- DVCS – additional problem of analytical continuation of $H(x, \xi)$
- Solved by using of Double Distributions Radon transform

$$H(z, \xi) = \int_{-1}^1 dx \int_{|x|-1}^{1-|x|} dy (F(x, y) + \xi G(x, y)) \delta(z - x - \xi y)$$

Double distributions and their integration

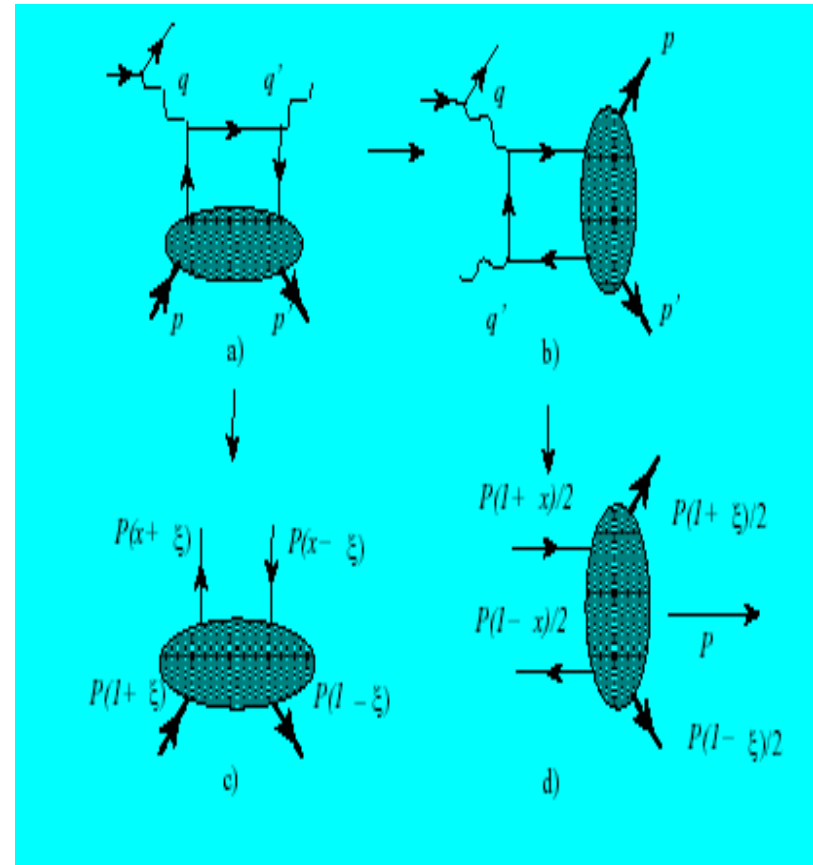
- Slope of the integration line-skewness
- Kinematics of DIS: $\xi = 0$
("forward") - vertical line (1)
- Kinematics of DVCS: $\xi < 1$
- line 2
- Line 3: $\xi > 1$ unphysical region - required to restore DD by inverse Radon transform: tomography



$$\begin{aligned}
 f(x, y) &= -\frac{1}{2\pi^2} \int_0^\infty \frac{dp}{p^2} \int_0^{2\pi} d\phi |\cos\phi| (H(p/\cos\phi + x + ytg\phi, t g\phi) - H(x + ytg\phi, t g\phi)) = \\
 &= -\frac{1}{2\pi^2} \int_{-\infty}^\infty \frac{dz}{z^2} \int_{-\infty}^\infty d\xi (H(z + x + y\xi, \xi) - H(x + y\xi, \xi))
 \end{aligned}$$

Crossing for DVCS and GPD

- DVCS \rightarrow hadron pair production in the collisions of real and virtual photons
- GPDs \rightarrow Generalized Distribution Amplitudes (Diehl, Gousset, Pire, OT'98)
- Straightforward generalization of DAs. Pioneered for pion FFs (Efremov & Radyushkin, Brodsky & Lepage '79)
- Duality between s and t channels (Polyakov, Shuvaev, Guzey, Vanderhaeghen, Pasquini...)



GDA -> back to unphysical regions for DIS and DVCS

- Recall DIS

$$H(x_B) = - \int_{-1}^1 dx \sum_{n=0}^{\infty} H(x) \frac{x^n}{x_B^{n+1}}$$

- Non-positive powers of x_B

- DVCS

$$H(\xi) = - \int_{-1}^1 dx \sum_{n=0}^{\infty} H(x, \xi) \frac{x^n}{\xi^{n+1}}$$

- Polynomiality (general property of Radon transforms): moments - integrals in x weighted with x^n - are polynomials in ξ of power $n+1$
- As a result, analyticity is preserved: only non-positive powers of ξ appear



Holographic property (OT'05)

■ *Factorization
Formula*

->

- *Analyticity ->
Imaginary part ->
Dispersion relation:*

$$\mathcal{H}(\xi) = \int_{-1}^1 dx \frac{H(x, \xi)}{x - \xi + i\epsilon}$$

$$\mathcal{H}(\xi) = \int_{-1}^1 dx \frac{H(x, x)}{x - \xi + i\epsilon}$$

$$\Delta\mathcal{H}(\xi) \equiv \int_{-1}^1 dx \frac{H(x, x) - H(x, \xi)}{x - \xi + i\epsilon}$$

- *"Holographic"
equation*

$$= \sum_{n=1}^{\infty} \frac{1}{n!} \frac{\partial^n}{\partial \xi^n} \int_{-1}^1 H(x, \xi) dx (x - \xi)^{n-1} = \text{const}$$



Holographic property - II

- Directly follows from double distributions

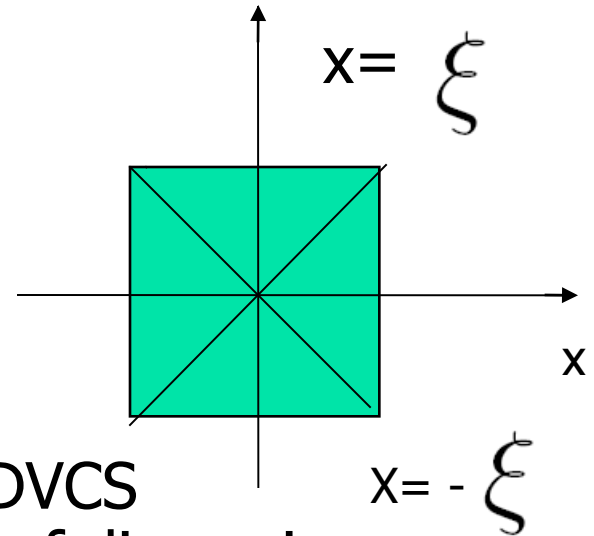
$$H(z, \xi) = \int_{-1}^1 dx \int_{|x|-1}^{1-|x|} dy (F(x, y) + \xi G(x, y)) \delta(z - x - \xi y)$$

- Constant is the SUBTRACTION one - due to the (generalized) Polyakov-Weiss term $G(x, y)$

$$\begin{aligned} \Delta \mathcal{H}(\xi) &= \int_{-1}^1 dx \int_{|x|-1}^{1-|x|} dy \frac{G(x, y)}{1 - y} \\ &= \int_{-\xi}^{\xi} dx \frac{D(x/\xi)}{x - \xi + i\epsilon} = \int_{-1}^1 dz \frac{D(z)}{z - 1} = \text{const} \end{aligned}$$

Holographic property - III

- 2-dimensional space \rightarrow 1-dimensional section!
- Momentum space: any relation to holography in coordinate space ?!
- Strategy (now adopted) of GPD's studies: start at diagonals
(through SSA due to imaginary part of DVCS amplitude) and restore by making use of dispersion relations + subtraction constants



Analyticity of Compton amplitudes in energy plane (Anikin, OT'07)

- Finite subtraction implied

$$\operatorname{Re}\mathcal{A}(\nu, Q^2) = \frac{\nu^2}{\pi} \mathcal{P} \int_{\nu_0}^{\infty} \frac{d\nu'^2}{\nu'^2} \frac{\operatorname{Im}\mathcal{A}(\nu', Q^2)}{(\nu'^2 - \nu^2)} + \Delta \quad \Delta = 2 \int_{-1}^1 d\beta \frac{D(\beta)}{\beta - 1}$$

$$\Delta_{\text{CQM}}^p(2) = \Delta_{\text{CQM}}^n(2) \approx 4.4, \quad \Delta_{\text{latt}}^p \approx \Delta_{\text{latt}}^n \approx 1.1$$

- Numerically close to Thomson term for real proton (but NOT neutron) Compton Scattering!
- Duality (sum of squares vs square of sum; proton: $4/9+4/9+1/9=1$)?!



From D-term to pressure

- *Inverse -> 1st moment (model)*
- *Kinematical factor – moment of pressure $D \sim -\langle p r^A \rangle$
($\langle p r^2 \rangle = 0$) M.Polyakov (2003)*

$$T_{\mu\nu}^Q(\vec{r}, \vec{s}) = \frac{1}{2E} \int \frac{d^3\Delta}{(2\pi)^3} e^{i\vec{r}\cdot\vec{\Delta}} \langle p', S' | \hat{T}_{\mu\nu}^Q(0) | p, S \rangle$$

$$T_{ij}(\vec{r}) = s(r) \left(\frac{r_i r_j}{r^2} - \frac{1}{3} \delta_{ij} \right) + p(r) \delta_{ij}$$

- *Possible justification: Born gravitational scattering*
- *Stable equilibrium $D < 0$: Holds for quarks (or leptons) in photon*

Pressure in hadron pairs production

- Back to GDA region
- -> moments of $H(x,x)$ - define the coefficients of powers of cosine! - $1/\xi$
- Higher powers of cosine in t-channel – threshold in s-channel (OT'10)
- Larger for pion than for nucleon pairs because of less fast decrease at $x \rightarrow 1$
- Large ξ limit – access to D-term

$$\begin{aligned} \mathcal{H}(\xi) &= - \int_{-1/\xi}^{1/\xi} dx \sum_{n=0}^{\infty} H(x, \xi) \frac{x^n}{\xi^{n+1}} \\ &= - \int_{-1/\xi}^{1/\xi} dx \sum_{n=0}^{\infty} H(x, x) \frac{x^n}{\xi^{n+1}} + \Delta \mathcal{H}. \end{aligned}$$

Gravitational FFs from Belle data on GDAs

S. Kumano, Qin-Tao Song and O. Teryaev, PRD 97 (2018) 014020.

Gravitational FFs are related to twist-2

GDAs

$$A_{\lambda_1 \lambda_2} = T_{\mu\nu} \varepsilon^\mu(\lambda_1) \varepsilon^\nu(\lambda_2) / e^2$$

$$A_{++} = \sum_q \frac{e_q^2}{2} \int_0^1 dz \frac{2z-1}{z(1-z)} \Phi^q(z, \xi, W^2)$$

$$\int dz (2z-1) \Phi_q^+(z, \xi, W^2) = \frac{2}{(P^+)^2} \langle \pi^+(p_1) \pi^-(p_2) | T_q^{++}(0) | 0 \rangle$$

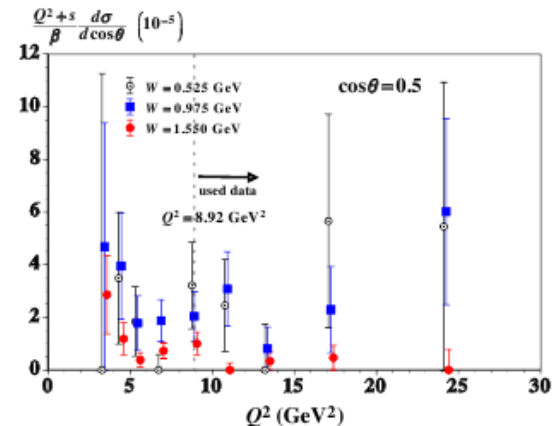
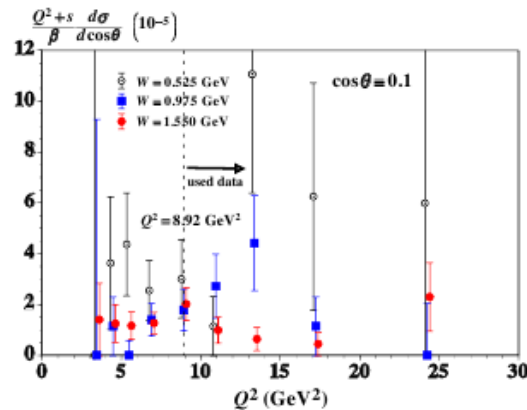
$$\langle \pi^0(p_1) \pi^0(p_2) | T^{\mu\nu}(0) | 0 \rangle = \frac{1}{2} \left[(sg^{\mu\nu} - P^\mu P^\nu) \Theta_1 + \Delta^\mu \Delta^\nu \Theta_2 \right]$$

$$P = p_1 + p_2, \Delta = p_1 - p_2$$

M. Masuda et al. [Belle Collaboration], PRD 93 (2016), 032003

Belle data and scaling : $W=0.525, 0.975, 1.55$ GeV

$$\frac{(Q^2+s)d\sigma}{\beta d|\cos\theta|} \propto \left| \Phi^{s^*s^*}(z, \cos\theta, W, Q) \right|^2$$



Phase shifts and resonances

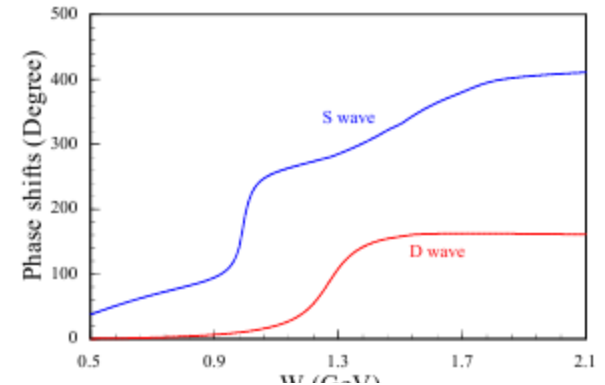
Leading harmonics

$$\begin{aligned} \sum_q \Phi_q^+(z, \xi, W^2) &= 18n_f z(1-z)(2z-1)[B_{10}(W) + B_{12}(W)P_2(2\xi-1)] \\ &= 18n_f z(1-z)(2z-1)[\tilde{B}_{10}(W) + \tilde{B}_{12}(W)P_2(\cos\theta)] \end{aligned}$$

$$\tilde{B}_{10}(W) = \bar{B}_{10}(W)e^{i\delta_0}, \quad \tilde{B}_{12}(W) = \bar{B}_{12}(W)e^{i\delta_2}$$

S/D shifts

$f_0(500)$, $f_2(1270)$
contributions



$$\bar{B}_{12}(W) = \beta^2 \frac{10g_{f_2\pi\pi} f_{f_2} M_{f_2}^2}{9\sqrt{2}\sqrt{(M_{f_2}^2 - W^2)^2 - \Gamma_{f_2}^2 M_{f_2}^2}}$$

$$\bar{B}_{10}(W) = \frac{5g_{f_0\pi\pi} f_{f_0}}{3\sqrt{2}\sqrt{(M_{f_0}^2 - W^2)^2 - \Gamma_{f_0}^2 M_{f_0}^2}}$$

Fits and results

Collection

$$\Phi_q^+(z, \xi, W^2) = N_h z^\alpha (1-z)^\alpha (2z-1) [\tilde{B}_{10}(W) + \tilde{B}_{12}(W) P_2(\cos \theta)]$$

$$\tilde{B}_{10}(W) = \left[\frac{-3 + \beta^2}{2} \frac{5R_\pi}{9} F_h(W^2) + \frac{5g_{f_0\pi\pi} f_{f_0}}{3\sqrt{2} \sqrt{(M_{f_0}^2 - W^2)^2 - \Gamma_{f_0}^2 M_{f_0}^2}} \right] e^{i\delta_0}$$

$$\tilde{B}_{12}(W) = \left[\beta^2 \frac{5R_\pi}{9} F_h(W^2) + \beta^2 \frac{10g_{f_2\pi\pi} f_{f_2} M_{f_2}^2}{9\sqrt{2} \sqrt{(M_{f_2}^2 - W^2)^2 - \Gamma_{f_2}^2 M_{f_2}^2}} \right] e^{i\delta_2}$$

$$F_h(W^2) = \frac{1}{\left[1 + \frac{W^2 - 4m_\pi^2}{\Lambda^2} \right]^{n-1}}$$

Best fit with (2) and without (1) f_0

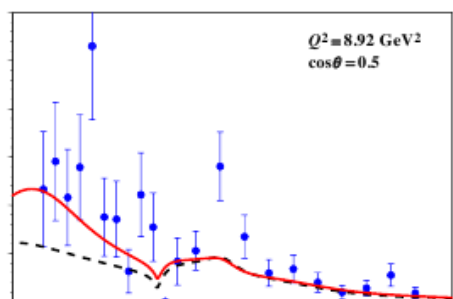
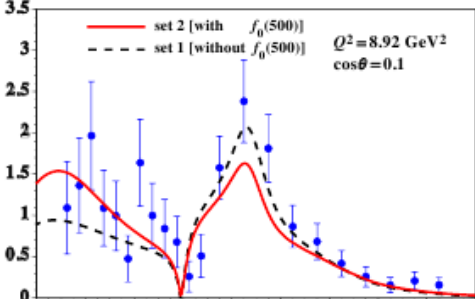
	Set 1	Set 2
α	0.801±0.042	1.157± 0.132
Λ	1.602±0.109	1.928±0.213
a	3.878± 0.165	3.800± 0.170
b	0.382± 0.040	0.407± 0.041
f_{f_0}	-----	0.0184± 0.034

$$\frac{\chi^2}{NOF} = 1.22$$

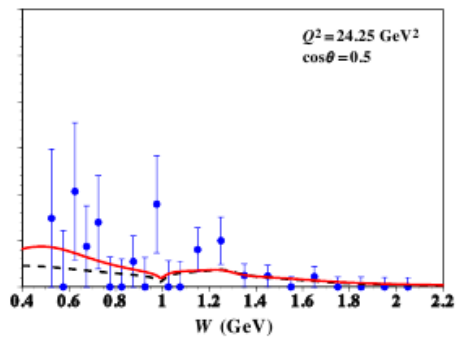
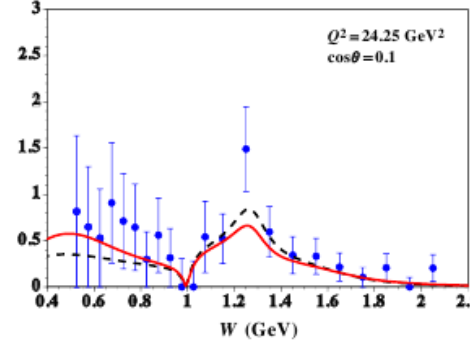
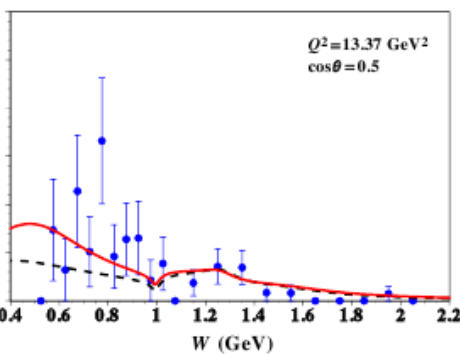
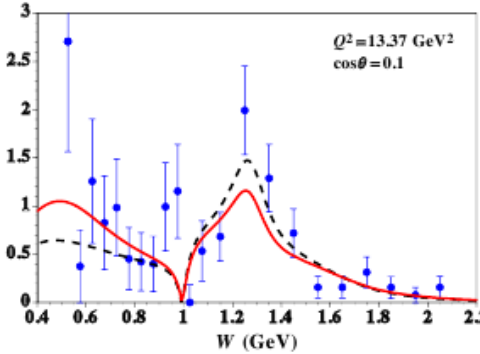
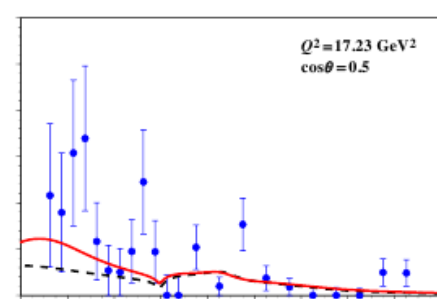
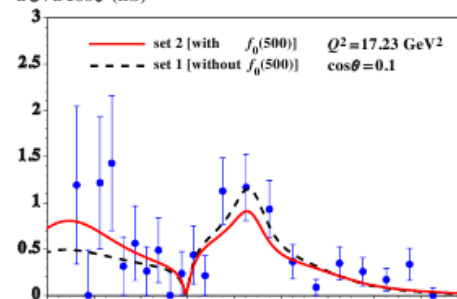
$$\frac{\chi^2}{NOF} = 1.09$$

Description of data

$d\sigma/d\cos\theta$ (nb)



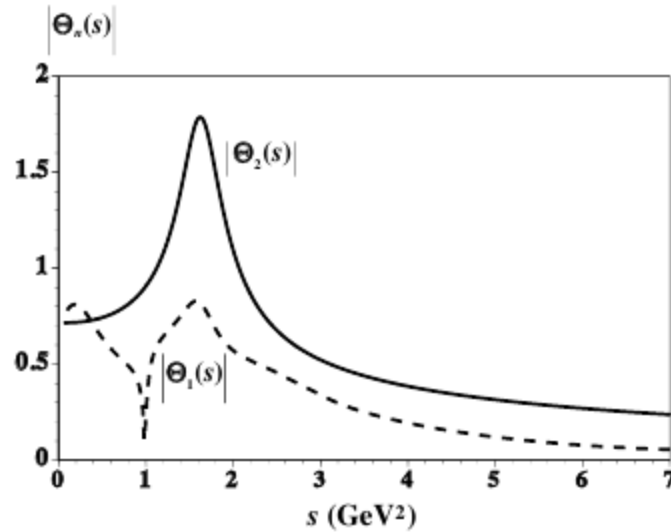
$d\sigma/d\cos\theta$ (nb)



Formfactors

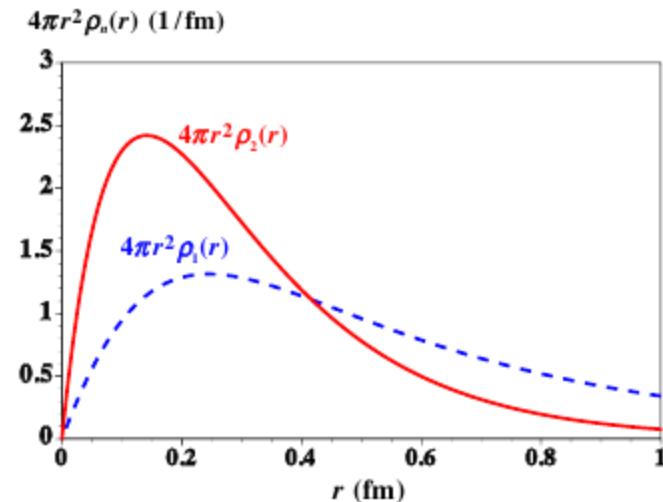
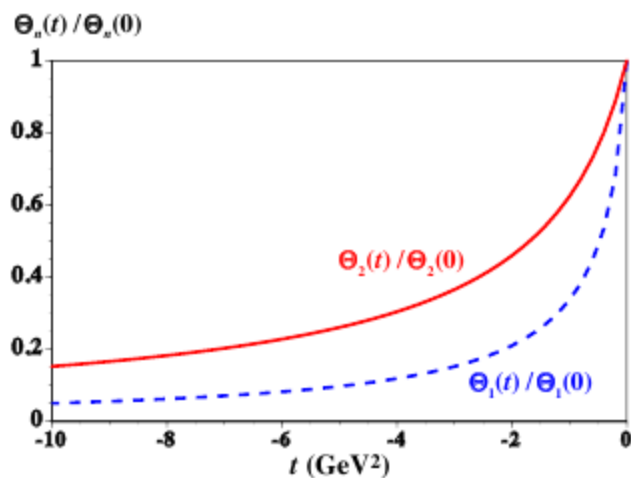
Resonance structure in pressure –related

Θ_1



Time-like \rightarrow space-like

Dispersion relation and Fourier transform



Mass radius

$$\langle r^2 \rangle = 6 \int_{4m^2}^{\infty} \frac{\text{Im}(F(s))}{s^2}$$

$$\sqrt{\langle r^2 \rangle} = 0.69 \text{ fm for } \Theta_2$$

Shear – natural counterpart of pressure

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Forces inside the nucleon on the light front from 3D Breit frame force distributions: Abel tomography case

Julia Yu. Panteleeva¹ and Maxim V. Polyakov^{1,2}

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²*Petersburg Nuclear Physics Institute, Gatchina, 188300 St. Petersburg, Russia*



(Received 4 March 2021; accepted 15 June 2021; published 9 July 2021)

3D <-> 2D relations



Shear viscosity

From spherically symmetric object to fluid (EoS!)

$$T^{\mu\lambda} = (e+p) v^\mu v^\lambda - p g^{\mu\lambda}$$

$V^\mu = P^\mu/M$: correct normalization but no coordinate dependence

Another suggestion (OT'19):

$$V^\mu = (P^\mu + a(t) k_T^\mu) / (M^2 + a^2(t) k_T^2)^{1/2}$$

Viscosity: $\eta dv^\mu/dx_T^\lambda \sim E_\eta p^{[\mu} \Delta^{\lambda]}$

Naïve T-oddness: phases

NO such term in total EMT (but can be for quarks separately)

Phases \leftrightarrow dissipation: polarization in pionic superfluidity model
(V. I. Zakharov, OT' 17)



Viscosity in GDA channel

Viscosity: will correspond to **Exotic $J^{PC}=1^{-+}$** meson
(already studied : Anikin, Pire, Szymanowski, OT,
Wallon'06)

Spin: related to structure of matrix element: One index
of EMT (0^{th} in rest frame) is carried by momentum
and other by polarization vector - just what we need
for **viscosity**

NO for conserved EM: zero coupling for (G)DA!

$\pi\eta$ pairs observation instead of $\pi\pi$ required

Smallness of viscosity: related to smallness of exotic
GDAs and ExEP violation?!

Exotic hybrid meson production

On exotic hybrid meson production in $\gamma^*\gamma$ collisions

I.V. Anikin¹, B. Pire^{2,a}, L. Szymanowski^{3,4,5}, O.V. Teryaev¹, S. Wallon⁵

Eur. Phys. J. C 47, 71–79 (2006)

Digital Object Identifier (DOI) 10.1140/epjc/s2006-02533-7

Possible candidate
 $\pi_1(1400)$

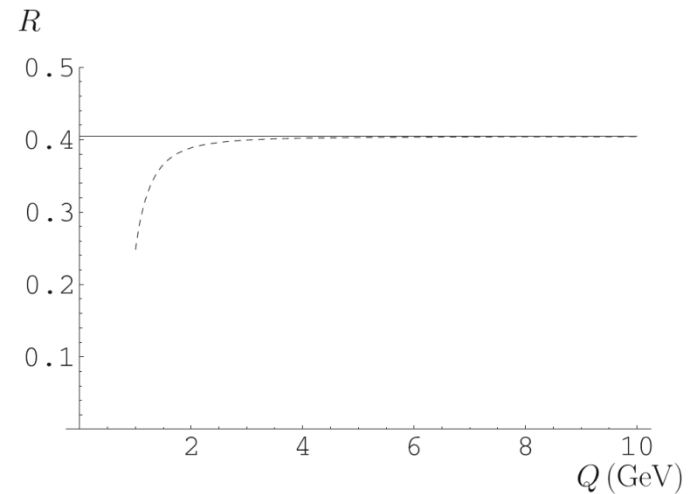


Fig. 2. The ratio $R(Q^2)$ of the squared amplitudes for H and π^0 production in $\gamma^*\gamma$ collisions at leading twist and zero-th order in α_s (solid line) and including twist three contributions in the numerator (dashed line)



Estimate of viscosity

Terms in EMT:

$$(e+p) v^\mu v^\lambda \sim A P^\mu P^\lambda$$

$$\eta dv^\mu/dx_T^\lambda \sim E_\eta p^{[\mu} \Delta^{\lambda]}$$

TD: $e+p \rightarrow Ts$

$$\eta/s (> 1/(4\pi)) \sim E_\eta T / AM$$

Correct dependence on Planck constant
recovered via $\Delta^\lambda \rightarrow -i\hbar d/dx_T^\lambda$

(cf K. Trachenko et al.)



Conclusions/Outlook

- Time-like Gravitational FFs: may be studied in meson pairs production
- Explore the s-t duality and analyticity
- Exotic hybrid mesons: access to shear viscosity and interplay between hadronic and heavy-ion physics
- Holographic bound: related to smallness of exotic GDA and violation of ExEP?