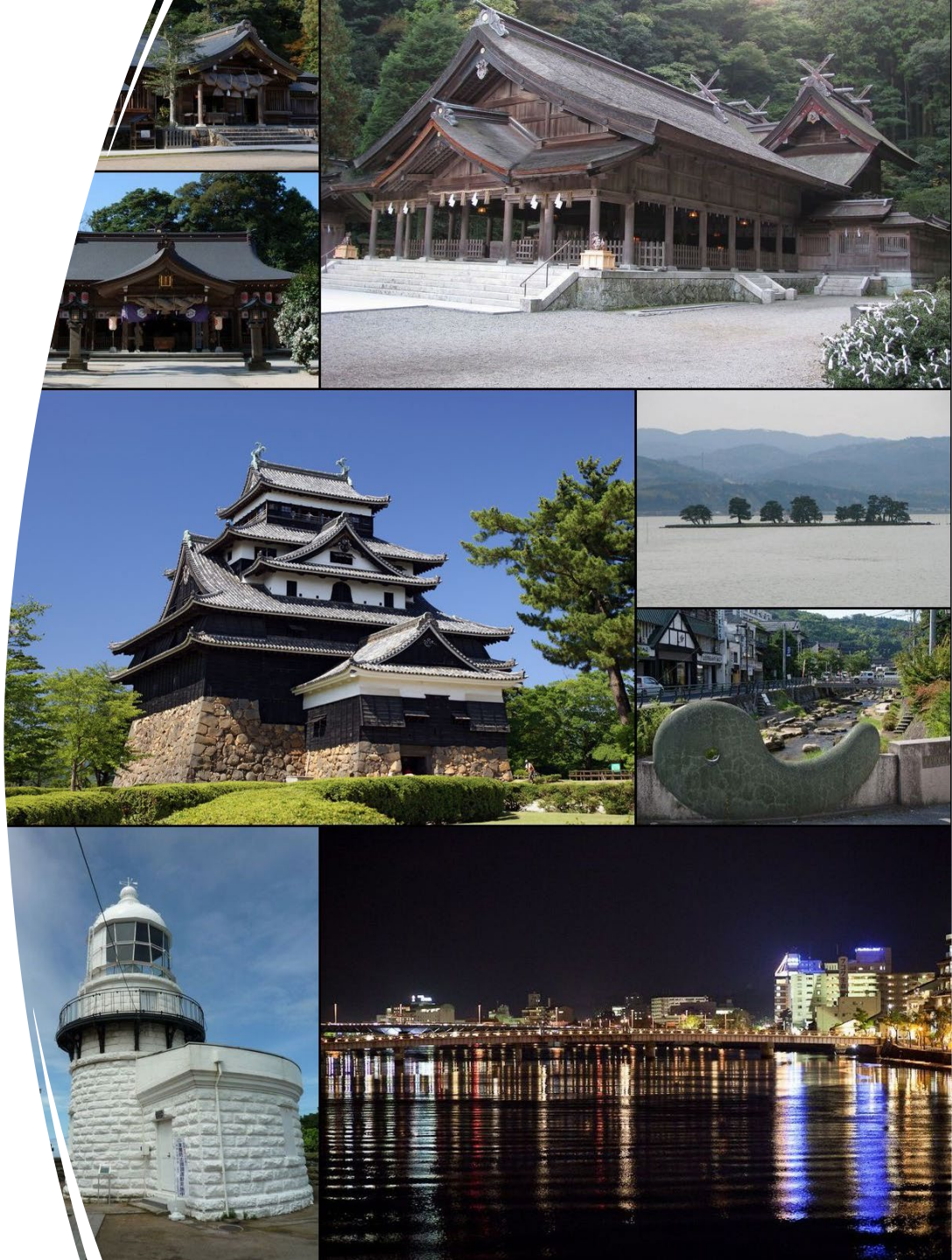


Transverse spin sum rules

Xiangdong Ji

University of Maryland

Spin2021: the 24th
International Spin
Symposium, Oct, 2021



Outline

- Introduction: good and bad news about the longitudinal spin sum rules
- Transverse-spin sum rules
 - Twist-2 case: a unique simple partonic spin sum rule
 - Twist-3 case: a rotation of Jaffe-Manohar sum rule
- Conclusions

Ref: Ji, Zhao, and Yuan, What we know and what we don't know about the proton spin after 30 years

Nature Reviews Physics 3, 65 (2021)

Introduction: good news
news about longitudinal
spin sum rules

Longitudinal spin sum rules

- The first longitudinal spin sum rule in QCD was put forward by Jaffe and Manohar (1990)

$$\frac{1}{2} = \frac{1}{2} \Delta\Sigma + \Delta G + \ell_q^Z + \ell_g^Z$$

- It is a sum rule in the infinite momentum frame $P^Z = \infty$ or partonic sum rule.
- **Good news**: it involves “simplest” experimentally accessible quantities:

$\Delta\Sigma$: total quark helicity

ΔG : total gluon helicity

Progress in measuring ΔG

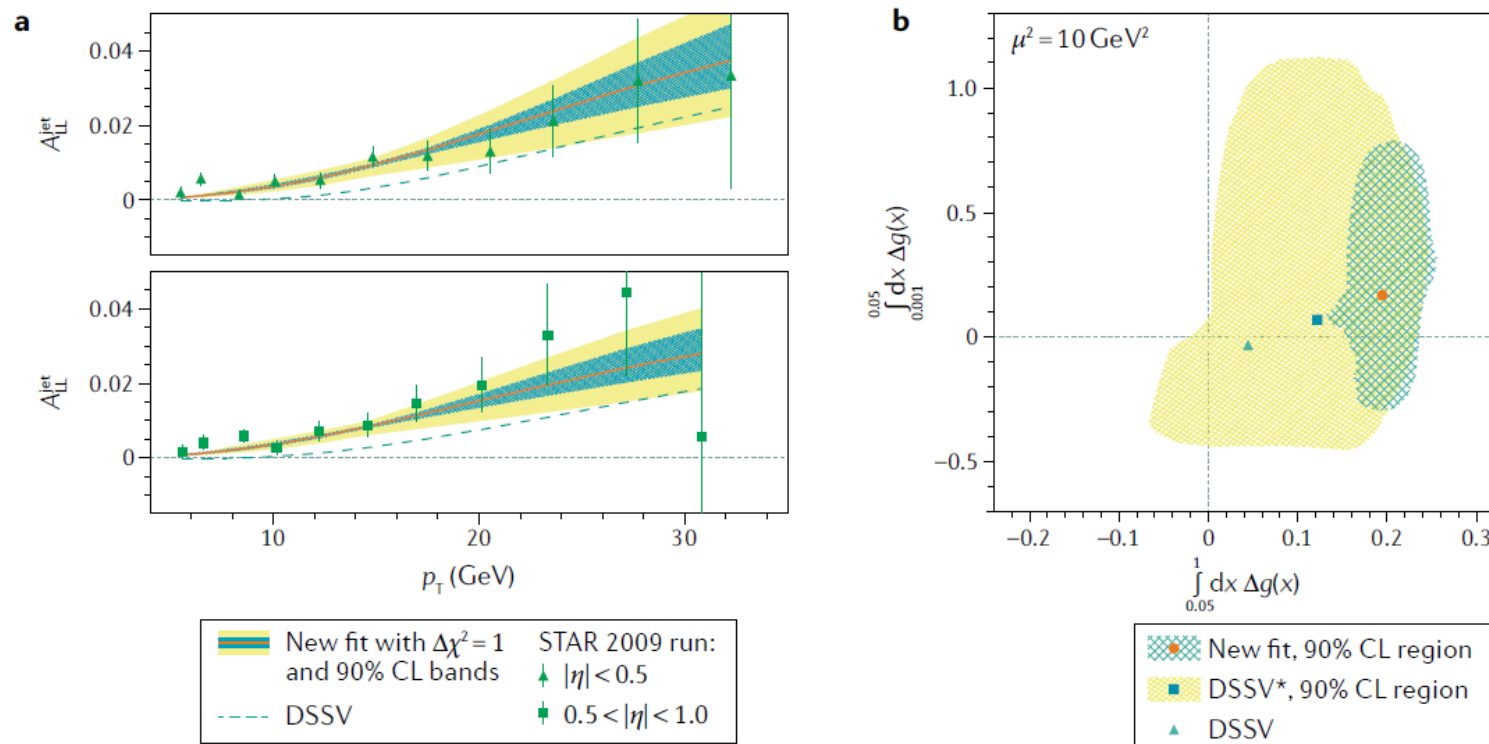
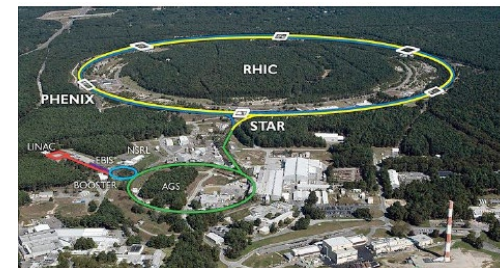


Fig. 2 | The Relativistic Heavy Ion Collider at Brookhaven National Laboratory provides strong evidence for the gluon helicity contribution to the proton spin. **a** | Double spin asymmetry A_{LL}^{jet} in inclusive jet production measured

Progress in understanding ΔG

- There is no gauge-invariant local axial vector corresponding to the gluon spin.
- There is a gauge-dependent candidate for ΔG

$$\vec{S}_g = \vec{E}_a \times \vec{A}_a$$

- In the infinite-momentum limit, the gauge-dependent part of the above operator vanishes

X. Ji, J. Zhang & Y. Zhao Phys.Rev.Lett. 111 (2013)

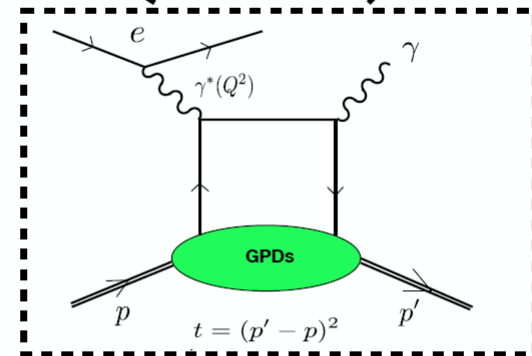
- This provides a recipe for lattice QCD calculation

Bad news about Jaffe-Manohar sum rules

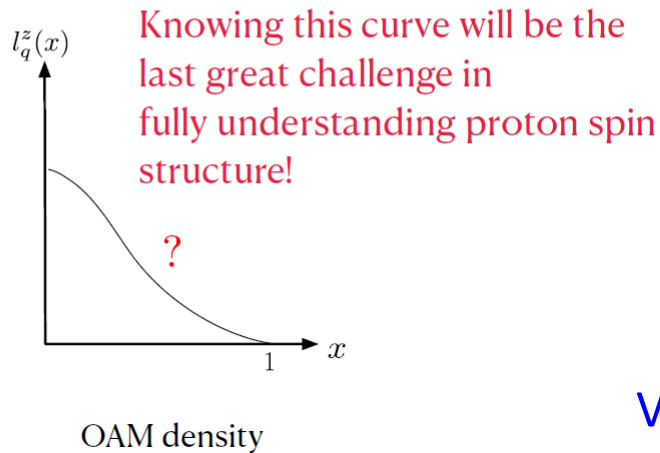
- The helicity of the nucleon, $\frac{1}{2}$, does not grow at the nucleon's momentum gets large. It stays as constant. Thus it is a twist-3 quantity.
- But it so happens that $\Delta\Sigma$ & ΔG are twist-2 quantities,
- OAM l_q and l_g are twist-3!
- The only way we know how to access them experimentally is through twist-three GPDs

Parton OAM & twist-3 GPD

$$l_q^z = \int dx l_q^z(x) \longrightarrow l_q^z(x) = \int dy \left[G_{q,D,3}(x, y) + \mathcal{P} \frac{1}{y-x} G_{q,F,3}(x, y) \right]$$



Twist-3 GPDs extractable from DVCS & DVMP



Very very challenging! It will be a while before we learn how to measure twist-3 GPDs

A frame-independent sum rule

- Frame-independent longitudinal spin sum rule (ji, 1996)

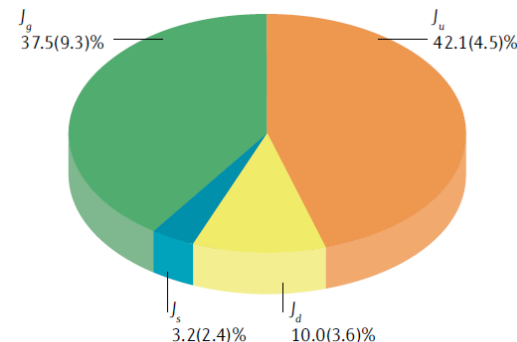
$$\frac{1}{2} = J_q + J_g = \frac{1}{2} \Delta\Sigma + L_q^Z + J_g$$

- J_q & J_g are related to the EMT form factor

$$J_{q,g} = \frac{1}{2} [A_{q,g}(0) + B_{q,g}(0)]$$

- They can be calculated using the standard lattice QCD approach

ETMC collaboration

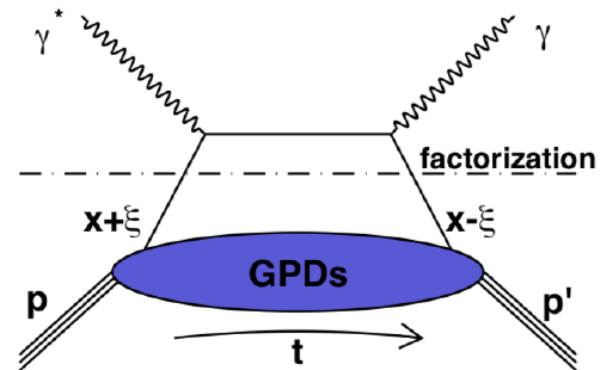


Twist-2 GPD sum rule and DVCS etc

- GPDs were introduced to extract the form factors of EMT through the sum rules

$$\int_{-1}^1 dx \, x H(x, \xi, t) = A(t) + \xi^2 C(t) ,$$
$$\int_{-1}^1 dx \, x E(x, \xi, t) = B(t) - \xi^2 C(t) .$$

- To extract H & E from DVCS and similar processes **model-independently** are somewhat challenging



Problem with frame-independed sum rule

- It does not have simple parton interpretation when considered in IMF.
- Even though H and E come from twist-two operators, but in the longitudinal polarized state, its parton structure is complicated!
- Consider a twist-2 operator

$$\bar{\psi} \gamma^{(\mu_1} i D^{\mu_2} \dots i D^{\mu_n)} \psi \quad (\text{all } \mu \text{ arbitray})$$

Its matrix element has only simple partonic interpretation when all indices +.

Transverse spin sum
rules: twist-2 case

Transverse spin is a twist-2 quantity!

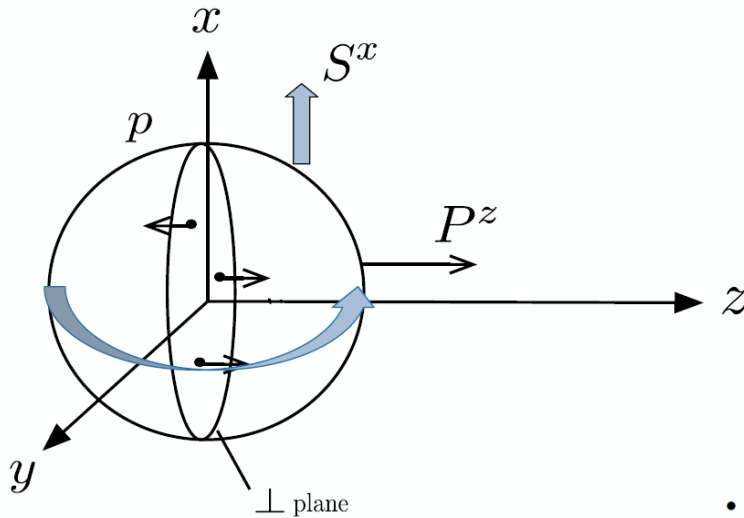
- Transverse AM grows with the external momentum of the hadron
- Consider the second order tensor $J^{\mu\nu} = (\vec{J}, \vec{K})$
- It transforms under Lorentz trans. Like (\vec{B}, \vec{E})

Under Lorentz trans., B_{\perp} increase by factor of γ , thus behave like a twist-2 quantity.

$$\text{transverse spin: } \frac{\hbar}{2} \rightarrow \frac{\gamma\hbar}{2}$$

ji & Yuan *Phys.Lett.B* 810 (2020) 135786

Transversely Polarized Proton:



$$J^x = \sum_i r_i^y p_i^z - r_i^z p_i^y$$

Kinematics:

$$\begin{array}{cc} \sim \gamma & \sim \gamma^{-1} \\ \text{Tw-2} & \text{Tw-4} \end{array}$$

$$\gamma = (1 - \beta^2)^{-1/2}$$

$$[J^{x,y}, K^z] \neq 0$$

- This means that J^x is P^z (frame)-dependent



- Less studied than Longitudinal case because:
 1. It is frame-dependent with non-trivial boost properties
 2. A key issue is separating intrinsic contributions from CM ones
- This has led to ~~some~~ controversy in previous works

a lot of

Frame dependence

Transverse Polarization Sum Rules:

- Let's look again at transverse AM, but **split in terms of its 2 contributions**:

$$J^x = \underbrace{\left(\gamma - \frac{1}{2\gamma}\right) J^{x(2)}}_{\xi^y T^{0z}(\xi)} + \underbrace{\frac{1}{2\gamma} J^{x(3)}}_{\xi^z T^{0y}(\xi)} = \gamma \frac{\hbar}{2}$$

Note: equal in rest frame by rotational symmetry

$J^{x(2)}$ & $J^{x(3)}$ ~ proton matrix elements of q and g fields

- Solution:

$$J^{x(2)} = J^{x(3)} = \frac{\hbar}{2}$$

Twist-2 Ji Sum Rule Novel Twist-3 Sum Rule

Spin sum rules are
 γ -independent!



2 Sum Rules:

$$J_q^{x(2)} + J_g^{x(2)} = \frac{\hbar}{2}$$

Twist-2

$$J_q^{x(3)} + J_g^{x(3)} = \frac{\hbar}{2}$$

Twist-3

Simple parton sum rule for transverse spin

- Spin sum rule as an infinite momentum frame
simple parton sum rule for transverse spin

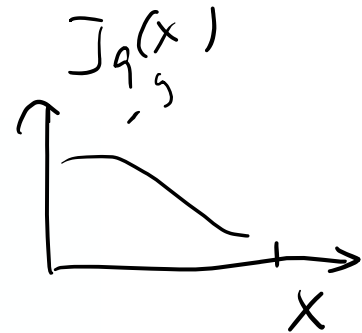
Twist-2 Sum Rule:

$$J_q^{x(2)} + J_g^{x(2)} = \frac{\hbar}{2}$$



$$J_q^{x(2)}(x) = \frac{x}{2} \left(H_q(x) + E_q(x) \right)$$

$$J_g^{x(2)}(x) = \frac{x}{2} \left(H_g(x) + E_g(x) \right)$$



$$J_{q,s} = \int_0^1 J_{q,s}(x) dx$$

X. Ji, PRL 78 (1997) 610

X. Ji, X. Xiong, and F. Yuan, PLB 717, 214 (2012)

- Complete twist-2 partonic sum rule!

Transverse spin sum
rules: twist-3 case

Twist-3 transverse spin sum rule

- Rotated version of the Jaffe Manohar Longitudinal spin sum rule

X. Ji, Y. Guo & K. Shiells, Nuc. Phys. B 969, 115440 (2021)

Transverse Polarization:
$$\frac{1}{2}\Delta q_T + \Delta G_T + l_q^{x(3)} + l_g^{x(3)} = \frac{\hbar}{2}$$

Δq_T , ΔG_T

Involve measurable PDFs in DIS and correspond to spin

$l_q^{x(3)}$, $l_g^{x(3)}$

Involve twist-3 GPDs and correspond to canonical OAM

- Δq_T is related to g_2 structure function
- ΔG_T is a twist-three, spin-dependent gluon distribution
- Twist-3 OAM are a challenge to measure

g_2 structure function

- Can be measured with a transversely polarized target.
- Jlab 6 GeV, 12 GeV
- Moments have been calculated on lattice in the past
- X distribution can be calculated using large momentum effective theory (talks in the previous session)

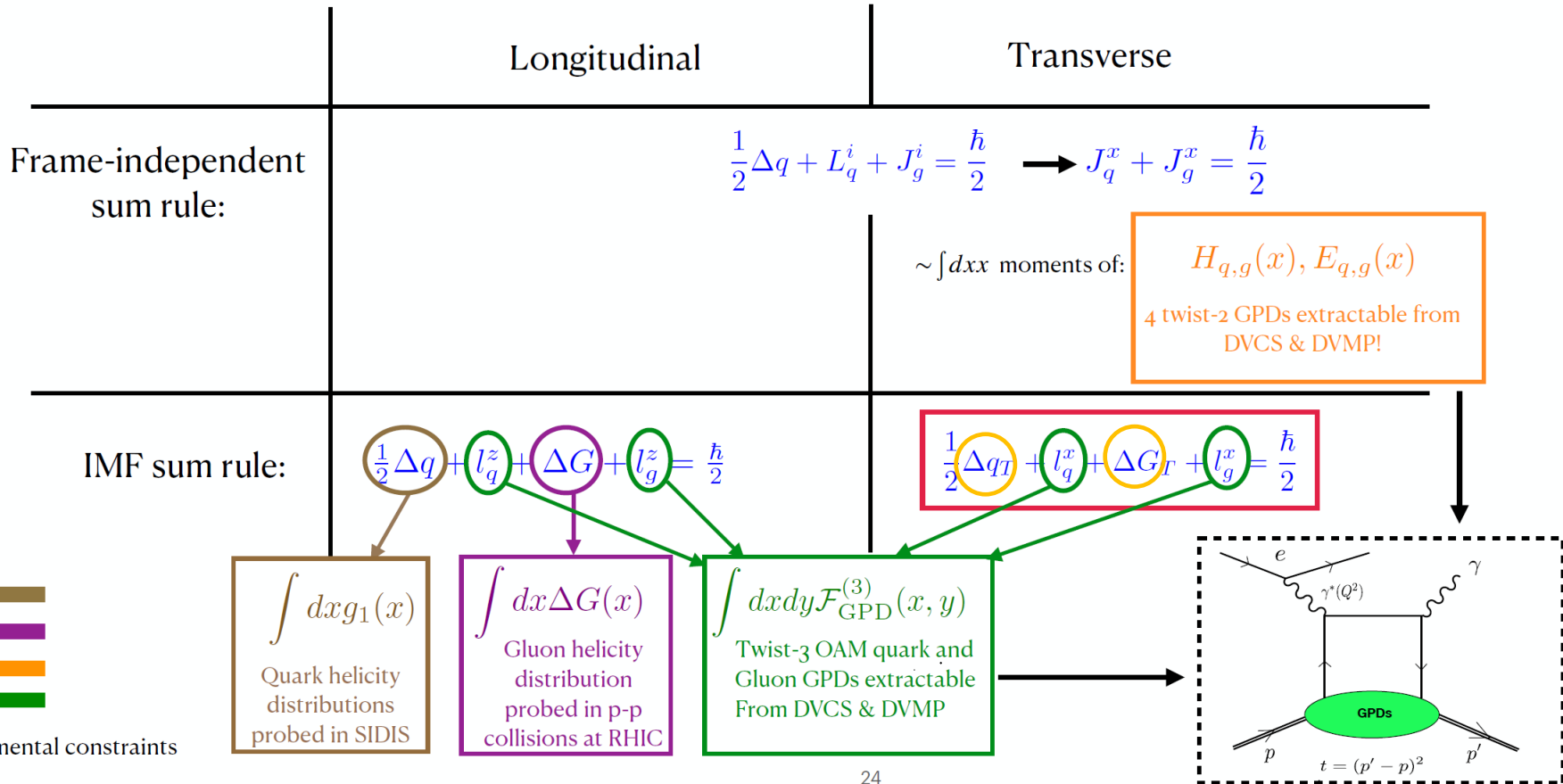
ΔG_T : twist-3 gluon polarization

- When the target is transversely polarized, there is a twist-3 gluon density.
- This mixes with g_2 in factorization formula.
- No experimental data.
- Can be calculated on lattice.

$$\langle PS | G_{(0)}^{+i} \tilde{G}_{(0)}^{\perp i} | P \rangle$$

$\int \mathcal{L}$

Experimental Roadmap for spin sums



less experimental constraints

$g_1, g_T = g_1 + g_2$

Conclusion

- The simplest spin sum rule is twist-2 transverse spin sum rule
- For long. pol., Jaffe-Manohar sum rule is a twist-3 one.
- For long. Pol, the frame independent sum rule does not have a simple parton interpretation,
- There is a frame-rotated version of transverse spin sum rule which involves g_2 and ΔG_T
- All spin sum rules involve GPDs