

# Transverse polarization of hyperons in SIDIS from twist-3 gluon fragmentation functions

↔ SPIN 2021

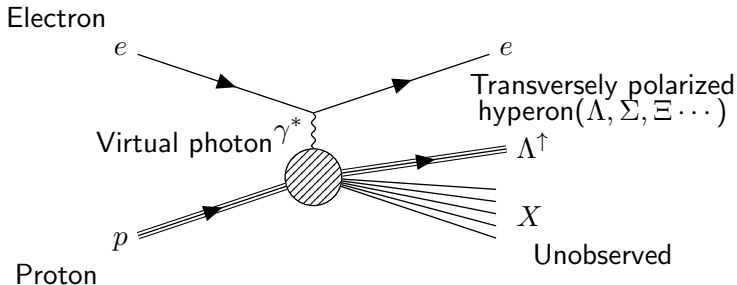
○ R. Ikarashi, Y. Koike<sup>A</sup>,  
K. Yabe, S. Yoshida<sup>B</sup>

- Grad. Sch. Sci. Tech., Niigata Univ.  
Dept. Phys., Niigata Univ.<sup>A</sup>  
South China Normal Univ.<sup>B</sup>

20 Sep 2021

## PURPOSE OF THIS TALK

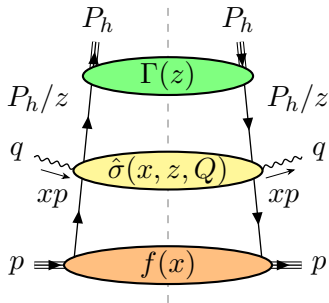
We have calculated the twist-3 gluon FF contribution to  $ep \rightarrow e\Lambda^\uparrow X$ ,  
which is relevant for the future EIC experiment.



# INTRODUCTION

→ QCD factorization for SIDIS

$$p(p) + \gamma^*(q) \rightarrow \Lambda^\uparrow(P_h) + X$$



- **Distribution fn.**  $f(x)$  with momentum fraction  $x$
- **Hard parts**  $\hat{\sigma}(x, z, Q)$  Perturbative parton interaction
- **Fragmentation fn.**  $\Gamma(z)$  with momentum fraction  $z$

Convolution with respect to momentum fractions  $x$  and  $z$ .

$$\text{Cross section (hadronic tensor)} \quad \sigma \sim f(x) \otimes \hat{\sigma}(x, z, Q) \otimes \Gamma(z)$$

# INTRODUCTION

3 types of the twist-3 contributions to  $ep \rightarrow e\Lambda^\uparrow X$ :

**A** twist-3 unpolarized PDF  $\otimes$  Twist-2 transversity FF [COMPLETED]

**B** twist-2 unpolarized PDF  $\otimes$  Twist-3 quark FF [COMPLETED]

formalism by [K. Kanazawa, Y. Koike, PRD88(2013)] for  $ep^\uparrow \rightarrow e\pi X$

**C** twist-2 unpolarized PDF  $\otimes$  Twist-3 gluon FF [THIS WORK]

formalism by [Y. Koike, K. Yabe, S. Yoshida PRD104(2021)]

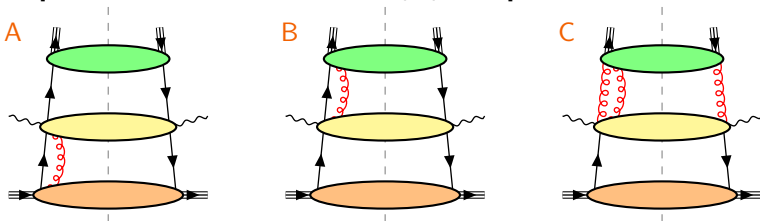
(previous talk by Y. Koike)

---

A,B  $\rightarrow$  [Y. Koike, K. Takada, S. Usui, K. Yabe, S. Yoshida, in preparation]

(talk # 88 by K. Takada)

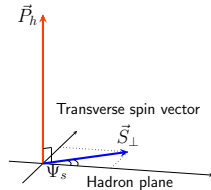
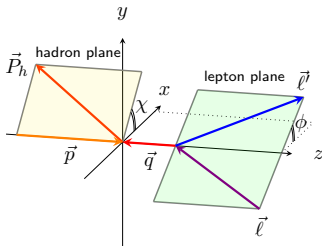
C  $\rightarrow$  [RI, Y. Koike, K. Yabe, S. Yoshida, in preparation]



# KINEMATICS

↪ Hadron frame

$$e(\vec{\ell}) + p(\vec{p}) \rightarrow e'(\vec{\ell}') + \Lambda^\uparrow(\vec{P}_h) + X$$



## 5 Lorentz invariants

1.  $S_{ep} = (p + l)^2 \simeq 2p \cdot l$
2.  $Q^2 = -q^2 = -(\ell - \ell')^2 > 0$
3.  $x_{bj} = Q^2/2p \cdot q$
4.  $z_f = p \cdot P_h/p \cdot q$
5.  $q_T = \sqrt{-q_t^2}, \quad q_t^\mu \equiv q^\mu - \frac{P_h \cdot q}{p \cdot P_h} p^\mu - \frac{p \cdot q}{p \cdot P_h} P_h^\mu$

- The cross section described by 5 Lorentz invariants

$$\frac{d^6\sigma}{dx_{bj}dQ^2dz_fdq_T^2d\phi d\chi}$$

**Leptonic tensor**

**Hadronic tensor**

$$= \frac{\alpha_{em}^2 z_f}{128\pi^4 x_{bj}^2 S_{ep}^2 Q^2}$$

$$L^{\rho\sigma}(\ell, \ell')$$

$$\int \frac{dx}{x} f_1(x) w_{\rho\sigma}(xp, q, P_h)$$

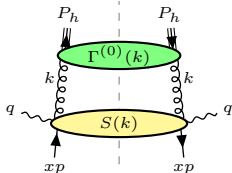
**Twist-2 PDF**

where  $L^{\rho\sigma} = 2(\ell^\rho \ell'^\sigma + \ell^\sigma \ell'^\rho) - Q^2 g^{\mu\nu}$

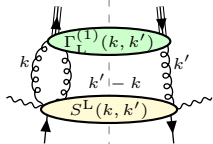
- $w_{\rho\sigma}$  consists of **FFs** and **hard parts**.
- Calculation of  $w_{\rho\sigma}$  is essential.

# HADRONIC TENSOR

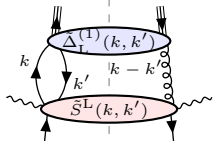
$w_{\rho\sigma} \equiv w_{\rho\sigma}^{(a)} + w_{\rho\sigma}^{(b)} + w_{\rho\sigma}^{(c)} + w_{\rho\sigma}^{(d)} + w_{\rho\sigma}^{(e)}$   
 (c) and (e) are of mirror of (b) and (d), respectively.



$$w^{(a)} = \int \frac{d^4k}{(2\pi)^4} \Gamma_{ab}^{(0)\mu\nu}(k) S_{\mu\nu}^{ab}(k)$$



$$w^{(b)} = \frac{1}{2} \iint \frac{d^4k}{(2\pi)^4} \frac{d^4k'}{(2\pi)^4} \Gamma_{Labc}^{(1)\mu\nu\lambda}(k, k') S_{\mu\nu\lambda}^{Labc}(k, k')$$

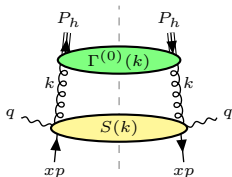


$$w^{(d)} = \text{Tr} \iint \frac{d^4k}{(2\pi)^4} \frac{d^4k'}{(2\pi)^4} \tilde{\Delta}_{La}^{(1)\alpha}(k, k') \tilde{S}_{\alpha}^{La}(k, k')$$

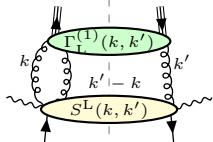
# HADRONIC TENSOR

↪ Definition of Fragmentation matrix elements

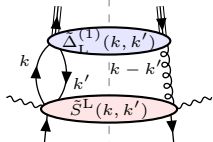
**NOT** color gauge invariant. (mirror diagrams also considered)



$$\Gamma_{ab}^{(0)\mu\nu}(k) \equiv \frac{1}{N} \sum_X \int d^4\xi e^{-ik\cdot\xi} \times \langle 0 | A_b^\nu(0) | \Lambda^\uparrow X \rangle \langle \Lambda^\uparrow X | A_a^\mu(\xi) | 0 \rangle$$



$$\Gamma_{Labc}^{(1)\mu\nu\lambda}(k, k') \equiv \frac{1}{N} \sum_X \iint d^4\xi d^4\eta e^{-ik\cdot\xi} e^{-i(k'-k)\cdot\eta} \times \langle 0 | A_b^\nu(0) | \Lambda^\uparrow X \rangle \langle \Lambda^\uparrow X | A_a^\mu(\xi) g A_c^\lambda(\eta) | 0 \rangle$$



$$\tilde{\Delta}_{La,ij}^{(1)\alpha}(k, k') \equiv \frac{1}{N} \sum_X \iint d^4\xi d^4\eta e^{-ik\cdot\xi} e^{-i(k'-k)\cdot\eta} \times \langle 0 | g A_a^\alpha(\eta) | \Lambda^\uparrow X \rangle \langle \Lambda^\uparrow X | \psi_i(0) \bar{\psi}_j(\xi) | 0 \rangle$$



## HADRONIC TENSOR

- Collinear expansion with respect to  $k$  and  $k'$  around  $P_h$ :

$$k^\alpha = P_h^\alpha/z + \Omega^\alpha_\beta k^\beta \quad k'^\alpha = P_h^\alpha/z' + \Omega^\alpha_\beta k'^\beta$$

$$\text{where } \Omega^\alpha_\beta = g^\alpha_\beta - P_h^\alpha w_\beta$$

- Ward-Takahashi identity for the hard parts:

$$k^\mu S_{\mu\nu}^{ab}(k) = 0,$$

$$k^\mu S_{\mu\nu\lambda}^{Labc}(k, k') = \frac{if^{abc}}{N^2 - 1} \delta^{a'b'} S_{\lambda\nu}^{a'b'}(k')$$

$$k'^\nu S_{\mu\nu\lambda}^{Labc}(k, k') = 0$$

$$(k' - k)^\lambda S_{\mu\nu\lambda}^{Labc}(k, k') = \frac{-if^{abc}}{N^2 - 1} \delta^{a'b'} S_{\mu\nu}^{a'b'}(k')$$

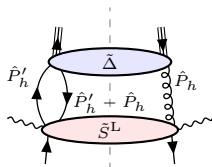
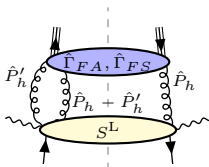
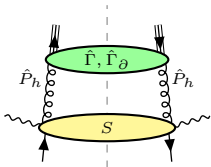
$$(k - k')^\alpha \tilde{S}_\alpha^{La}(k, k') = 0$$

After very lengthy calculation .....

# HADRONIC TENSOR

↪ Color gauge invariant FFs

$$\begin{aligned}
 w_{\rho\sigma} = & \Omega^\alpha_\mu \Omega^\beta_\nu \int d\left(\frac{1}{z}\right) z^2 \hat{\Gamma}^{\mu\nu}(z) S_{\alpha\beta,\rho\sigma}(1/z) \\
 & - i\Omega^\alpha_\mu \Omega^\beta_\nu \Omega^\gamma_\lambda \int d\left(\frac{1}{z}\right) z^2 \hat{\Gamma}_\partial^{\mu\nu\lambda}(z) \left. \frac{\partial S_{\alpha\beta,\rho\sigma}(k)}{\partial k^\gamma} \right|_{k \rightarrow P_h/z (\equiv \hat{P}_h)} \\
 & + \Re \left[ i\Omega^\alpha_\mu \Omega^\beta_\nu \Omega^\gamma_\lambda \iint d\left(\frac{1}{z}\right) d\left(\frac{1}{z'}\right) z z' \frac{1}{1/z - 1/z'} \right. \\
 & \times \left( -\frac{if^{abc}}{N} \hat{\Gamma}_{FA}^{\mu\nu\lambda}\left(\frac{1}{z'}, \frac{1}{z}\right) + \frac{Nd^{abc}}{N^2 - 4} \hat{\Gamma}_{FS}^{\mu\nu\lambda}\left(\frac{1}{z'}, \frac{1}{z}\right) \right) S_{\alpha\beta\gamma,\rho\sigma}^{Labc}\left(\frac{1}{z'}, \frac{1}{z}\right) \left. \right] \\
 & - \Im \Omega^\alpha_\mu \iint d\left(\frac{1}{z}\right) d\left(\frac{1}{z'}\right) z \text{Tr} \left[ \tilde{\Delta}^\mu\left(\frac{1}{z'}, \frac{1}{z'} - \frac{1}{z}\right) \tilde{S}_{\alpha,\rho\sigma}^L\left(\frac{1}{z'}, \frac{1}{z'} - \frac{1}{z}\right) \right]
 \end{aligned}$$



## ■ Intrinsic FF

$$\begin{aligned} \hat{\Gamma}^{\alpha\beta}(z) &= \frac{1}{N^2 - 1} \sum_X \int \frac{d\lambda}{2\pi} e^{-i\lambda/z} \langle 0 | ([\infty w, 0] F^{w\beta}(0))_a | \Lambda^\uparrow X \rangle \langle \Lambda^\uparrow X | (F^{w\alpha}(\lambda w) [\lambda w, \infty w])_a | 0 \rangle \\ &= M_h \epsilon^{P_h w S_\perp \{\alpha w\beta\}} \underbrace{\Delta \hat{G}_{3T}(z)}_{\text{wavy}} + \dots \end{aligned}$$

## ■ Kinematical FF

$$\begin{aligned} \hat{\Gamma}_\theta^{\alpha\beta\gamma}(z) &= \frac{1}{N^2 - 1} \sum_X \int \frac{d\lambda}{2\pi} e^{-i\lambda/z} \langle 0 | ([\infty w, 0] F^{w\beta}(0))_a | \Lambda^\uparrow X \rangle \langle \Lambda^\uparrow X | (F^{w\alpha}(\lambda w) [\lambda w, \infty w])_a | 0 \rangle \overleftarrow{\partial}^\gamma \\ &= -i \frac{M_h}{2} g_\perp^{\alpha\beta} \epsilon^{P_h w S_\perp \gamma} \underbrace{\hat{G}_T^{(1)}(z)}_{\text{wavy}} - i \frac{M_h}{8} \left( \epsilon^{P_h w S_\perp \{\alpha g_\perp^\beta\} \gamma} + \epsilon^{P_h w \gamma \{\alpha S_\perp^\beta\}} \right) \underbrace{\Delta \hat{H}_T^{(1)}(z)}_{\text{wavy}} + \dots \end{aligned}$$

where  $[\lambda w, \infty w]$

gauge link in the adjoint representation connecting  $\lambda w$  and  $\infty w$ .

- NOT gluon field  $A_a^\mu$ , BUT strength tensor  $F_a^{\mu\nu}$ .
- $w_{\rho\sigma}$  is color gauge invariant in  $\mathcal{O}(g)$  accuracy.

## ■ Dynamical FF

$$\begin{aligned}
 & \hat{\Gamma}_{FA}^{\alpha\beta\gamma} \left( \frac{1}{z}, \frac{1}{z'} \right) \\
 &= \frac{-if_{abc}}{N^2 - 1} \sum_X \iint \frac{d\lambda}{2\pi} \frac{d\mu}{2\pi} e^{-i\lambda/z} e^{-i\mu(1/z' - 1/z)} \langle 0 | F_b^{w\beta}(0) | \Lambda^\uparrow X \rangle \langle \Lambda^\uparrow X | F_a^{w\alpha}(\lambda w) g F_c^{w\gamma}(\mu w) | 0 \rangle \\
 &= -M_h \left( \underbrace{\hat{N}_1 \left( \frac{1}{z}, \frac{1}{z'} \right) g_{\perp}^{\alpha\gamma} \epsilon^{PhwS\perp\beta}}_{\text{wavy}} + \underbrace{\hat{N}_2 \left( \frac{1}{z}, \frac{1}{z'} \right) g_{\perp}^{\beta\gamma} \epsilon^{PhwS\perp\alpha}}_{\text{wavy}} - \underbrace{\hat{N}_2 \left( \frac{1}{z'} - \frac{1}{z}, \frac{1}{z'} \right) g_{\perp}^{\alpha\beta} \epsilon^{PhwS\perp\gamma}}_{\text{wavy}} \right)
 \end{aligned}$$

$$\begin{aligned}
 & \hat{\Gamma}_{FS}^{\alpha\beta\gamma} \left( \frac{1}{z}, \frac{1}{z'} \right) \\
 &= \frac{d_{abc}}{N^2 - 1} \sum_X \iint \frac{d\lambda}{2\pi} \frac{d\mu}{2\pi} e^{-i\lambda/z} e^{-i\mu(1/z' - 1/z)} \langle 0 | F_b^{w\beta}(0) | \Lambda^\uparrow X \rangle \langle \Lambda^\uparrow X | F_a^{w\alpha}(\lambda w) F_c^{w\gamma}(\mu w) | 0 \rangle \\
 &= -M_h \left( \underbrace{\hat{O}_1 \left( \frac{1}{z}, \frac{1}{z'} \right) g_{\perp}^{\alpha\gamma} \epsilon^{PhwS\perp\beta}}_{\text{wavy}} + \underbrace{\hat{O}_2 \left( \frac{1}{z}, \frac{1}{z'} \right) g_{\perp}^{\beta\gamma} \epsilon^{PhwS\perp\alpha}}_{\text{wavy}} + \underbrace{\hat{O}_2 \left( \frac{1}{z'} - \frac{1}{z}, \frac{1}{z'} \right) g_{\perp}^{\alpha\beta} \epsilon^{PhwS\perp\gamma}}_{\text{wavy}} \right)
 \end{aligned}$$

$$\begin{aligned}
 \tilde{\Delta}^{\alpha} \left( \frac{1}{z}, \frac{1}{z'} \right) &= \frac{1}{N} \sum_X \iint \frac{d\lambda}{2\pi} \frac{d\mu}{2\pi} e^{-i\lambda/z} e^{-i\mu(1/z' - 1/z)} \langle 0 | F_a^{w\alpha}(\mu w) | \Lambda^\uparrow X \rangle \langle \Lambda^\uparrow X | \bar{\psi}_j(\lambda w) t^a \psi_i(0) | 0 \rangle \\
 &= M_h \left( \underbrace{\epsilon^{\alpha PhwS\perp} (\mathcal{P}_h)_{ij} \tilde{D}_{FT} \left( \frac{1}{z}, \frac{1}{z'} \right)}_{\text{wavy}} + i S_{\perp}^{\alpha} (\gamma_5 \mathcal{P}_h)_{ij} \tilde{G}_{FT} \left( \frac{1}{z}, \frac{1}{z'} \right) \right)
 \end{aligned}$$

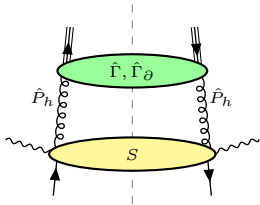
- Twist-3 gluon FFs contributing to  $ep \rightarrow e\Lambda^\dagger X$

$$\left\{ \Delta \hat{G}_{3T}(z), \hat{G}_T^{(1)}(z), \Delta \hat{H}_T^{(1)}(z), \mathfrak{S} \hat{N}_1 \left( \frac{1}{z}, \frac{1}{z'} \right), \mathfrak{S} \hat{N}_2 \left( \frac{1}{z}, \frac{1}{z'} \right), \mathfrak{S} \hat{O}_1 \left( \frac{1}{z}, \frac{1}{z'} \right), \mathfrak{S} \hat{O}_2 \left( \frac{1}{z}, \frac{1}{z'} \right), \mathfrak{S} \tilde{D}_{FT} \left( \frac{1}{z}, \frac{1}{z'} \right), \mathfrak{S} \tilde{G}_{FT} \left( \frac{1}{z}, \frac{1}{z'} \right) \right\}$$

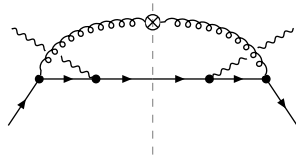
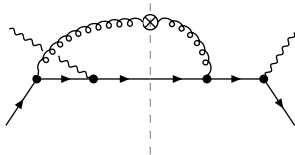
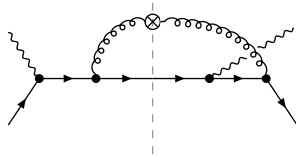
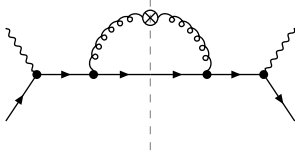
- Exact relations among FFs from QCD equation of motion and Lorentz invariance  
[Y. Koike, K. Yabe and S. Yoshida, PRD 101 (2020)]

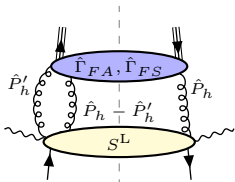
$$\left\{ \cancel{\Delta \hat{G}_{3T}(z)}, \hat{G}_T^{(1)}(z), \Delta \hat{H}_T^{(1)}(z), \mathfrak{S} \hat{N}_1 \left( \frac{1}{z}, \frac{1}{z'} \right), \mathfrak{S} \hat{N}_2 \left( \frac{1}{z}, \frac{1}{z'} \right), \mathfrak{S} \hat{O}_1 \left( \frac{1}{z}, \frac{1}{z'} \right), \mathfrak{S} \hat{O}_2 \left( \frac{1}{z}, \frac{1}{z'} \right), \mathfrak{S} \tilde{D}_{FT} \left( \frac{1}{z}, \frac{1}{z'} \right), \mathfrak{S} \tilde{G}_{FT} \left( \frac{1}{z}, \frac{1}{z'} \right) \right\}$$

# HARD PARTS

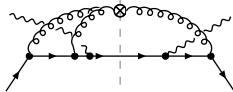
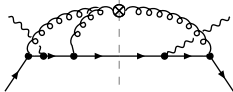
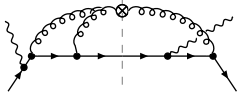
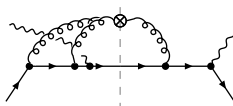
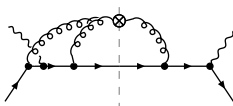
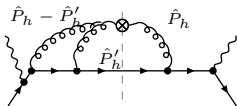
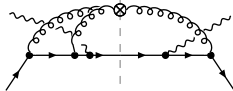
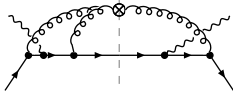
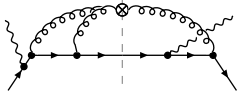
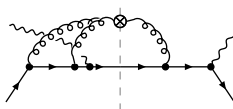
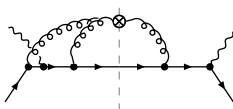
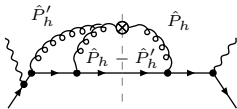


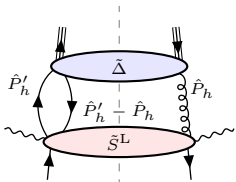
Regarding  $\hat{\Gamma}^{\mu\nu}(z)$  and  $\hat{\Gamma}_\partial^{\mu\nu\lambda}(z)$



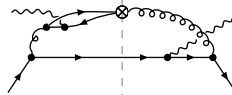
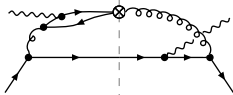
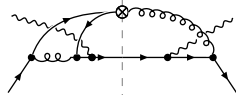
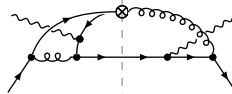
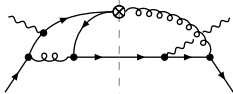
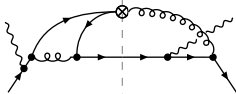
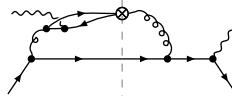
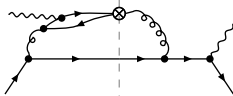
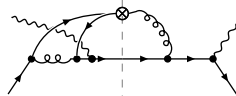
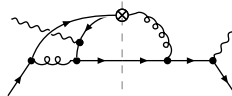
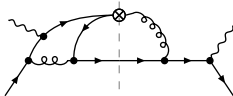
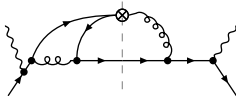


Regarding  $\hat{\Gamma}_{FA}^{\mu\nu\lambda} \left( \frac{1}{z'}, \frac{1}{z} \right)$  and  $\hat{\Gamma}_{FS}^{\mu\nu\lambda} \left( \frac{1}{z'}, \frac{1}{z} \right)$





Regarding  $\tilde{\Delta}^\mu \left( \frac{1}{z'}, \frac{1}{z'} - \frac{1}{z} \right)$





# CROSS SECTION

$$\begin{aligned}
 & \frac{d^6\sigma}{dx_{bj}dQ^2dz_fdq_T^2d\phi d\chi} \quad \hat{x} \equiv x_{bj}/x, \quad \hat{z} \equiv z_f/z \\
 &= \frac{\alpha_{em}^2 \alpha_s M_h}{16\pi^2 x_{bj}^2 S_{ep}^2 Q^2} \sum_k \mathcal{A}_k(\varphi) \mathcal{S}_k(\Psi_s) \iint dx d\left(\frac{1}{z}\right) \frac{z^3}{x} f_1(x) \delta\left(\frac{q_T^2}{Q^2} - \left(1 - \frac{1}{\hat{x}}\right) \left(1 - \frac{1}{\hat{z}}\right)\right) \\
 & \left\{ \hat{G}_T^{(1)}(z) \hat{\sigma}_G^k + \Delta \hat{H}_T^{(1)}(z) \hat{\sigma}_H^k \right. \\
 & + \int d\left(\frac{1}{z'}\right) \left[ \frac{1}{1/z - 1/z'} \mathfrak{S}\left(\hat{N}_1\left(\frac{1}{z'}, \frac{1}{z}\right) \hat{\sigma}_{N1}^k + \hat{N}_2\left(\frac{1}{z'}, \frac{1}{z}\right) \hat{\sigma}_{N2}^k + \hat{N}_2\left(\frac{1}{z} - \frac{1}{z'}, \frac{1}{z}\right) \hat{\sigma}_{N3}^k\right) \right. \\
 & + \frac{1}{z} \left(\frac{1}{1/z - 1/z'}\right)^2 \mathfrak{S}\left(\hat{N}_1\left(\frac{1}{z'}, \frac{1}{z}\right) \hat{\sigma}_{DN1}^k + \hat{N}_2\left(\frac{1}{z'}, \frac{1}{z}\right) \hat{\sigma}_{DN2}^k + \hat{N}_2\left(\frac{1}{z} - \frac{1}{z'}, \frac{1}{z}\right) \hat{\sigma}_{DN3}^k\right) \\
 & + \{N \rightarrow O\} \\
 & + \int d\left(\frac{1}{z'}\right) \frac{2}{C_F} \left[ \mathfrak{S} \tilde{D}_{FT}\left(\frac{1}{z'}, \frac{1}{z'} - \frac{1}{z}\right) \left( \hat{\sigma}_{DF1'}^k + \frac{1}{z} \frac{1}{1/z - 1/z'} \hat{\sigma}_{DF2}^k + \frac{z'}{z} \hat{\sigma}_{DF3}^k \right. \right. \\
 & \left. \left. + \frac{1}{1 - (1 - q_T^2/Q^2) z_f/z'} \hat{\sigma}_{DF4}^k + \frac{1}{1 - (1 - q_T^2/Q^2) z_f(1/z - 1/z')} \hat{\sigma}_{DF5}^k \right) \right. \\
 & \left. + \{D \rightarrow G\} \right\}.
 \end{aligned}$$

# CROSS SECTION

$$\frac{d^6\sigma}{dx_{bj}dQ^2dz_fdq_T^2d\phi d\chi}$$

$$\sim \sum_k \mathcal{A}_k(\varphi) \mathcal{S}_k(\Psi_s) \boxed{f_1(x)} \otimes$$

**Twist-2 PDF**

**Twist-3 gluon FFs**

$$\left\{ \begin{array}{ll} \hat{G}_T^{(1)}, & \Delta \hat{H}_T^{(1)} \\ \mathfrak{S} \hat{N}_i, & \mathfrak{S} \hat{O}_i \\ \mathfrak{S} \tilde{D}_{FT}, & \mathfrak{S} \tilde{G}_{FT} \end{array} \right\} \otimes$$

$$\boxed{\{\hat{\sigma}^k\}}$$

**Hard parts**

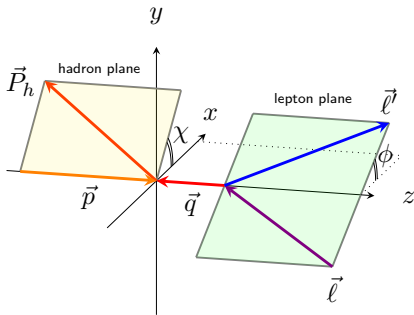
$$\varphi \equiv \phi - \chi$$

$\Psi_s$ : azimuthal angle of  $S_\perp$

$\hat{\sigma}^k$  depend on  $\hat{x}, \hat{z}, Q$  and  $q_T$ ,

where  $\hat{x} \equiv x_{bj}/x, \hat{z} \equiv z_f/z$ .

$\mathcal{A}_k$  and  $\mathcal{S}_k$  are defined  
in the next slide



## CROSS SECTION

### Angle dependencies

$$\begin{aligned}\mathcal{A}_1(\varphi) &= 1 + \cosh^2 \psi, & \mathcal{A}_2(\varphi) &= -2 \\ \mathcal{A}_3(\varphi) &= -\cos \varphi \sinh 2\psi, & \mathcal{A}_4(\varphi) &= \cos 2\varphi \sinh^2 \psi \\ \mathcal{A}_8(\varphi) &= -\sin \varphi \sinh 2\psi, & \mathcal{A}_9(\varphi) &= \sin 2\varphi \sinh^2 \psi\end{aligned}$$

[R. Meng, F. I. Olness and D. E. Soper, NPB371 (1992)]

---

$$\cosh \psi \equiv 2x_{bj} S_{ep} / Q^2 - 1$$

---

$$\mathcal{S}_{1,2,3,4} \equiv \sin \Psi_s, \mathcal{S}_{8,9} \equiv \cos \Psi_s.$$

### ■ 5 structure functions

$$\begin{aligned}& \frac{d^6\sigma}{dx_{bj}dQ^2dz_fdq_T^2d\phi d\chi} \\ &= F_0(\sin \Psi_s) + F_1(\sin \Psi_s \cos \varphi) + F_2(\sin \Psi_s \cos 2\varphi) \\ &+ F_3(\cos \Psi_s \sin \varphi) + F_4(\cos \Psi_s \sin 2\varphi)\end{aligned}$$

## CONCLUSION

- We have derived the cross-section formula ( $\mathcal{O}(\alpha_s)$ ) of  $ep \rightarrow e\Lambda^\uparrow X$  process related to the twist-3 gluon fragmentation functions in the collinear framework.

- A Twist-3 distribution effect
- B Twist-3 quark fragmentation effect )←[talk #88 by K. Takada]
- C Twist-3 gluon fragmentation effect [THIS WORK]

- Twist-3 gluon FF contribution could play an important role, since gluons are ample in the collision environment.
- Satisfying EM and color gauge invariances.
- Classified into 5 structure functions.

### Future outlook

- Numerical estimate of each effect.
- Inclusion of the NLO correction,  $\mathcal{O}(\alpha_s^2)$ .

\*NLO contribution has large correction for other unpolarized cross section [Hinderer,Schlegel, Vogelsang PRD92(2015),erratum[PRD93(2016)]]