


- Cammarota, Gamberg, Kang, Miller, Pitonyak, Prokudin, Rogers, Sato, Phys.Rev.D 102 (2020) 5, 054002
- Gamberg, Kang, Pitonyak, Prokudin, Sato, Seidl, Phys.Lett.B 816 (2021) 136255

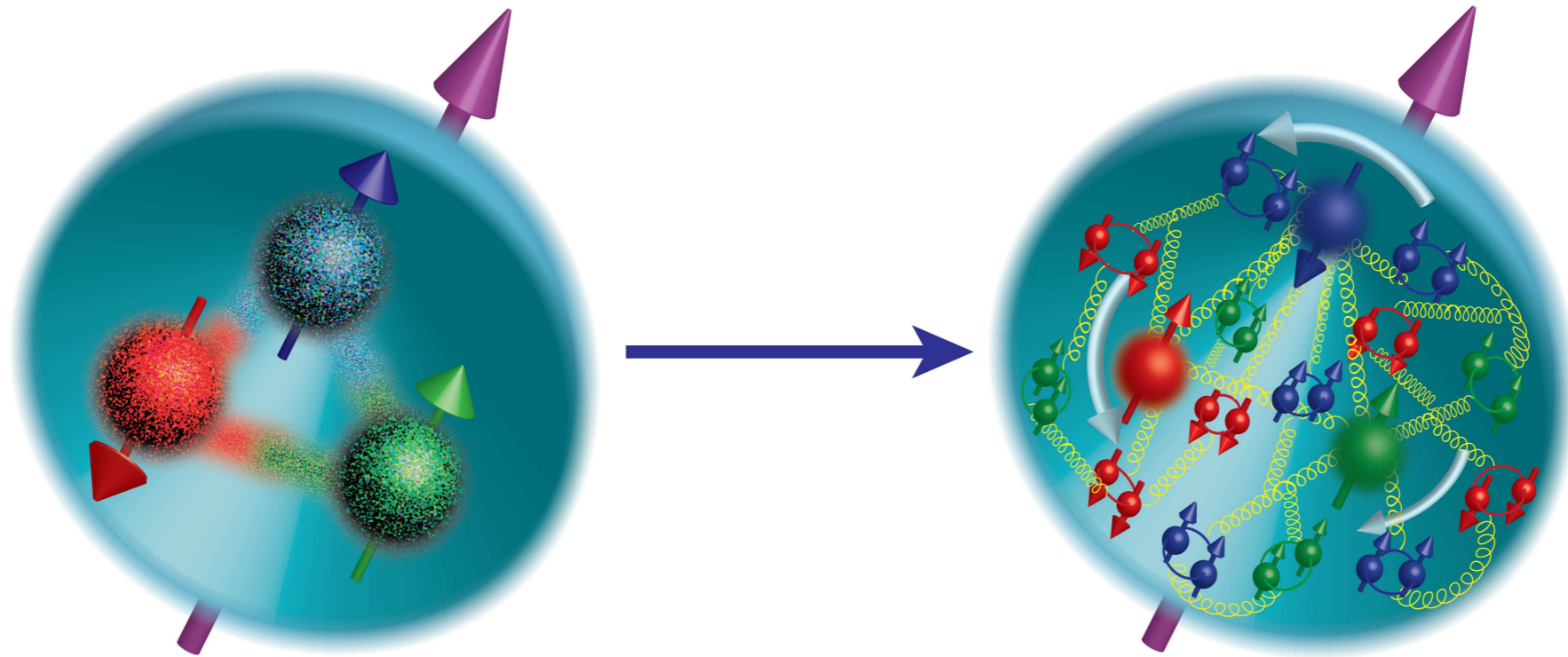


# Global analysis of SSAs and the EIC impact on the tensor charge

Alexei Prokudin

# EVOLUTION OF OUR UNDERSTANDING OF THE SPIN STRUCTURE

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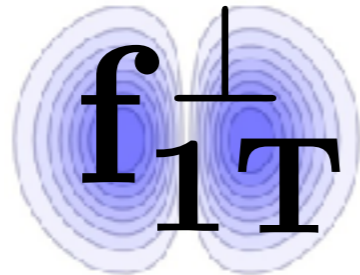


1980' - the spin of the nucleon is due to the valence quarks

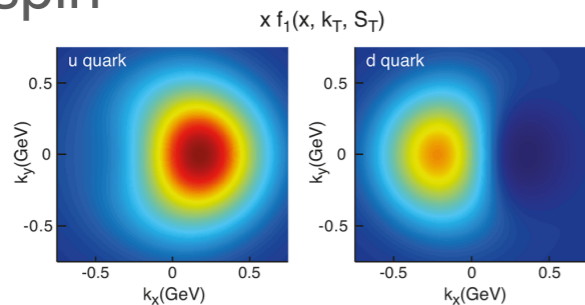
Modern concept: valence quarks, sea quarks, and gluons together with orbital angular momentum are contributing

# POLARIZED TMD FUNCTIONS

## Sivers function

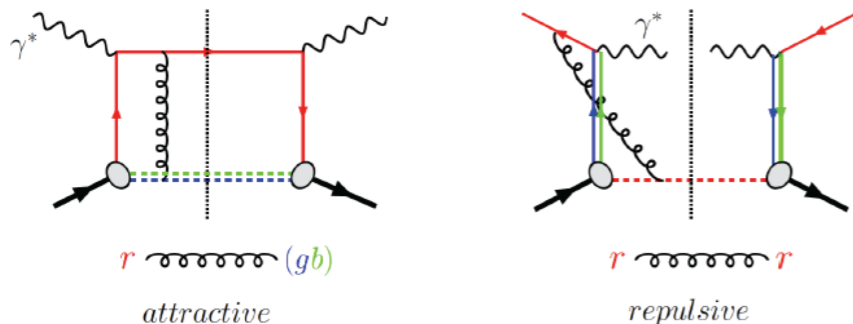


- Describes unpolarized quarks inside of transversely polarized nucleon
- Encodes the correlation of orbital motion with the spin



- Sign change of Sivers function is fundamental consequence of QCD

Brodsky, Hwang, Schmidt (2002), Collins (2002)



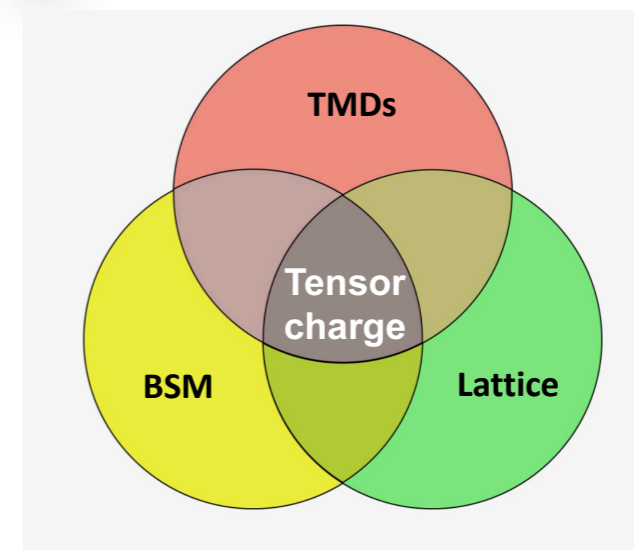
$$f_{1T}^\perp \text{SIDIS} = -f_{1T}^\perp \text{DY}$$

## Transversity



- The only source of information on tensor charge of the nucleon
- Couples to Collins fragmentation function or di-hadron interference fragmentation functions in SIDIS

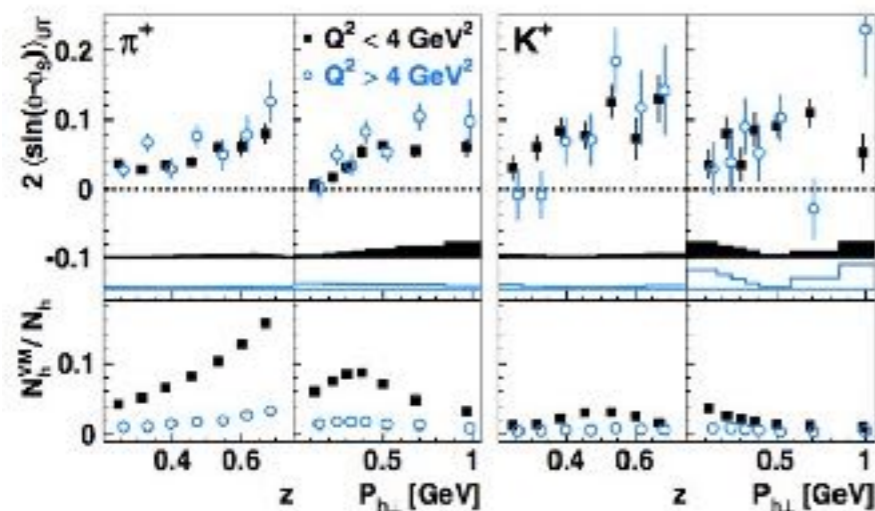
$$\delta q \equiv g_T^q = \int_0^1 dx [h_1^q(x, Q^2) - h_1^{\bar{q}}(x, Q^2)]$$



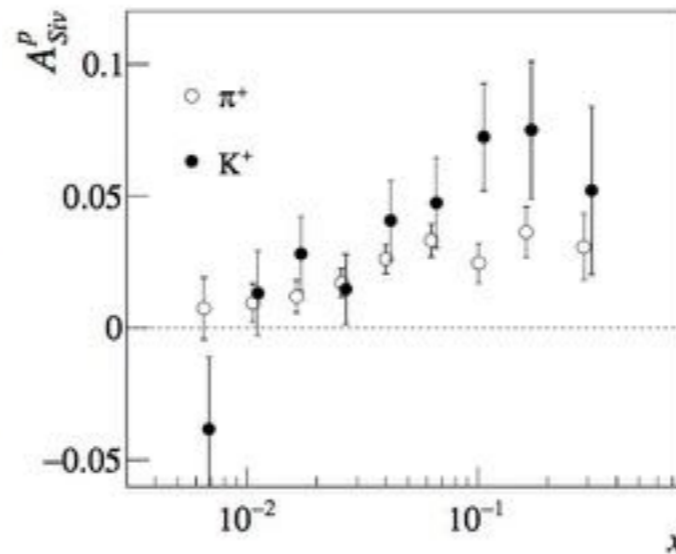
# TRANSVERSE SPIN ASYMMETRIES

Transverse Single Spin Asymmetries (SSAs) have been observed in a variety of processes

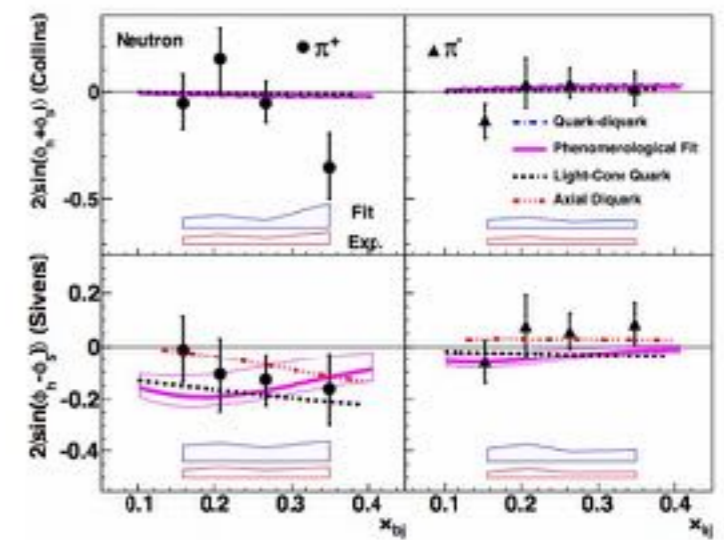
## Sivers asymmetry in SIDIS



HERMES (09)



COMPASS (15)



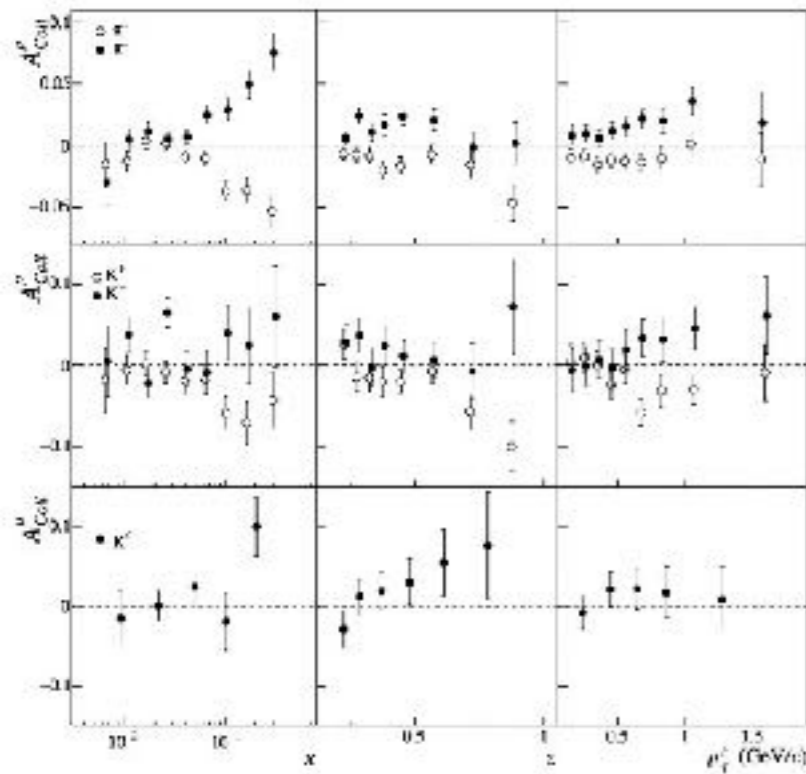
JLAB (11)

$$F_{UT}^{\sin(\phi_h - \phi_S)} = C \left[ -\frac{\hat{h} \cdot \vec{k}_T}{M} f_{1T}^\perp D_1 \right]$$

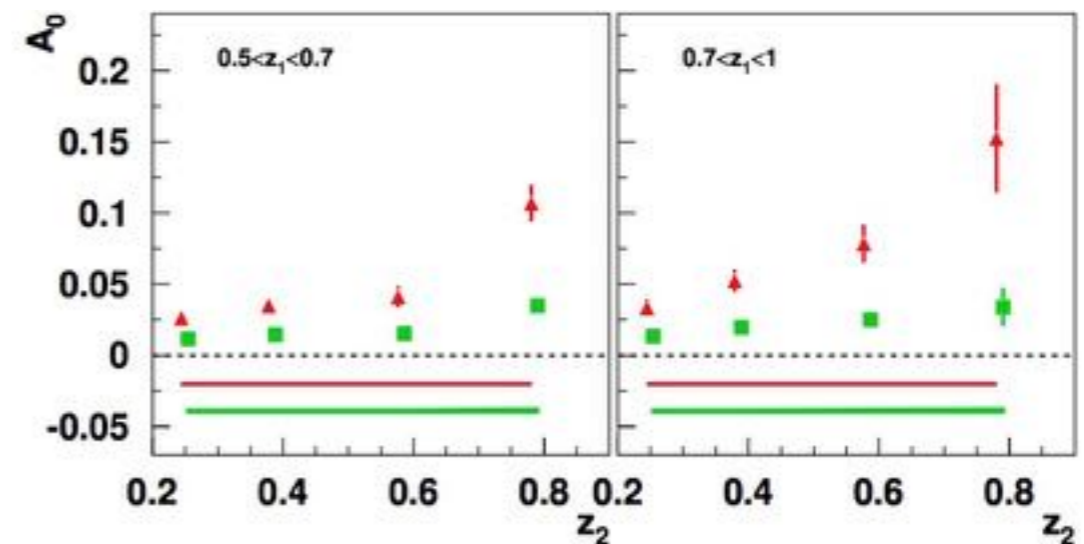
# TRANSVERSE SPIN ASYMMETRIES

Transverse Single Spin Asymmetries (SSAs) have been observed in a variety of processes

Collins asymmetry in SIDIS and  $e^+e^-$



COMPASS (15),  
also HERMES (05,10, 20), JLab (11,14)



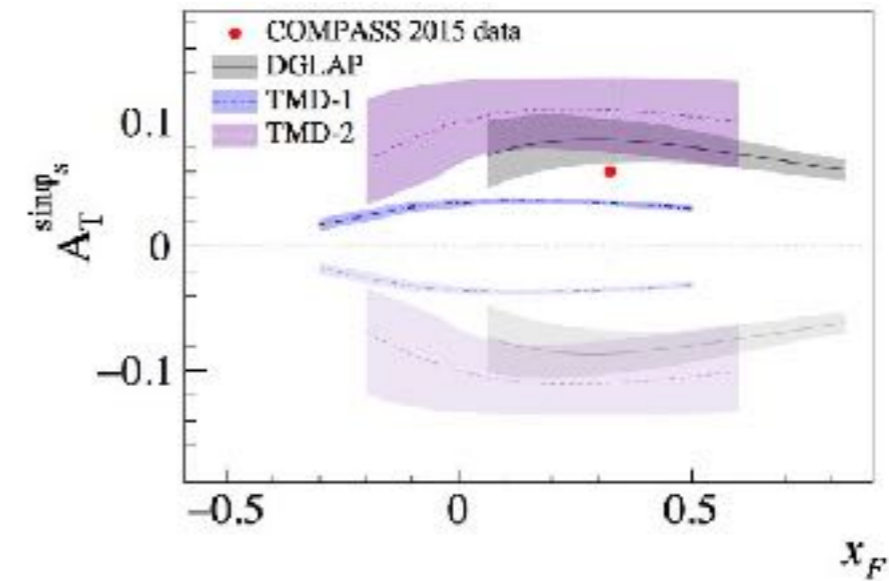
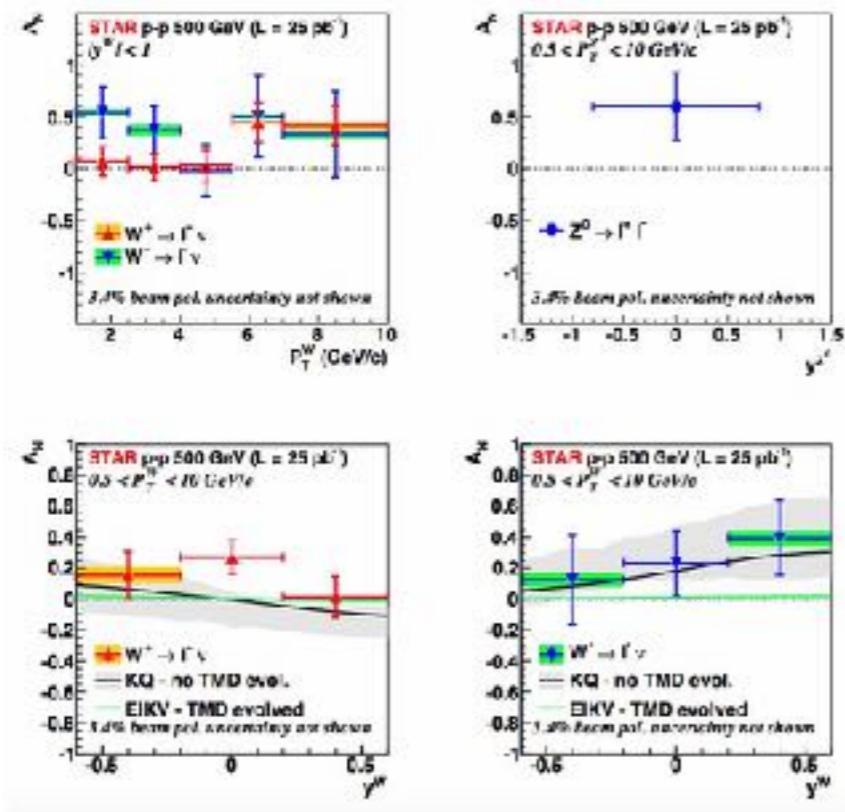
BELLE (08),  
also BaBar (14), BESIII (16)

$$F_{UT}^{\sin(\phi_h + \phi_S)} = C \left[ -\frac{\hat{h} \cdot \vec{p}_\perp}{M_h} h_1 H_1^\perp \right] \quad F_{UU}^{\cos(2\phi_0)} = C \left[ \frac{2\hat{h} \cdot \vec{p}_{a\perp} \hat{h} \cdot \vec{p}_{b\perp} - \vec{p}_{a\perp} \cdot \vec{p}_{b\perp}}{M_a M_b} H_1^\perp \bar{H}_1^\perp \right]$$

# TRANSVERSE SPIN ASYMMETRIES

Transverse Single Spin Asymmetries (SSAs) have been observed in a variety of processes

## Sivers effect in Drell-Yan

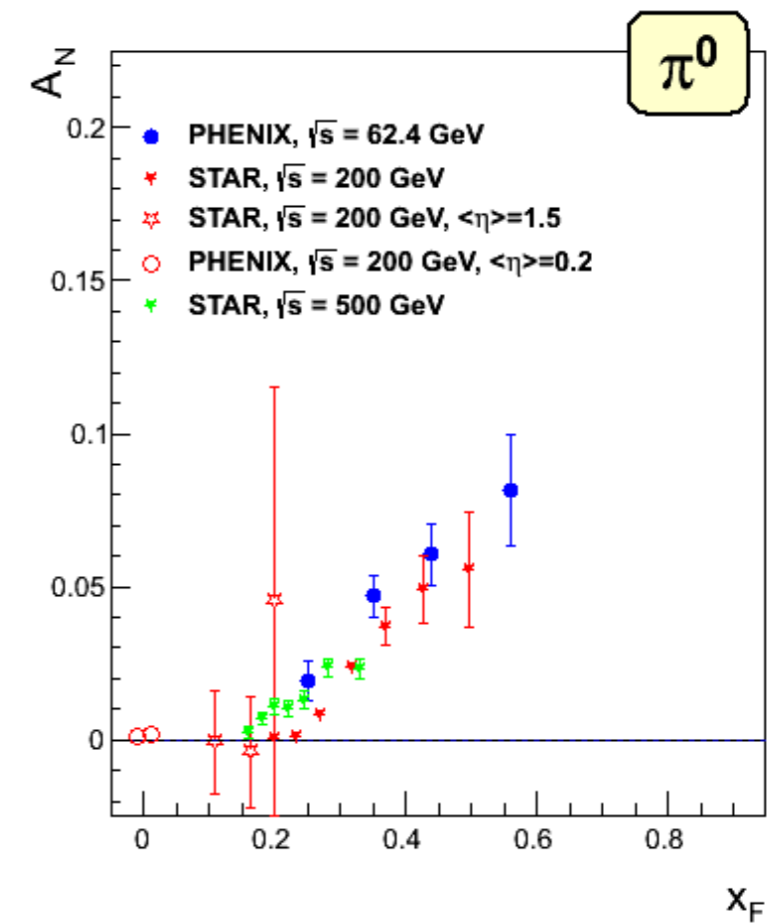
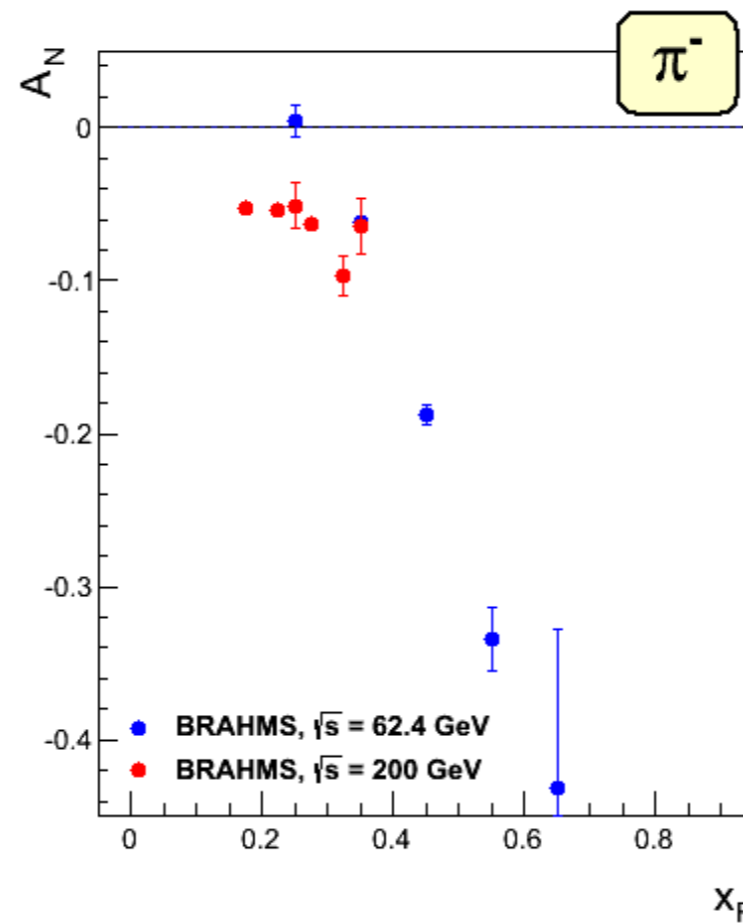
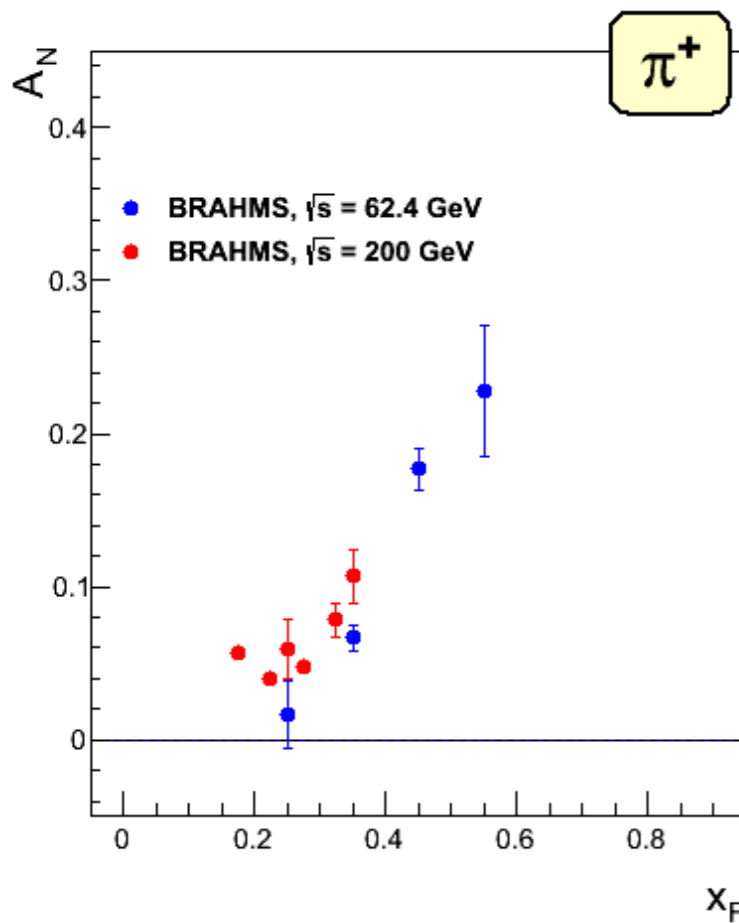
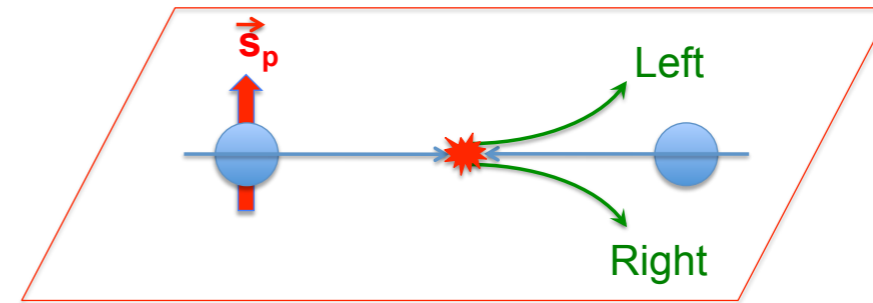


$$F_{TU}^1 = C \left[ -\frac{\vec{h} \cdot \vec{k}_{aT}}{M_a} \mathbf{f}_{1T}^\perp \bar{f}_1 \right]$$

# TRANSVERSE SPIN ASYMMETRIES

$A_N$  in pp scattering

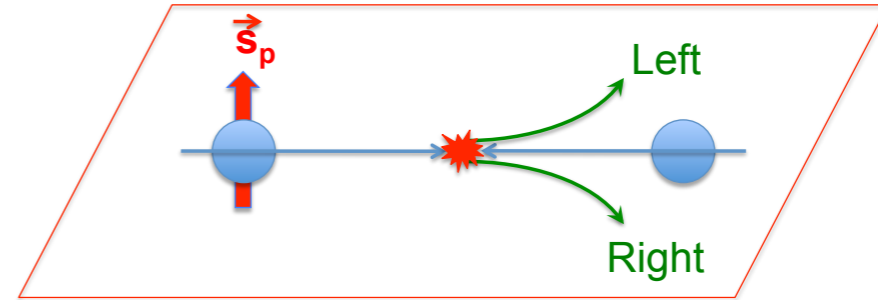
RHIC: STAR, BRAHMS, PHENIX



*“The RHIC SPIN Program: Achievements and Future Opportunities”, Aschenauer et al (15)*

# TRANSVERSE SPIN ASYMMETRIES

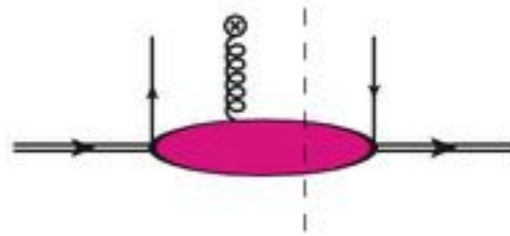
$A_N$  in pp scattering is related to collinear twist-3 (CT3) factorization



$$d\Delta\sigma(S_T) \sim \underbrace{H_{QS} \otimes f_1 \otimes \mathbf{F}_{FT} \otimes D_1}_{\text{Qiu-Sterman term}} + H_F \otimes f_1 \otimes \mathbf{h}_1 \otimes \left( H_1^{\perp(1)}, \tilde{H} \right)$$

Qiu-Sterman term

$\mathbf{F}_{FT} \sim$



quark-gluon-quark correlator

Qiu, Sterman (99), Kouvaris, et al (06)

$$\pi \mathbf{F}_{FT}(\mathbf{x}, \mathbf{x}) = \int d^2 \vec{k}_T \frac{k_T^2}{2M^2} \mathbf{f}_{1T}^{\perp}(\mathbf{x}, k_T^2) \equiv f_{1T}^{\perp(1)}(\mathbf{x}) \quad \text{the first moment of Sivers function}$$

Boer, et al (03)

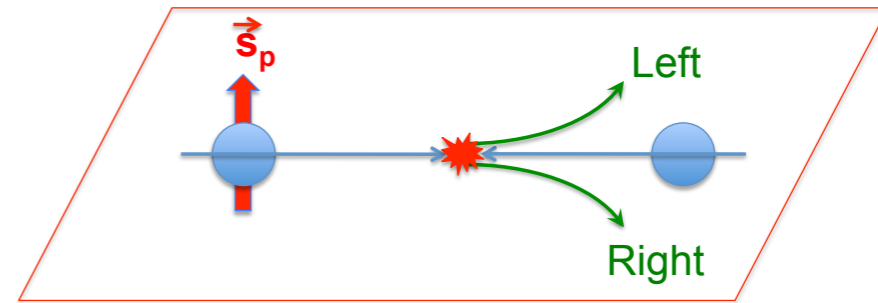
TMD and CT3 factorization agree in their overlapping region of applicability

Ji, et al (06); Koike, et a. (08); Zhou, et al (08, 10); Yuan and Zhou (09)



# TRANSVERSE SPIN ASYMMETRIES

$A_N$  in pp scattering is related to collinear twist-3 (CT3) factorization

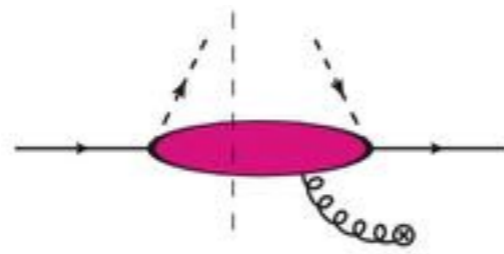


$$d\Delta\sigma(S_T) \sim H_{QS} \otimes f_1 \otimes \mathbf{F}_{FT} \otimes D_1 + H_F \otimes f_1 \otimes \mathbf{h}_1 \otimes \left( \mathbf{H}_1^{\perp(1)}, \tilde{\mathbf{H}} \right)$$

Fragmentation term

$\mathbf{h}_1$  collinear transversity

$\mathbf{H}_1^{\perp(1)}$   $\tilde{\mathbf{H}}$



Kanazawa, Koike, Metz, Pitonyak, Schlegel, (16)

quark-gluon-quark fragmentation functions

$$\mathbf{H}_1^{\perp(1)}(z) \equiv z^2 \int d^2\vec{p}_\perp \frac{p_\perp^2}{2M_h^2} \mathbf{H}_1^\perp(z, z^2 p_\perp^2)$$

the first moment of Collins FF

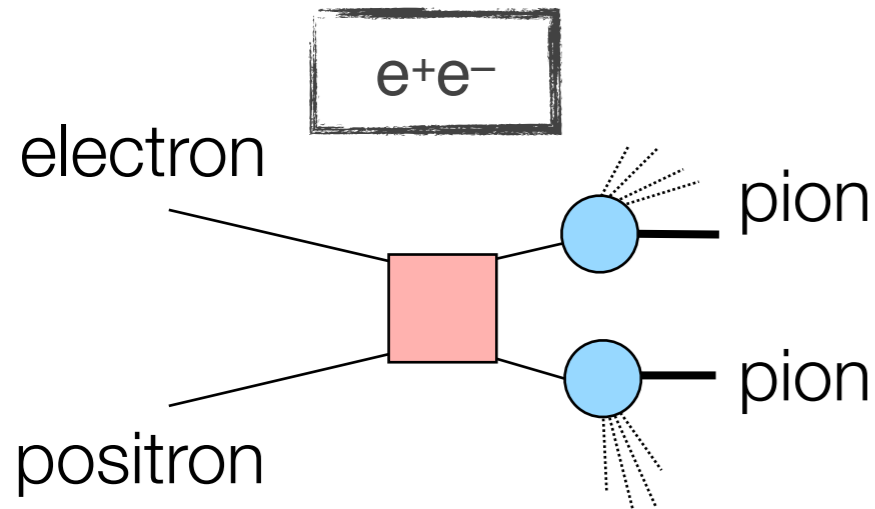
$$F_{UT}^{\sin\phi_S} \sim \sum_a e_a^2 \frac{2M_h}{Q} \mathbf{h}_1^a(x) \frac{\tilde{\mathbf{H}}(z)}{z}$$

Mulders, Tangerman (96); Bacchetta, et al (07)

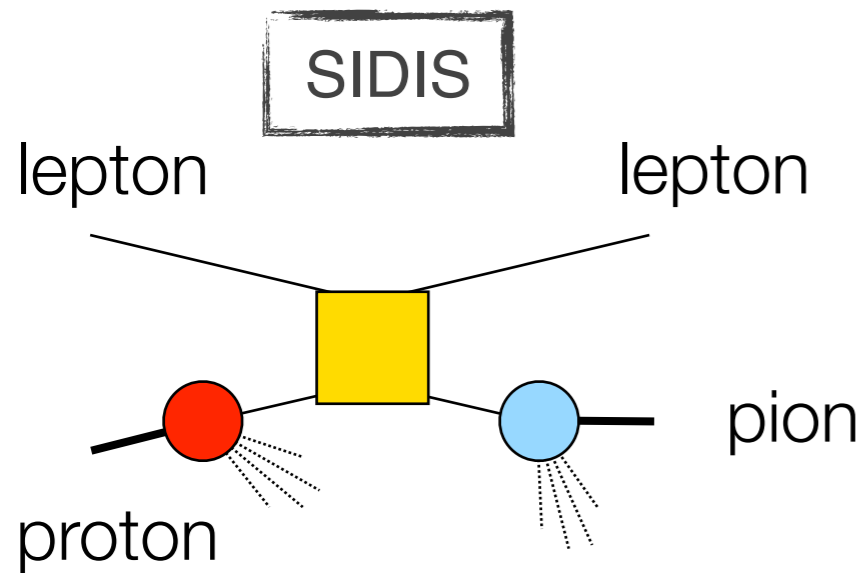
# JAM20 ANALYSIS

# UNIVERSAL GLOBAL FIT 2020

*Cammarota, Gamberg, Kang, Miller, Pitonyak, Prokudin, Rogers, Sato (2020)*

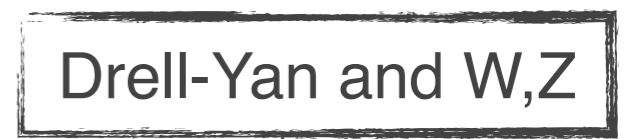


*Collins asymmetries  
BELLE, BaBar, BESIII data*

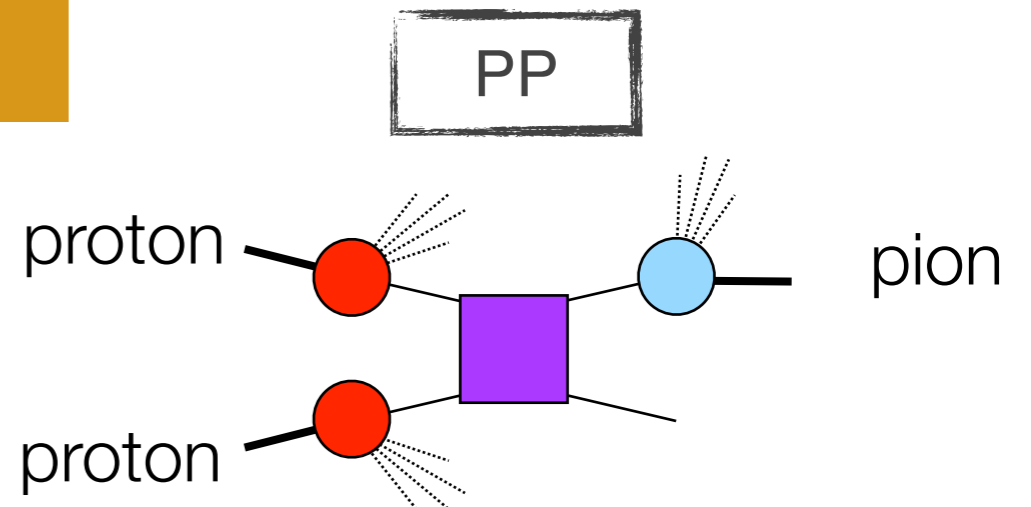


*Sivers, Collins asymmetries  
COMPASS, HERMES, JLab data*

To demonstrate the common origin of SSAs in various processes, we combined all available data and extracted a universal set of non perturbative functions that describes all of them



*Sivers asymmetries  
COMPASS, STAR data*



*$A_N$  asymmetry  
STAR, PHENIX, BRAHMS data*

# UNIVERSAL GLOBAL FIT 2020

Jefferson Lab Angular Momentum Collaboration

<https://www.jlab.org/theory/jam>

| Observable                      | Reactions   | Non-Perturbative Function(s)   | $\chi^2/N_{\text{pts.}}$ |
|---------------------------------|---|--|--------------------------|
| $A_{\text{SIDIS}}^{\text{Siv}}$ | $e + (p, d)^{\uparrow} \rightarrow e + (\pi^+, \pi^-, \pi^0) + X$ | $f_{1T}^{\perp}(x, k_T^2)$   | 150.0/126 = 1.19         |
| $A_{\text{SIDIS}}^{\text{Col}}$ | $e + (p, d)^{\uparrow} \rightarrow e + (\pi^+, \pi^-, \pi^0) + X$ | $h_1(x, k_T^2), H_1^{\perp}(z, z^2 p_{\perp}^2)$                               | 111.3/126 = 0.88         |
| $A_{\text{SIA}}^{\text{Col}}$   | $e^+ + e^- \rightarrow \pi^+ \pi^- (UC, UL) + X$                  | $H_1^{\perp}(z, z^2 p_{\perp}^2)$  | 154.5/176 = 0.88         |
| $A_{\text{DY}}^{\text{Siv}}$    | $\pi^- + p^{\uparrow} \rightarrow \mu^+ \mu^- + X$                | $f_{1T}^{\perp}(x, k_T^2)$   | 5.96/12 = 0.50           |
| $A_{\text{DY}}^{\text{Siv}}$    | $p^{\uparrow} + p \rightarrow (W^+, W^-, Z) + X$                  | $f_{1T}^{\perp}(x, k_T^2)$   | 31.8/17 = 1.87           |
| $A_N^h$                         | $p^{\uparrow} + p \rightarrow (\pi^+, \pi^-, \pi^0) + X$          | $h_1(x), F_{FT}(x, x) = \frac{1}{\pi} f_{1T}^{\perp(1)}(x), H_1^{\perp(1)}(z)$ | 66.5/60 = 1.11           |

*Cammarota, Gamberg, Kang, Miller, Pitonyak, Prokudin, Rogers, Sato (2020)*

► 18 observables and 6 non-perturbative functions (Sivers up/down; transversity up/down; Collins favored/unfavored)

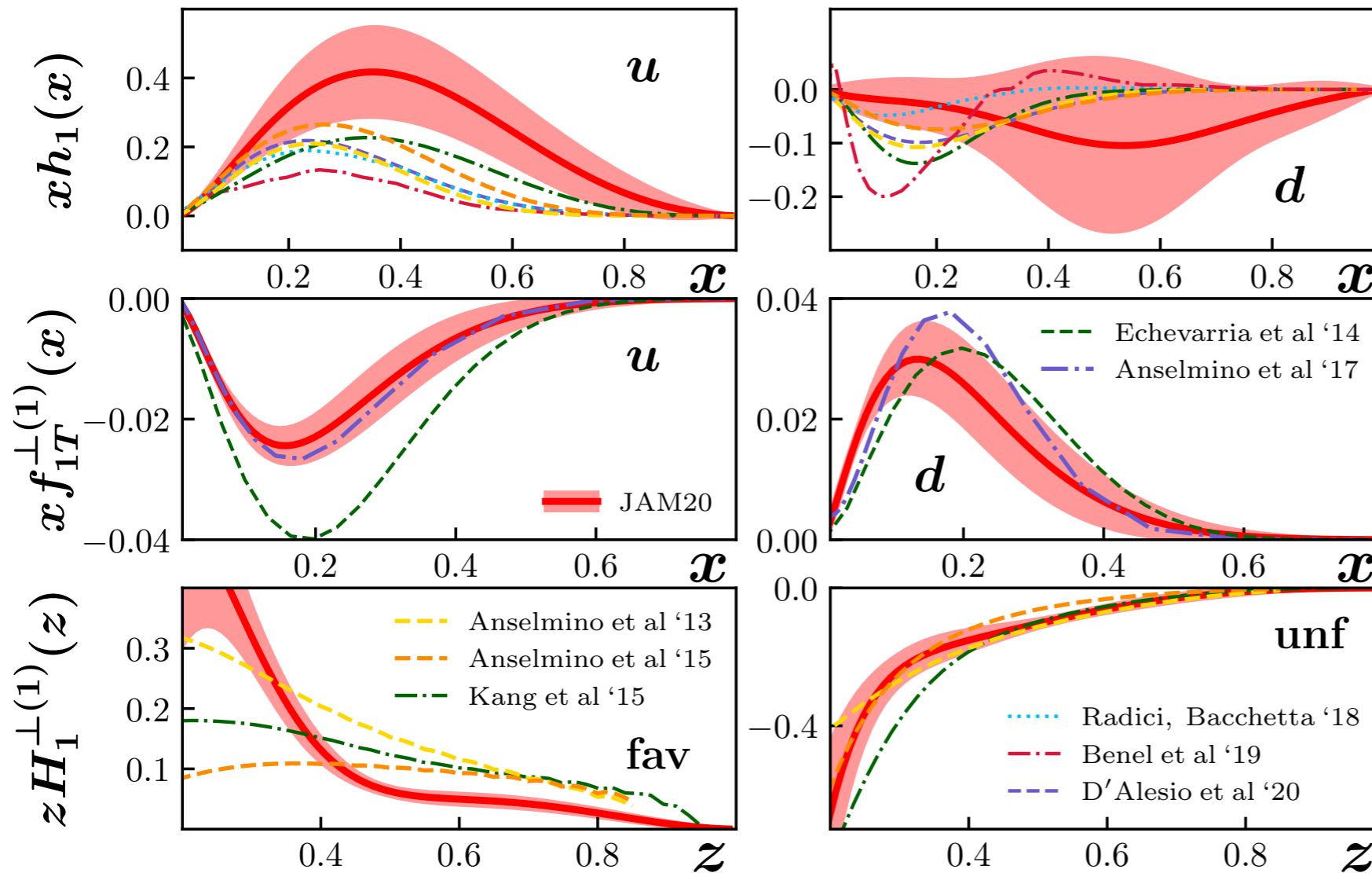
$$h_1(x), F_{FT}(x, x), H_1^{\perp(1)}(z), \hat{H}(z)$$

► Broad kinematical coverage to test universality

► The analysis is performed at parton level leading order, gaussian model is used for TMDs, and DGLAP-type evolution is implemented

# UNIVERSAL GLOBAL FIT 2020

*Cammarota, Gamberg, Kang, Miller, Pitonyak, Prokudin, Rogers, Sato (2020)*



Transversity

$$h_1(x)$$

Sivers

$$f_{1T}^{\perp(1)}(x)$$

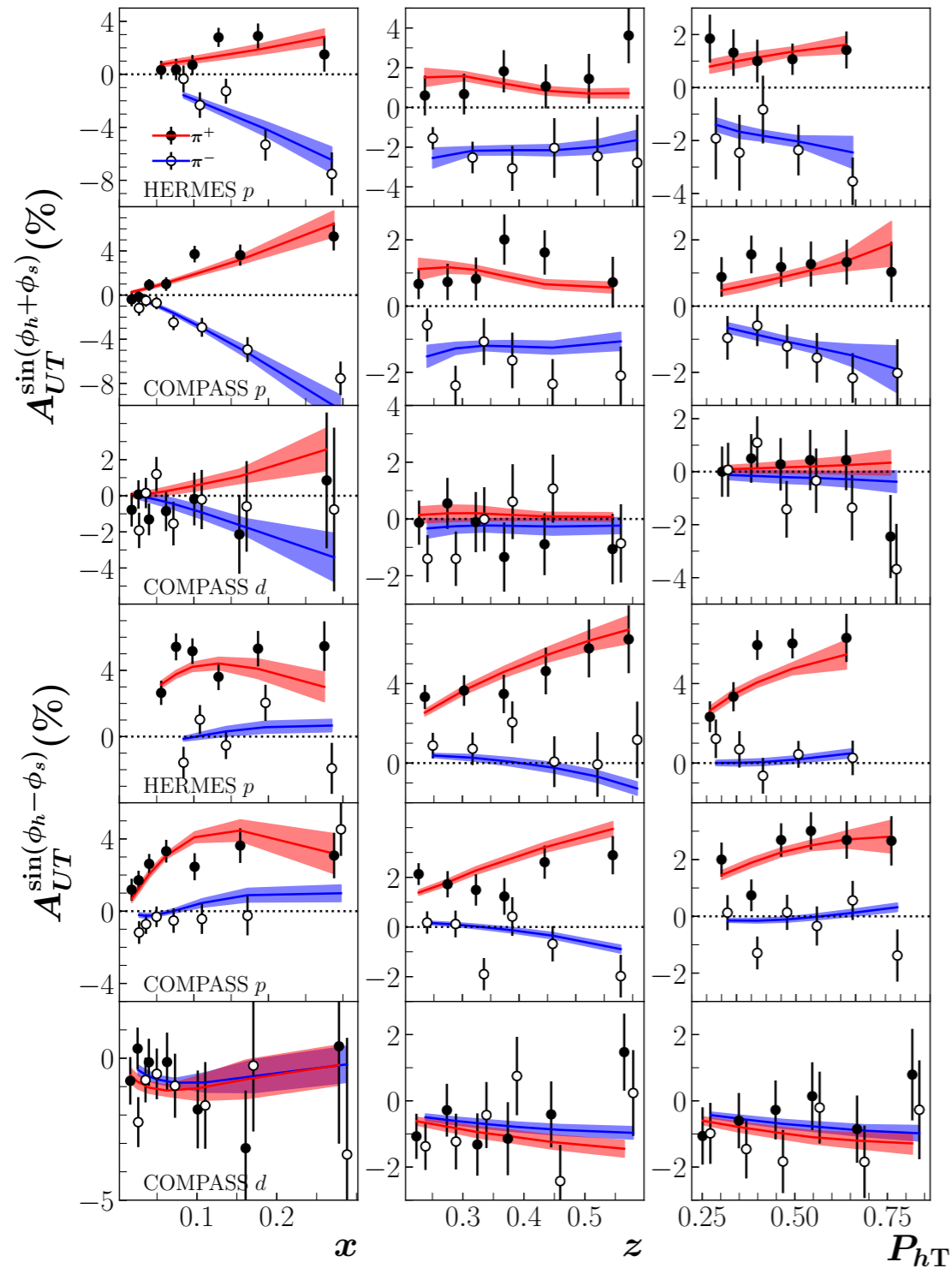
Collins FF

$$H_1^{\perp(1)}(z)$$

# UNIVERSAL GLOBAL FIT 2020

*Cammarota, Gamberg, Kang, Miller, Pitonyak, Prokudin, Rogers, Sato (2020)*

## SIDIS



## Collins asymmetry

$$\frac{\chi^2}{npoints} = \frac{107.1}{126} = 0.85$$

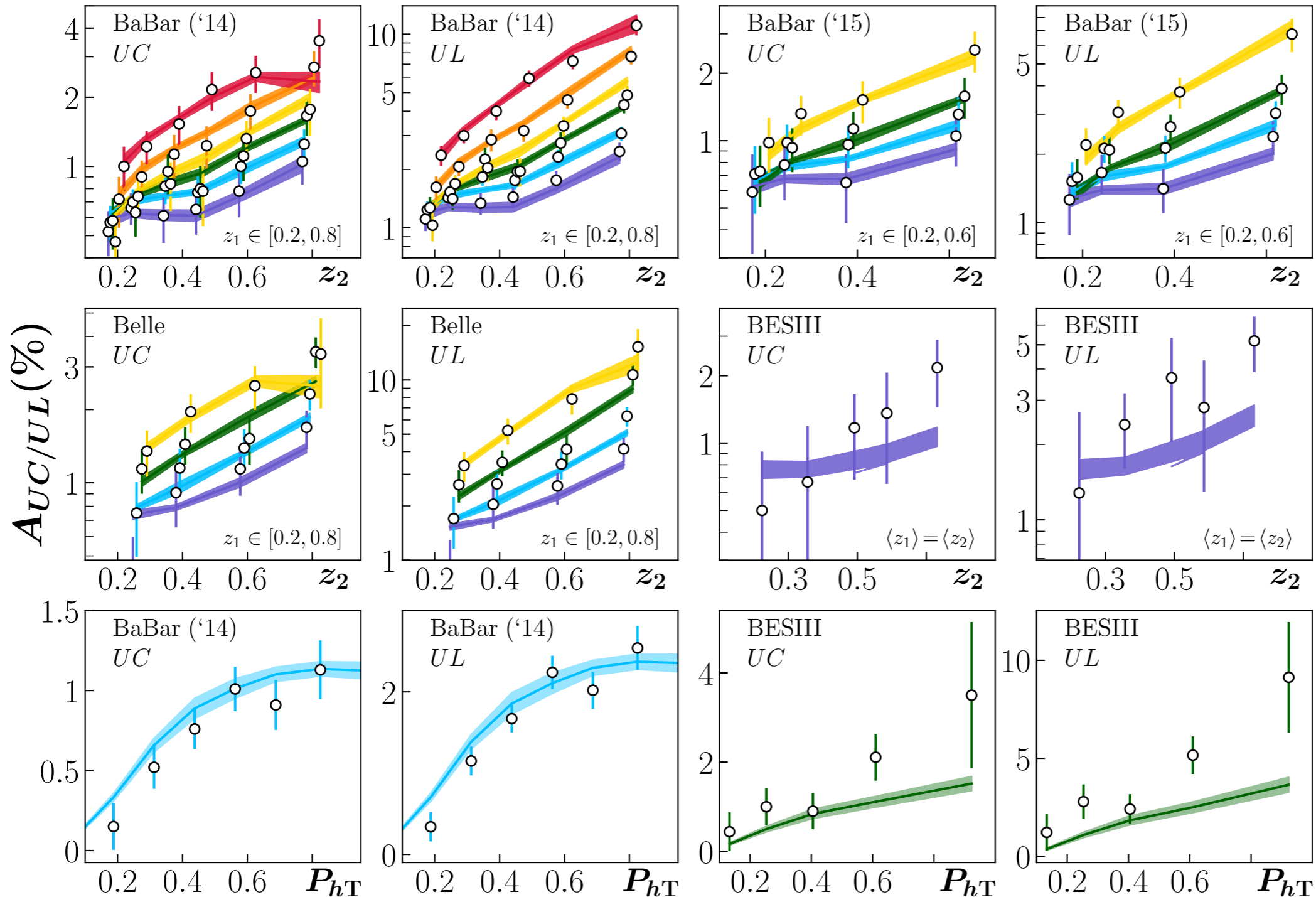
## Sivers asymmetry

$$\frac{\chi^2}{npoints} = \frac{85.4}{88} = 0.97$$

# UNIVERSAL GLOBAL FIT 2020

*Cammarota, Gamberg, Kang, Miller, Pitonyak, Prokudin, Rogers, Sato (2020)*

$e^+e^-$

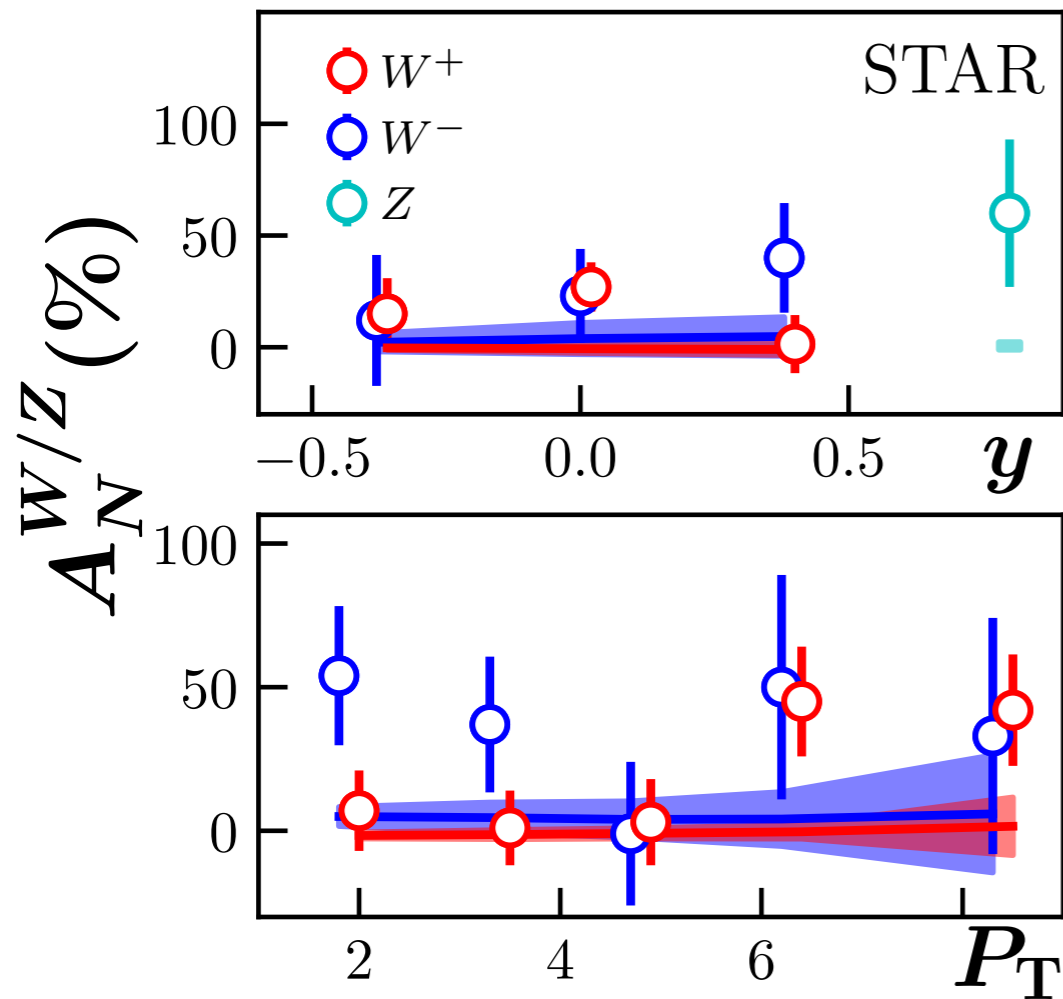


$$\frac{\chi^2}{npoints} = \frac{149.6}{176} = 0.85$$

# UNIVERSAL GLOBAL FIT 2020

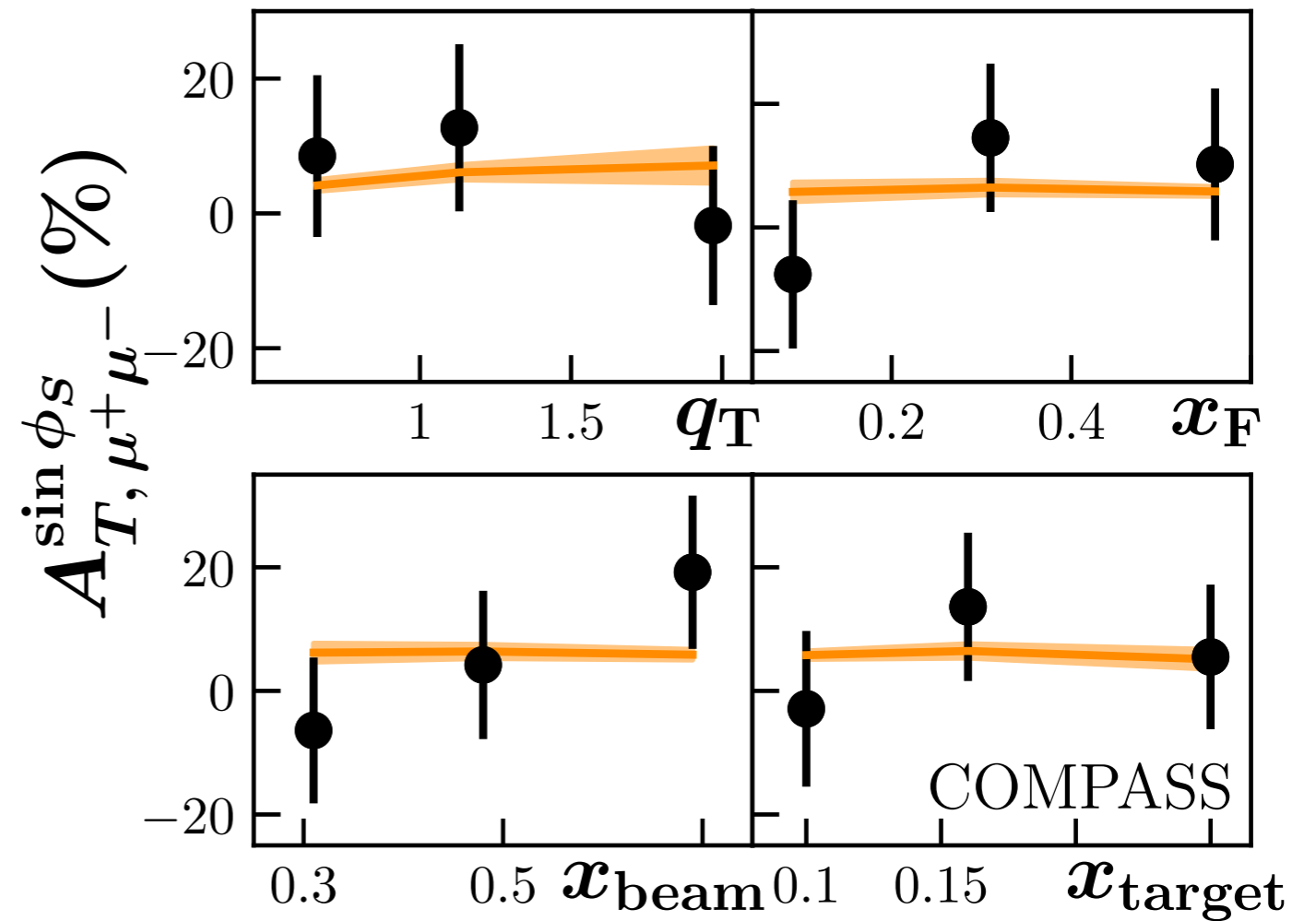
*Cammarota, Gamberg, Kang, Miller, Pitonyak, Prokudin, Rogers, Sato (2020)*

## Drell-Yan



$$\frac{\chi^2}{n_{\text{points}}} = \frac{29.8}{17} = 1.75$$

STAR



$$\frac{\chi^2}{n_{\text{points}}} = \frac{7.6}{12} = 0.63$$

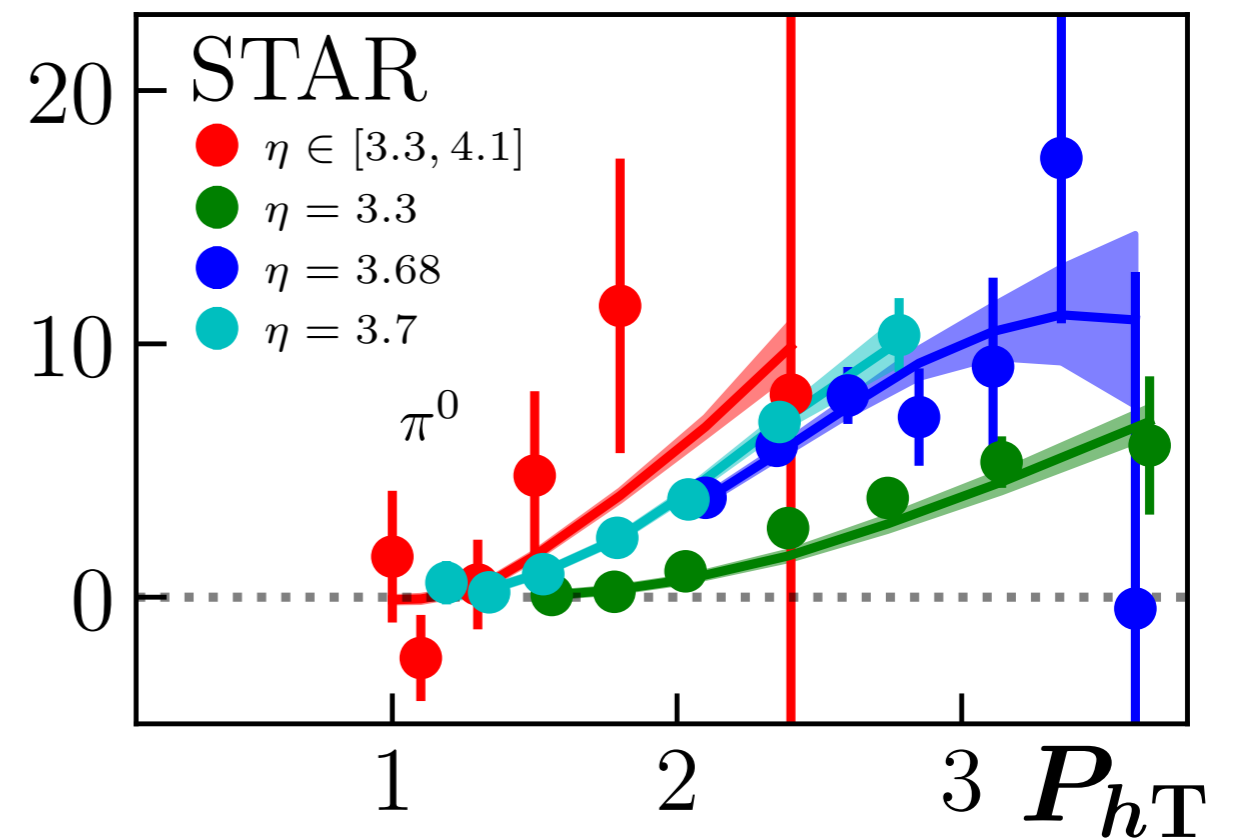
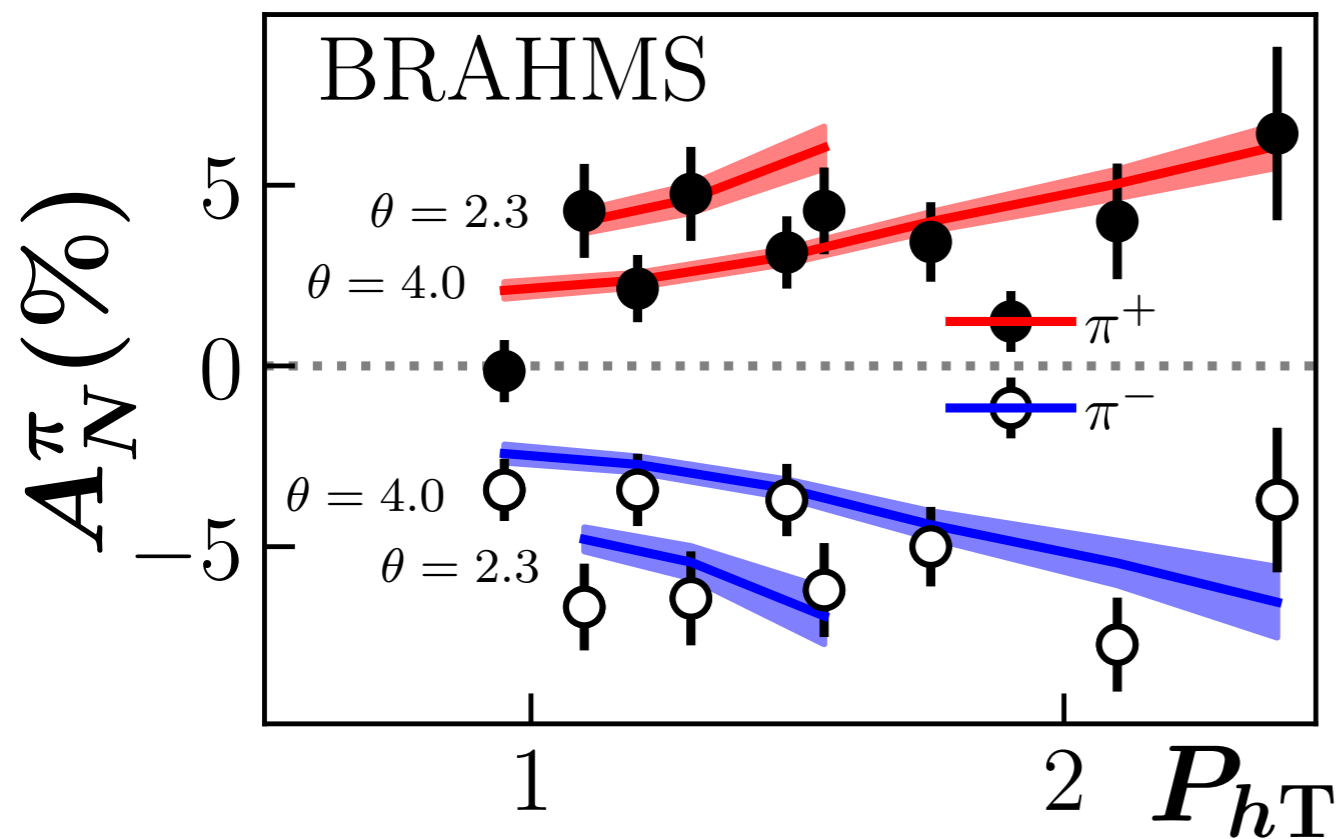
COMPASS DY



# UNIVERSAL GLOBAL FIT 2020

*Cammarota, Gamberg, Kang, Miller, Pitonyak, Prokudin, Rogers, Sato (2020)*

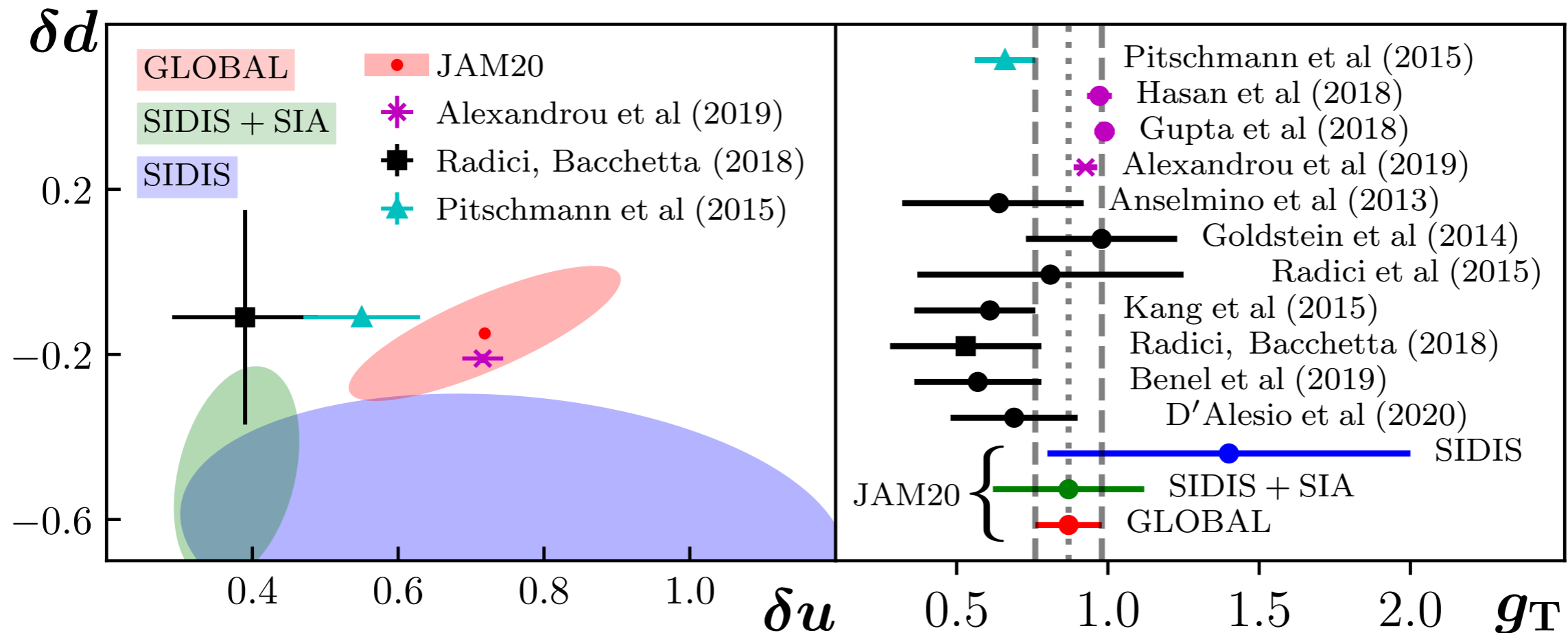
proton-proton  $A_N$



$$\frac{\chi^2}{npoints} = \frac{72.0}{60} = 1.2$$

# UNIVERSAL GLOBAL FIT 2020

Cammarota, Gamberg, Kang, Miller, Pitonyak, Prokudin, Rogers, Sato (2020)



Tensor charge from up and down quarks is constrained and compatible with lattice results

Isovector tensor charge  $g_T = \delta u - \delta d$

$g_T = 0.89 \pm 0.12$  compatible with lattice results

$\delta u$  and  $\delta d$   $Q^2=4$   $\text{GeV}^2$

$\delta u = 0.65 \pm 0.22$

$\delta d = -0.24 \pm 0.2$

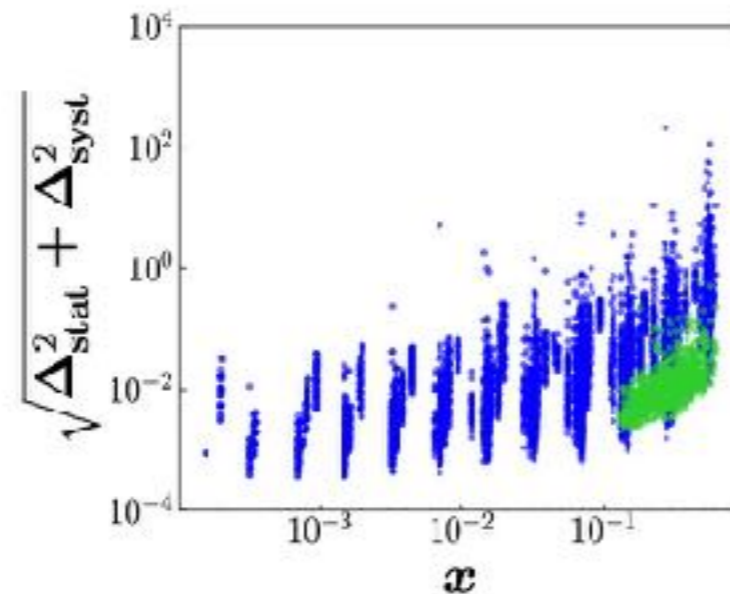
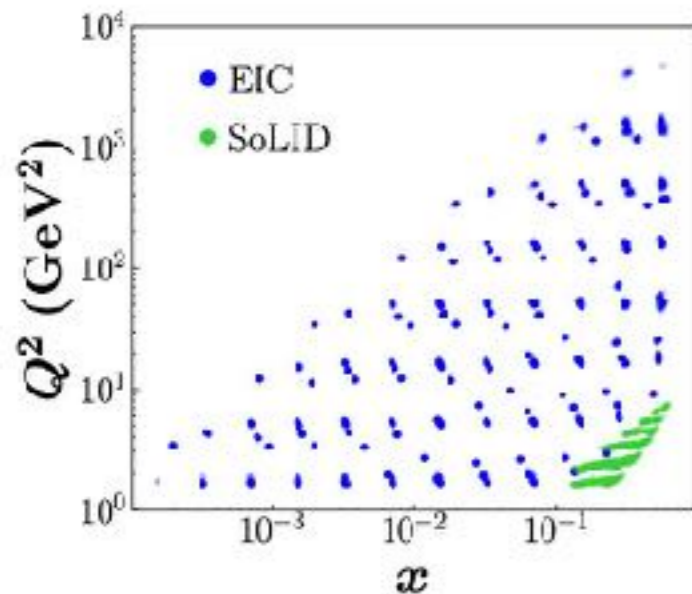
# **EIC IMPACT STUDY**

# TENSOR CHARGE AND FUTURE FACILITIES

Gamberg, Kang, Pitonyak, Prokudin, Sato, Seidl (21)

| EIC Pseudo-data |  |   |                                    |
|-----------------|--|---|------------------------------------|
| Observable      | Reactions  | CM Energy ( $\sqrt{S}$ )                      | $N_{\text{pts.}}$                  |
| Collins (SIDIS) | $e + p^\uparrow \rightarrow e + \pi^\pm + X$             | 141 GeV                                       | 756 ( $\pi^+$ )<br>744 ( $\pi^-$ ) |
|                 |  | 63 GeV  | 634 ( $\pi^+$ )<br>619 ( $\pi^-$ ) |
|                 |  | 45 GeV  | 537 ( $\pi^+$ )<br>556 ( $\pi^-$ ) |
|                 |  | 29 GeV  | 464 ( $\pi^+$ )<br>453 ( $\pi^-$ ) |
|                 | $e + {}^3\text{He}^\uparrow \rightarrow e + \pi^\pm + X$ | 85 GeV  | 647 ( $\pi^+$ )<br>650 ( $\pi^-$ ) |
|                 |  | 63 GeV  | 622 ( $\pi^+$ )<br>621 ( $\pi^-$ ) |
|                 |  | 29 GeV  | 461 ( $\pi^+$ )<br>459 ( $\pi^-$ ) |
|                 |  | <b>Total EIC <math>N_{\text{pts.}}</math></b> | <b>8223</b>                        |

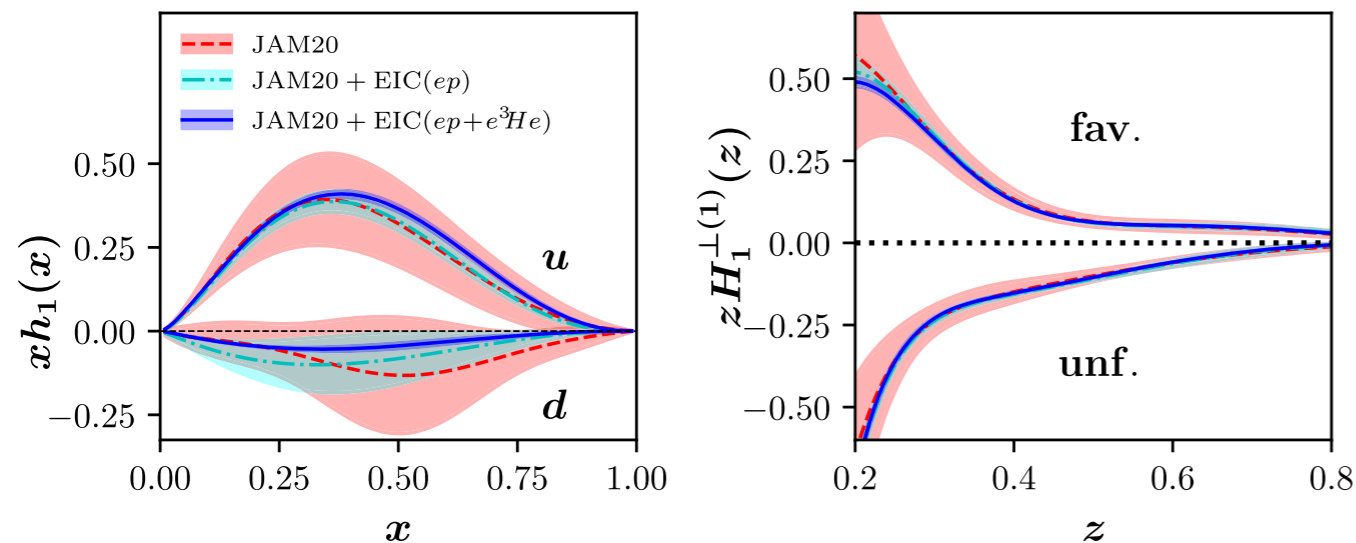
Assumed accumulated luminosities of  $10 \text{ fb}^{-1}$ , 70% polarization, conservatively accounted for detector smearing and acceptance effects



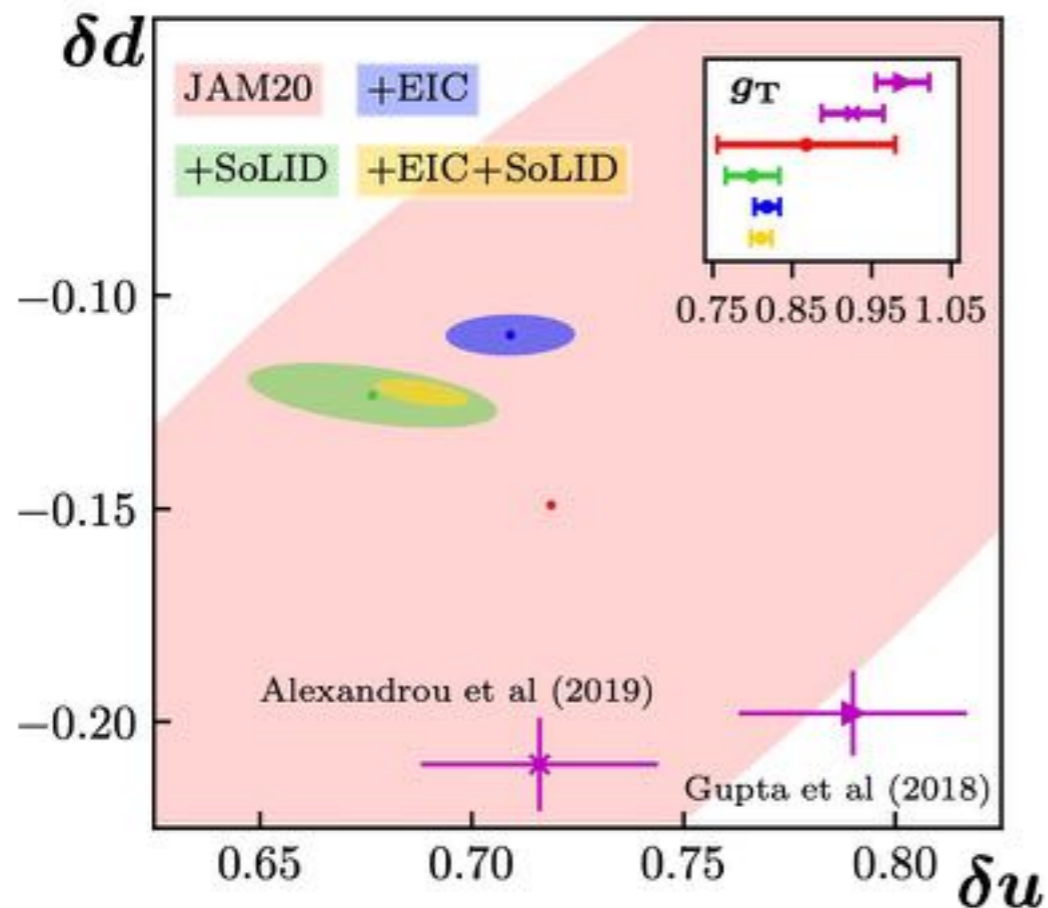
SoLID at JLab covers a complimentary region at higher  $x$  and lower  $Q^2$  with much greater luminosity – important to explore the effect of multiple measurements in different kinematic regions

# TENSOR CHARGE AND FUTURE FACILITIES

Gamberg, Kang, Pitonyak, Prokudin, Sato, Seidl (21)



Accuracy vs. precision – a precise measurement cannot always guarantee a very accurate extraction of the distributions, and multiple experiments, such as EIC and SoLID, should be performed in a wide kinematical region in order to minimize bias and expose any potential tensions between data sets (also one reason to have IR2@EIC)



The combined fit that includes both EIC and SoLID pseudo-data provides the best constraint on transversity and the tensor charges, with the results more precise than current lattice calculations

**TOWARDS JAM21**

# JAM21

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- Collins and Sivers effects (3D binned) SIDIS data from HERMES (2020) data



- $A_{UT}^{\sin \phi_S}$  (x and z projections only) from HERMES (2020)  $\tilde{H}$

- $A_N$  ( $x_F$  and  $p_T$  dependent) data from STAR (2021)

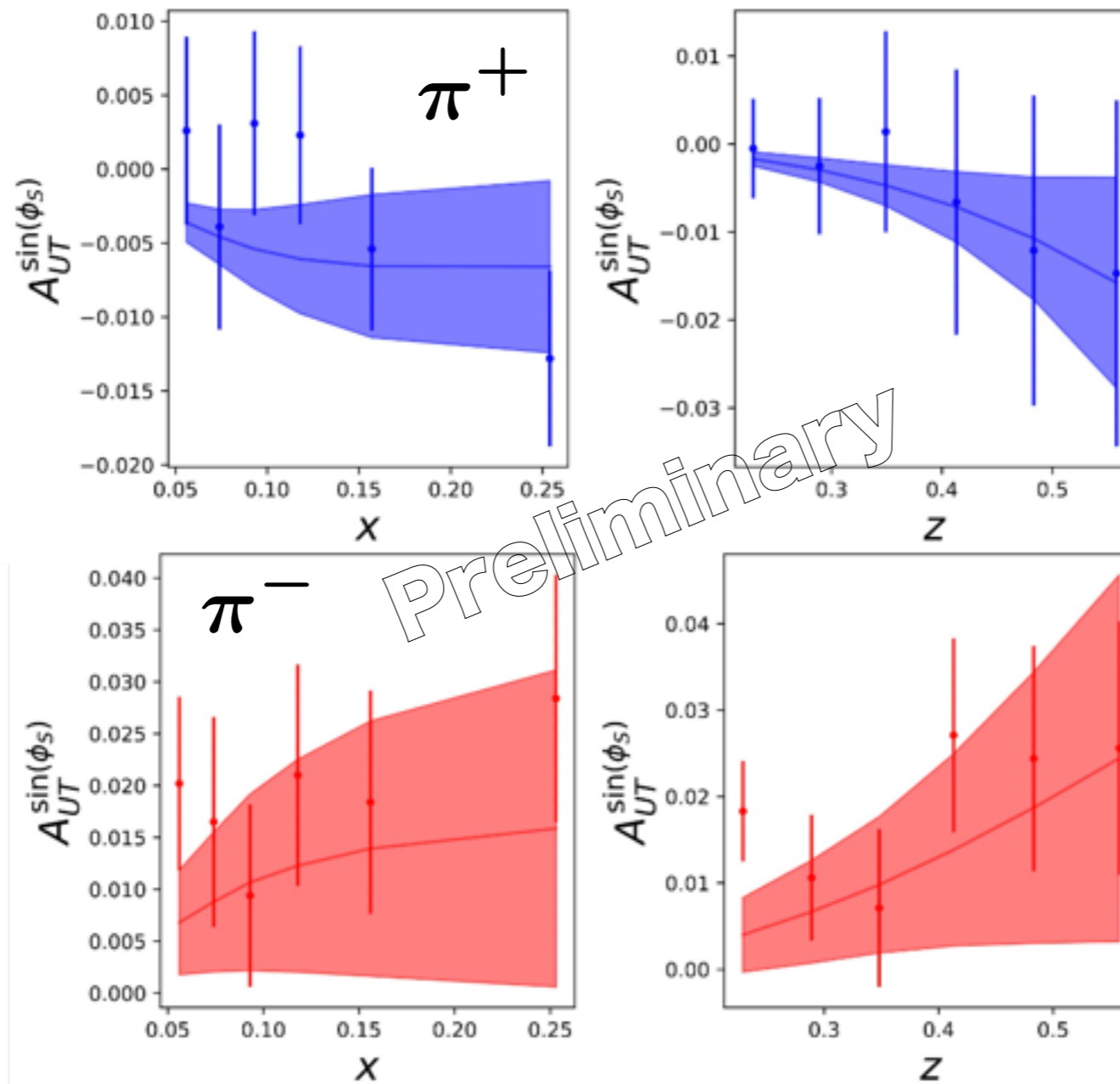


- Lattice data on  $\delta u$ ,  $\delta d$ ,  $g_T$  at the physical pion mass from Alexandrou, et al. (2019)

- Imposing the Soffer bound on transversity  $|h_1^q(x)| \leq \frac{1}{2}(f_1^q(x) + g_1^q(x))$

# JAM21

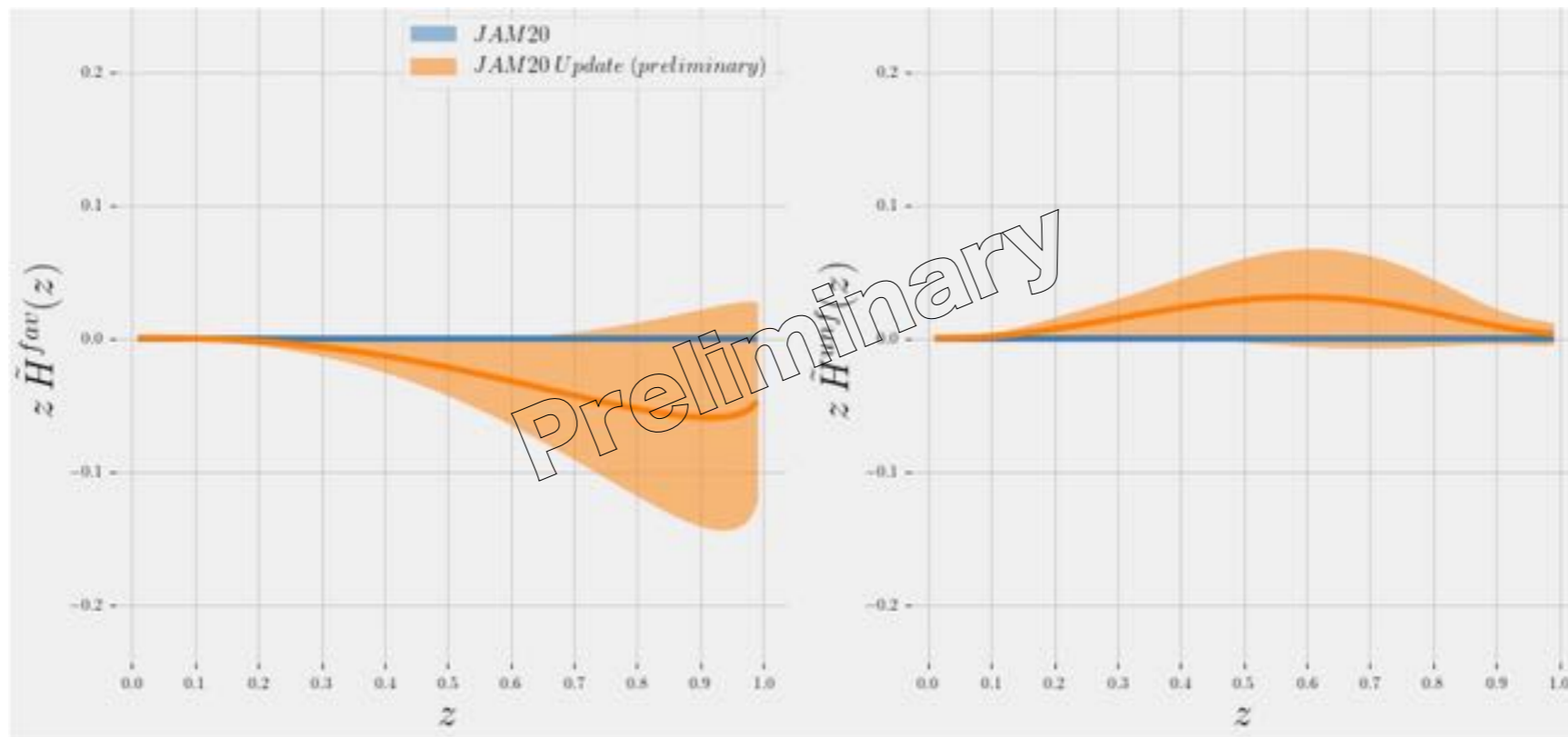
The first insight to the twist-3 fragmentation function  $\tilde{H}$





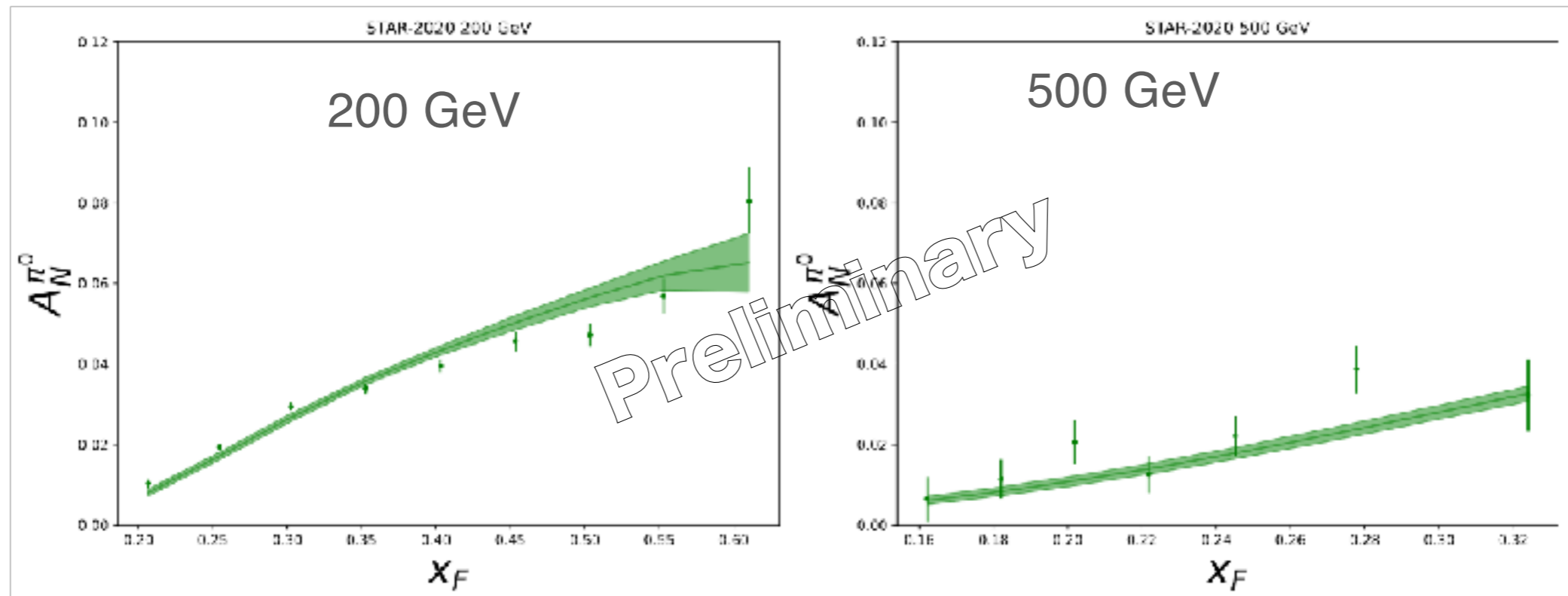
# JAM21

The first insight to the twist-3 fragmentation function  $\tilde{H}$



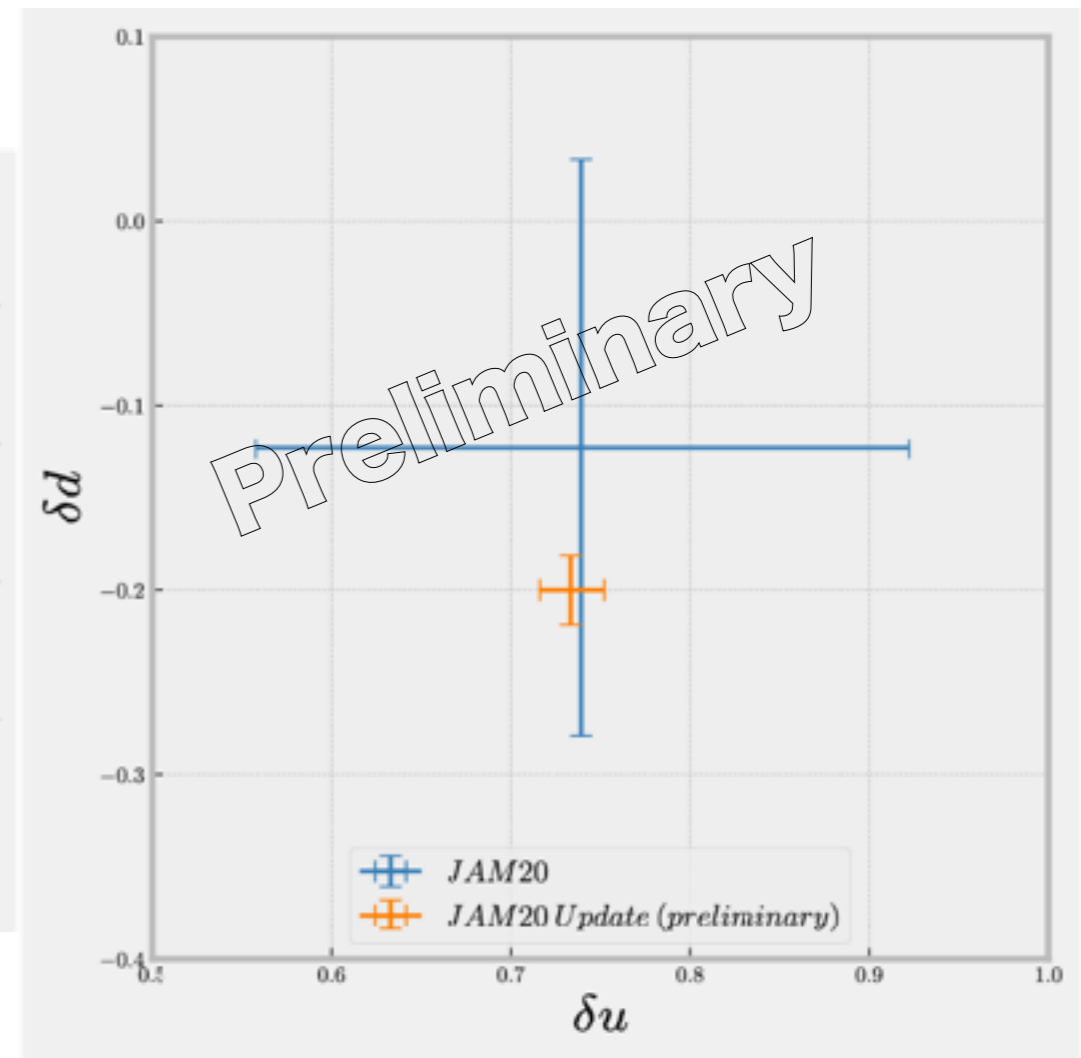
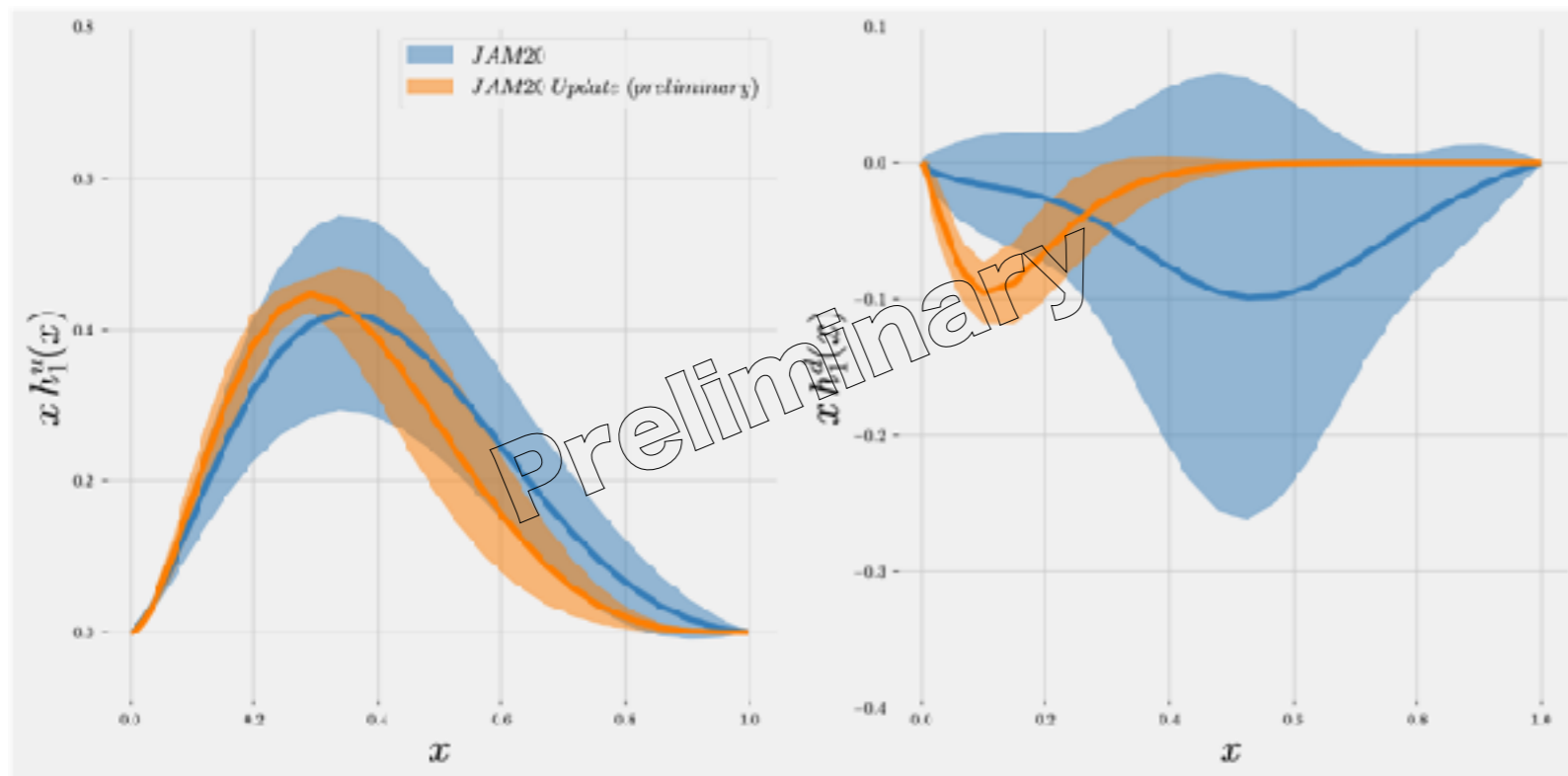
# JAM21

Description of the new  $A_N$  STAR data



# JAM21

Improved extraction of transversity and the tensor charge.



# CONCLUSIONS

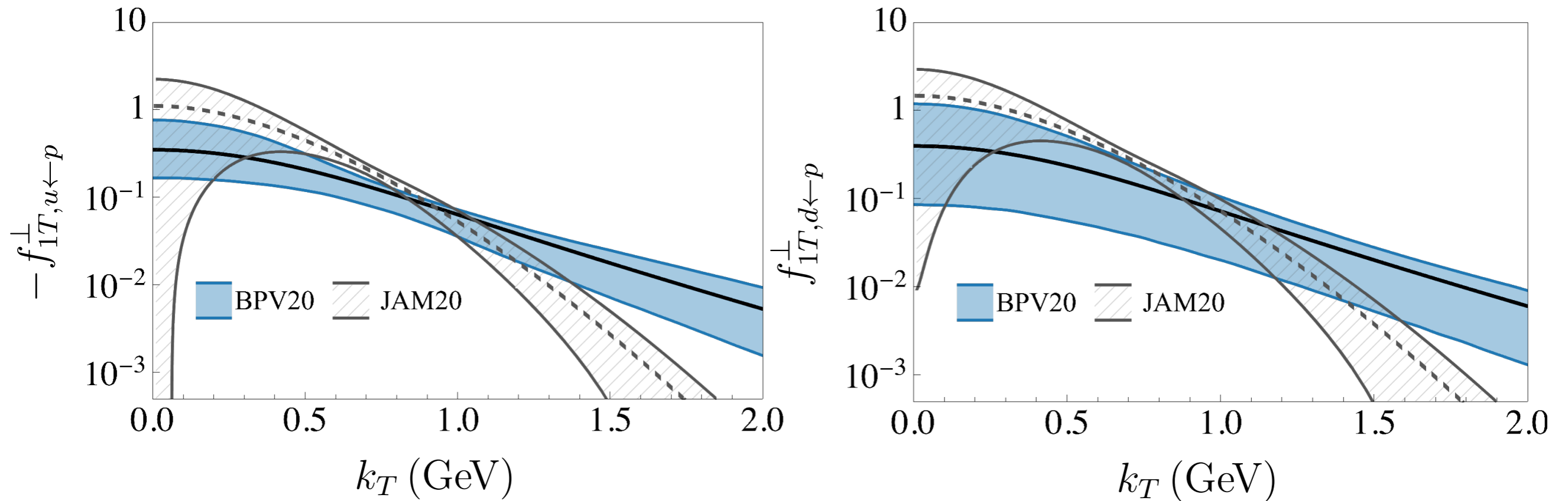
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- The transverse spin asymmetries in a variety of processes SIDIS, Drell-Yan,  $e^+e^-$ , and proton proton scattering have the same origin: (multi) parton correlation functions
- These effects have predominantly non perturbative origin and are universal
- Phenomenological results are consistent with lattice QCD in extraction of the isovector tensor charge  $g_T$  and individual contributions from up and down quarks
- The future EIC and JLab data will allow to have tensor charge extraction at the precision level of lattice QCD calculations

**BACKUP SLIDES**

# SIVERS FUNCTION IN THE MOMENTUM SPACE

Bury, Prokudin, Vladimirov (2021)



► Comparison to Jam20 (LO) analysis

Jam20: Cammarota, Gamberg, Kang, Miller, Pitonyak, Prokudin, Rogers, Sato (2020)