



Fondazione
di Sardegna



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**GENERAL HELICITY FORMALISM
FOR TWO-HADRON PRODUCTION IN
 e^+e^- COLLISIONS AND THE Λ
POLARIZING FRAGMENTATION
FUNCTION**

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in collaboration with
F. Murgia & M. Zaccheddu

PRD **102**, 054001 (2020)

JHEP **10**, 078 (2021)

OUTLINE

□ $e^+e^- \rightarrow h_1 h_2 + X$: theory

- ✓ TMD helicity formalism & FFs for Spin-½ hadrons
- ✓ Transverse Λ polarization

□ Phenomenology

- ✓ Fit of Belle data and extraction of the polarizing FF
- ✓ Predictions for SIDIS@EIC

□ Concluding remarks



$$e^+ e^- \rightarrow h_1(P_1) h_2(P_2) + X$$

- Complete results for the azimuthal and polarization observables
Boer, Jakob, Mulders 1997 - Pitonyak, Schlegel, Metz 2014
- TMD factorization for small relative transverse momenta w.r.t. Q^2
Collins 2011 - Echevarria, Idilbi, Scimemi 2012
- Formulation in the helicity formalism: partonic interpretation
UD, Murgia, Zaccheddu 2021



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- $e^+ e^- \rightarrow h_1(P_1)$ jet X
 - ✓ TMD factorization, role of soft factors, universality *Kang, Shao, Zhao 2020 - Boglione, Simonelli 2021 - Gumberg, Kang, Shao, Terry, Zhao 2021*
 - ✓ access to the intrinsic transverse momentum within a simplified TMD model
(see also *Anselmino, Kishore, Mukherjee 2019*)

➤ Improved TMD analysis *UD, Gumberg, Murgia, Zaccheddu 2021*
[talk by M. Zaccheddu]



$$e^+ e^- \rightarrow h_1(P_1) h_2(P_2) + X$$

Master formula [TMD factorization]

Helicity density matrix

$$\begin{aligned} & \rho_{\lambda_{h_1}, \lambda'_{h_1}}^{h_1, S_1} \rho_{\lambda_{h_2}, \lambda'_{h_2}}^{h_2, S_2} \frac{d\sigma^{e^+ e^- \rightarrow h_1 h_2 X}}{d \cos \theta dz_1 d^2 \mathbf{p}_{\perp 1} dz_2 d^2 \mathbf{p}_{\perp 2}} \\ &= \sum_{q_c, \bar{q}_d} \sum_{\{\lambda\}} \frac{1}{32\pi s} \frac{1}{4} \hat{M}_{\lambda_c \lambda_d, \lambda_a \lambda_b} \hat{M}_{\lambda'_c \lambda'_d, \lambda_a \lambda_b}^* \hat{D}_{\lambda_c, \lambda'_c}^{\lambda_{h_1}, \lambda'_{h_1}}(z_1, \mathbf{p}_{\perp 1}) \hat{D}_{\lambda_d, \lambda'_d}^{\lambda_{h_2}, \lambda'_{h_2}}(z_2, \mathbf{p}_{\perp 2}) \end{aligned}$$



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$$= \sum_{q_c, \bar{q}_d} \sum_{\{\lambda\}} \frac{1}{32\pi s} \frac{1}{4} \hat{M}_{\lambda_c \lambda_d, \lambda_a \lambda_b} M_{\lambda'_c \lambda'_d, \lambda_a \lambda_b}^* \hat{D}_{\lambda_c, \lambda'_c}^{\lambda_{h_1}, \lambda'_{h_1}}(z_1, \mathbf{p}_{\perp 1}) \hat{D}_{\lambda_d, \lambda'_d}^{\lambda_{h_2}, \lambda'_{h_2}}(z_2, \mathbf{p}_{\perp 2})$$

Unpol x-sec



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Unpol x-sec

Helicity scatt.
Amplitude
(hard)



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Unpol x-sec

Fragm. Funct.
(soft)

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$$e^+ e^- \rightarrow h_1(P_1) h_2(P_2) + X$$

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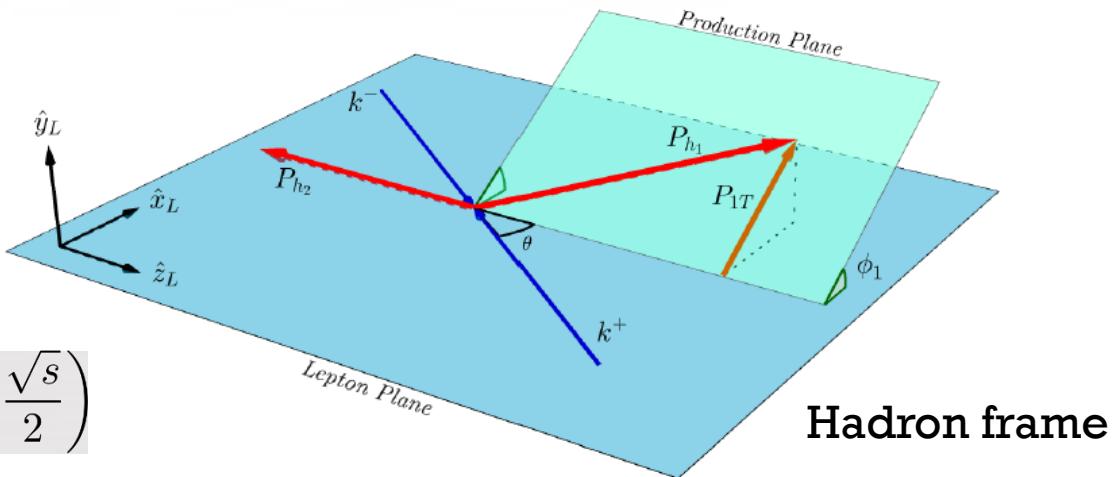
Scaling variables

$$z_p = 2|\mathbf{P}_h|/\sqrt{s} \quad z_h = 2E_h/\sqrt{s}$$

quark/antiquark momenta

$$q_2 \simeq \left(\frac{\sqrt{s}}{2}, -\frac{p_{\perp 2}}{z_{p_2}} \cos \varphi_2, -\frac{p_{\perp 2}}{z_{p_2}} \sin \varphi_2, -\frac{\sqrt{s}}{2} \right)$$

$$\mathbf{q}_1 = -\mathbf{q}_2$$



$$e^+ e^- \rightarrow h_1(P_1) h_2(P_2) + X$$

$$\begin{aligned} & \rho_{\lambda_{h_1}, \lambda'_{h_1}}^{h_1, S_1} \rho_{\lambda_{h_2}, \lambda'_{h_2}}^{h_2, S_2} \frac{d\sigma^{e^+ e^- \rightarrow h_1 h_2 X}}{d \cos \theta dz_1 d^2 \mathbf{p}_{\perp 1} dz_2 d^2 \mathbf{p}_{\perp 2}} \\ &= \sum_{q_c, \bar{q}_d} \sum_{\{\lambda\}} \frac{1}{32\pi s} \frac{1}{4} \hat{M}_{\lambda_c \lambda_d, \lambda_a \lambda_b} \hat{M}_{\lambda'_c \lambda'_d, \lambda_a \lambda_b}^* \hat{D}_{\lambda_c, \lambda'_c}^{\lambda_{h_1}, \lambda'_{h_1}}(z_1, \mathbf{p}_{\perp 1}) \hat{D}_{\lambda_d, \lambda'_d}^{\lambda_{h_2}, \lambda'_{h_2}}(z_2, \mathbf{p}_{\perp 2}) \end{aligned}$$

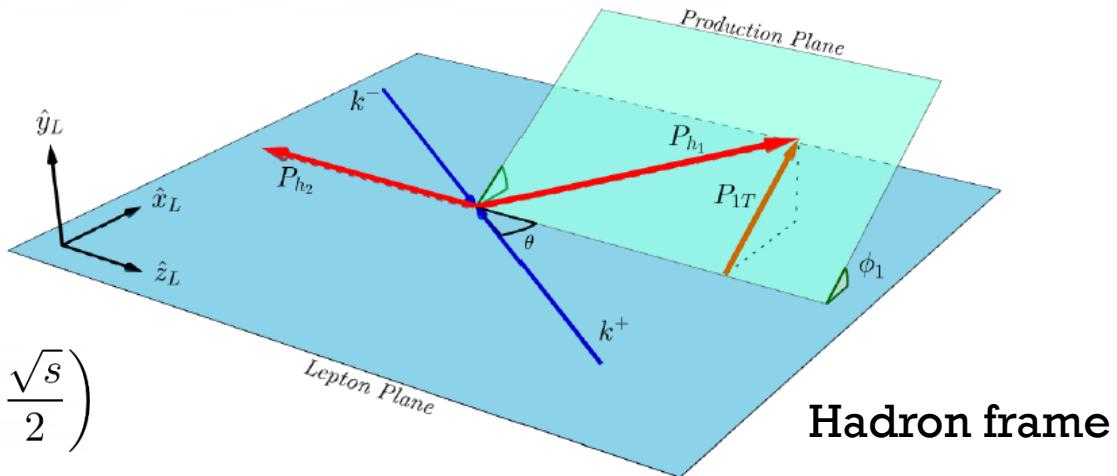
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$$\frac{d\sigma^{e^+ e^- \rightarrow h_1(S_1) h_2(S_2) X}}{d \cos \theta dz_1 d^2 \mathbf{p}_{\perp 1} dz_2 d^2 \mathbf{p}_{\perp 2}} = \int d^2 \mathbf{p}_{\perp 2} d^2 \mathbf{p}_{\perp 1} \delta^2(\mathbf{p}_{\perp 1} - \mathbf{P}_1 + \mathbf{p}_{\perp 2} z_{p_1}/z_{p_2}) \frac{d\sigma^{e^+ e^- \rightarrow h_1(S_1) h_2(S_2) X}}{d \cos \theta dz_1 d^2 \mathbf{p}_{\perp 1} dz_2 d^2 \mathbf{p}_{\perp 2}}$$



QUARK TMD-FFS FOR SPIN-1/2 HADRONS

Hadron helicity
density matrix

$$\rho_{\lambda_h, \lambda'_h}^{h, S_h} \hat{D}_{h/q, s_q}(z, \mathbf{p}_\perp) = \sum_{\lambda_q, \lambda'_q} \rho_{\lambda_q, \lambda'_q}^q \hat{D}_{\lambda_q, \lambda'_q}^{\lambda_h, \lambda'_h}(z, \mathbf{p}_\perp)$$

$$\rho_{\lambda_h, \lambda'_h}^{h, S_h} = \frac{1}{2} \begin{pmatrix} 1 + P_Z^h & P_X^h - iP_Y^h \\ P_X^h + iP_Y^h & 1 - P_Z^h \end{pmatrix}$$



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$$\mathbf{P}^h = (P_T^h \cos \phi_h, P_T^h \sin \phi_h, P_Z^h)$$

polarization components
in the **hadron helicity frame**



QUARK TMD-FFS FOR SPIN-1/2 HADRONS

Hadron helicity density matrix

$$\rho_{\lambda_h, \lambda'_h}^{h, S_h} \hat{D}_{h/q, s_q}(z, \mathbf{p}_\perp) = \sum_{\lambda_q, \lambda'_q} \rho_{\lambda_q, \lambda'_q}^q \hat{D}_{\lambda_q, \lambda'_q}^{\lambda_h, \lambda'_h}(z, \mathbf{p}_\perp)$$

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$$(P_J^h \hat{D}_{h/q, s_T}) = \Delta \hat{D}_{S_J/s_T}^{h/q} = \hat{D}_{S_J/s_T}^{h/q} - \hat{D}_{-S_J/s_T}^{h/q} \equiv \Delta \hat{D}_{S_J/s_T}^{h/q}(z, \mathbf{p}_\perp)$$

$$(P_J^h \hat{D}_{h/q, s_z}) = \Delta \hat{D}_{S_J/s_z}^{h/q} = \hat{D}_{S_J/+}^{h/q} - \hat{D}_{-S_J/+}^{h/q} \equiv \Delta \hat{D}_{S_J/+}^{h/q}(z, \mathbf{p}_\perp)$$

$$\hat{D}_{h/q, s_T} = \hat{D}_{h/q}(z, p_\perp) + \frac{1}{2} \Delta \hat{D}_{h/s_T}(z, \mathbf{p}_\perp)$$

8 TMD-FFs



QUARK TMD-FFS FOR SPIN-1/2 HADRONS

quark helicity
density matrix

$$\rho_{\lambda_h, \lambda'_h}^{h, S_h} \hat{D}_{h/q, s_q}(z, \mathbf{p}_\perp) = \sum_{\lambda_q, \lambda'_q} \rho_{\lambda_q, \lambda'_q}^q \hat{D}_{\lambda_q, \lambda'_q}^{\lambda_h, \lambda'_h}(z, \mathbf{p}_\perp)$$

$$\rho_{\lambda_q, \lambda'_q}^q = \frac{1}{2} \begin{pmatrix} 1 + P_L^q & P_T^q e^{-i\phi_{s_q}} \\ P_T^q e^{i\phi_{s_q}} & 1 - P_L^q \end{pmatrix}$$



QUARK TMD-FFS FOR SPIN-1/2 HADRONS

$$\rho_{\lambda_h, \lambda'_h}^{h, S_h} \hat{D}_{h/q, s_q}(z, \mathbf{p}_\perp) = \sum_{\lambda_q, \lambda'_q} \rho_{\lambda_q, \lambda'_q}^q \hat{D}_{\lambda_q, \lambda'_q}^{\lambda_h, \lambda'_h}(z, \mathbf{p}_\perp)$$

quark helicity density matrix product of fragm. amplitudes

$$\sum_X \hat{D}_{\lambda_h, \lambda_X; \lambda_q}(z, \mathbf{p}_\perp) \hat{D}_{\lambda'_h, \lambda_X; \lambda'_q}^*(z, \mathbf{p}_\perp)$$

The diagram illustrates the decomposition of the quark helicity density matrix $\rho_{\lambda_h, \lambda'_h}^{h, S_h}$ into a sum of products of quark fragmentation amplitudes and quark helicity density matrices. The equation is:

$$\rho_{\lambda_h, \lambda'_h}^{h, S_h} \hat{D}_{h/q, s_q}(z, \mathbf{p}_\perp) = \sum_{\lambda_q, \lambda'_q} \rho_{\lambda_q, \lambda'_q}^q \hat{D}_{\lambda_q, \lambda'_q}^{\lambda_h, \lambda'_h}(z, \mathbf{p}_\perp)$$

Below this, another sum is shown:

$$\sum_X \hat{D}_{\lambda_h, \lambda_X; \lambda_q}(z, \mathbf{p}_\perp) \hat{D}_{\lambda'_h, \lambda_X; \lambda'_q}^*(z, \mathbf{p}_\perp)$$

Annotations with arrows point to specific terms:

- A green arrow points to the term $\rho_{\lambda_q, \lambda'_q}^q$, labeled "quark helicity density matrix".
- A red arrow points to the term $\hat{D}_{\lambda_q, \lambda'_q}^{\lambda_h, \lambda'_h}(z, \mathbf{p}_\perp)$, labeled "product of fragm. amplitudes".
- A blue arrow points to the term $\hat{D}_{\lambda_h, \lambda_X; \lambda_q}(z, \mathbf{p}_\perp)$.

At the bottom right, a summary diagram shows a quark line q, λ_q entering a brown oval labeled \hat{D} . From the oval, three arrows emerge: one labeled h, λ_h and two labeled λ_X .



QUARK TMD-FFS FOR SPIN-1/2 HADRONS

$$\rho_{\lambda_h, \lambda'_h}^{h, S_h} \hat{D}_{h/q, s_q}(z, \mathbf{p}_\perp) = \sum_{\lambda_q, \lambda'_q} \rho_{\lambda_q, \lambda'_q}^q \hat{D}_{\lambda_q, \lambda'_q}^{\lambda_h, \lambda'_h}(z, \mathbf{p}_\perp)$$

real complex real (q) / imm (g)

8 real quantities
by parity



QUARK TMD-FFS FOR SPIN-1/2 HADRONS

$$\rho_{\lambda_h, \lambda'_h}^{h, S_h} \hat{D}_{h/q, s_q}(z, \mathbf{p}_\perp) = \sum_{\lambda_q, \lambda'_q} \rho_{\lambda_q, \lambda'_q}^q \hat{D}_{\lambda_q, \lambda'_q}^{\lambda_h, \lambda'_h}(z, \mathbf{p}_\perp)$$

\$D_{++}^{++}, D_{--}^{++}\$
\$D_{+-}^{++}, D_{++}^{+-}\$
\$D_{+-}^{+-}, D_{-+}^{+-}\$
real complex real (\$q\$) / imm (\$g\$)

**8 real quantities
by parity**

$\hat{D}_{h/q, s_q}$	$= (D_{++}^{++} + D_{--}^{++}) + 2P_T^q \text{Im} D_{+-}^{++} \sin(\phi_{s_q} - \phi_h)$
$P_Z^h \hat{D}_{h/q, s_q}$	$= P_L^q (D_{++}^{++} - D_{--}^{++}) + 2P_T^q \text{Re} D_{+-}^{++} \cos(\phi_{s_q} - \phi_h)$
$P_X^h \hat{D}_{h/q, s_q}$	$= 2P_L^q \text{Re} D_{++}^{+-} + P_T^q (D_{+-}^{+-} + D_{-+}^{+-}) \cos(\phi_{s_q} - \phi_h)$
$P_Y^h \hat{D}_{h/q, s_q}$	$= -2\text{Im} D_{++}^{+-} + P_T^q (D_{+-}^{+-} - D_{-+}^{+-}) \sin(\phi_{s_q} - \phi_h)$



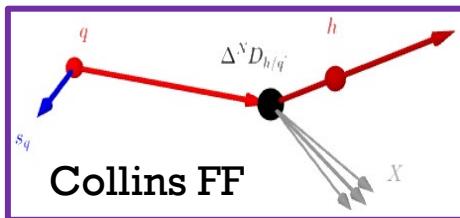
8 independent TMD Fragmentation Functions

		Hadron		
	Pol. States	U	L	T
Q u a r k	U	$\hat{D}_{h/q}$		$\Delta \hat{D}_{h^\uparrow/q}$
	L		$\Delta \hat{D}_{S_Z/s_L}^{h/q}$	$\Delta \hat{D}_{S_X/s_L}^{h/q}$
	T	$\Delta \hat{D}_{h/q^\uparrow}$	$\Delta \hat{D}_{S_Z/s_T}^{h/q}$	$\Delta \hat{D}_{S_X/s_T}^{h/q} / \Delta^- \hat{D}_{S_Y/s_T}^{h/q}$

Surviving in the collinear limit

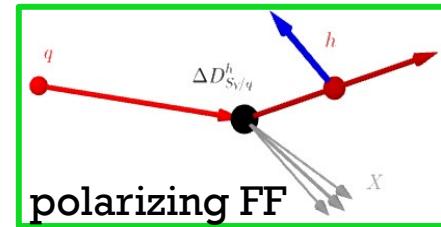


8 independent TMD Fragmentation Functions



T-odd & chiral-odd

UNIVERSAL



T-odd & chiral-even

		Hadron		
		U	L	T
Q u a r k	U	$\hat{D}_{h/q}$		$\Delta \hat{D}_{h^\uparrow/q}$
	L		$\Delta \hat{D}_{S_Z/s_L}^{h/q}$	$\Delta \hat{D}_{S_X/s_L}^{h/q}$
	T	$\Delta \hat{D}_{h/q^\uparrow}$	$\Delta \hat{D}_{S_Z/s_T}^{h/q}$	$\Delta \hat{D}_{S_X/s_T}^{h/q} / \Delta^- \hat{D}_{S_Y/s_T}^{h/q}$



OBSERVABLES (EXAMPLES)

$$\begin{aligned}
 P_X^{h_1} \frac{d\sigma^{e^+ e^- \rightarrow h_1 h_2 X}}{d \cos \theta dPS_{12}} &= \frac{3\pi\alpha^2}{4s} \sum_q e_q^2 \sin^2 \theta \Delta D_{S_X/s_T}^{h_1/q}(z_1, p_{\perp 1}) \Delta^N D_{h_2/\bar{q}\uparrow}(z_2, p_{\perp 2}) \sin(2\varphi_2 + \phi_{h_1}) \\
 P_Y^{h_1} \frac{d\sigma^{e^+ e^- \rightarrow h_1 h_2 X}}{d \cos \theta dPS_{12}} &= \frac{3\pi\alpha^2}{2s} \sum_q e_q^2 \left\{ (1 + \cos^2 \theta) \Delta D_{S_Y/q}^{h_1}(z_1, p_{\perp 1}) D_{h_2/\bar{q}}(z_2, p_{\perp 2}) \right. \\
 &\quad \left. + \frac{1}{2} \sin^2 \theta \Delta^- D_{S_Y/s_T}^{h_1}(z_1, p_{\perp 1}) \Delta^N D_{h_2/\bar{q}\uparrow}(z_2, p_{\perp 2}) \cos(2\varphi_2 + \phi_{h_1}) \right\} \\
 P_Z^{h_1} \frac{d\sigma^{e^+ e^- \rightarrow h_1 h_2 X}}{d \cos \theta dPS_{12}} &= \frac{3\pi\alpha^2}{4s} \sum_q e_q^2 \sin^2 \theta \Delta D_{S_Z/s_T}^{h_1/q}(z_1, p_{\perp 1}) \Delta^N D_{h_2/\bar{q}\uparrow}(z_2, p_{\perp 2}) \sin(2\varphi_2 + \phi_{h_1})
 \end{aligned}$$

$$\cos \phi_{h_1} \simeq \frac{P_{1T}}{p_{\perp 1}} \cos(\phi_1 - \varphi_2) - \frac{z_{p_1}}{z_{p_2}} \frac{p_{\perp 2}}{p_{\perp 1}} \quad \sin \phi_{h_1} = \frac{P_{1T}}{p_{\perp 1}} \sin(\phi_1 - \varphi_2)$$



OBSERVABLES (EXAMPLES)

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 \end{aligned}$$

convolutions

$$\mathcal{C}[wD\bar{D}] = \sum_q e_q^2 \int d^2 \mathbf{p}_{\perp 1} d^2 \mathbf{p}_{\perp 2} \delta^{(2)}(\mathbf{p}_{\perp 1} - \mathbf{P}_{1T} + \mathbf{p}_{\perp 2} z_{p_1}/z_{p_2}) w(\mathbf{p}_{\perp 2}, \mathbf{P}_{1T}) D(z_1, p_{\perp 1}) \bar{D}(z_2, p_{\perp 2})$$



OBSERVABLES (EXAMPLES)

$$\begin{aligned}
 P_X^{h_1} \frac{d\sigma^{e^+ e^- \rightarrow h_1 h_2 X}}{d \cos \theta dPS_{12}} &= \frac{3\pi\alpha^2}{4s} \sum_q e_q^2 \sin^2 \theta \Delta D_{S_X/s_T}^{h_1/q}(z_1, p_{\perp 1}) \Delta^N D_{h_2/\bar{q}\uparrow}(z_2, p_{\perp 2}) \sin(2\varphi_2 + \phi_{h_1}) \\
 P_Y^{h_1} \frac{d\sigma^{e^+ e^- \rightarrow h_1 h_2 X}}{d \cos \theta dPS_{12}} &= \frac{3\pi\alpha^2}{2s} \sum_q e_q^2 \left\{ (1 + \cos^2 \theta) \Delta D_{S_Y/q}^{h_1}(z_1, p_{\perp 1}) D_{h_2/\bar{q}}(z_2, p_{\perp 2}) \right. \\
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 P_Z^{h_1} \frac{d\sigma^{e^+ e^- \rightarrow h_1 h_2 X}}{d \cos \theta dPS_{12}} &= \frac{3\pi\alpha^2}{4s} \sum_q e_q^2 \sin^2 \theta \Delta D_{S_Z/s_T}^{h_1/q}(z_1, p_{\perp 1}) \Delta^N D_{h_2/\bar{q}\uparrow}(z_2, p_{\perp 2}) \sin(2\varphi_2 + \phi_{h_1})
 \end{aligned}$$

TENSORIAL ANALYSIS

ϕ_1 dependence – S_i scalar functions

Tensors

$$\begin{aligned}
 T^i &= \frac{1}{P_{1T}} \int d^2 \mathbf{p}_{\perp 2} p_{\perp 2}^i \Delta D^{h_1}(z_1, p_{\perp 1}) \Delta D^{h_2}(z_2, p_{\perp 2}) \\
 T^{ij} &= \frac{1}{P_{1T}^2} \int d^2 \mathbf{p}_{\perp 2} p_{\perp 2}^i p_{\perp 2}^j \Delta D^{h_1}(z_1, p_{\perp 1}) \Delta D^{h_2}(z_2, p_{\perp 2})
 \end{aligned}$$

$$\begin{aligned}
 T^i &= \frac{P_{1T}^i}{P_{1T}} S_1(P_{1T}) \\
 T^{ij} &= \frac{P_{1T}^i P_{1T}^j}{P_{1T}^2} S_2(P_{1T}) + \delta^{ij} S_3(P_{1T})
 \end{aligned}$$



STRUCTURE FUNCTIONS

(EXAMPLES)

Unpol.
cross section

$$\frac{d\sigma^{e^+ e^- \rightarrow h_1 h_2 X}}{d\cos\theta dz_1 dz_2 d^2\mathbf{P}_{1T}} = \frac{3\pi\alpha^2}{2s} \left\{ (1 + \cos^2\theta) F_{UU} + \sin^2\theta \cos(2\phi_1) \boxed{F_{UU}^{\cos(2\phi_1)}} \right\}$$



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(EXAMPLES)

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$$\frac{d\sigma^{e^+ e^- \rightarrow h_1 h_2 X}}{d \cos \theta dz_1 dz_2 d^2 \mathbf{P}_{1T}} = \frac{3\pi\alpha^2}{2s} \left\{ (1 + \cos^2 \theta) F_{UU} + \sin^2 \theta \cos(2\phi_1) \boxed{F_{UU}^{\cos(2\phi_1)}} \right\}$$

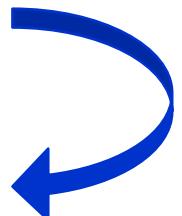
$$\cos(2\phi_1) F_{UU}^{\cos(2\phi_1)}$$

$$= \sum_q e_q^2 \int d^2 \mathbf{p}_{\perp 2} \frac{1}{4} \left[\frac{P_{1T}}{p_{\perp 1}} \cos(\phi_1 + \varphi_2) - \frac{z_{p_1}}{z_{p_2}} \frac{p_{\perp 2}}{p_{\perp 1}} \cos(2\varphi_2) \right] \Delta^N D_{h_1/q^\uparrow} \Delta^N D_{h_2/\bar{q}^\uparrow}$$

$$= \cos(2\phi_1) \mathcal{C} \left[\frac{1}{4} \left\{ \frac{P_{1T}}{p_{\perp 1}} \hat{\mathbf{p}}_{\perp 2} \cdot \hat{\mathbf{P}}_{1T} - \frac{z_{p_1}}{z_{p_2}} \frac{p_{\perp 2}}{p_{\perp 1}} \left[2(\hat{\mathbf{p}}_{\perp 2} \cdot \hat{\mathbf{P}}_{1T})^2 - 1 \right] \right\} \Delta^N D_{h_1/q^\uparrow} \Delta^N D_{h_2/\bar{q}^\uparrow} \right]$$

$$= \cos(2\phi_1) \mathcal{C} \left[\left\{ \frac{\mathbf{p}_{\perp 2} \cdot \mathbf{P}_{1T}}{z_1 z_2} - \left[2 \left(\frac{\mathbf{p}_{\perp 2} \cdot \hat{\mathbf{P}}_{1T}}{z_2} \right)^2 - \frac{p_{\perp 2}^2}{z_2^2} \right] \right\} \frac{H_1^\perp \bar{H}_1^\perp}{M_{h_1} M_{h_2}} \right],$$

Tensorial analysis



**Access to
the Collins function**

in agreement with *Boer, Jakob, Mulders 1997 - Pitonyak, Schlegel, Metz 2014*



TRANSVERSE POLARIZATION

$$\mathbf{P}_T^{h_1} = P_X^{h_1} \hat{\mathbf{X}}_{h_1} + P_Y^{h_1} \hat{\mathbf{Y}}_{h_1} = P_{x_L}^{h_1} \hat{\mathbf{x}}_L + P_{y_L}^{h_1} \hat{\mathbf{y}}_L = P_T^{h_1} (\cos \phi_{S_1}^L \hat{\mathbf{x}}_L + \sin \phi_{S_1}^L \hat{\mathbf{y}}_L)$$

From the helicity to the lab frame

$$\begin{aligned} P_T^{h_1} \frac{d\sigma^{e^+ e^- \rightarrow h_1 h_2 X}}{d\cos\theta dz_1 dz_2 d^2\mathbf{P}_{1T}} &= \frac{3\pi\alpha^2}{4s} \left\{ \left(1 + \cos^2\theta\right) \sin(\phi_1 - \phi_{S_1}^L) F_{TU}^{\sin(\phi_1 - \phi_{S_1}^L)} \right. \\ &+ \left. \sin^2\theta \left(\sin(\phi_1 + \phi_{S_1}^L) F_{TU}^{\sin(\phi_1 + \phi_{S_1}^L)} + \sin(3\phi_1 - \phi_{S_1}^L) F_{TU}^{\sin(3\phi_1 - \phi_{S_1}^L)} \right) \right\} \end{aligned}$$



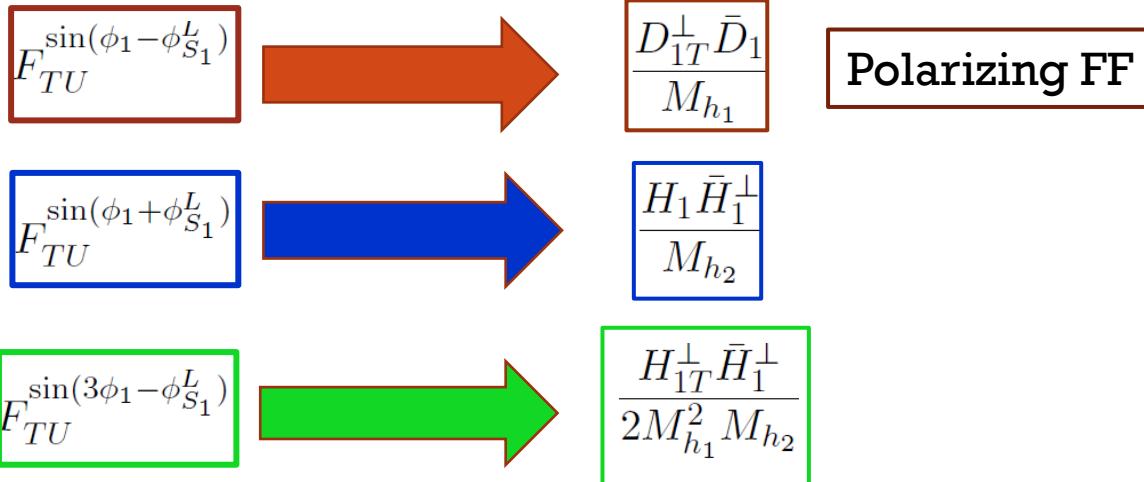
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$$+ \sin^2\theta \left(\sin(\phi_1 + \phi_{S_1}^L) F_{TU}^{\sin(\phi_1 + \phi_{S_1}^L)} + \sin(3\phi_1 - \phi_{S_1}^L) F_{TU}^{\sin(3\phi_1 - \phi_{S_1}^L)} \right) \left. \right\}$$



TRANSVERSE POLARIZATION

Projection along the normal to the hadron plane [lab frame] (phenom.)

$$\hat{\mathbf{n}} \equiv (\cos \phi_n, \sin \phi_n, 0) = \frac{-\mathbf{P}_2 \times \mathbf{P}_1}{|\mathbf{P}_2 \times \mathbf{P}_1|} = -\sin \phi_1 \hat{\mathbf{x}}_L + \cos \phi_1 \hat{\mathbf{y}}_L, \quad P_n^{h_1} \equiv \mathbf{P}^{h_1} \cdot \hat{\mathbf{n}} = \mathbf{P}_T^{h_1} \cdot \hat{\mathbf{n}}$$

$$\phi_n = \phi_1 + \frac{\pi}{2}$$



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$$\boxed{\phi_n = \phi_1 + \frac{\pi}{2}} \quad \rightarrow \quad \boxed{\phi_{S_1}^L = \phi_n}$$

$$\begin{aligned} P_n^{h_1} \frac{d\sigma^{e^+ e^- \rightarrow h_1 h_2 X}}{d \cos \theta dz_1 dz_2 d^2 \mathbf{P}_{1T}} &= \frac{3\pi\alpha^2}{4s} \left\{ \left(1 + \cos^2 \theta\right) \left[-F_{TU}^{\sin(\phi_1 - \phi_{S_1}^L)} \right] \right. \\ &+ \left. \sin^2 \theta \left(F_{TU}^{\sin(\phi_1 + \phi_{S_1}^L)} - F_{TU}^{\sin(3\phi_1 - \phi_{S_1}^L)} \right) \cos(2\phi_1) \right\} \end{aligned}$$



TRANSVERSE POLARIZATION

Projection along the normal to the hadron plane [lab frame] (phenom.)

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By integrating over \mathbf{P}_{1T} , the $\cos(2\phi_1)$ term is washed out, leaving



TRANSVERSE POLARIZATION

Projection along the normal to the hadron plane [lab frame] (phenom.)

$$P_n^{h_1} \frac{d\sigma^{e^+ e^- \rightarrow h_1 h_2 X}}{d \cos \theta dz_1 dz_2} = \frac{3\pi\alpha^2}{4s} \left\{ (1 + \cos^2 \theta) \left[-F_{TU}^{\sin(\phi_1 - \phi_{S_1}^L)} \right] \right.$$



$$P_n^{h_1}(z_1, z_2) = -\frac{F_{TU}^{\sin(\phi_1 - \phi_{S_1}^L)}}{F_{UU}}.$$

Direct access to
the polarizing FF



PHENOMENOLOGY

- First attempt to describe $P_T(\Lambda)$ data in unpolarized hadron-hadron collisions within a phenomenological TMD model *Anselmino, Boer, UD, Murgia 2001*
- Twist-three approach: *Kanazawa, Koike 2001*



- Belle data for $P_T(\Lambda)$ in $e^+e^- \rightarrow \Lambda^\uparrow h + X$ and $e^+e^- \rightarrow \Lambda^\uparrow(\text{jet}) + X$
Guan et al. (Belle Coll.) 2019



GAUSSIAN MODEL

$$D_{h/q}(z, p_\perp) = D_{h/q}(z) \frac{e^{-p_\perp^2/\langle p_\perp^2 \rangle}}{\pi \langle p_\perp^2 \rangle} \quad \langle p_\perp^2 \rangle_{\text{pol}} = \frac{M_{\text{pol}}^2}{M_{\text{pol}}^2 + \langle p_\perp^2 \rangle} \langle p_\perp^2 \rangle$$
$$\Delta D_{h^\uparrow/q}(z, p_\perp) = \Delta D_{h^\uparrow/q}(z) \frac{\sqrt{2e} p_\perp}{M_{\text{pol}}} \frac{e^{-p_\perp^2/\langle p_\perp^2 \rangle_{\text{pol}}}}{\pi \langle p_\perp^2 \rangle}$$

$$\Delta D_{\Lambda^\uparrow/q}(z) = N_q z^{a_q} (1-z)^{b_q} \frac{(a_q + b_q)^{(a_q + b_q)}}{a_q^{a_q} b_q^{b_q}} D_{\Lambda/q}(z)$$



FIT OF BELLE DATA

No isospin symmetry constraint

$$N_u = 0.67^{+0.33}_{-0.32} \quad N_d = -0.65^{+0.32}_{-0.35}$$

$$N_s = -1.00^{+0.02}_{-0.00} \quad N_{\text{sea}} = -0.40^{+0.18}_{-0.28}$$

$$a_s = 2.17^{+0.76}_{-0.71}$$

$$b_u = 3.36^{+3.15}_{-2.31} \quad b_{\text{sea}} = 2.14^{+3.06}_{-1.92}$$

$$\langle p_\perp^2 \rangle_{\text{pol}} = 0.056^{+0.026}_{-0.019} \text{ GeV}^2$$

$\Lambda + \pi/K$ data fit : $\chi^2_{\text{dof}} = 1.26$

- ✓ 8 parameter fit
- ✓ No TMD evolution
(fixed-scale analysis)

Similar analysis

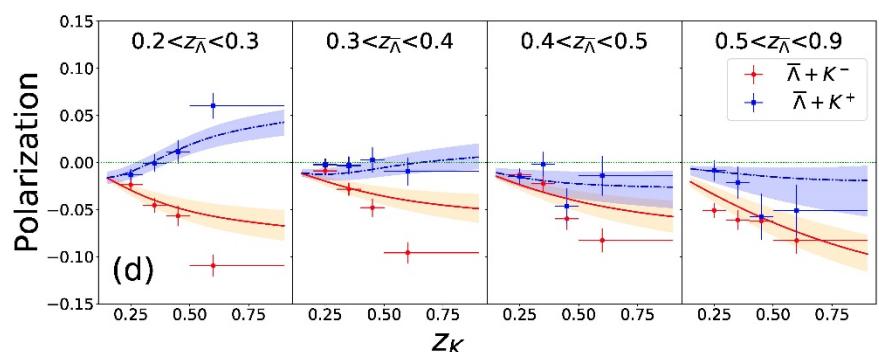
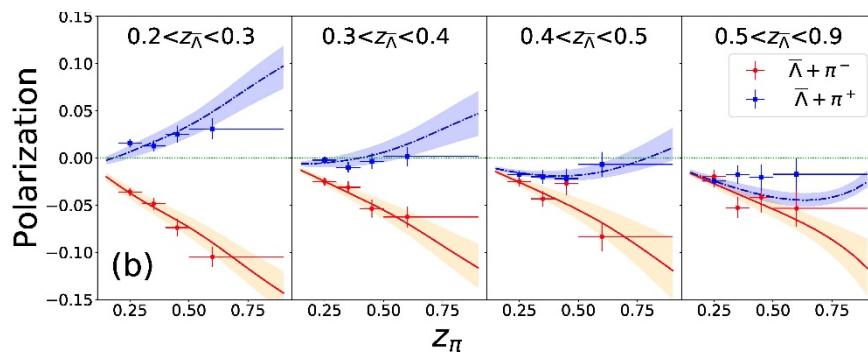
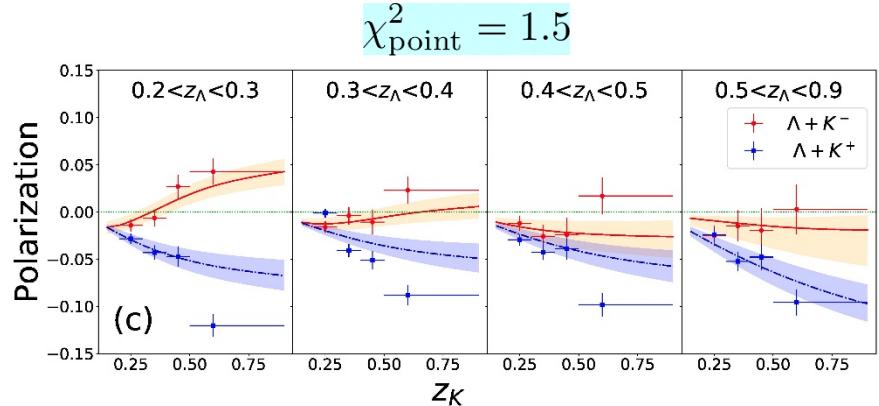
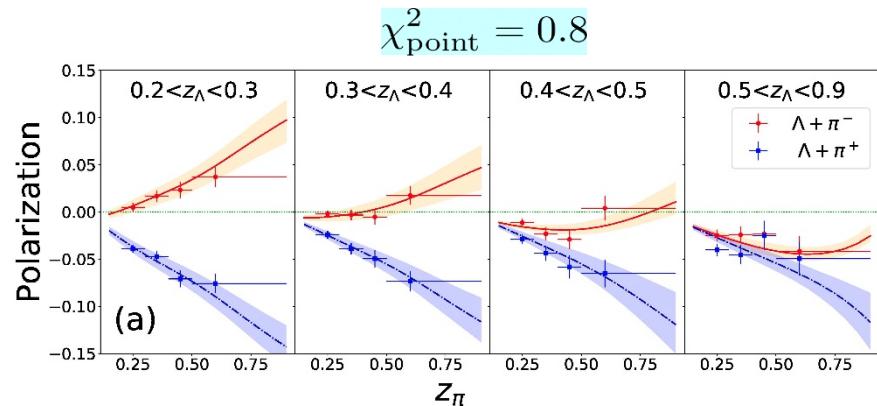
Callos, Kang, Terry 2020

isospin symmetry constraint

Chen, Liang, Pan, Song, Wei 2021



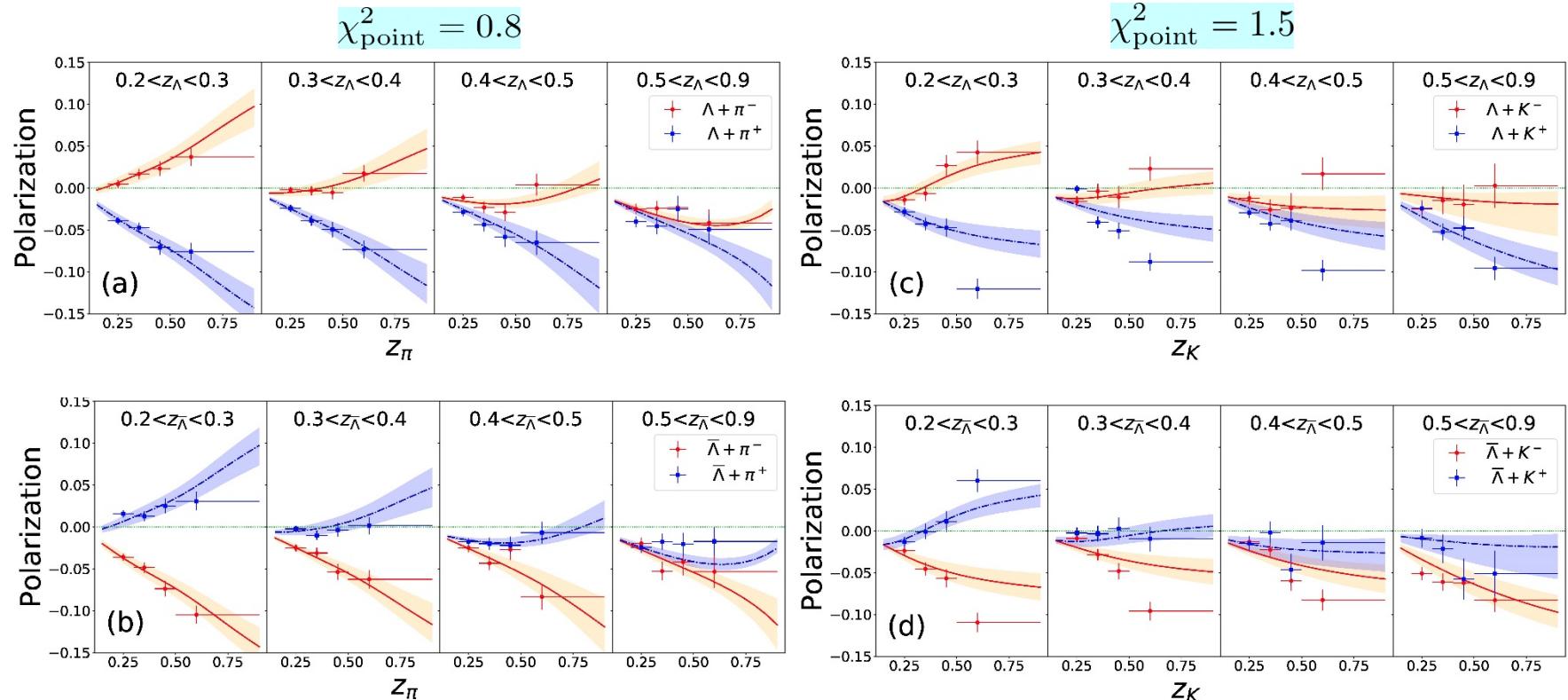
Associated production



- simpler fits with only two pFFs ($u=d, s$ and/or no sea) $\rightarrow \chi^2 \sim 2$
- kaon, some tension (unpol FFs?, parametriz. of polFF?)



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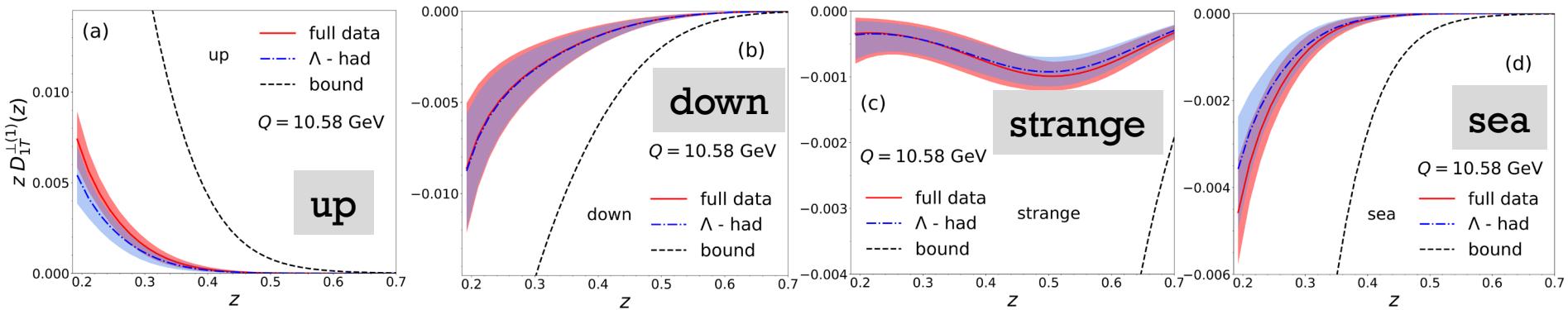


- simpler fits with only two pFFs ($u=d$, s and/or no sea) $\rightarrow \chi^2 \sim 2$
- kaon, some tension (unpol FFs?, parametriz. of polFF?)
- u, d, s **AND isospin symmetry** $\rightarrow \chi^2 \sim 2$
- u, d, s **plus charm** w/wo isospin symm. $\rightarrow \chi^2 \sim 1.27$ [& better $\chi^2(K)$]



First moments of the polarizing FF

$$D_{1T}^{\perp(1)}(z) = \int d^2\mathbf{P}_\perp \frac{p_\perp}{2zm_h} \Delta D_{h^\dagger/q}(z, p_\perp)$$



Full-data (including Λ (jet)) and Λ -had extractions
Consistency: central lines and (overlapping) uncertainty bands!!!

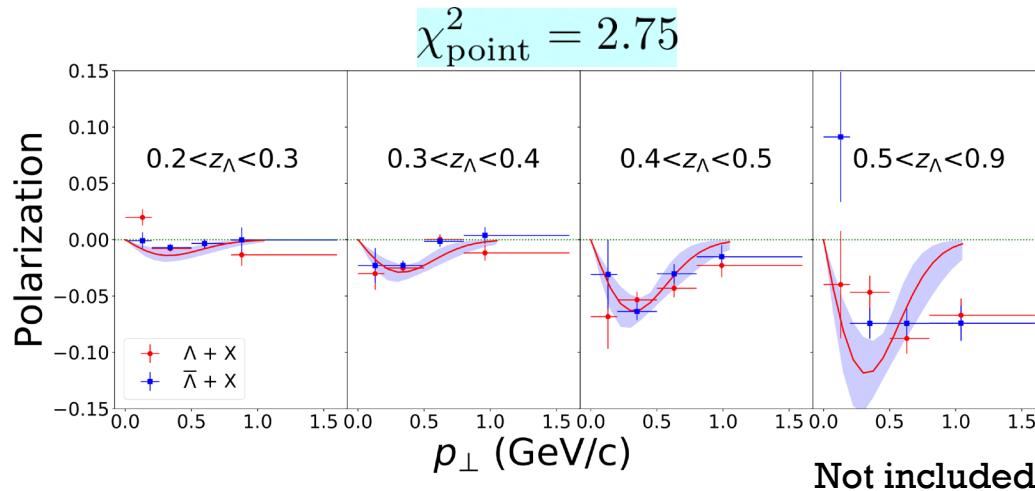
- $e^+e^- \rightarrow h_1(P_1) \text{ jet } X$ within a simplified TMD scheme

$$\mathcal{P}_T(z_1, p_{\perp 1}) = \frac{\sum_q e_q^2 \Delta D_{h_1^\dagger/q}(z_1, p_{\perp 1})}{\sum_q e_q^2 D_{h_1/q}(z_1, p_{\perp 1})}$$

orthogonal to the thrust plane



Inclusive Λ production [full-data analysis]



- Less good description...but not so bad!
- Theory
 - ✓ Polarization = 0 at $p_\perp = 0$
 - ✓ $P(\Lambda) = P(\bar{\Lambda})$
- not *clearly* visible in the data
- **Reanalysis within TMD factorization** [next talk by M. Zaccheddu]



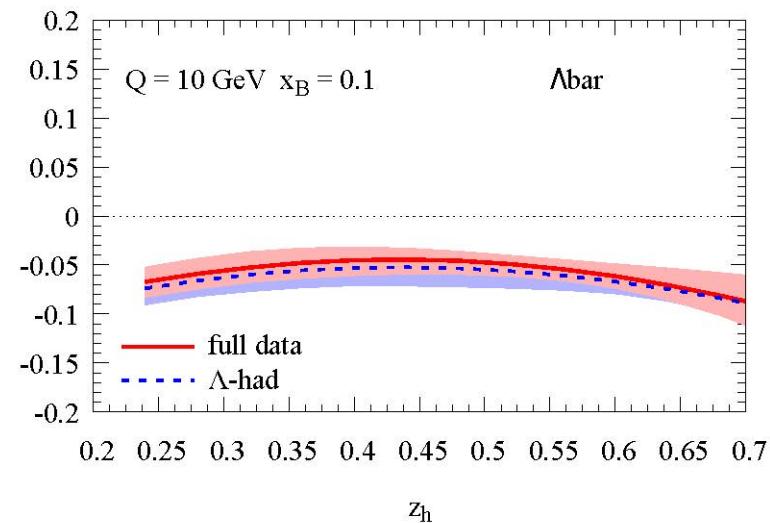
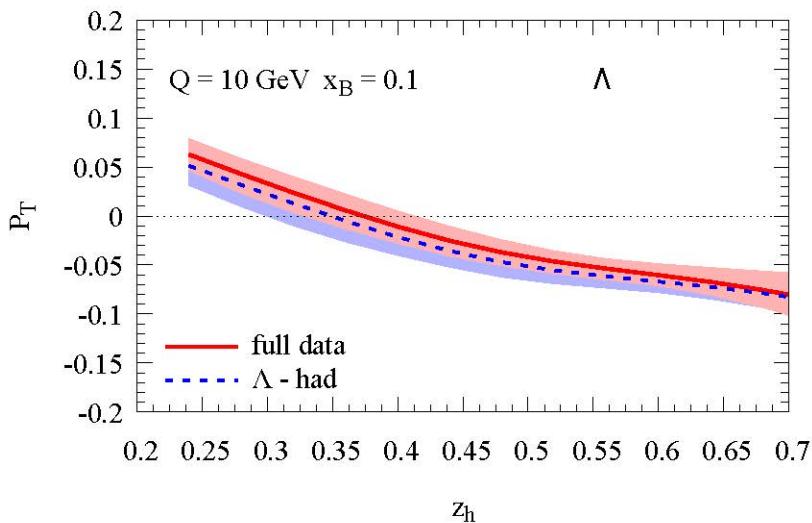
TRANSVERSE Λ POLARIZATION IN SIDIS

$$P_T(x_B, z_h) = \frac{\sqrt{2e\pi}}{2M_p} \frac{\langle p_\perp^2 \rangle_{\text{pol}}^2}{\langle p_\perp^2 \rangle} \frac{1}{\sqrt{\langle p_\perp^2 \rangle_{\text{pol}} + \xi_p^2 \langle k_\perp^2 \rangle}} \times \frac{\sum_q e_q^2 f_{q/p}(x_B) \Delta D_{\Lambda^\uparrow/q}(z_h)}{\sum_q e_q^2 f_{q/p}(x_B) D_{\Lambda/q}(z_h)}$$

Transverse w.r.t.
the target- Λ plane

$$\xi_p = z_h \left(1 - \frac{m_h^2}{z_h^2 Q^2} \frac{x_B}{1 - x_B} \right)$$

EIC kinematics



CONCLUDING REMARKS

- ❑ General helicity formalism for $e^+e^- \rightarrow h_1 h_2 + X$ AND LT quark (and gluon) TMD-FFs for spin-1/2 hadrons
- ❑ Fit of Belle e^+e^- data on transverse Λ polarization and extraction of the polarizing FF
 - ❑ Polarizing FF
 - ✓ three different valence polFFs (up, down and strange) + sea
 - ✓ relative sign between up and down polFFs
- ❑ SIDIS measurements @EIC important to complement these findings and to check the universality of the polFF



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NEXT

- ❑ TMD analysis of inclusive production and role of charm and isospin symmetry (ongoing)

THANKS for the ATTENTION



BACK-UP SLIDES



FIT OF BELLE DATA

Data selection:

- $\Lambda + \pi/K$: $z_{\pi.K} = [0.5 - 0.9]$ excluded $\rightarrow 96/128$ data points
- $\Lambda(\text{jet})$: $z_\Lambda = [0.5 - 0.9]$ excluded $\rightarrow 24/32$ data points

- **Unpol. FFs:** DSS07 for π/K , AKK08 for Λ

$$\langle p_\perp^2 \rangle = 0.2 \text{ GeV}^2$$

$$D_{\Lambda/\bar{q}}(z_p) = (1 - z_p) D_{\Lambda/q}(z_p)$$

flavour separation

- **Polarizing FF**

$$\Delta D_{\Lambda^\uparrow/q}(z) = N_q z^{a_q} (1 - z)^{b_q} \frac{(a_q + b_q)^{(a_q + b_q)}}{a_q^{a_q} b_q^{b_q}} D_{\Lambda/q}(z)$$

$$|N_q| \leq 1 \quad \text{Positivity bound}$$

- ✓ No TMD evolution (fixed-scale analysis)



CORRESPONDENCE WITH AMSTERDAM NOTATION

Projection of the
correlator $\Delta(z, \mathbf{k}_T)$

$$(-z \mathbf{k}_T) = \mathbf{p}_\perp$$

$$D_1(z, p_\perp) = D_{++}^{++} + D_{--}^{++} = D_{h/q}$$

$$\frac{p_\perp}{zM_h} D_{1T}^\perp(z, p_\perp) = -2 \operatorname{Im} D_{++}^{+-} = \Delta D_{S_Y/q}^h = \Delta D_{h^\uparrow/q}$$

$$G_{1L}(z, p_\perp) = D_{++}^{++} - D_{--}^{++} = \Delta D_{S_Z/s_L}^{h/q}$$

$$\frac{p_\perp}{zM_h} G_{1T}(z, p_\perp) = 2 \operatorname{Re} D_{+-}^{++} = \Delta D_{S_X/s_L}^{h/q}$$

$$\frac{p_\perp}{zM_h} H_{1L}^\perp(z, p_\perp) = 2 \operatorname{Re} D_{++}^{+-} = \Delta D_{S_Z/s_T}^{h/q}$$

$$\frac{2p_\perp}{zM_h} H_1^\perp(z, p_\perp) = 4 \operatorname{Im} D_{+-}^{++} = \Delta D_{h/q^\uparrow}$$

$$H_1(z, p_\perp) = D_{+-}^{+-}$$

$$\frac{p_\perp^2}{z^2 M_h^2} H_{1T}^\perp(z, p_\perp) = 2 D_{-+}^{+-}$$

Polarizing FF

Collins FF

Extension of the
Trento Convention results
Bacchetta, UD, Diehl, Miller 2004

