Lattice results on twist-3 parton distributions

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In collaboration with

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- Hadron structure described by a tower of partonic distributions

$$f = f^{(0)} + \frac{1}{Q}f^{(1)} + \frac{1}{Q^2}f^{(2)} + \dots$$

twist-2   twist-3

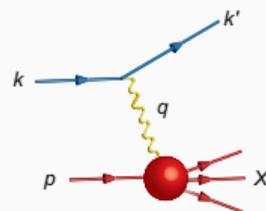

- f: classified according to their **twist** - order in $1/Q$ at which $f^{(i)}$ appear in factorization theorem of structure functions

$$\sigma_{DIS}(x, Q^2) = \sum_i \left[H_{DIS}^i \otimes f_i \right](x, Q^2)$$

Twist-2 case:
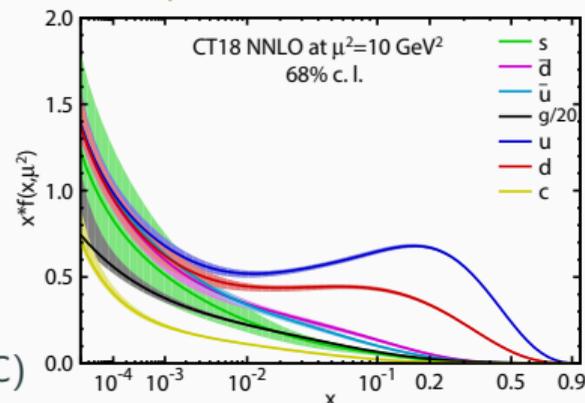* probabilistic interpretation
* some twist-2 PDFs are very well known

Twist-3 case:
* information about *qgq* correlations
* kinematically suppressed
* challenging to probe experimentally (JLAB 12 GeV upgrade, EIC)

Unpolarized twist-2 PDFs

[Kovarik et al., 2019, Rev.Mod.Phys. 92 (2020) 4, 045003]

- Isovector proton twist-3 PDFs: $g_T^{u-d}(x), h_L^{u-d}(x)$

Polarization	PDF	Operator
U	$e(x)$	1
L	$h_L(x)$	σ^{ij}
T	$g_T(x)$	$\gamma_5 \gamma_i$

* $g_T^{u-d}(x)$: ⊥ force acting on the active quark in DIS off a transversely polarized nucleon after it has absorbed the virtual photon [M.Burkardt, PRD 88 (2013) 114502]

* $h_L^{u-d}(x)$: chiral odd, accessible in e.g. Drell-Yan process [R. Jaffe, PRL 67 (1991) 552-555], proton-proton collisions [Y. Koike et al., PLB 759 (2016) 75,], . . .

Interesting aspects to address:

1. Are twist-3 effects sizeable?

2. Wandzura-Wilczek approximation

[S. Wandzura and F. Wilczek, Phys. Lett.72B, 195(1977)][R.L. Jaffe and X. Ji, PRL 67 (1991) 552-555]

E.g.: $g_T(x) \stackrel{?}{=} \int_x^1 dy \, \frac{g_1(y)}{y}$ $g_1(y)$: helicity twist-2 PDF

- Twist-3 GPD $\tilde{G}_2(x, \xi, t)$

$$[\tilde{H} + \tilde{G}_2](x, \xi = 0, Q^2 = 0) = g_T(x)$$

[D. Kiptily and M. Polyakov, Eur. Phys. J. C37 (2004) 105]

- PDFs computed through purely space-like correlation functions

$$\tilde{\Phi}^{\Gamma}(x, \mu, P_3) = \int \frac{dz}{4\pi} e^{-ixP_3 z} \langle N(P_3)|\underbrace{\bar{\psi}(0)\Gamma W(0,z)\psi(z)|N(P_3)}\rangle_\mu$$

<div align="center">
matter elements
of fast moving nucleons
</div>

* μ: renormalization scale   * $z = (0,0,0,z)$ * $\vec{P} = (0,0,P_3)$: proton boost

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- For sufficiently large P_3 $(P_3 >> M)$, quasi-PDFs matched to light-cone PDFs

  <div align="right">LaMET</div>

  $$\tilde{\Phi}^{\Gamma}(x, \mu, P_3) = \int_{-1}^{+1} \frac{dy}{|y|} C\left(\frac{x}{y}, \mu\right) \Phi(y, \mu) + \mathcal{O}\left(\frac{1}{P_3^2}\right)$$

  * <u>twist-2 case</u>: factorization to all orders
    many works between 2013-2020

  * <u>twist-3 case</u>: proved factorization to 1-loop order

    [S.Bhattacharya, K.Cichy, M.Constantinou, A.Metz, AS, F. Steffens,
    Phys.Rev.D 102 (2020) 3, 034005, Phys.Rev.D 102 (2020) 114025]

    [V.M. Braun et al., 2021, JHEP 05 (2021) 086]

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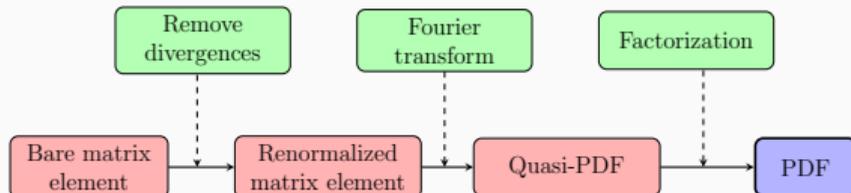 [V.M. Braun et al., 2021, JHEP 05 (2021) 086]

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→ S.Bhattacharya's talk

→ matching kernel

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  matching kernel

- Road map

3 / 13

Lattice setup

- Configurations of $N_f = 2 + 1 + 1$ flavors & clover term [ETM collaboration]

N_f	$L^3 \times T$	lattice spacing a	m_π	$m_\pi L$
4	$32^3 \times 64$	0.093 fm	270 MeV	4

Lattice setup

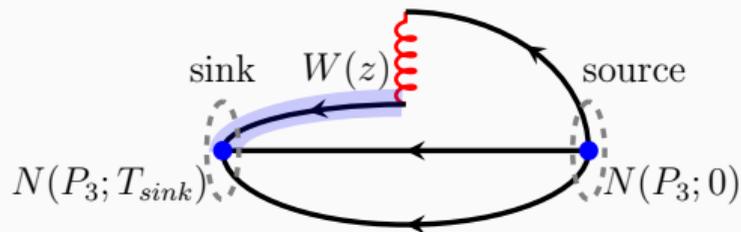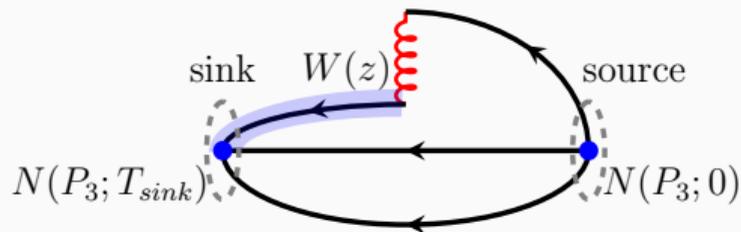- Configurations of $N_f = 2 + 1 + 1$ flavors & clover term [ETM collaboration]

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* $\Gamma = \gamma_x \gamma_5, \gamma_y \gamma_5$ for $g_T(x)$
* $\Gamma = \sigma^{12}$ for $h_L(x)$

sink $W(z)$ source

$N(P_3; T_{sink})$ $N(P_3; 0)$

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Lattice techniques:

* sequential inversions for all-to-all propagator (here $T_{sink} = 12a \simeq 1.12$ fm)
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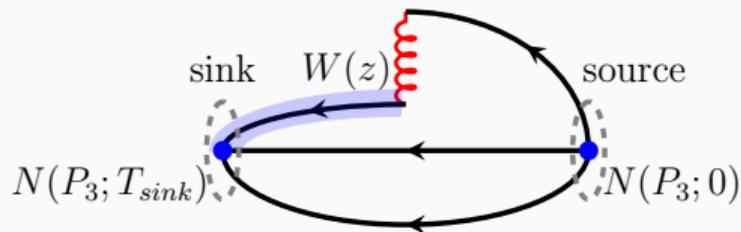
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- Momentum dependence of the lattice PDFs studied at three proton boosts

P_3[GeV]	N_{conf}	N_{meas}
0.83	194	1552
1.25	731	11696
1.67	1644	105216
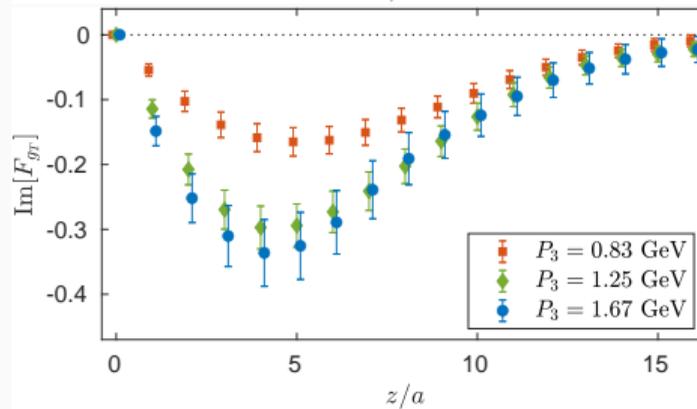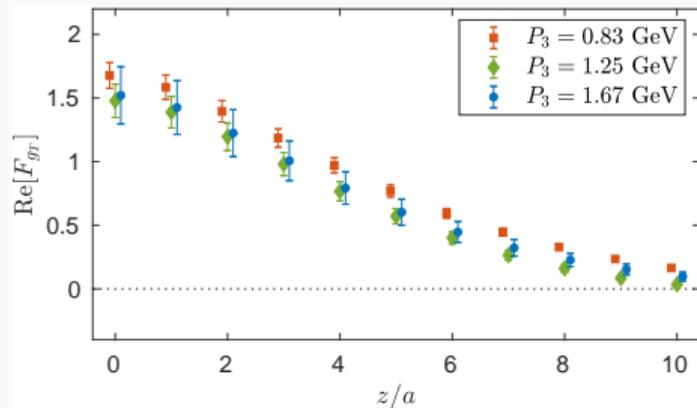
Results: $g_T(x)$

Bare matrix elements

Bare matrix elements
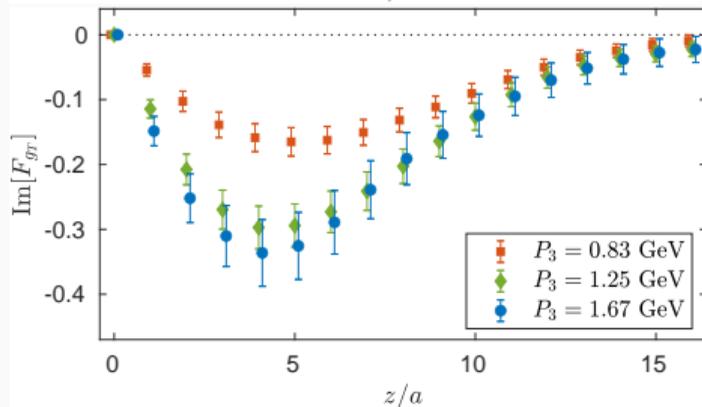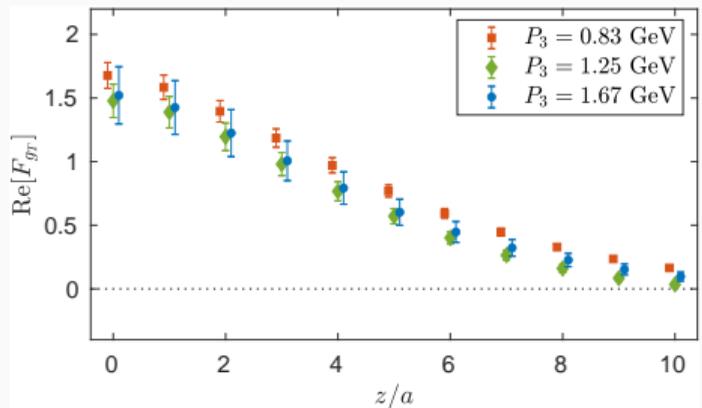
Reconstructed distributions
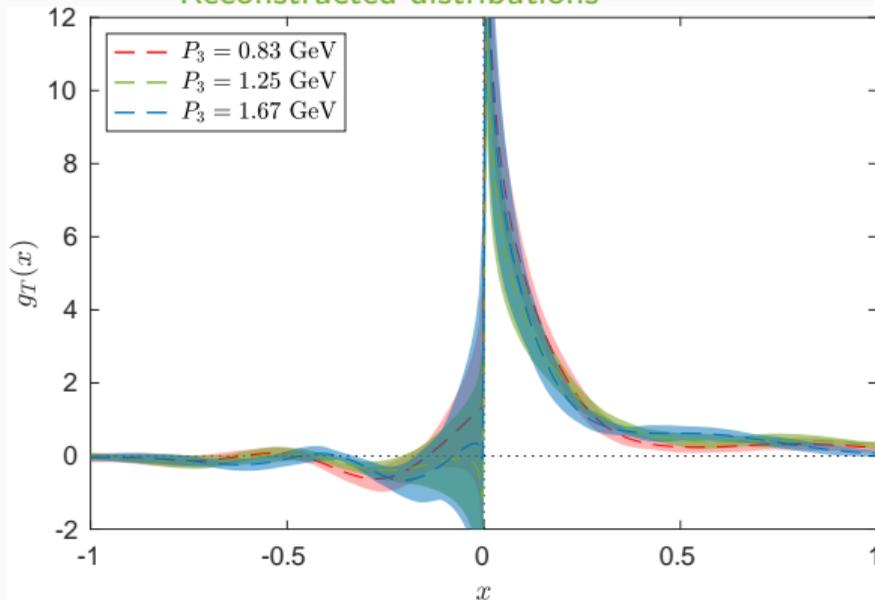
(Upon renormalization, F.T and matching)

[S.Bhattacharya, K.Cichy, M. Constantinou,A. Metz, AS, F. Steffens, Phys.Rev.D 102 (2020) 11]

# Comparison between $g_1^{u-d}(x)$ and $g_T^{u-d}(x)$

[S.Bhattacharya, K.Cichy, M. Constantinou, A. Metz, AS, F. Steffens, Phys.Rev.D 102 (2020) 11]

$P_3 = 1.67$ GeV
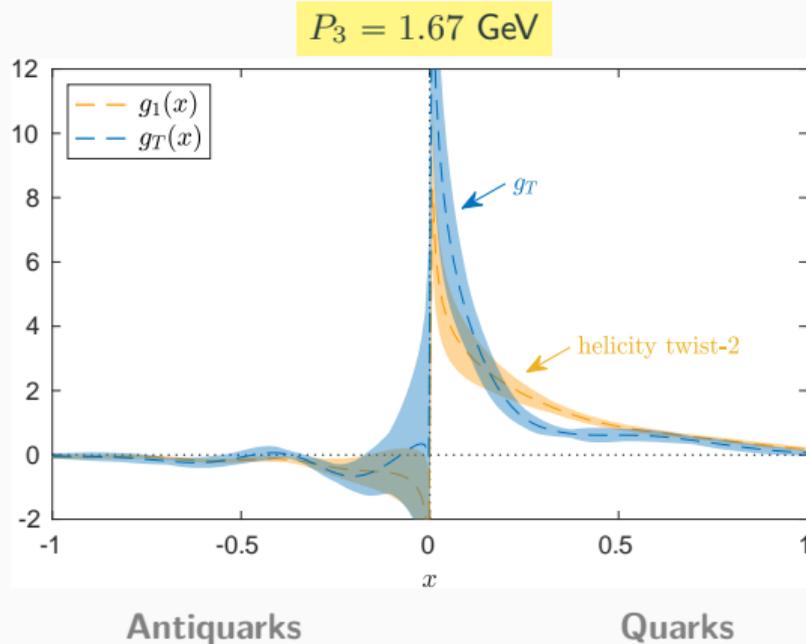
- $g_1(x)$ extracted on the same gauge ensemble

- g_1 and g_T compatible for antiquarks, $\bar{u} - \bar{d}$

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  $\rightarrow$ twist-3 contributions may be sizable

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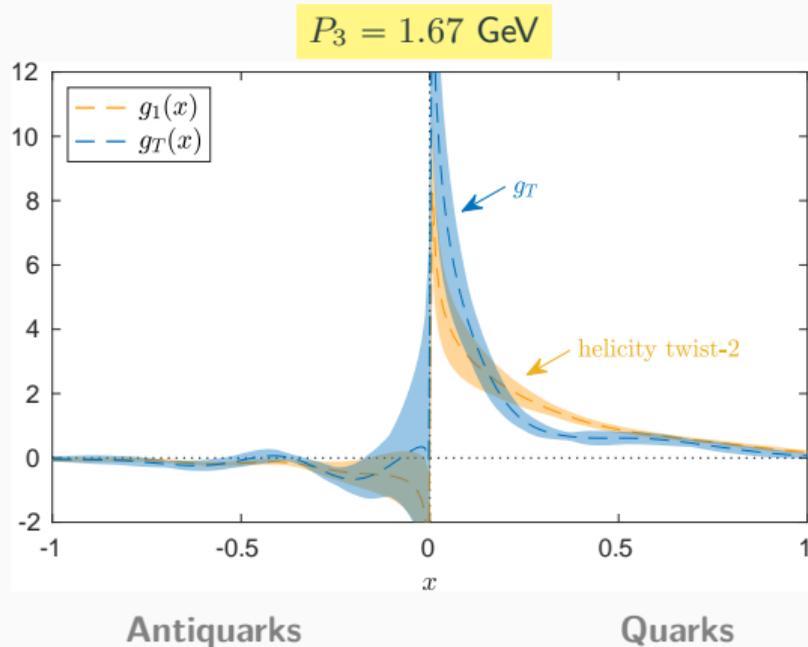
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Burkhardt-Cottingham sum rule:
[H. Burkhardt and W. N. Cottingham, Annals Phys.56,453 (1970)]

$$\int_{-1}^{1} dx\, g_1(x) = \int_{-1}^{1} dx\, g_T(x)$$

(lattice) $\int_{-1}^{1} dx\, g_1(x) - \int_{-1}^{1} dx\, g_T(x) = 0.01(20)$

Wandzura-Wilczek approximation for g_T^{u-d}

- Wandzura-Wilczek relation:

$$g_T(x) = \int_x^1 dy \, \frac{g_1(y)}{y} + g_T^{twist-3}(x) \quad \text{(on the light cone)}$$

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- Approximation: $\boxed{g_T(x) \stackrel{?}{=} \int_x^1 dy \, \frac{g_1(y)}{y} = g_T^{WW}(x)}$

$$d_2 = \int dx \, 3x^2[g_T(x) - g_T^{WW}(x)] \sim \mathcal{O}(10^{-3})$$

[D. Flay et al., PRD 94, 5 (2016) 052003]

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- agreement between g_T and g_T^{WW} for $x \lesssim 0.5$ within uncertainties
- still violations are possible up to $30 - 40\%$
- similar violations are observed in experimental analysis at JLab [A. Accardi et al., JHEP11 (2009) 093]

$P_3 = 1.67$ GeV

$g_T^{WW}(x)$ lattice
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$g_T^{WW}(x)$ NNPDF1.1pol
$g_T^{WW}(x)$ JAM17
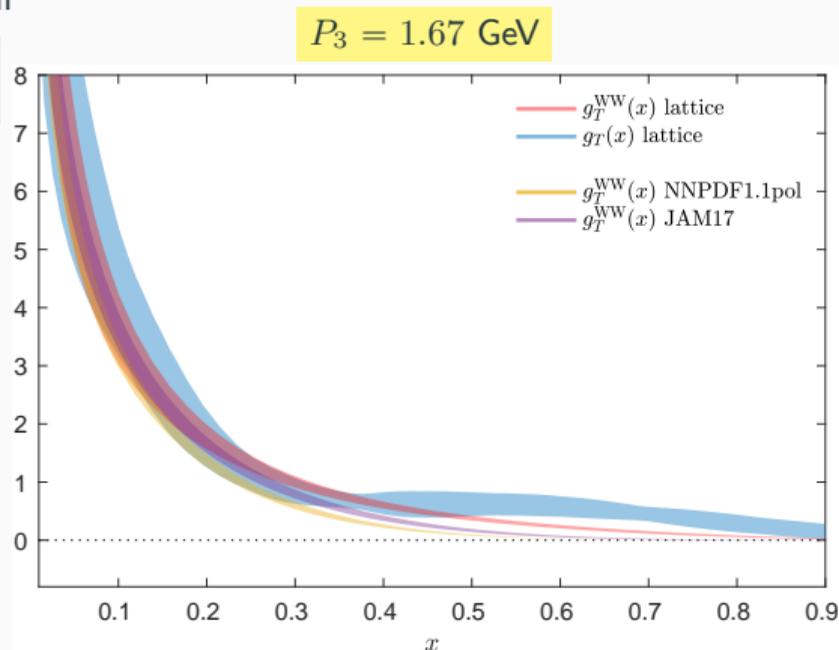
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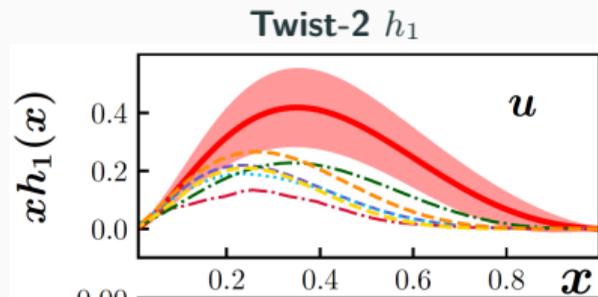**Beyond $g_T(x)$... The $h_L(x)$ case**
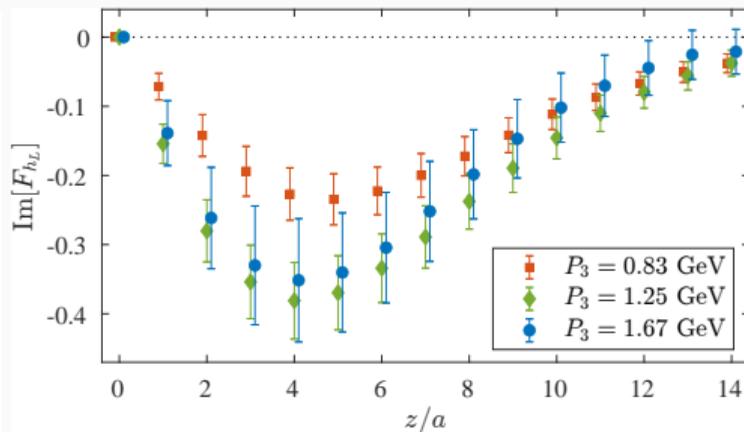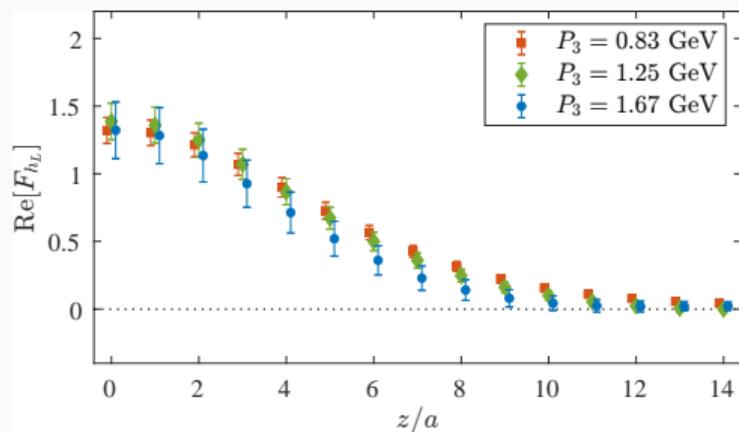
Extraction of $h_L(x)$ PDF

- $h_L(x)$ contains the leading twist-2 h_1 transversity PDF
- No experimental data on $h_L(x)$
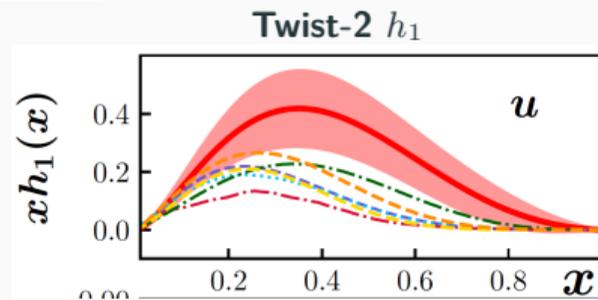
Twist-2 h_1

[JAM, Phys.Rev.D 102 (2020) 5, 054002]

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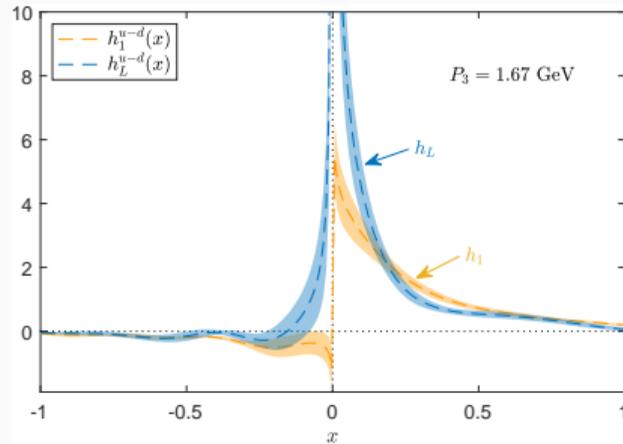- No experimental data on $h_L(x)$

Twist-2 h_1

[JAM, Phys.Rev.D 102 (2020) 5, 054002]

h_L from quasi-PDFs $\longrightarrow$ $F_{h_L}(P_3, z) = \langle N(P_3) | \bar{\psi}(0) \sigma^{12} \tau_3 W(0, z) \psi(z) | N(P_3) \rangle$

- h_1^{u-d} dominant only for $0.3 \lesssim x \lesssim 0.5$

Results for $h_L(x)$ and twist-2 $h_1(x)$
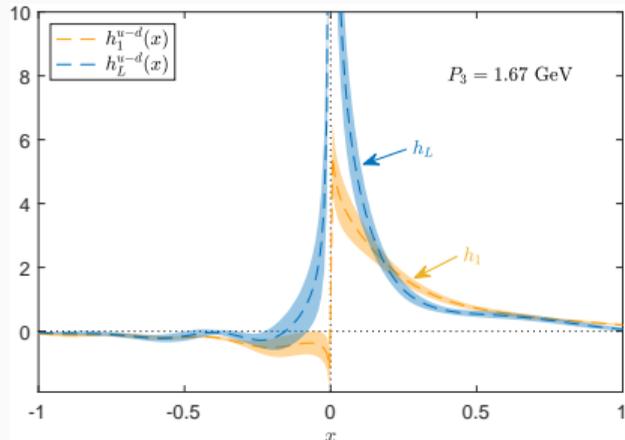

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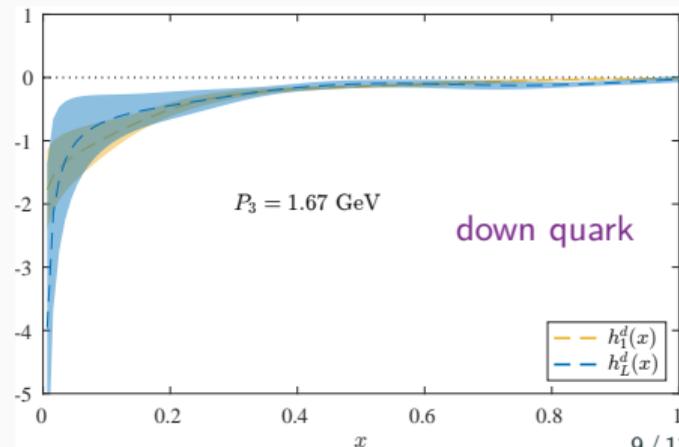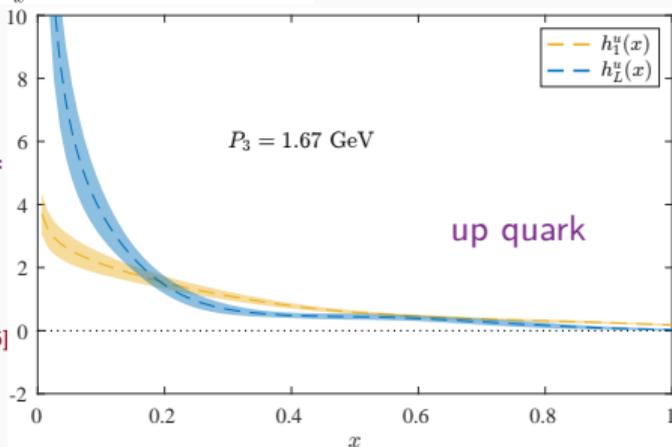
- h_1^{u-d} dominant only for $0.3 \lesssim x \lesssim 0.5$

- up-quark especially dominant for $x < 0.5$

- differences between h_L and h_1 mostly come from the up quark

What about for the single quark flavors?*

*quark loops here neglected - found very small for $h_1(x)$
[C.Alexandrou at al.,arXiv:2106.16065]

$$h_L(x) = 2x \int_x^1 \frac{h_1(y)}{y^2} dy + h_L^{twist-3}(x)$$

[R.L. Jaffe and Xiangdong Ji, Nucl. Phys. B 375 (1992) 527—560]

Test: $\boxed{\; \mathsf{h}_L(x) \overset{?}{=} 2x \int_x^1 dy \, \frac{h_1(y)}{y^2} \equiv h_L^{WW}(x) \;}$

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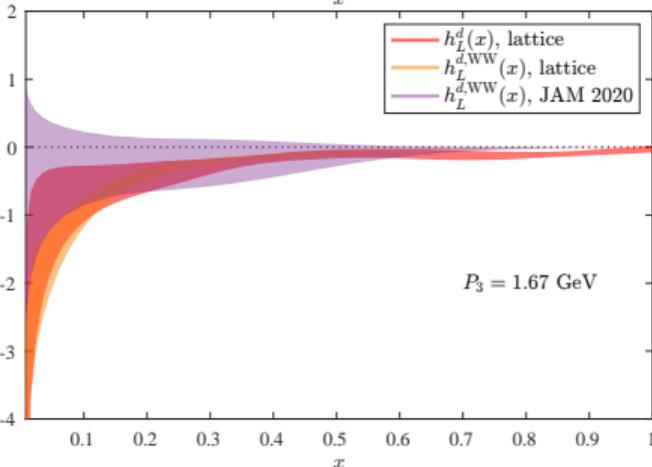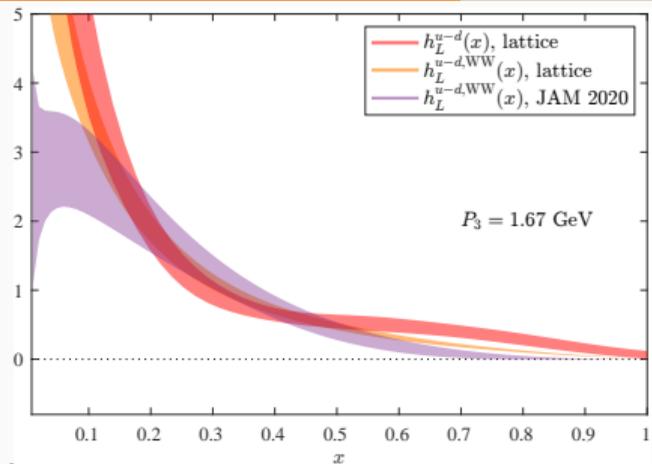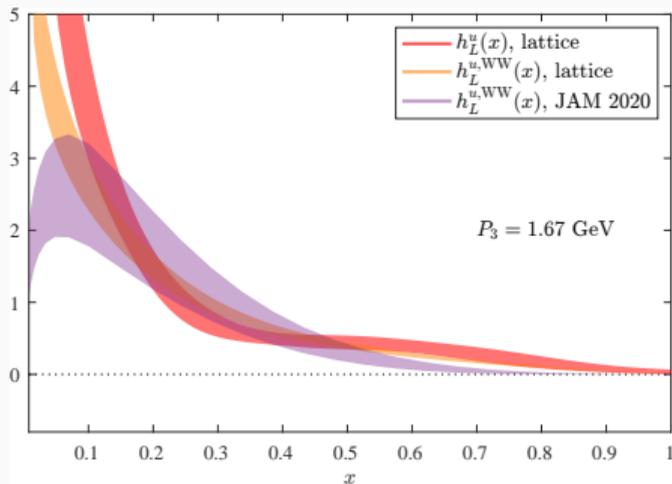Test: $\boxed{h_L(x) \overset{?}{=} 2x \int_x^1 dy \frac{h_1(y)}{y^2} \equiv h_L^{WW}(x)}$

- h_L seems to be largely determined by its twist-2 counterpart within current uncertainties

From twist-3 PDFs...to twist-3 GPDs

Twist-3 GPDs (axial-vector case)

Theoretical framework

(DVCS at twist-3 level)

$$\tilde{F}_q^\mu(x,\xi,\Delta) = \int_{-\infty}^{\infty} \frac{d\lambda}{2\pi} e^{-i\lambda x} \langle p'|\bar{q}(\tfrac{\lambda}{2}n)\gamma^\mu\gamma_5 W(\tfrac{\lambda}{2}n,-\tfrac{\lambda}{2}n)q(-\tfrac{\lambda}{2}n)|p\rangle$$

$$\tilde{F}^\mu = P^\mu \frac{\tilde{h}^+}{P^+}\tilde{H} + P^\mu\frac{\tilde{e}^+}{P^+}\tilde{E} + \Delta_\perp^\mu \frac{\tilde{b}}{2m}(\tilde{E} + \boxed{\tilde{G}_1}) + \tilde{h}_\perp^\mu(\tilde{H} + \boxed{\tilde{G}_2}) + \Delta_\perp^\mu\frac{\tilde{h}^+}{P^+}\boxed{\tilde{G}_3} + \tilde{\Delta}_\perp^\mu\frac{h^+}{P^+}\boxed{\tilde{G}_4}$$

$$\Delta = p' - p, P = \frac{p'+p}{2}$$
$$\xi = -\frac{\Delta^+}{2P^+}, P^+ = \frac{P^0 + P^3}{2}$$

4 twist-3 GPDs to disentangle

[D. Kiptily and M. Polyakov, Eur. Phys. J. C37 (2004) 105]
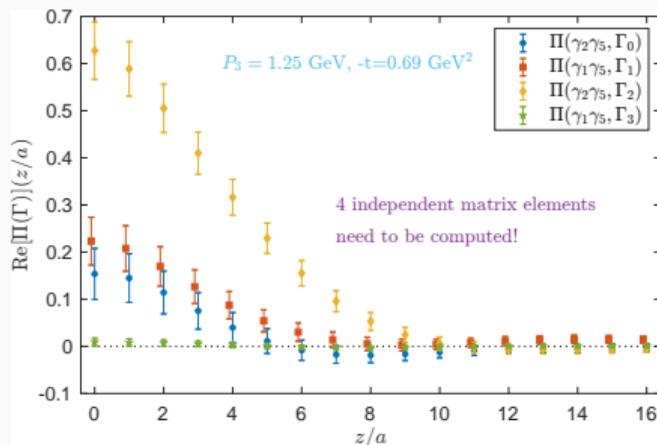
[F. Aslan et a., Phys. Rev. D 98, 014038 (2018)]

Lattice setup

Higher computational cost than PDFs!

- Breit frame
- Same ensemble as twist-3 PDFs

P_3[GeV]	$\vec{Q} \times \frac{L}{2\pi}$	$-t$[GeV^2]	ξ	N_{meas}
1.25	(2,0,0)	0.69	0	4288
1.25	(2,2,0)	1.08	0	8704

many more Q-values will be added [in progress]

$P_3 = 1.25$ GeV, -t=0.69 GeV^2

4 independent matrix elements need to be computed!

x-dependence of twist-3 GPDs ($\xi = 0$)
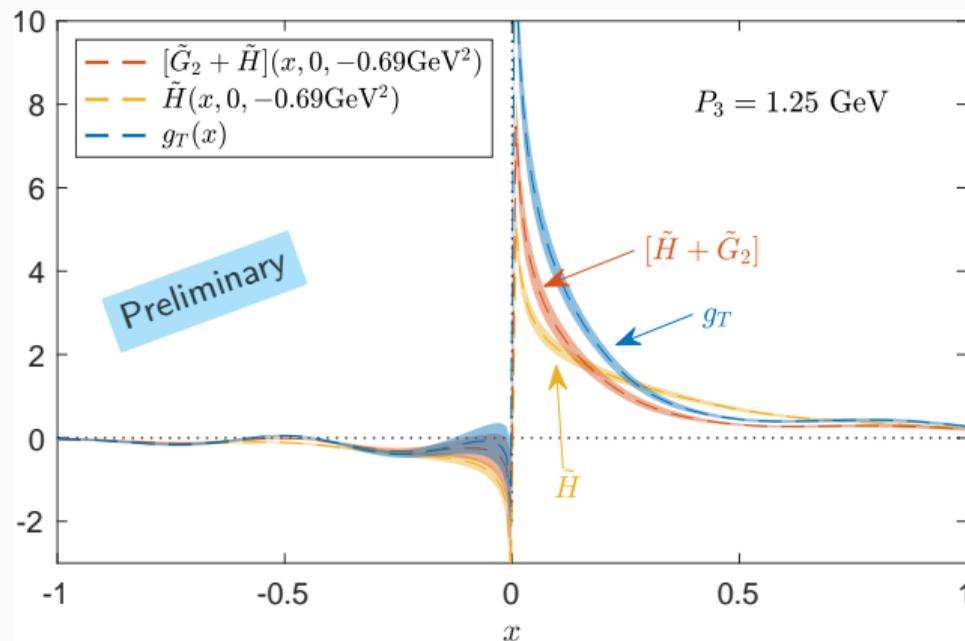
[S. Bhattacharya, K. Cichy, M. Constantinou, J. Dodson, A. Metz, AS, F. Steffens, (DIS2021) arXiv:2107.12818]

- g_T dominant as expected ($[\tilde{H} + \tilde{G}_2](x, 0, 0) = g_T(x)$)
- $[\tilde{H} + \tilde{G}_2]$ enhanced for small-x with respect to $\tilde{H}$
  (similar behavior as for g_T vs. g_1)

Summary & Outlook

Study of twist-3 PDFs & GPDs

- Lattice investigations can be pursued using quasi-distribution formalism (and also other approaches)

- For the first time, a qualitative comparison between twist-2 and twist-3 PDFs can be made

  $\rightarrow$ test of the Wandzura-Wilczek approximation for g_T and h_L for $x > 0$

- Several challenges in front of us:

  * can we disentangle the twist-2 and twist-3 parts of g_T and h_L?
    $\rightarrow$ calculation of quark-gluon-quark matrix elements
    $\rightarrow$ to what extent numerical results are affected?
    [V.M. Braun, Y. Ji , A. Vladimirov, JHEP 05 (2021) 086]

  * Systematics on the lattice: finite lattice spacing and volume effects, pion mass dependence, . . .

- Lots of work, but of great impact in phenomenology in the next few years

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Thank you!