

# Gluon TMDs and inelastic $J/\psi$ lepton production at the EIC



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Paper in preparation



Fondazione  
di Sardegna



# Overview

- Quarkonium as a Lab for Gluon TMDs and the gluon Sivers function
- $J/\psi$  leptoproduction at small-intermediate transverse momentum
- Interplay TMD [CGI-GPM] and NRQCD approaches: Challenging but intriguing!
- Comparison with H1 (unp. cross sect.) and COMPASS (Sivers asymmetry) data
- Estimates for EIC kinematics and concluding remarks

# Gluon TMDs at leading twist

gluon pol.

|   | U                  | Circularly | Linearly                  |
|---|--------------------|------------|---------------------------|
| U | $f_1^g$            |            | $h_1^{\perp g}$           |
| L |                    | $g_{1L}^g$ | $h_{1L}^{\perp g}$        |
| T | $f_{1T}^{\perp g}$ | $g_{1T}^g$ | $h_1^g, h_{1T}^{\perp g}$ |

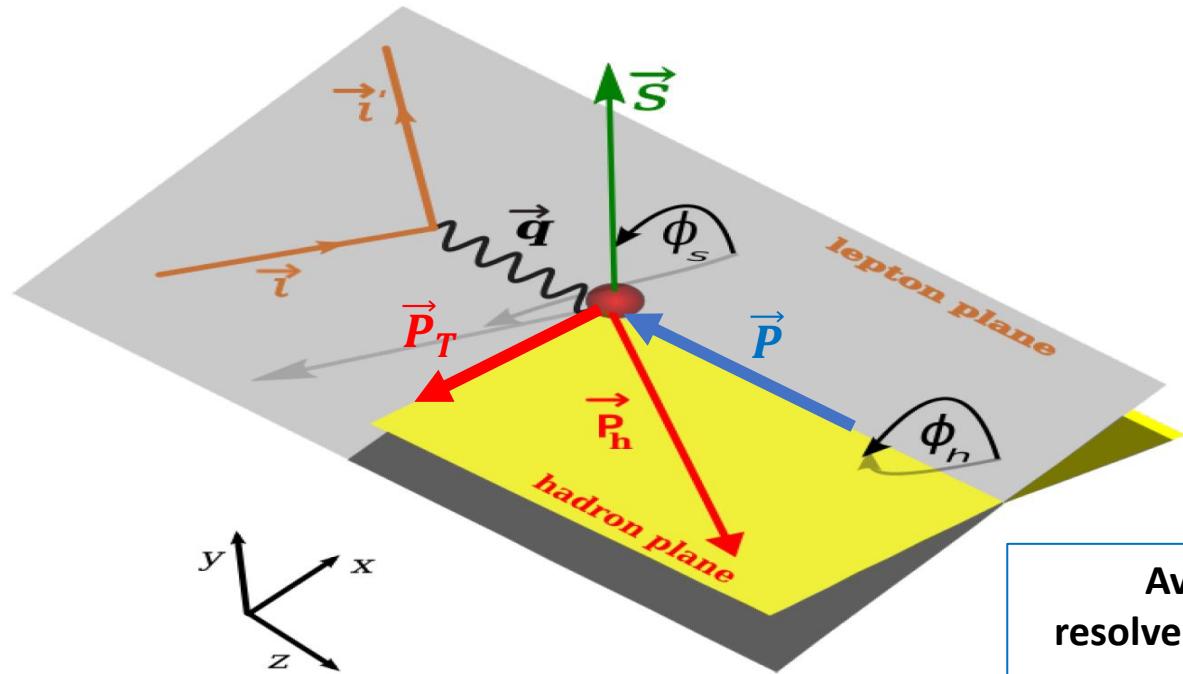
$$\begin{array}{cc} \xrightarrow{} & \xrightarrow{} \\ f_{1T}^{\perp g(d)} & f_{1T}^{\perp g(f)} \end{array}$$

Mulders & Rodrigues PRD 63 094021 (2001)

- Eight LT Gluon TMDs for spin  $\frac{1}{2}$  hadrons
- Analogous to quarks but linear vs transverse pol.
- More complicate gauge link structure
- Two independent gluon Sivers functions
- Different properties and features
- Different relevance in different processes
- Phenomenology needs accurate fits of as much as possible observables

Boer, Lorcé, Pisano, Zhou  
Adv. High Energy Phys. 2015 371396 (2015)

# Quarkonium lepto production: $\ell p^\uparrow \rightarrow \ell' J/\psi + X$



$$s = (\ell + P)^2 \quad W_{\gamma p}^2 = (q + P)^2 \quad Q^2 = -q^2$$

$$x_B = \frac{Q^2}{2P \cdot q} \quad y = \frac{P \cdot q}{P \cdot \ell} \quad z = \frac{P \cdot P_h}{P \cdot q}$$

$$\hat{s} = (k_a + q)^2 = x_a y S - Q^2 \left( 1 + \frac{k_{\perp a}^2}{x_a y S} \right),$$

$$\hat{t} = (q - P_h)^2 = (z - 1) Q^2 + M^2 - \frac{M^2 + P_T^2}{z},$$

$$\hat{u} = (k_a - P_h)^2 = M^2 - x_a y z S - \frac{k_{\perp a}^2}{x_a y S} \frac{M^2 + P_T^2}{z} + 2 k_{\perp a} P_T \cos(\phi_a - \phi_h)$$

$0.3 < z < 0.9$

Avoid resolved-photon region

$P_T \leq 8 \text{ GeV}$

Avoid collinear divergences at  $z=1$

$P_T \leq \text{few GeV: TMD regime}$   
 $P_T \text{ up to } 8 \text{ GeV: Matches LO collinear scheme}$

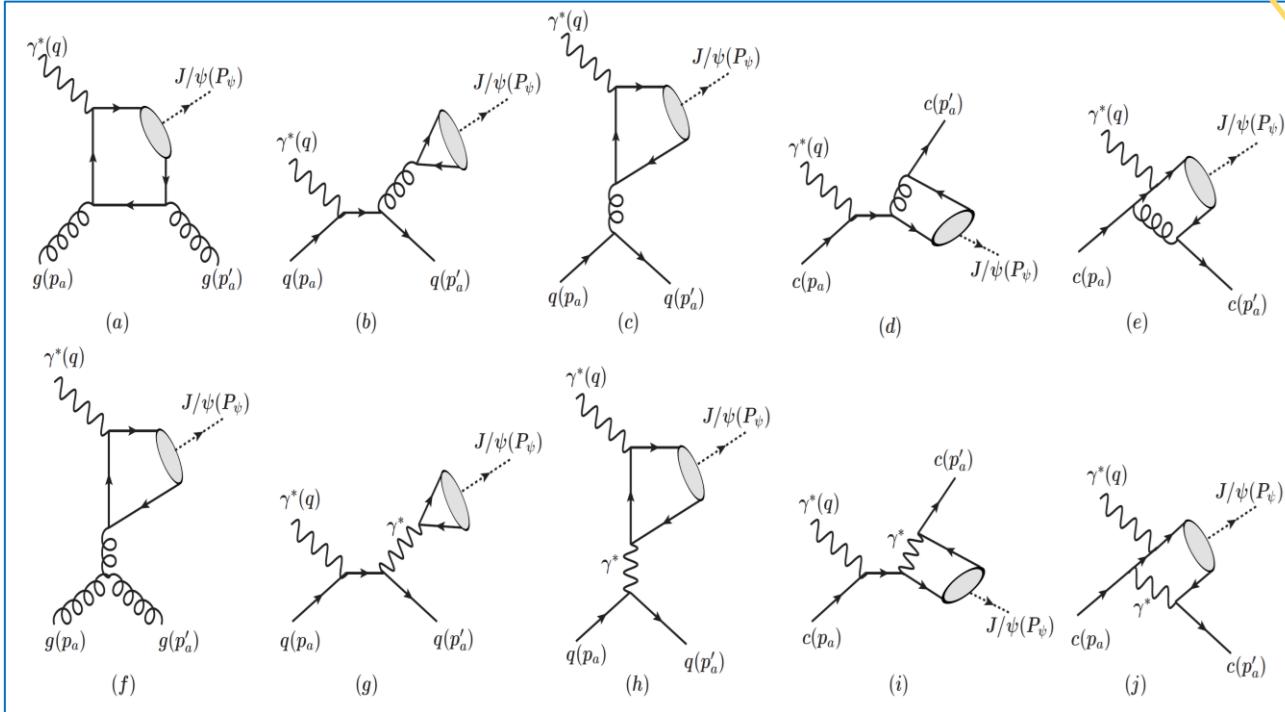
For  $\ell p^\uparrow \rightarrow \ell' J/\psi + X$  at  $z = 1$  see:  
A. Mukherjee & S. Rajesh EPJC 77 854 (2017)  
A. Bacchetta, Boer, Pisano, Taels EPJC 80 72 (2020)

# Unpolarized cross section in a TMD GPM + NRQCD scheme

$$\frac{d\sigma}{dQ^2 dy d^2 \mathbf{P}_T dz} = \frac{1}{2S} \frac{2}{(4\pi)^4 z} \sum_a \int \frac{dx_a}{x_a} d^2 \mathbf{k}_{\perp a} \delta(\hat{s} + \hat{t} + \hat{u} - M^2 + Q^2)$$

$$\times \sum_n \frac{1}{Q^4} f_{a/p}(x_a, k_{\perp a}) L^{\mu\nu} H_{\mu\nu}^{a,U}[n] \langle 0 | \mathcal{O}^{J/\psi}(n) | 0 \rangle ,$$

**QCD @ $\alpha\alpha_s^2$  + QED @ $\alpha^3$  + intrinsic charm**

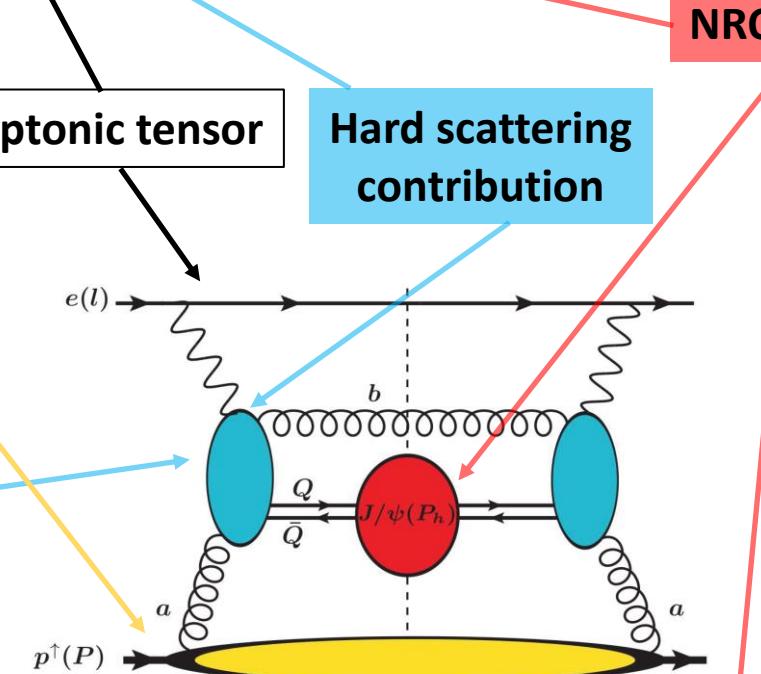


**TMD PDFs**

**Leptonic tensor**

**Hard scattering contribution**

**NRQCD LDMEs**



$$n = {}^3S_1^{(1,8)}, {}^1S_0^{(8)}, {}^3P_J^{(8)}, J = 0, 1, 2$$

Bodwin, Braaten, Lepage PRD 51 1125 (1995)

# Unpolarized cross section in a TMD GPM + NRQCD scheme

- We consider two sets of LDMEs:
  - Butenschoen & Kniehl PRD 84 051501 (2011) [BK11]
  - Sun, Yuan & Yuan PRD 88 054008 (2013) [SYY13] [ no CS]
  
- All partonic channels @ $\alpha\alpha_s^2$ :  $\gamma^* + g \rightarrow J/\psi + g$ ,  $\gamma^* + q(\bar{q}) \rightarrow J/\psi + q(\bar{q})$  [including intrinsic charm]
- Direct  $c(\bar{c})$  fragmentation into  $J/\psi$  negligible in the low  $P_T$  region
  
- QED contributions @ $\alpha^3$  also included
  
- Feed-down contributions:
  - $\psi(2S)$  included; LDME set: Sharma & Vitev PRC 87 044905 (2013)
  - $\chi_c$  and  $b$ -quark contributions negligible in the low  $P_T$  region

TABLE III: Numerical values of LDMEs for  $J/\psi$ .

|           | $\langle \mathcal{O}_1^{J/\psi}(^3S_1) \rangle$<br>GeV <sup>3</sup> | $\langle \mathcal{O}_8^{J/\psi}(^3S_1) \rangle$<br>$\times 10^{-2}$ GeV <sup>3</sup> | $\langle \mathcal{O}_8^{J/\psi}(^1S_0) \rangle$<br>$\times 10^{-2}$ GeV <sup>3</sup> | $\langle \mathcal{O}_8^{J/\psi}(^3P_0) \rangle$<br>$\times 10^{-2}$ GeV <sup>5</sup> |
|-----------|---|--|--|--|
| BK11 [5]  | 1.32  | +0.224   | +4.97  | -1.61  |
| SYY13 [6] |   | -0.93  | +14.23   | -1.75 $m_c^2$  |

TABLE IV: Numerical values of LDMEs for  $\psi(2S)$ .

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|---------|---|--|--|--|
| Ref.[7] | 0.76  | +0.33  | +0.8   | +0.8 $m_c^2$   |

Quarkonium overviews:

J.P. Lansberg Phys. Rep. 889 1 (2020)  
N. Brambilla et al EPJC 71 1534 (2011)

# Unpolarized cross section in a TMD GPM + NRQCD scheme

Unp. TMD

$$f(x_a, \mathbf{k}_{\perp a}^2, \mu) = f(x_a, \mu) \frac{1}{\pi \langle k_{\perp a}^2 \rangle} e^{-\mathbf{k}_{\perp a}^2 / \langle k_{\perp a}^2 \rangle}$$

$$\langle k_{\perp q}^2 \rangle = 0.25 \text{ GeV}^2 \quad \langle k_{\perp g}^2 \rangle = 1.0 \text{ GeV}^2 \quad \mu^2 = M^2 + P_T^2 \quad \text{collinear PDFs CTEQ-L1}$$

Anselmino, Boglione, D'Alesio  
 Kotzinian, FM, Prokudin  
 PRD 71 074006 (2005)

U. D'Alesio, FM, C. Pisano, S. Rajesh  
 EPJC 79 12 (2019)

## NRQCD LDMEs

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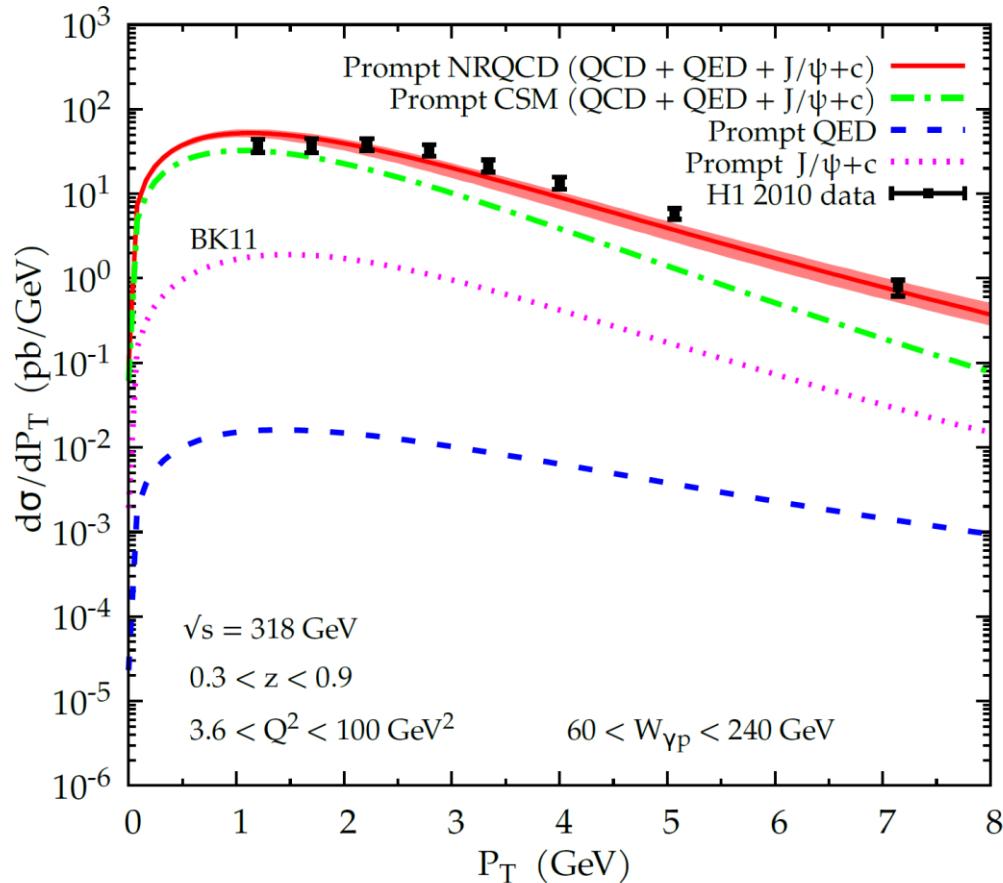
Butenschoen & Kniehl PRD 84 051501 (2011) [BK11]  
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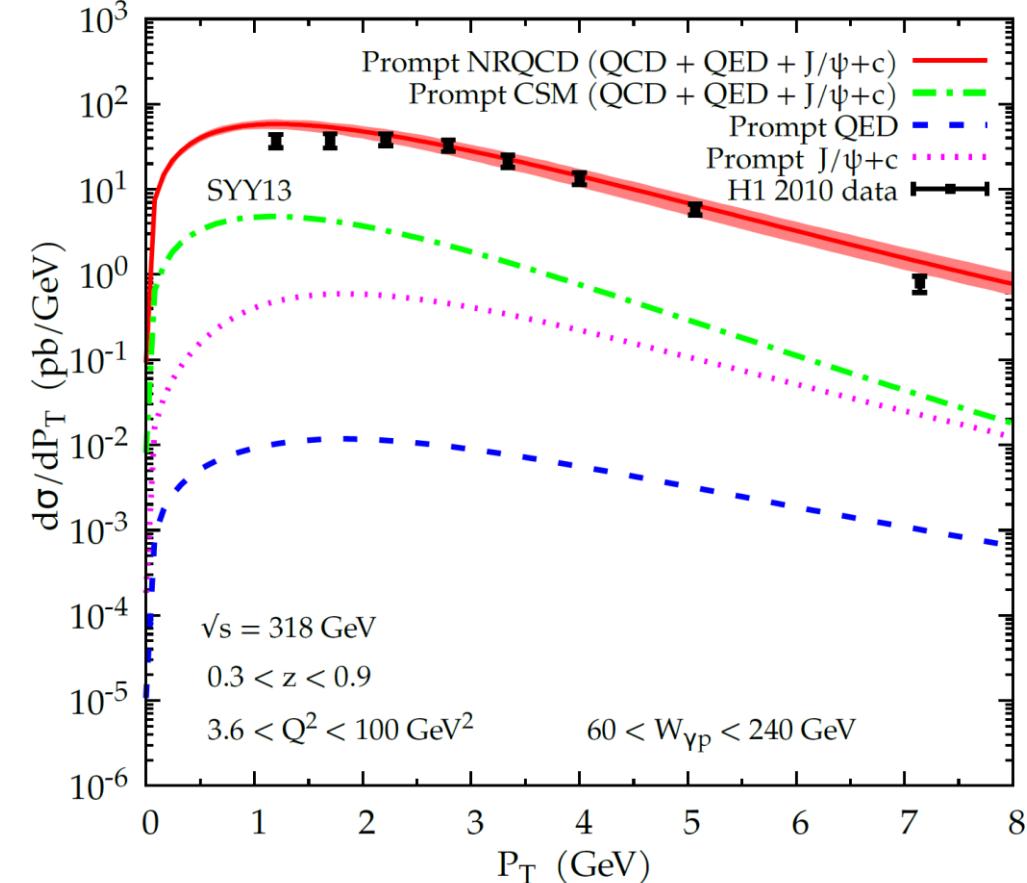
|         | $\langle \mathcal{O}_1^{\psi(2S)}(^3S_1) \rangle$<br>GeV <sup>3</sup> | $\langle \mathcal{O}_8^{\psi(2S)}(^3S_1) \rangle$<br>$\times 10^{-2}$ GeV <sup>3</sup> | $\langle \mathcal{O}_8^{\psi(2S)}(^1S_0) \rangle$<br>$\times 10^{-2}$ GeV <sup>3</sup> | $\langle \mathcal{O}_8^{\psi(2S)}(^3P_0) \rangle$<br>$\times 10^{-2}$ GeV <sup>5</sup> |
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Sharma & Vitev PRC 87 044905 (2013)

# Unpolarized cross section in a TMD GPM + NRQCD scheme vs. H1 data



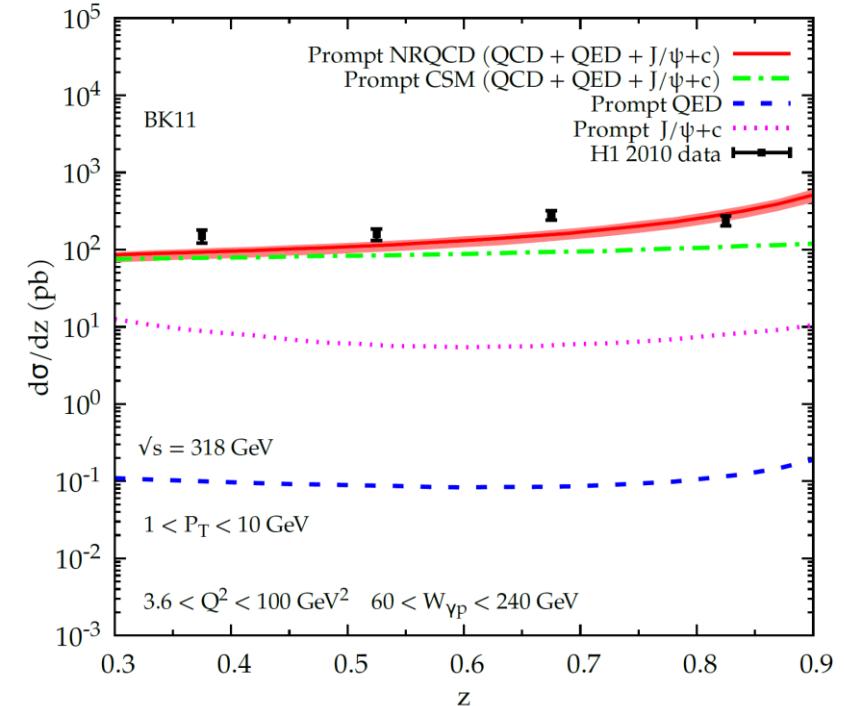
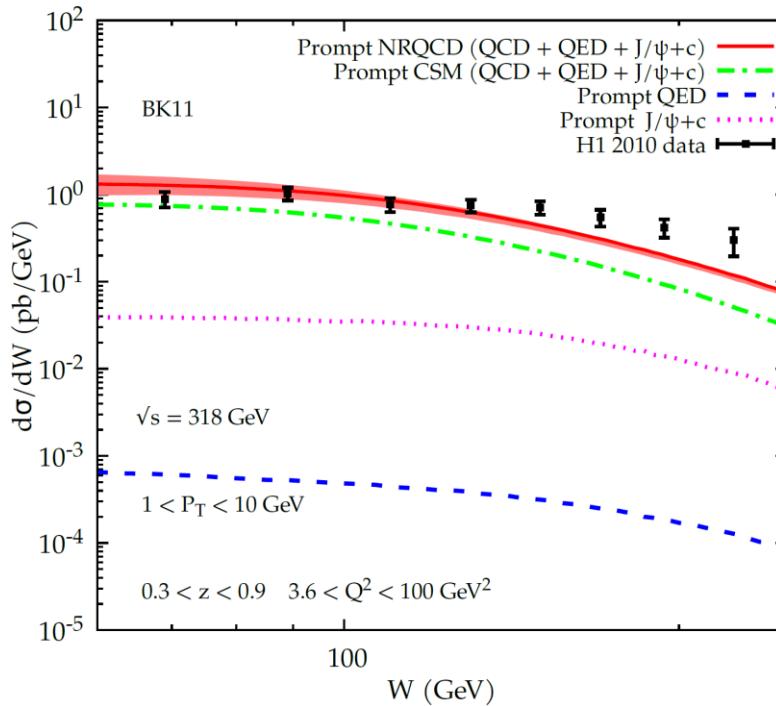
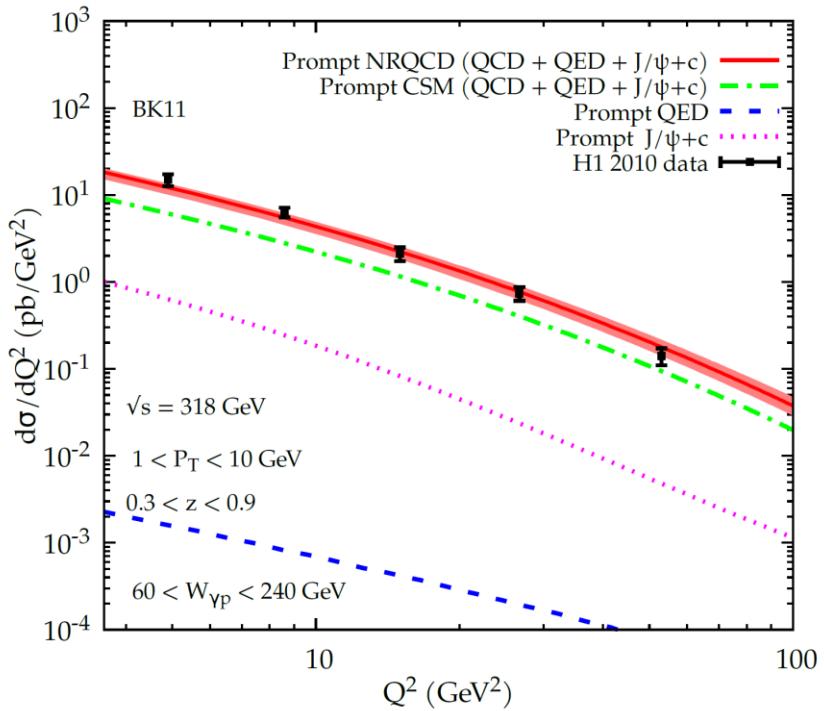
$\mu^2 = M^2 + P_T^2$       collinear PDFs CTEQ-L1



Data: H1 Collab. EPJC 68 401 (2010)

Uncertainty bands: only scale dependence  $\in [\mu/2 \div 2\mu]$

# Unpolarized cross section in a TMD GPM + NRQCD scheme vs. H1 data



$$\mu^2 = M^2 + P_T^2$$

collinear PDFs CTEQ-L1

Data: H1 Collab. EPJC 68 401 (2010)

Uncertainty bands: only scale dependence  $\in [\mu/2 \div 2\mu]$

# Sivers Asymmetry in $J/\psi$ lepto production

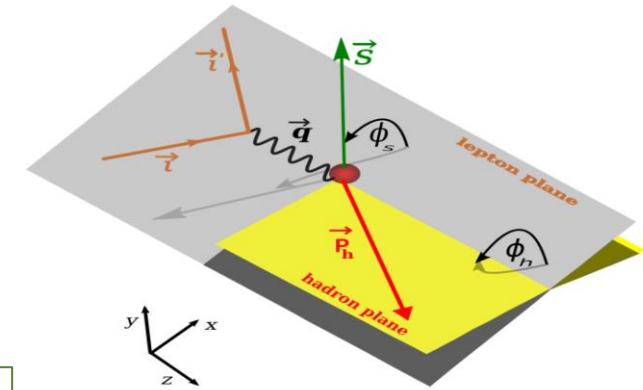
$$A_N^{\sin(\phi_h - \phi_s)} \equiv 2 \frac{\int d\phi_s d\phi_h \sin(\phi_h - \phi_s) (d\sigma^\uparrow - d\sigma^\downarrow)}{\int d\phi_s d\phi_h (d\sigma^\uparrow + d\sigma^\downarrow)} \equiv 2 \frac{\int d\phi_s d\phi_h \sin(\phi_h - \phi_s) d\Delta\sigma(\phi_s, \phi_h)}{\int d\phi_s d\phi_h 2d\sigma}$$

$$\begin{aligned} d\Delta\sigma^{\text{GPM}} = & \frac{1}{2S} \frac{2}{(4\pi)^4 z} \sum_a \int \frac{dx_a}{x_a} d^2 \mathbf{k}_{\perp a} \delta(\hat{s} + \hat{t} + \hat{u} - M^2 + Q^2) \left( -2 \frac{k_{\perp a}}{M_p} \right) \sin(\phi_s - \phi_a) \\ & \times \sum_n \frac{1}{Q^4} f_{1T}^{\perp a}(x_a, k_{\perp a}) L^{\mu\nu} H_{\mu\nu}^{a,U}[n] \langle 0 | \mathcal{O}^{J/\psi}(n) | 0 \rangle. \end{aligned}$$

$$\begin{aligned} d\Delta\sigma^{\text{CGI-GPM}} = & \frac{1}{2S} \frac{2}{(4\pi)^4 z} \int \frac{dx_a}{x_a} d^2 \mathbf{k}_{\perp a} \delta(\hat{s} + \hat{t} + \hat{u} - M^2 + Q^2) \left( -2 \frac{k_{\perp a}}{M_p} \right) \sin(\phi_s - \phi_a) \\ & \times \sum_n \frac{1}{Q^4} L^{\mu\nu} \left\{ \sum_q f_{1T}^{\perp q}(x_a, k_{\perp a}) H_{\mu\nu}^{q,\text{Inc}}[n] + f_{1T}^{\perp(f)g}(x_a, k_{\perp a}) H_{\mu\nu}^{g,\text{Inc}(f)}[n] \right\} \langle 0 | \mathcal{O}^{J/\psi}(n) | 0 \rangle \end{aligned}$$

$$H_{\mu\nu}^{a,\text{Inc}}[n] = \frac{C^{\text{Inc}}}{C_U} H_{\mu\nu}^{a,U}[n] = \frac{C_I + C_F}{C_U} H_{\mu\nu}^{a,U}[n]$$

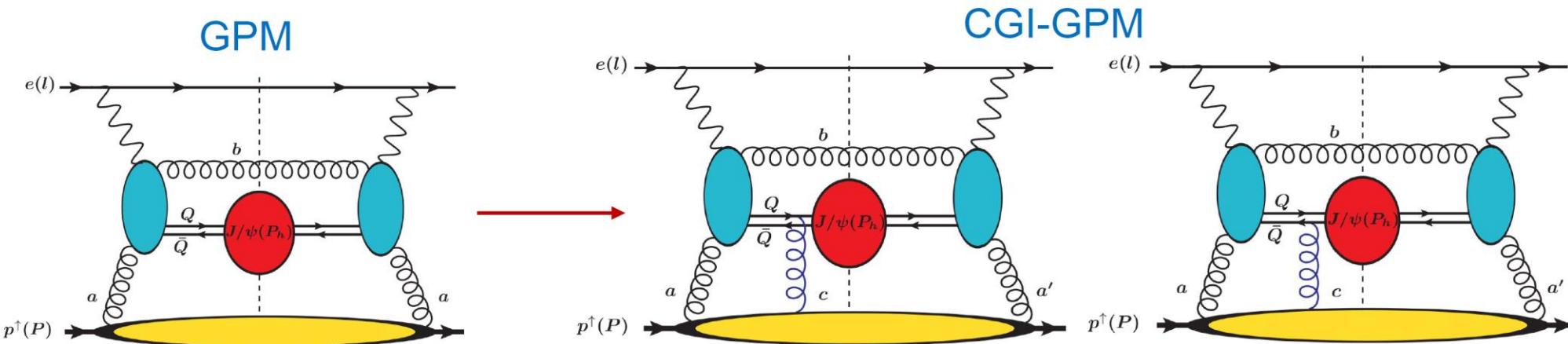
Process-dependent color factors  
can be absorbed into modified hard cross sections  
Two universal GSFs



GPM: TMDs hypothetically  
UNIVERSAL  
one single GSF

CGI-GPM:  
TMDs process dependent  
via ISIs and FSIs  
Two (f- and d-type) GSFs

# Sivers Asymmetry in $J/\psi$ lepto production



$$H^{Inc(f/d)} \equiv \frac{C_I^{(f/d)} + C_F^{(f/d)}}{C_U} H^U$$

No ISIs (colorless photon)  $\Rightarrow C_I^{(f/d)} = 0$

$C_F^{(f/d)} = 0$  for CS state, only CO states contribute

$C_F^{(d)} = 0 \Rightarrow$  only  $f$ -type GSF contributes to the SSA

More details in:  
 Gamberg & Kang PLB 696 109 (2011)  
 U. D'Alesio, L. Maxia, FM, C. Pisano, S. Rajesh  
 PRD 102 094011 (2020)

# Sivers Asymmetry in $J/\psi$ lepto production

Sivers TMD  $\Delta^N f_{a/p^\dagger}(x_a, k_{\perp a}, \mu) = 2\mathcal{N}_a(x_a) f_{a/p}(x_a, \mu) \frac{\sqrt{2e}}{\pi} \sqrt{\frac{1-\rho}{\rho}} k_{\perp g} \frac{e^{-k_{\perp a}^2/\rho \langle k_{\perp a}^2 \rangle}}{\langle k_{\perp a}^2 \rangle^{3/2}}$

$$\mathcal{N}_a(x_a) = N_a x_a^\alpha (1-x_a)^\beta \frac{(\alpha+\beta)^{(\alpha+\beta)}}{\alpha^\alpha \beta^\beta}$$

$$\rho = \frac{\langle k_\perp^2 \rangle_S}{\langle k_\perp^2 \rangle + \langle k_\perp^2 \rangle_S}$$

**Maximizing the Asymmetry:**  $\mathcal{N}_q(x) = +1$ ,  $\mathcal{N}_g^{(f/d)}(x) = +1$ ,  $\rho = 2/3$

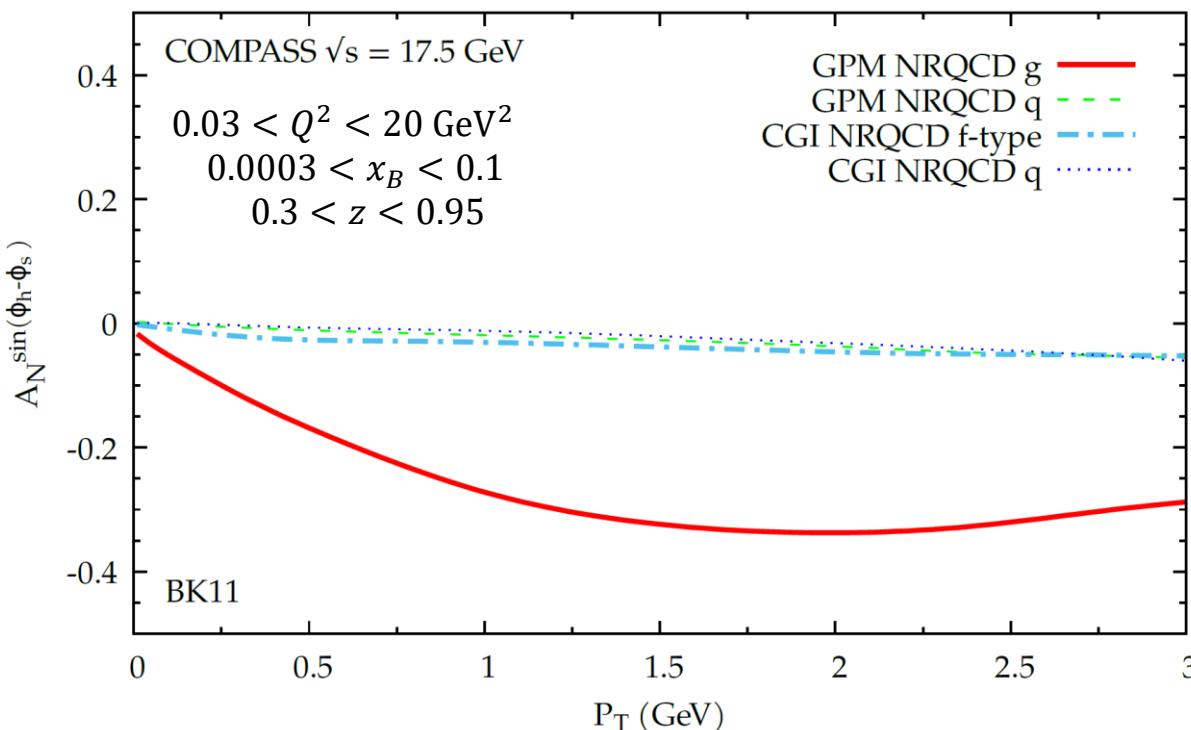
**Cancellations among different partonic contributions  
are still possible due to relative signs of NRQCD LDMEs**

TABLE III: Numerical values of LDMEs for  $J/\psi$ .

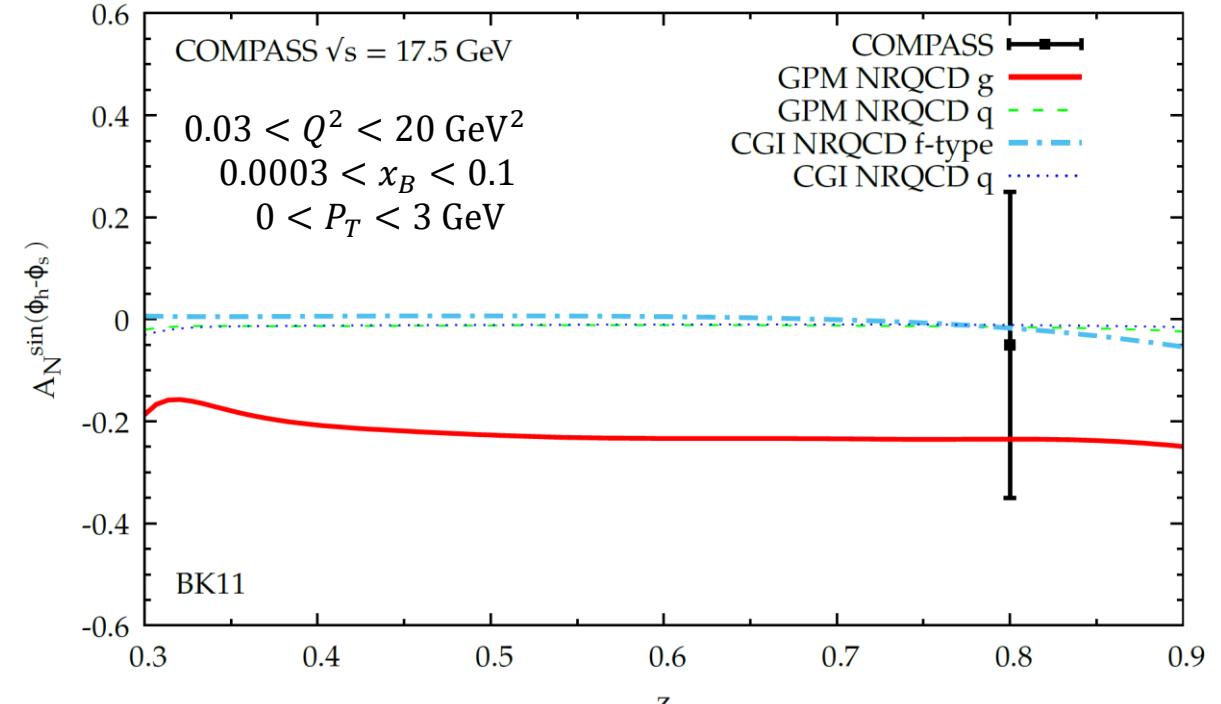
|           | $\langle \mathcal{O}_1^{J/\psi}(^3S_1) \rangle$<br>GeV <sup>3</sup> | $\langle \mathcal{O}_8^{J/\psi}(^3S_1) \rangle$<br>$\times 10^{-2}$ GeV <sup>3</sup> | $\langle \mathcal{O}_8^{J/\psi}(^1S_0) \rangle$<br>$\times 10^{-2}$ GeV <sup>3</sup> | $\langle \mathcal{O}_8^{J/\psi}(^3P_0) \rangle$<br>$\times 10^{-2}$ GeV <sup>5</sup> |
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# Maximized Sivers Asymmetry @ COMPASS

$e p^\uparrow \rightarrow e + J/\psi + X$



$e p^\uparrow \rightarrow e + J/\psi + X$



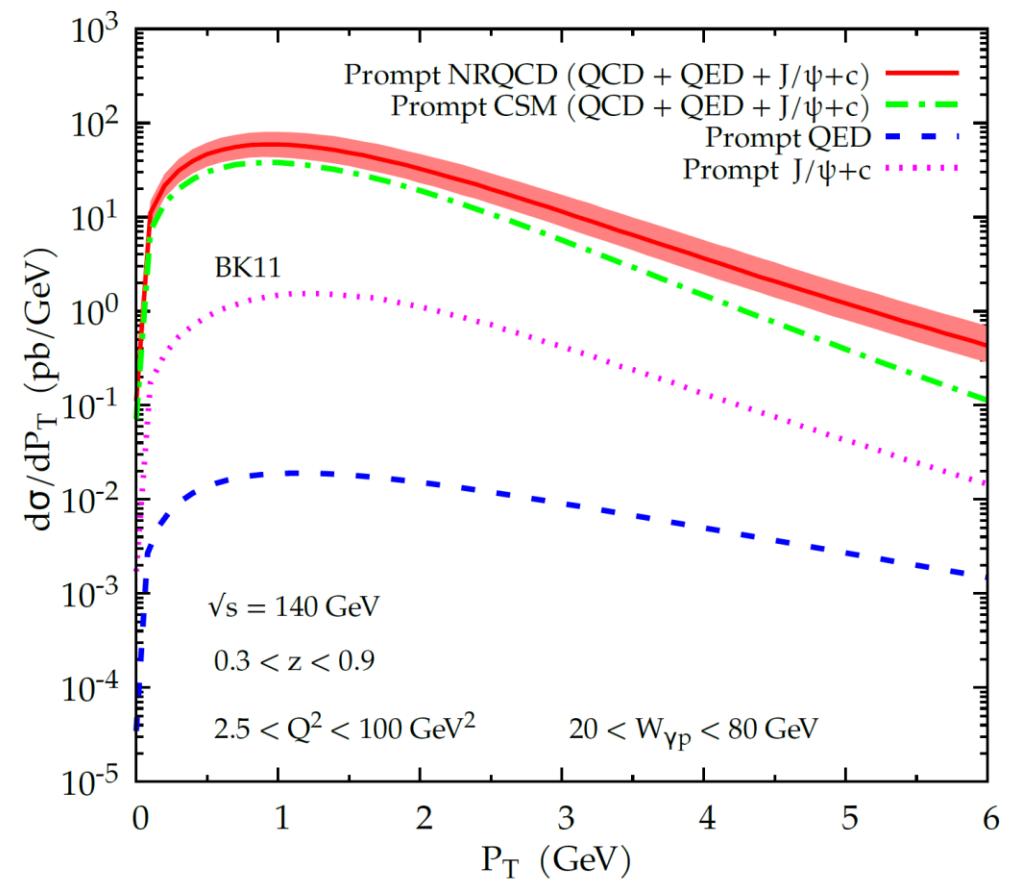
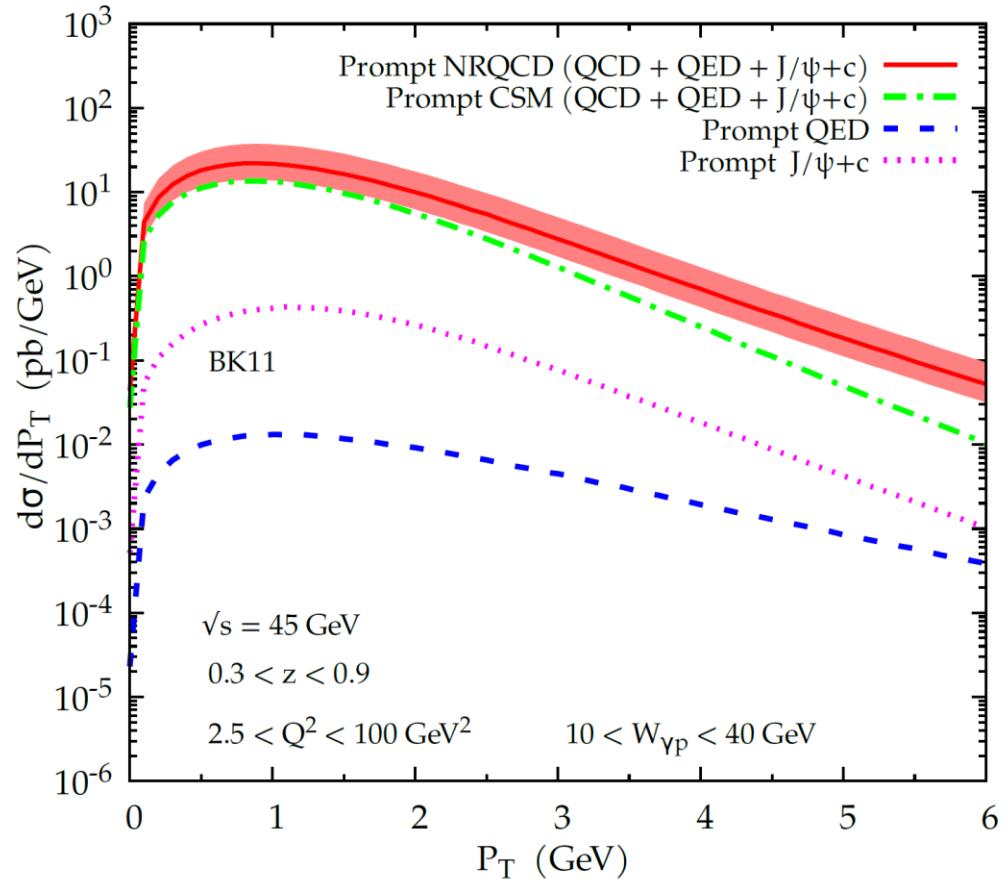
**SSA Data point: COMPASS Collab. J. Phys. Conf. Ser. 678 012050 (2017)**

Negative gluon Sivers asymmetry measured by COMPASS  
Monte Carlo analysis with enhanced photon-gluon contribution  
COMPASS Collab. PLB 772 854 (2017)

| Subprocess | Deuteron data |                   |                        | Proton data |                   |                    |
|------------|---------------|-------------------|------------------------|-------------|-------------------|--------------------|
|            | Asymmetry     | Statistical error | Systematic uncertainty | Asymmetry   | Statistical error | System uncertainty |
| PGF        | -0.14         | 0.15              | 0.10                   | -0.26       | 0.09              | 0.06               |

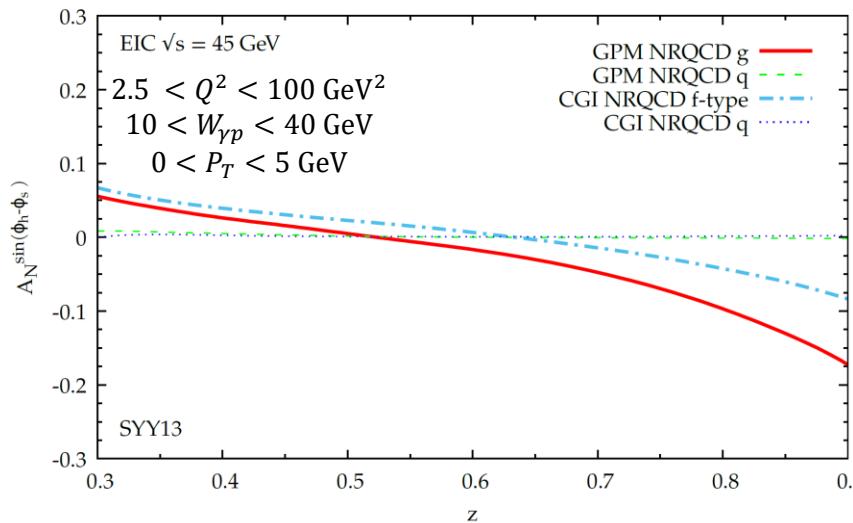
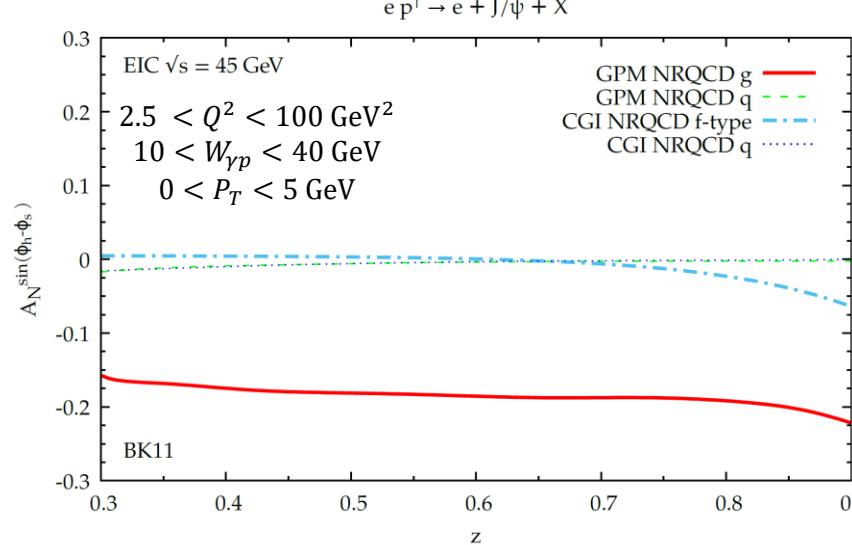
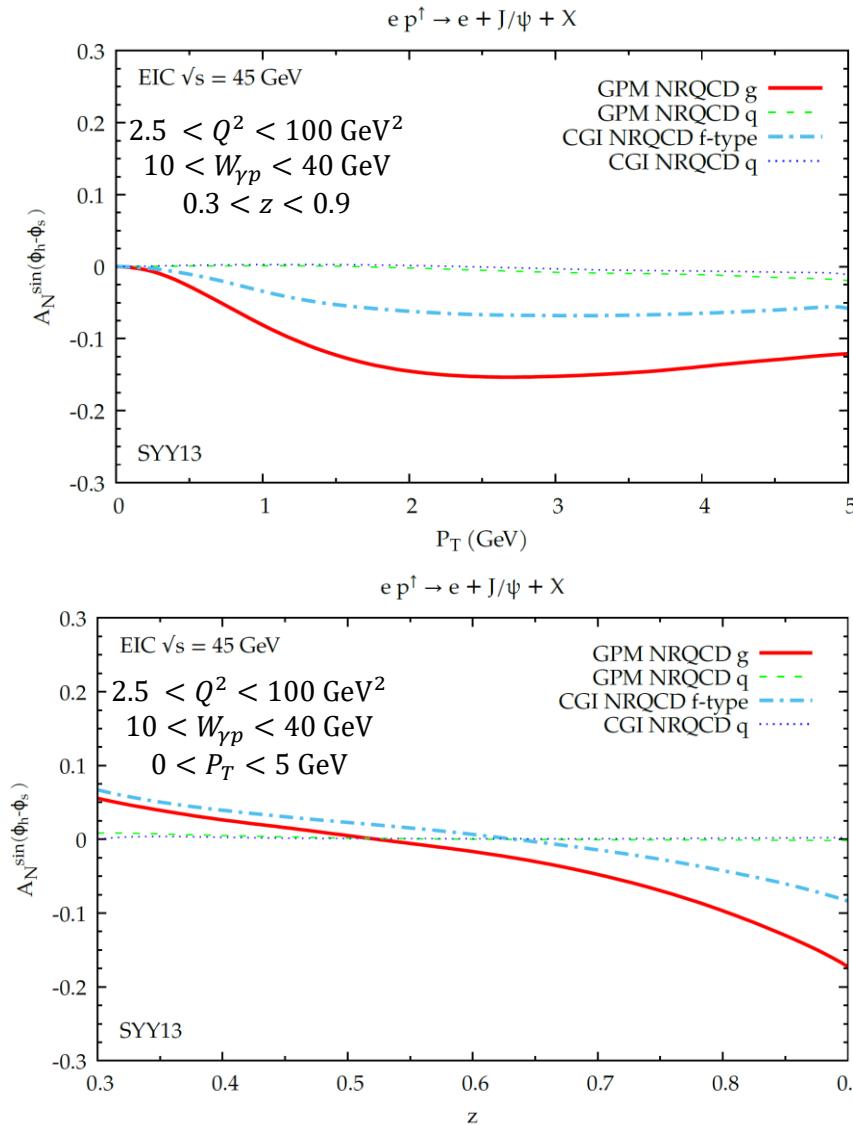
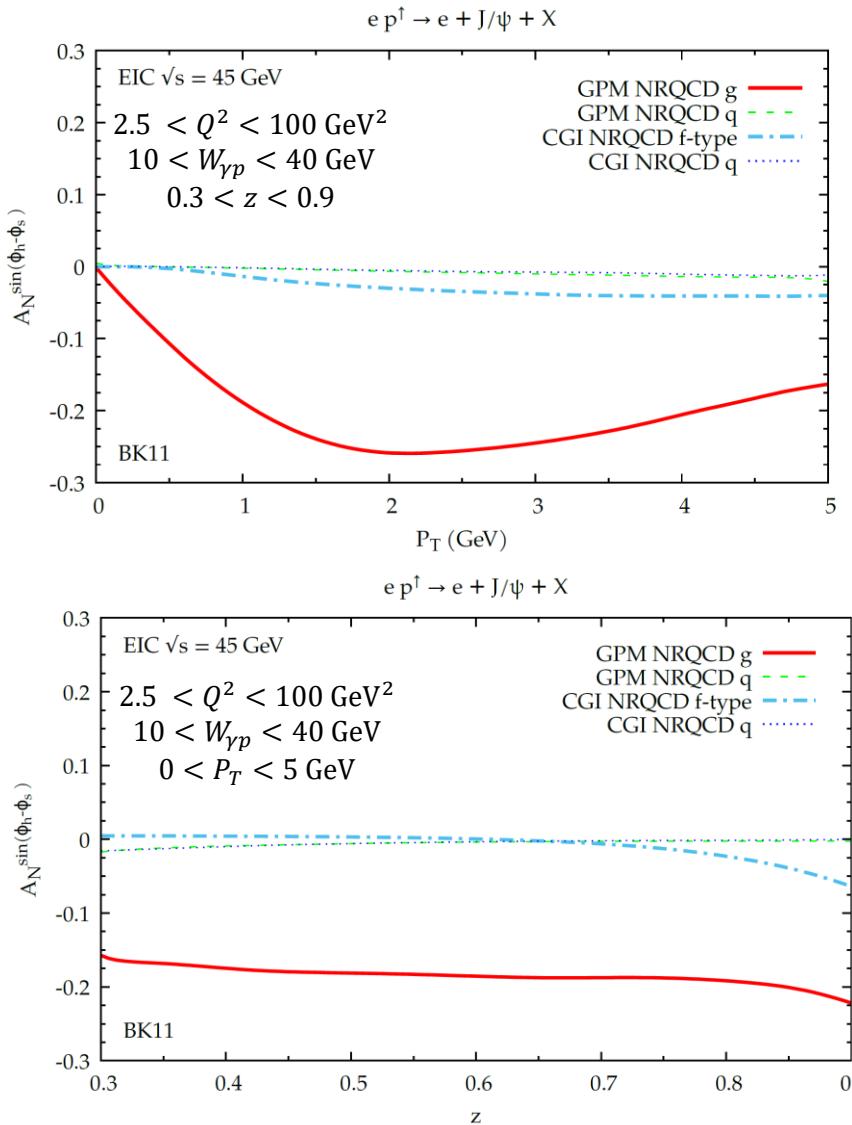
# Estimates for $J/\psi$ lepto production at the EIC

Unpolarized cross section at  $\sqrt{s} = 45 \text{ GeV}$  [left] and  $\sqrt{s} = 140 \text{ GeV}$  [right]



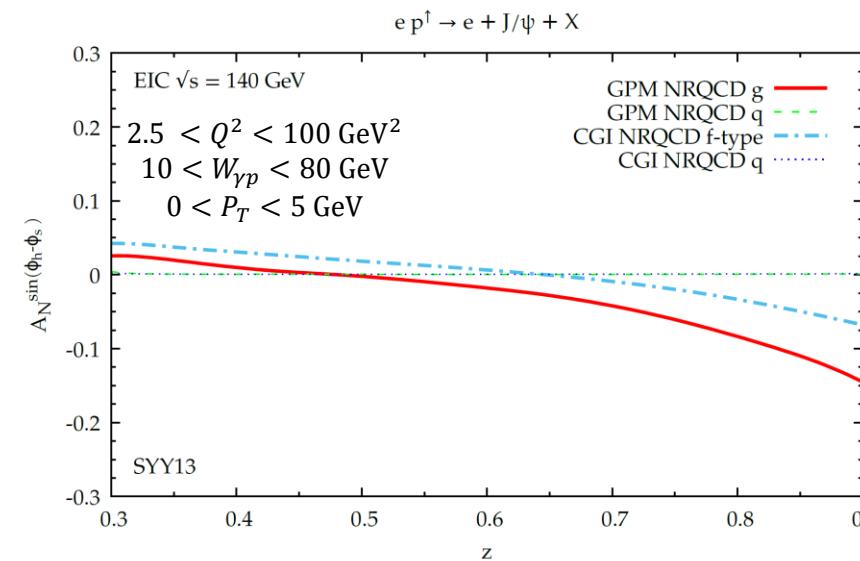
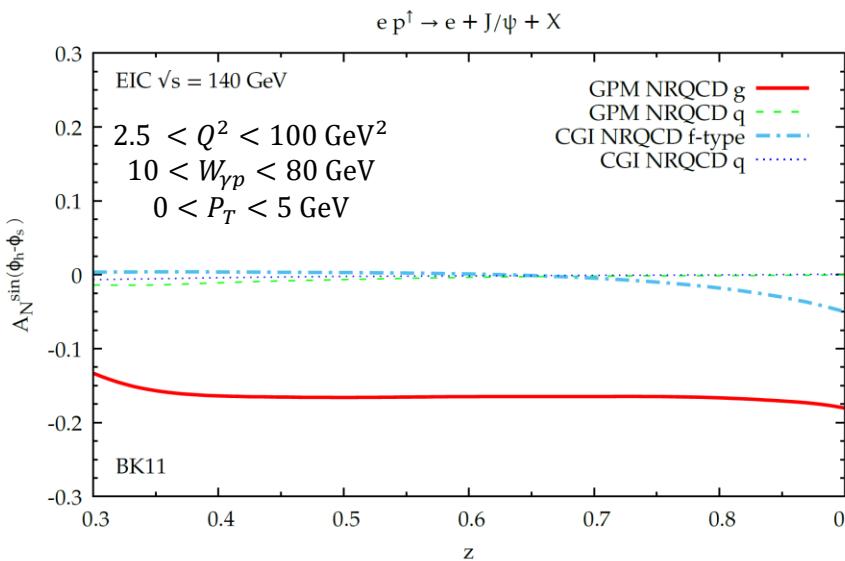
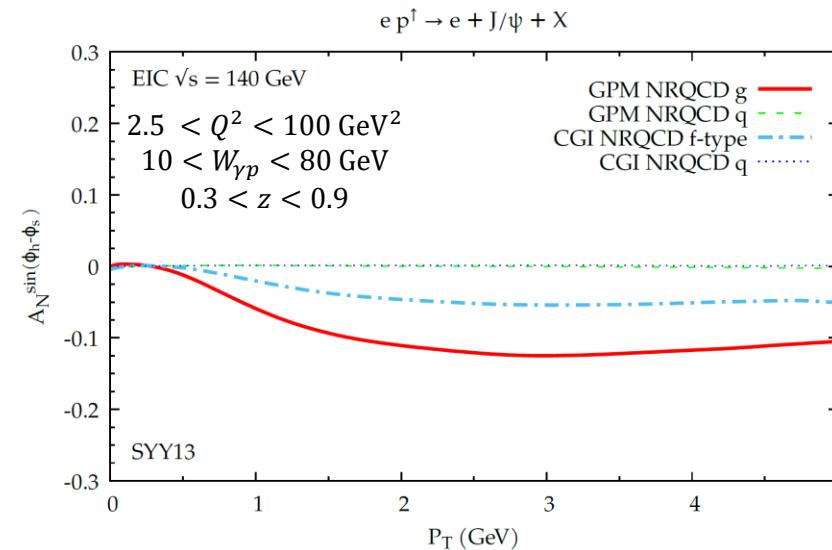
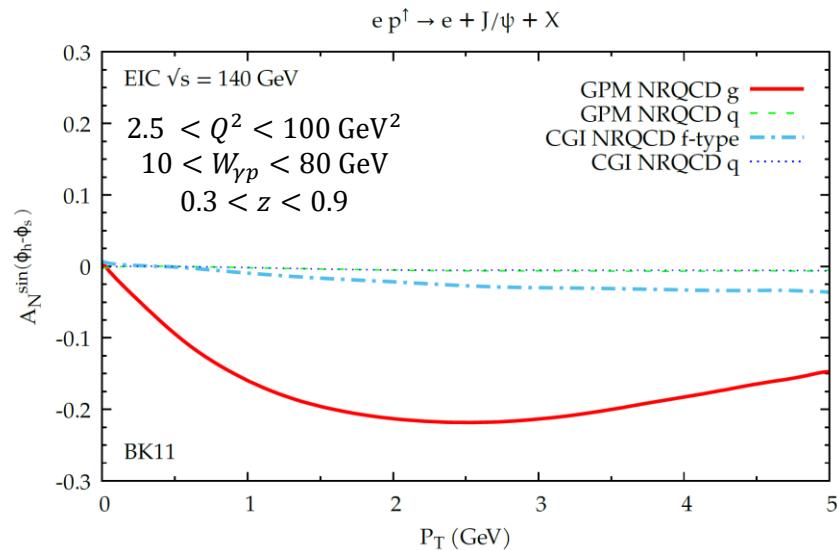
# Estimates for $J/\psi$ leptoproduction at the EIC

Maximized Sivers Asymmetry at  $\sqrt{s} = 45$  GeV for BK11 [left] and SYY13 [right] sets, vs.  $P_T$  [top] and  $z$  [bottom]



# Estimates for $J/\psi$ leptoproduction at the EIC

Maximized Sivers Asymmetry at  $\sqrt{s} = 140$  GeV for BK11 [left] and SYY13 [right] sets, vs.  $P_T$  [top] and  $z$  [bottom]



# Summary

- Inelastic  $J/\psi$  electroproduction at EIC as a tool for studying gluon TMDs and the Sivers asymmetry
- (Color-gauge-invariant) generalized parton model + NRQCD simultaneously at work
- Dependence on NRQCD LDMEs investigated by adopting two available sets for the  $P_T$  range considered
- Reasonable agreement with H1 unpolarized cross sections in the small-intermediate  $P_T$  region
- Only color-octet FSIs and  $f$ -type GSF are at work for Sivers SSA in  $\ell p^\uparrow \rightarrow \ell' + J/\psi + X$  useful for phenomenology
- Estimates for cross sections and maximized Sivers asymmetry at EIC and COMPASS kinematics

Thanks for your attention!