Accessing TMD fracture functions in backto-back hadron production at CLAS12

Timothy B. Hayward, Harut Avakian





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Accessing 3D nucleonic structure

- Decades of study have led to detailed mappings of the momentum distribution of the nucleon in terms of 1-D parton distribution functions (PDFs).
- Theoretical advances have led to a framework in which information on the confined motion of the partons inside a nucleon are matched to 3-D transverse momentum dependent parton distribution functions (TMDs).
- TMDs/PDFs
 Confined motion of quarks and gluons inside the nucleus
 Orbital motion of quarks, correlations between quarks and gluons
 Fragmentation Functions
 Probability for a quark to form particular final state particles
 Insight into transverse momenta and polarization

TMD-sensitive CLAS12 Measurements

CLAS12, with its high luminosity, wide acceptance and low Q² (for higher twist measurements) is an ideal place to study numerous physics observables which access these 3-D dependent quantities... Measurements traditionally focus on factorization theorems and assumption that hadrons are produced in current fragmentation.



S. Diehl, A. Kim, G. Angelini, K. Joo et al., (2021) hep:ex/2101.03544.

Fracture Functions

- Information on the momentum distribution of quarks and gluons are encoded in TMDs that can be measured in inclusive processes such as SIDIS.
- TMDs have been extensively studied via azimuthal modulations of a final state hadron generated in the fragmentation of a struck quark (CFR).
- Final state hadrons can also form from the left-over target remnant (TFR) whose partonic structure is defined by "fracture functions": the probability to form a certain hadron given a particular ejected quark.



Figure 1: The handbag diagram for the SIDIS hadronic tensor in the current fragmentation region (left) and in the target fragmentation region (right).

Phys. Lett. B. 699 (2011), 108-118, [hep-ph] 1102.4214

	U	\sum_{L}	T
U	\hat{u}_1	$\hat{l}_1^{\perp h}$	$\hat{t}_1^h, \hat{t}_1^\perp$
L	$\hat{u}_{1L}^{\perp h}$	\hat{l}_{1L}	$\hat{t}^h_{1L}, \hat{t}^\perp_{1L}$
Т	$\hat{u}^h_{1T}, \hat{u}^\perp_{1T}$	$\hat{l}^h_{1T}, \hat{l}^\perp_{1T}$	$\hat{t}_{1T}, \hat{t}_{1T}^{hh}, \hat{t}_{1T}^{\perp\perp}, \hat{t}_{1T}^{\perp h}$

Back-to-back Formalism

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Phys. Lett. B. 713 (2012), 317-320, [hep-ph] 1112.2604



When two hadrons are produced "backto-back" with one in the CFR and one in the TFR the structure function contains a convolution of a fracture function and a fragmentation function.



Figure 1: Lepto-production of two hadrons, one in the CFR and one in the TFR. Phys. Lett. B. 706 (2011), 46-52, [hep-ph] 1109.1132

Additional Modulations

$$\mathcal{F}_{LU} = \frac{|p_{\pi}^{\perp}||p_{P}^{\perp}|}{m_{p}m_{\pi}} \mathcal{C} \left[w_{5} \hat{l}_{1}^{\perp h} D_{1} \right] -$$

If the correlations are assumed small, the fracture functions can be expanded in powers of $k^{\perp} \cdot p_P^{\perp}$.

Structure functions carry a dependence on $|p_{\pi}^{\perp}||p_{P}^{\perp}|$ which introduces a dependence on $\cos \Delta \phi$.

$$\hat{l}_{1}^{\perp h}(x,\zeta,\mathbf{k}^{\perp 2},\mathbf{p}_{\mathbf{P}}^{\perp 2},\mathbf{k}^{\perp}\cdot\mathbf{p}_{\mathbf{P}}^{\perp}) \\\approx a(x,\zeta,\mathbf{k}^{\perp 2},\mathbf{p}_{\mathbf{P}}^{\perp 2}) \\+b(x,\zeta,\mathbf{k}^{\perp 2},\mathbf{p}_{\mathbf{P}}^{\perp 2})\mathbf{k}^{\perp}\cdot\mathbf{p}_{\mathbf{P}}^{\perp}$$

The term linear in $k^{\perp} \cdot p_P^{\perp}$ yields a $\cos \Delta \phi$ which when combined with the already existing $\sin \Delta \phi$ term results in a $\sin 2\Delta \phi$.

$$\mathcal{A}_{LU}(x,\zeta,\mathbf{k}^{\perp^2},\mathbf{p}_{\mathbf{P}}^{\perp^2},\Delta\phi) = A(x,\zeta,\mathbf{k}^{\perp^2},\mathbf{p}_{\mathbf{P}}^{\perp^2})\sin\Delta\phi + B(x,\zeta,\mathbf{k}^{\perp^2},\mathbf{p}_{\mathbf{P}}^{\perp^2})\sin(2\Delta\phi)$$

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Phys. Lett. B. 713 (2012), 317-320, [hep-ph] 1112.2604

CLAS12 Experimental Setup





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V. Burkert et al., Nucl. Instrum. Meth. A 959 (2020) 163419
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- Data taken during Fall 2018 and Spring 2019.
- 10.6 (2018) and 10.2 (2019) GeV longitudinally polarized electron beam with a H₂ target.
- Analyzed data is a small percentage of the total available data.

Particle ID

- Electron
 - Electromagnetic calorimeter.
 - Cherenkov detector.
 - Vertex and fiducial cuts.



Hadron

- β vs p comparison between vertex timing and event start time.
- Vertex and fiducial cuts.



Selecting "back-to-back" events

- Formalism requires one hadron in the CFR and one hadron in the TFR.
- A natural choice for a first analysis are events with a pion (CFR biased) and proton (TFR biased).

PLOTS REMOVED ONLINE (DISTRIBUTIONS OF M_X AND X_F OF PROTON AND PION)

Monte Carlo

- SIDIS MC "clasdis"¹ based on PEPSI² generator, the polarized version of the well-known LEPTO³ generator.
- Parameters changed to reproduce observed distributions include average transverse momentum, fraction of spin-1 light mesons and fraction of spin-1 strange mesons.
- CLAS12 detector system described in "GEMC^{"4,} a detailed GEANT4 simulation package.
- Excellent agreement between data and MC!

PLOTS REMOVED ONLINE (Data/MC Comparisons of Q², P^{\perp}, X_F of pion and $\Delta\phi$)

- 1. H. Avakian, "clasdis." https://github.com/JeffersonLab/clasdis, 2020.
- 2. L. Mankiewicz, A. Schafer, and M. Veltri, "Pepsi: A monte carlo generator for polarized leptoproduction," Comput. Phys. Commun., vol. 71, pp. 305–318, 1992.
- 3. G. Ingelman, A. Edin, and J. Rathsman, "LEPTO 6.5: A Monte Carlo generator for deep inelastic 912 lepton nucleon scattering," Comput. Phys. Commun., vol. 101, pp. 108–134, 1997.
- 4. M. Ungaro et al., "The CLAS12 Geant4 simulation," Nucl. Instrum. Meth. A, vol. 959, p. 163422, 2020.

Extracting A_{LU}

- Select $ep \rightarrow e'P \pi^+ + X$.
- Consider all possible hadron pairs.
- Amplitudes are extracted simultaneously via maximizing a likelihood function.
- Unbinned maximum likelihood method:

$$\frac{N_{\pm}(\Delta\phi)}{N(\Delta\phi)} = \frac{1}{2} \frac{L_{\pm}a_{\pm}(\Delta\phi)\eta_{\pm}(\Delta\phi)}{La(\Delta\phi)\eta(\Delta\phi)} \left[1 \pm PA_{LU}(\Delta\phi)\right]$$

- Use MINUIT to minimize the -log likelihood.
- Include relevant beam polarization (~85% at JLab).

Channel selection

- Q²>1.0 GeV²
- W>2.0 GeV
- M_{miss}>0.85 GeV
- y<0.75

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$$x_{F\pi} > C$$

• $z_{\pi} > 0.2$

Back-to-back results

- Observed linear dependence on the product of transverse momenta is consistent with expectations.
- Non-zero asymmetries are the first experimental observation of possible spin-orbit correlations between hadrons produced simultaneously in the CFR and TFR.

$$\mathcal{F}_{LU} = \frac{|p_{\pi^+}^{\perp}||p_P^{\perp}|}{m_p m_{\pi}} \mathcal{C} \left[w_5 \hat{l}_1^{\perp h} D_1 \right]$$



Access to unmeasured fracture function

- Fracture function $\hat{l}_1^{\perp h}$ depends on ζ .
- Fragmentation function D_1 depends on z_{π} .

$$\mathcal{F}_{LU}^{\sin(\Delta\phi)} = \frac{|p_{\pi^+}^{\perp}||p_P^{\perp}|}{m_P m_{\pi^+}} \mathcal{C} \left[w_5 \hat{l}_1^{\perp h} D_1 \right]$$
$$\zeta = E_P / E \quad z_{\pi} = E_{\pi} / \nu$$



Conclusions

- Significant single-spin asymmetries have been observed for the first time in back-to-back proton-pion electroproduction.
- Beam-spin asymmetries amplitudes indicate that spin-orbit correlations may exist between hadrons produced simultaneously in the target and current fragmentation regions.
- Analysis in final collaboration review with the aim of submission for publication this year.