# POLARZIATION TEST FOR QUSI-FREE KNOCKOUT OF NUCLEON FROM NUCLEAR SHORT-RANGE CORRELATED NN PAIR

Yu.N. UZIKOV (DLNP, JINR, Dubna)

- Main properties of SRC NN pairs in nuclei
- ${}^{12}C + p \rightarrow p + p + N + {}^{10}B$  at BM@N NICA
- ISI@FSI in <sup>12</sup>C(p,ppN)<sup>10</sup>A
- T<sub>20</sub> in pd-pX, pd-dp, ed-ed in the SRC region
- ISI@FSI in pd->ppn
- Conclusion

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- Dubna, M.G. Mesheryakov et al., (1957) p+<sup>12</sup>C → d+X at 670 MeV; quasi-elastic knock-out of the fast deuterons.
  - D.I. Blokhintsev (1957) : fluctons in nuclei Blokhintsev, Efremov, Lukjanov, Titov, 6q, 1978; Strikman, Frankfurt, 1978, SRC
  - Two nucleons at short distances r<sub>NN</sub> <0.5 fm with relative momentum q>1/r<sub>NN</sub> =0.4 GeV/c Repulsive core in NN-potential → high-momentum part of the w.f. of NN pair



$$\vec{k_1} + \vec{k_2} \approx 0$$
  
Two nucleon SRC  $k_1 > k_F$ 

- Summary of the main results on SRC of BNL, Jlab, SLAC
- \* High-momentum part (q >  $p_F$ ) accounts for 20% nucleons .
- \* pn- SRC pairs at 0.3 GeV/c < q < 0.7 GeV/c dominate by factor of 20 as compared to pp and nn due to the tensor forces.
- \* Factorization  $n(p_1, p_2) = C_A n_{c.m.}(k_{cm}) n_{rel}(q)$ .
- \* SRC are connected with **neutrino-nucleus interaction**, **neutron stars** structure, modification of the bound nucleon structure (**EMC effect**).
- \* 3N SRC

Cioffi degli Atti, Phys. Rep. 590 (2015) 1 O. Hen et al. Rev. Mod. Phys. 89 (2017) 045002.

#### M.Duer et al. PRL 122 (2019) 172502 (CLAS Collabotration)



## Study of SRC at BM@N

4=GeV/c /nucleon beam of <sup>12</sup>C and proton target in inverse kinematics



M. Patsyuk et al., Nature Phys. 17 (2021) 693; arXiv:2102.02626 [nucl-ex]

12C(p,2p)11B (g.s.), shell model is OK !

**12C(p,2pN)10A**, SCR: 10B 23 events (pn); 10Be – 2 event (pp)

**CONCLUSION**: unperturbed by ISI@FSI distributions are measured "All measured reactions are well described by theoretical calculations that do not contain ISI/FSI distortions"

# REACTION ${}^{12}C+p \rightarrow {}^{10}B+P+P+N$





- Matrix element of the  $p + \langle NN \rangle \rightarrow p + p + N$  ——

In the Light front dynamics

$$\begin{split} M_{fi}^{LFD}(p < NN > \to p_1 p_2) &= \frac{\Psi_d^{LFD}(\mathbf{k}_{\perp}, \xi)}{1 - \xi} M_{fi}(pN \to p_1 p_2), \\ & \mathbf{p} & \boldsymbol{\xi} = \frac{p_r^+}{p_r^+ + p_N^+}, \ \mathbf{q}_{\perp} = (1 - \xi) \mathbf{p}_{r\perp} - \xi \mathbf{p}_{N\perp}, \\ & M_{pN}^2 = \frac{m_p^2 + \mathbf{p}_N^2}{\xi(1 - \xi)}. \\ & \mathbf{M}_{pN}^2 = \frac{m_p^2 + \mathbf{p}_N^2}{\xi(1 - \xi)}. \end{split}$$

Yu.N. Uzikov, EPJ Web Conf. 222 (2019) 03027

#### **Transition matrix element**

$$T_{fi} = \binom{A}{x}^{1/2} \sum_{x' \ \nu\Lambda} \langle \psi_A | \psi_B \psi_{x'}, \psi_{\nu\Lambda} \rangle \Phi_{\nu\Lambda}(\mathbf{k}_B) T^{px' \to Nx}.$$

$$T^{px' \to Nx} = \langle \mathbf{k}_N \mathbf{k}_x \chi_N \psi_x | \tau (px' \to Nx) | \mathbf{k}_p, -\mathbf{k}_B \chi_p \psi_{x'} \rangle$$
$$N_A - N_B = N_x + \nu$$

Yu.N. Uzikov, Izv. RAN. Ser. Fiz. 84 (2020) 580; EPJ Web Conf. 222 (2019) 03027; Phys. Part. Nucl. 52 (2021) 652

How to take into account ISI and FSI?



## ISI@FSI ELASTIC RESCATTERINGS

Feynman graphs + generalized eikonal appr. L.L. Frankfurt et al. PRC 56(1997) 2752

T-operator formalism with eikonal appr. collinear kinematics. M.A. Zhusupov, Yu. N. Uzikov ,

Fiz. Elem. Chast. At. Yadr. 18(1987) 323





<NN>

f)

В

А



Feynman graphs + generalized eikonal approximation (GEA) L.L. Frankfurt et al. PRC 56(1997) 2752

$$\begin{aligned} \frac{1}{(p_1 + p_B - p_{B'})^2 - m^2 + i\varepsilon} &\cong \frac{1}{2p_{1z}} \frac{1}{p_{B'_z} - p_{B_z} + \Delta_1 + i\varepsilon};\\ \Delta_1 &= -\frac{\vec{p}_{1t}}{p_{1z}} (\vec{p}_{Bt} - \vec{p}_{B'_t}) + \frac{E_1}{p_{1z}} (E_B - E_{B'});\\ \Delta_1 &\neq 0 - > GEA \end{aligned}$$

T-operator formalism, eikonal appr.; kollinear kinematics. M.A. Zhusupov, Yu.N.U. Sov.J.Part.Nucl 18 (1987)

$$\begin{split} \Phi_{\nu\Lambda}(\mathbf{k}_B) &= \psi_{\nu\Lambda}(\mathbf{k}_B) + \frac{i}{4\pi k_{pA}} \int d^2 \mathbf{q}_p F_{pB}(\mathbf{q}_p) \psi_{\nu\Lambda}(\mathbf{k}_B - \mathbf{q}_p) + \\ \frac{i}{4\pi k_{xB}} \int d^2 \mathbf{q}_x F_{xB}(\mathbf{q}_x) \psi_{\nu\Lambda}(\mathbf{k}_B - \mathbf{q}_x) + \frac{i}{4\pi k_{NB}} \int d^2 \mathbf{q}_N F_{NB}(\mathbf{q}_N) \psi_{\nu\Lambda}(\mathbf{k}_B - \mathbf{q}_N) \\ &- \frac{1}{(4\pi)^2 k_{pA}} \int d^2 \mathbf{q}_p \ d^2 \mathbf{q}_N F_{pB}(\mathbf{q}_p) \ F_{NB}(\mathbf{q}_N) \ \psi_{\nu\Lambda}(\mathbf{k}_B - \mathbf{q}_N - \mathbf{q}_p) - \\ &- \frac{1}{(4\pi)^2 k_{xB}} \int d^2 \mathbf{q}_p \ d^2 \mathbf{q}_x F_{pB}(\mathbf{q}_p) \ F_{xB}(\mathbf{q}_x) \ \psi_{\nu\Lambda}(\mathbf{k}_B - \mathbf{q}_p - \mathbf{q}_x) - \\ &- \frac{1}{(4\pi)^2 k_{xB}} \int d^2 \mathbf{q}_x \ d^2 \mathbf{q}_N F_{xB}(\mathbf{q}_p) \ F_{NB}(\mathbf{q}_N) \ \psi_{\nu\Lambda}(\mathbf{k}_B - \mathbf{q}_p - \mathbf{q}_x) - \\ &- \frac{1}{(4\pi)^2 k_{xB}} \int d^2 \mathbf{q}_x \ d^2 \mathbf{q}_x F_{xB}(\mathbf{q}_p) \ F_{NB}(\mathbf{q}_N) \ \psi_{\nu\Lambda}(\mathbf{k}_B - \mathbf{q}_p - \mathbf{q}_N) - \\ &- \frac{i}{(4\pi)^3 k_{pA}} \ k_{NB} \ k_{xB}} \int d^2 \mathbf{q}_p \ d^2 \mathbf{q}_N F_{xN}(\mathbf{q}_x) \ \psi_{\nu\Lambda}(\mathbf{k}_B - \mathbf{q}_N - \mathbf{q}_p) - \\ &- \frac{i}{(4\pi)^3 k_{pA}} \ k_{NB} \ k_{xB}} \int d^2 \mathbf{q}_p \ d^2 \mathbf{q}_N F_{xN}(\mathbf{q}_N) \ \psi_{\nu\Lambda}(\mathbf{k}_B - \mathbf{q}_N - \mathbf{q}_N) - \\ &- \frac{i}{(4\pi)^3 k_{pA}} \ k_{NB} \ k_{xB} \ k_{xB}} \int d^2 \mathbf{q}_p \ d^2 \mathbf{q}_N d^2 \mathbf{q}_N \times \\ &\times F_{pB}(\mathbf{q}_p) \ F_{NB}(\mathbf{q}_N) \ F_{xN}(\mathbf{q}_x) \ \psi_{\nu\Lambda}(\mathbf{k}_B - \mathbf{q}_N - \mathbf{q}_p - \mathbf{q}_N). \end{split}$$



![](_page_12_Figure_0.jpeg)

# What about ISI@FSI in p<NN>-> p+N+N?

![](_page_13_Picture_1.jpeg)

L. Frankfurt et al, PRC 51 (1995) 890

![](_page_14_Figure_0.jpeg)

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# Polarization test for the IA in p<NN>->pNN and measurement of n<sub>rel</sub>(q<sub>rel</sub>)

$$p < NN > \rightarrow p + N + N;$$
  
 $p + \vec{d} \rightarrow p + p + n;$   
 $T_{20}$  – measurement

IA for 
$$T_{20}$$
  $\longrightarrow_{rel}$  (q)

A similar test was suggested for mechanisms of backward pd->{pp}s+n at 0.5-2 GeV at COSY(1999)

![](_page_16_Figure_0.jpeg)

![](_page_17_Figure_0.jpeg)

![](_page_17_Figure_1.jpeg)

Fig. 2.  $T_{20}$  versus k; (•) for backward elastic dp and (•) for inclusive breakup. Solid curve is IA using Paris deuteron wave function; straight line is asymptotic prediction [24].

![](_page_18_Figure_0.jpeg)

Fig. 3. Tensor analyzing power  $T_{20}$  for the backward elastic dp scattering versus k; (\*) Saclay data [2], ( $\circ$ ) 93-run ( $\Box$ ) 94-run: this experiment. Solid and broken curves are the ONE predictions

L. Azhgirey et al. PLB 387 (1996) 37

![](_page_18_Figure_3.jpeg)

Fig. 4.  $T_{20}$  data for H(d, p)X at 9 GeV/c (filled triangles), compared with  $T_{20}$  data for backward elastic  $dp \rightarrow pd$  from refs [19] (empty circles) and [22] (empty squares). The IA prediction that both reactions are closely related is not validated by these data. The curve represents the IA, as in Fig. 2.

![](_page_19_Figure_0.jpeg)

![](_page_20_Figure_0.jpeg)

# Impulse approximation dominates!

**<u>SUGGESTION for BM@N or SPD NICA</u>: In order to check validation of the IA for the quasi-elastic knock out of the nucleon from SRC NN pair</u>** 

## p+<NN>-p+N+N,(1)

it is of importance to measure the  $T_{20}$  for the free reaction p+d-> p+p+n (2) at the same kinematics as in the subprocess (1) of the reaction  ${}^{12}C+p -> p+N+N+ {}^{10}A$ 

BM@N result : 23 event for pn, and 2 for pp , while was expected ~400; for free deuteron one may expect ~ 100.

### In addition:

Testing rescattering dynamics (including color transparency effects - dashed curves)

L.L. Frankfurt et al. PRC 56 (1997)

$$T = \sigma^{DWIA} / \sigma^{IA}$$
 pd->ppn

![](_page_22_Figure_4.jpeg)

### P<sub>L</sub>=6 GeV/c is accessible at BM@N

 $\alpha_s = 1$  optimal for testing dynamics of multinucleon rescatterings

## **TO CONCLUSION**

The reaction p+d  $\rightarrow$ p+p+n has two aspects :

(i) Study of the deuteron wave function, u(q) and w(q) waves at high q
 (ii) Study of wave packets evolution over distances < 2 fm under interference between IA, single and double rescatterings. One can choose kinematics with minimal and maximal FSI-> CT.

## THANK YOU FOR ATTENTION!

# SUPPLEMENTARY SLIDES

 $T = \sigma(\gamma A) / \sigma(\gamma p)$ 

D. Dutta et al. / Progress in Particle and Nuclear Physics 69 (2013) 1–27

![](_page_25_Figure_2.jpeg)

**Fig. 13.** The nuclear transparency of  ${}^{4}\text{He}(\gamma, p\pi)$  at  $\theta_{cm}^{\pi} = 70^{\circ}$  (left) and  $\theta_{cm}^{\pi} = 90^{\circ}$  (right), as a function of momentum transfer square |t| [80]. The inner error bars shown are statistical uncertainties only, while the outer error bars are statistical and point-to-point systematic uncertainties (2.7%) added in quadrature. In addition there is a 4% normalization/scale systematic uncertainty which leads to a total systematic uncertainty of 4.8%.

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![](_page_26_Figure_0.jpeg)

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Possible studies at the first stage of the NICA collider operation with polarized and unpolarized proton and deuteron beams

V. V. Abramov<sup>1</sup>, A. Aleshko<sup>2</sup>, V.A. Baskov<sup>3</sup>, E. Boos<sup>2</sup>,
V. Bunichev<sup>2</sup>, O.D. Dalkarov<sup>3</sup>, R. El-Kholy<sup>4</sup>, A. Galoyan<sup>5</sup>, A.V. Guskov<sup>6</sup>,
V.T. Kim<sup>7,8</sup>, E. Kokoulina<sup>5,9</sup>, I.A. Koop<sup>10,11,12</sup>, B.F. Kostenko<sup>13</sup>,
A.D. Kovalenko<sup>5</sup>, V.P. Ladygin<sup>5</sup>, A. B. Larionov<sup>14,15</sup>, A.I. L'vov<sup>3</sup>, A.I. Milstein<sup>10,11</sup>,
V.A. Nikitin<sup>5</sup>, N. N. Nikolaev<sup>16,26</sup>, A. S. Popov<sup>10</sup>, V.V. Polyanskiy<sup>3</sup>,
J.-M. Richard<sup>17</sup>, S. G. Salnikov<sup>10</sup>, A.A. Shavrin<sup>18</sup>, P.Yu. Shatunov<sup>10,11</sup>,
Yu.M. Shatunov<sup>10,11</sup>, O.V. Selyugin<sup>14</sup>, M. Strikman<sup>19</sup>, E. Tomasi-Gustafsson<sup>20</sup>,
V.V. Uzhinsky<sup>13</sup>, Yu.N. Uzikov<sup>6,21,22,\*</sup>, Qian Wang<sup>23</sup>, Qiang Zhao<sup>24,25</sup>, A.V. Zelenov<sup>7</sup>

<sup>1</sup> NRC "Kurchatov Institute" - IHEP, Protvino 142281, Moscow region, Russia

<sup>2</sup> Skobeltsyn Institute of Nuclear Physics, MSU, Moscow, 119991 Russia

<sup>3</sup> P.N. Lebedev Physical Institute, Leninsky prospect 53, 119991 Moscow, Russia

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<sup>4</sup> Astronomy Department, Faculty of Science, Cairo University, Giza, Egypt, 12613