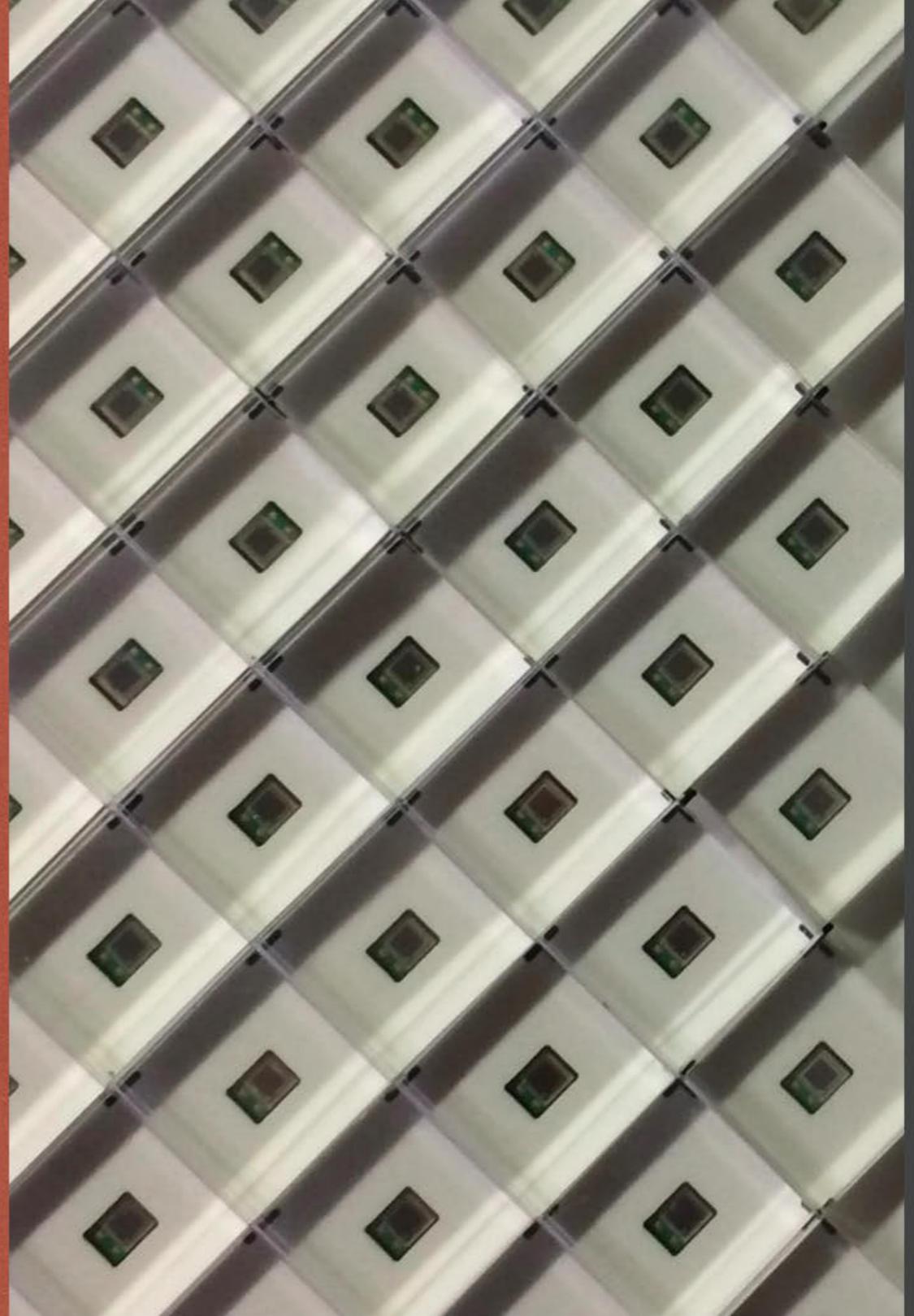


# Precision microwave spectroscopy of the ground-state hyperfine splitting in muonium atom



1

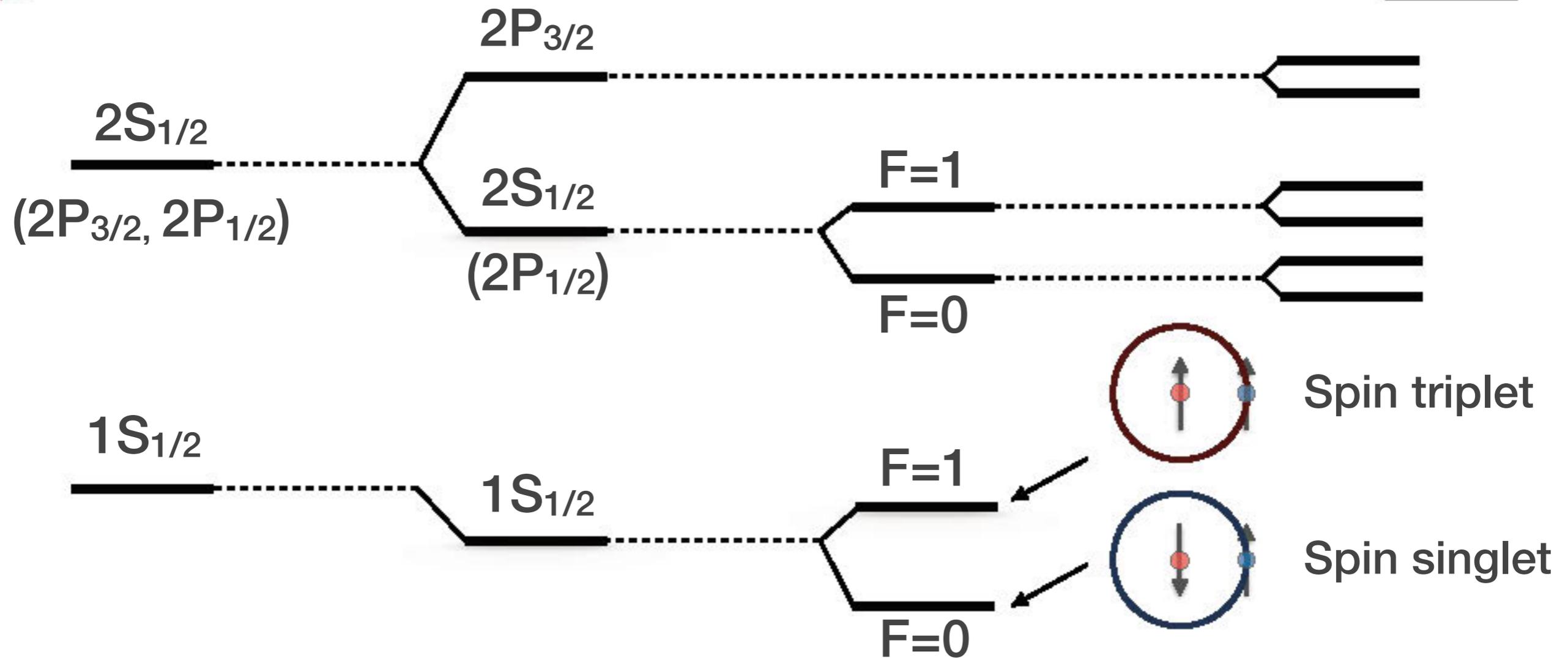
Sohtaro Kanda / 神田 聡太郎

2019/11/26

[sohtaro.kanda@riken.jp](mailto:sohtaro.kanda@riken.jp)

# Hydrogen Atom Spectroscopy

2



Lyman series (1906)

Bohr model (1913)

Hyperfine structure (1881)

Nuclear magnetic moment (1924)

Fine structure (1887)

Sommerfeld-Bohr model (1916)

Lamb shift (1947)

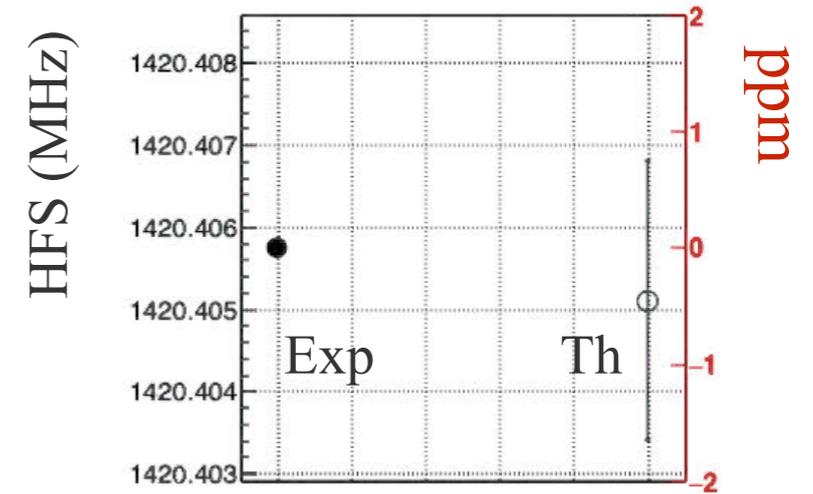
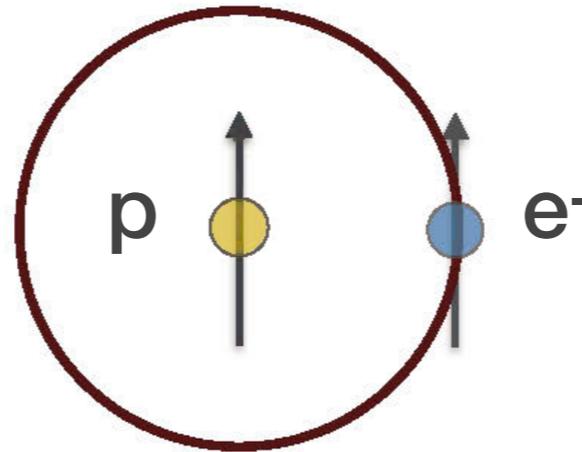
QED (1949)

- Spectroscopy of hydrogen atoms has led to the development of quantum mechanics and quantum field theory.

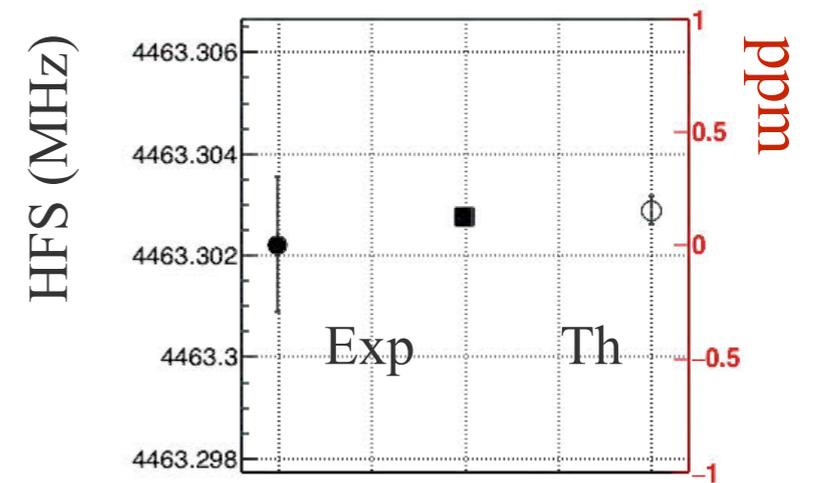
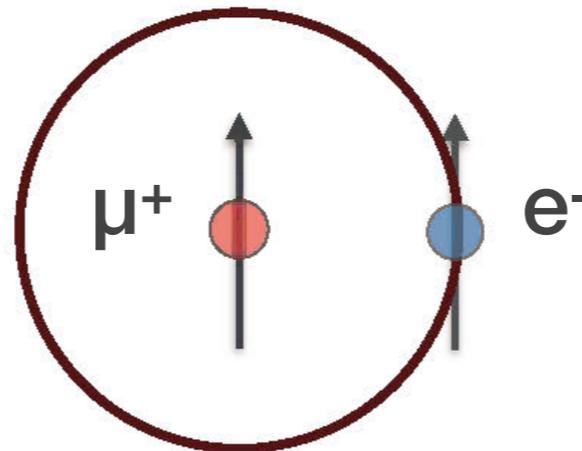
# HFS of Hydrogen-like Atoms

3

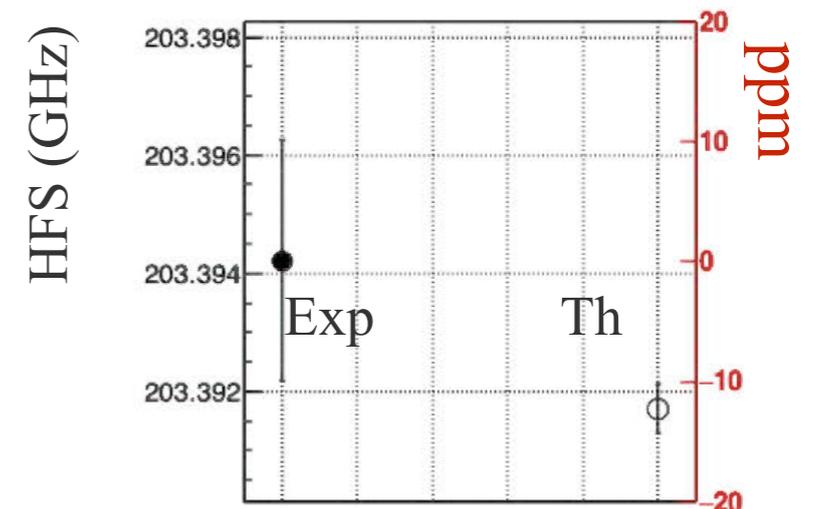
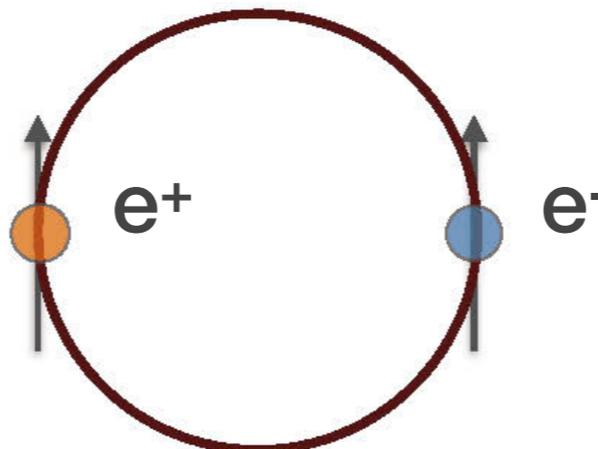
- Hydrogen
- 1420 MHz
- Atomic hydrogen maser



- Muonium
- 4464 MHz
- Decay positron asymmetry



- Positronium
- 203 GHz
- Annihilation gamma rays



- Muonium provides the most precise theory/experiment comparison.

- Theoretical prediction of muonium HFS

$$\Delta\nu_{\text{Mu}} = 4\,463\,302\,868(271) \text{ Hz}$$

- Contributions from each sector

$$\Delta\nu_{\text{Mu}} (\text{QED}) = 4\,463\,302\,700 (271) \text{ Hz}$$

$$\Delta\nu_{\text{Mu}} (\text{QCD}) = 232.7 (1.4) \text{ Hz}$$

$$\Delta\nu_{\text{Mu}} (\text{EW}) = -65 \text{ Hz}$$

- Theoretical uncertainty

| Term                        | Contribution (Hz) |
|-----------------------------|-------------------|
| Radiative correction        | 5                 |
| Recoil correction           | 64                |
| Radiative recoil correction | 55                |
| Hadronic correction         | 1.4               |
| Total                       | 85                |

- Dominant uncertainties are arising from errors in physical constants (fine structure constant, Rydberg constant, and muon mass).

- By comparing the theoretical expression and experimental result of the muonium HFS, the muon-to-electron mass ratio can be extracted.

$$\nu_{\text{theory}} = \frac{16}{3} \alpha^2 c R_{\infty} \frac{m_e}{m_{\mu}} \left[ 1 + \frac{m_e}{m_{\mu}} \right]^{-3} + \delta(\alpha, m_e/m_{\mu})$$

- The muon-to-electron mass ratio is one of the necessities to the experimental determination of the muon anomalous magnetic moment (muon  $g-2$ ).

$$a_{\mu}(\text{Exp.}) = 11\,659\,208.9(6.3)$$

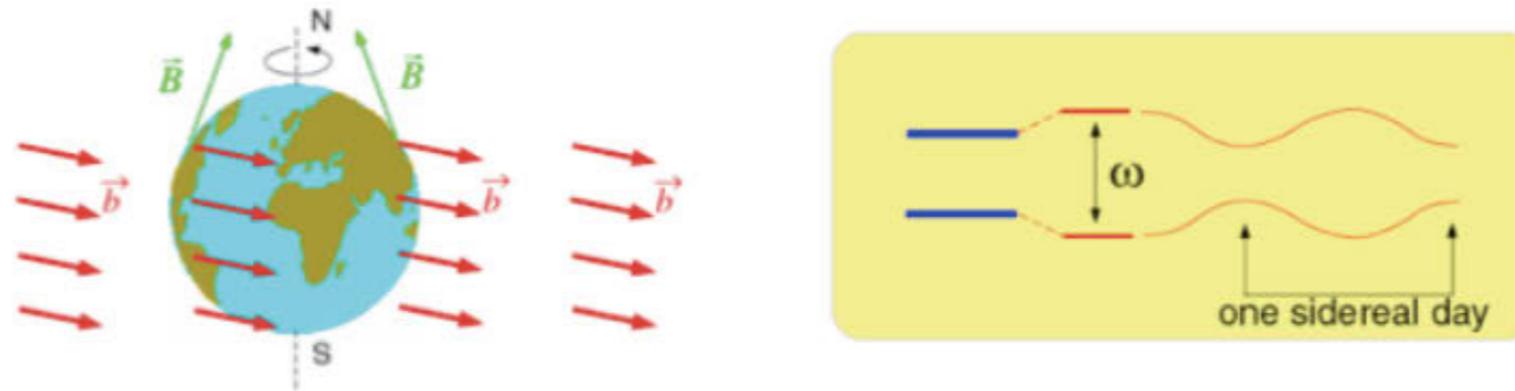
$$a_{\mu}(\text{Theory}) = 11\,659\,181.8(4.9)$$

$$a_{\mu} = \frac{g_e}{2} \frac{\omega_a}{\omega_p} \frac{\mu_p}{\mu_e} \frac{m_{\mu}}{m_e}$$

P. J. Mohr, D. B. Newell, and B. N. Taylor, "CODATA2014", Rev. Mod. Phys. 88, 035009 (2016).

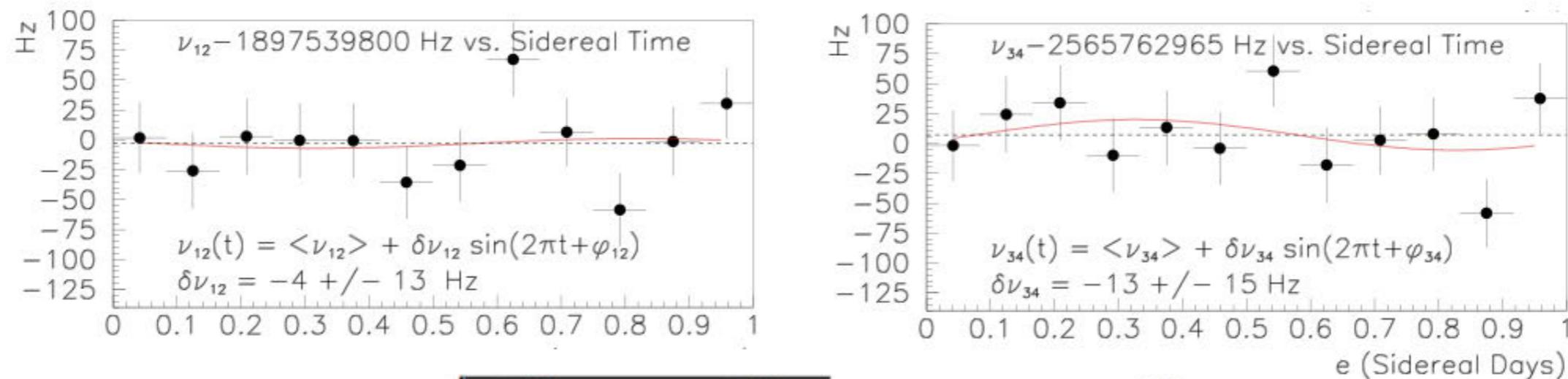
G.W. Bennett *et al.*, Phys. Rev. D73, 072003 (2006). K. Hagiwara *et al.*, JPHGB G38, 085003 (2011).

- Sidereal oscillation of the transition frequency means Lorentz invariance violation.



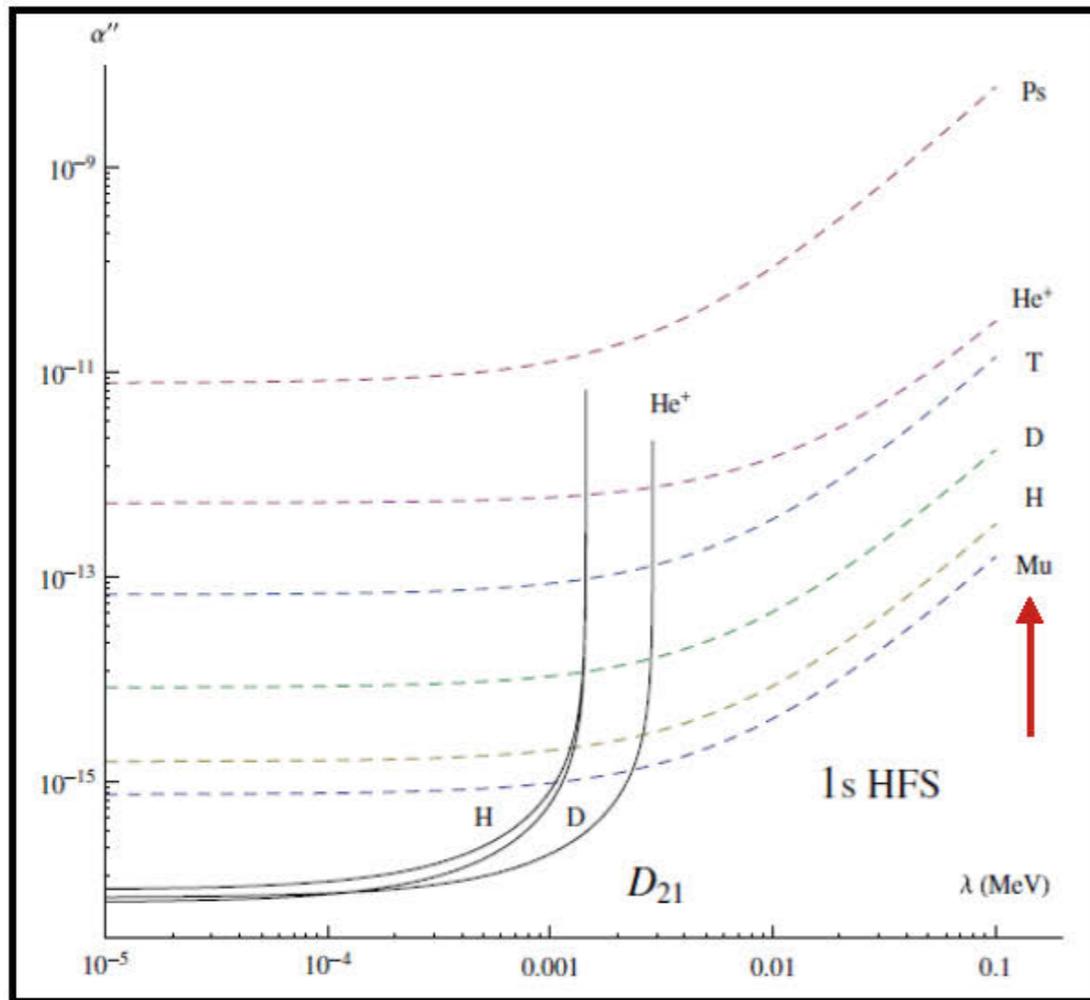
R. Bluhm, V.A. Kostelecký, and C. Da Lane, Phys. Rev. Lett. 84, 1098 (2000).

- Limits on the model parameters from the muonium HFS.



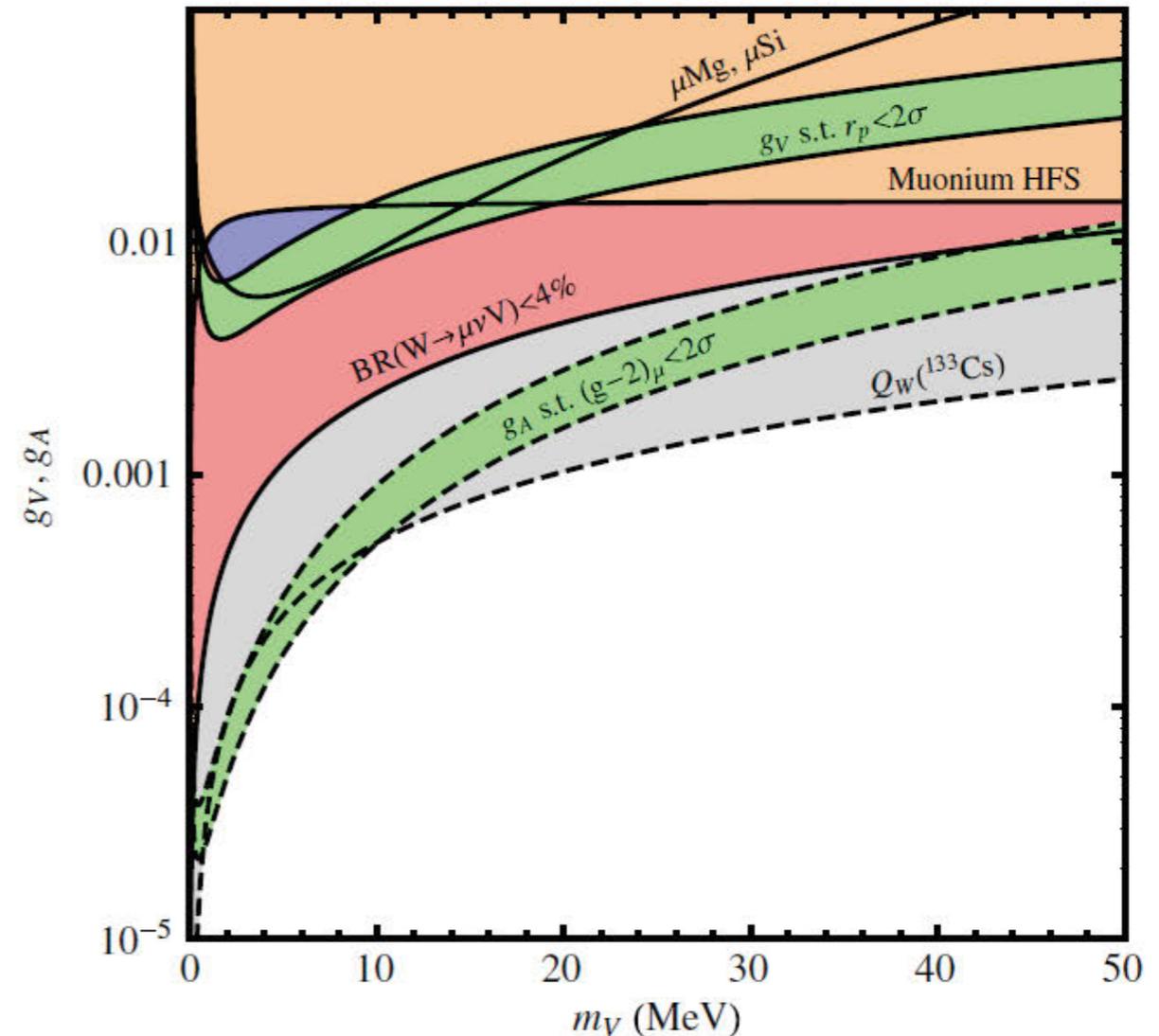
$$|\sin\chi| \sqrt{(\tilde{b}_X^\mu)^2 + (\tilde{b}_Y^\mu)^2} \lesssim 2 \times 10^{-22} \text{ GeV}$$

V.W. Hughes *et al.*, Phys.Rev.Lett.87, 111804 (2001).



## Light pseudo-vector boson

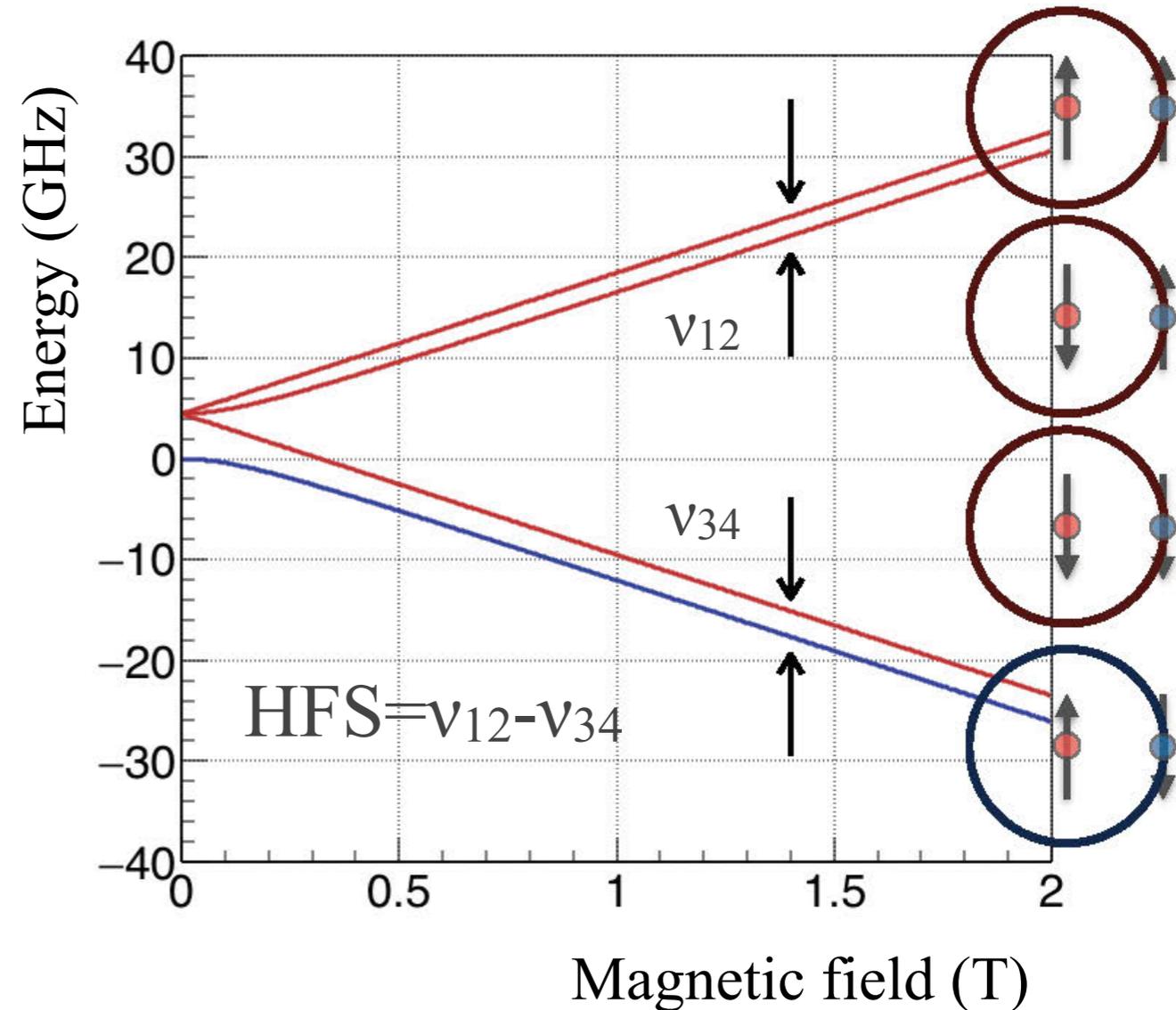
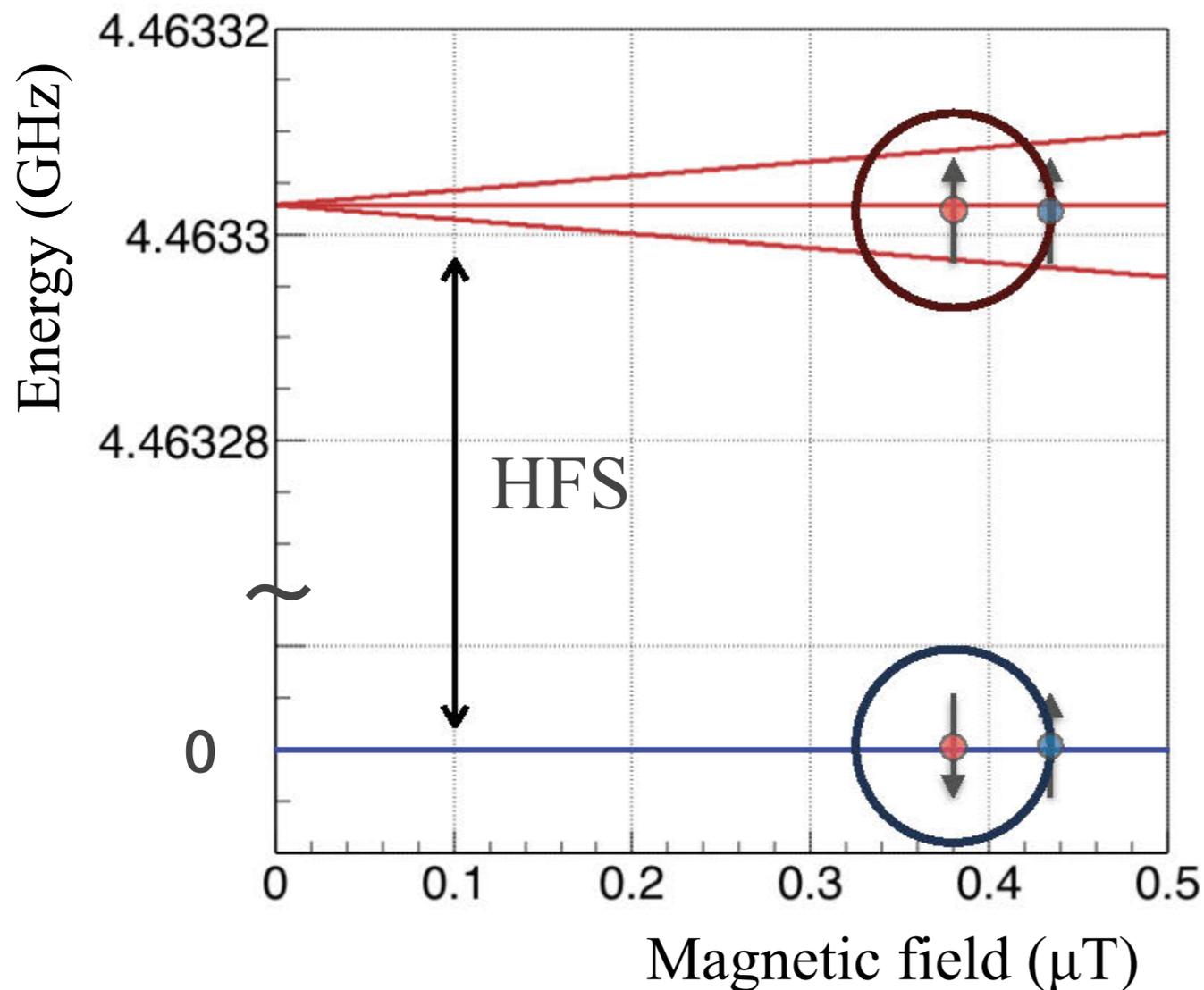
S. G. Karshenboim, PRL 104, 220406 (2010).



## Massive vector boson

S. G. Karshenboim, PRD 90, 073004 (2014).

- Muonium HFS constraints “dark” force carriers.
- Test of lepton universality via a search for muon-specific force.

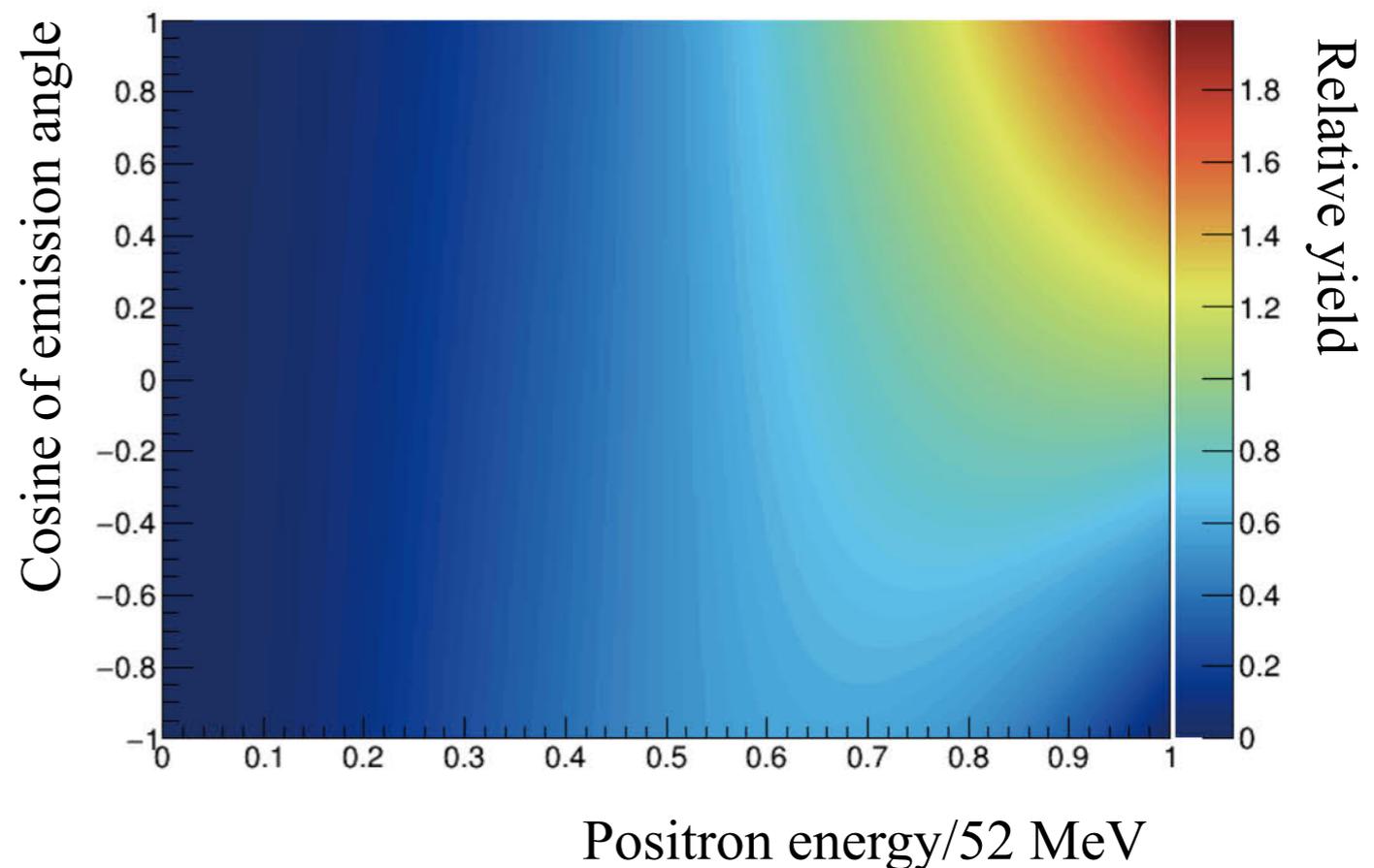
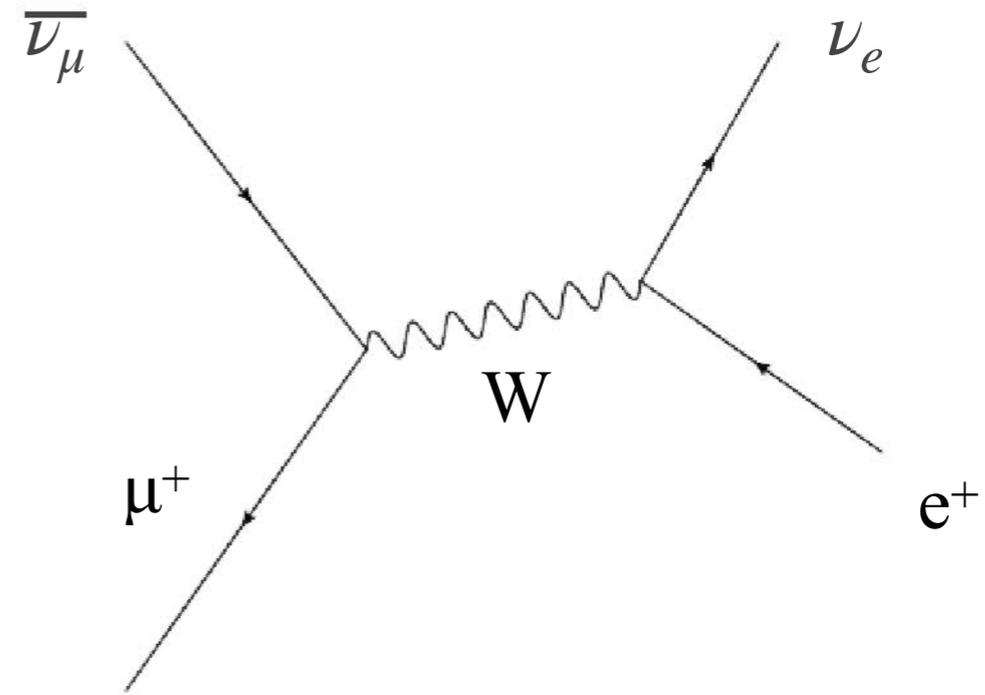


- Two independent methods for the hyperfine splitting measurement
  - Direct measurement at “zero” magnetic field (ZF, in progress).
  - Indirect measurement in a high magnetic field (HF, in preparation).
- **Our goal is ten-fold of improvements in both ZF and HF experiments.**

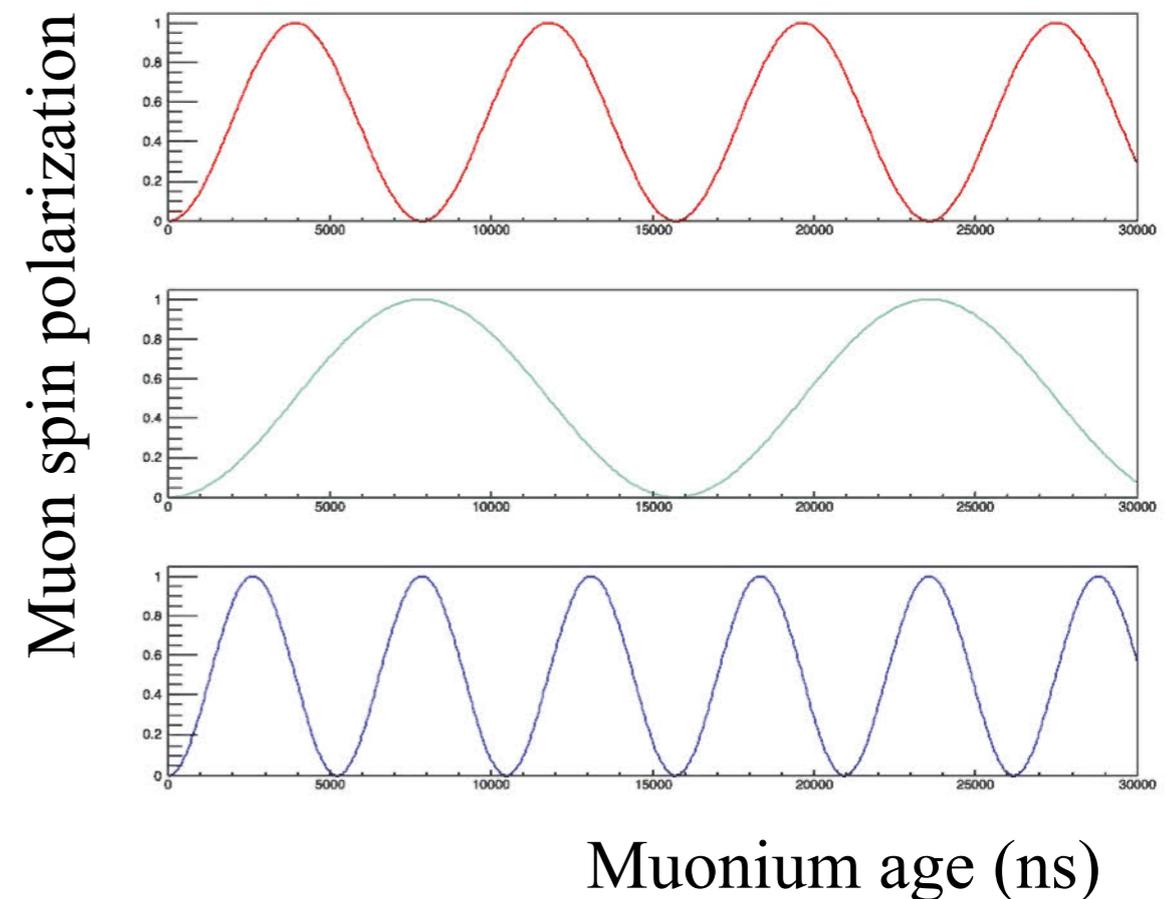
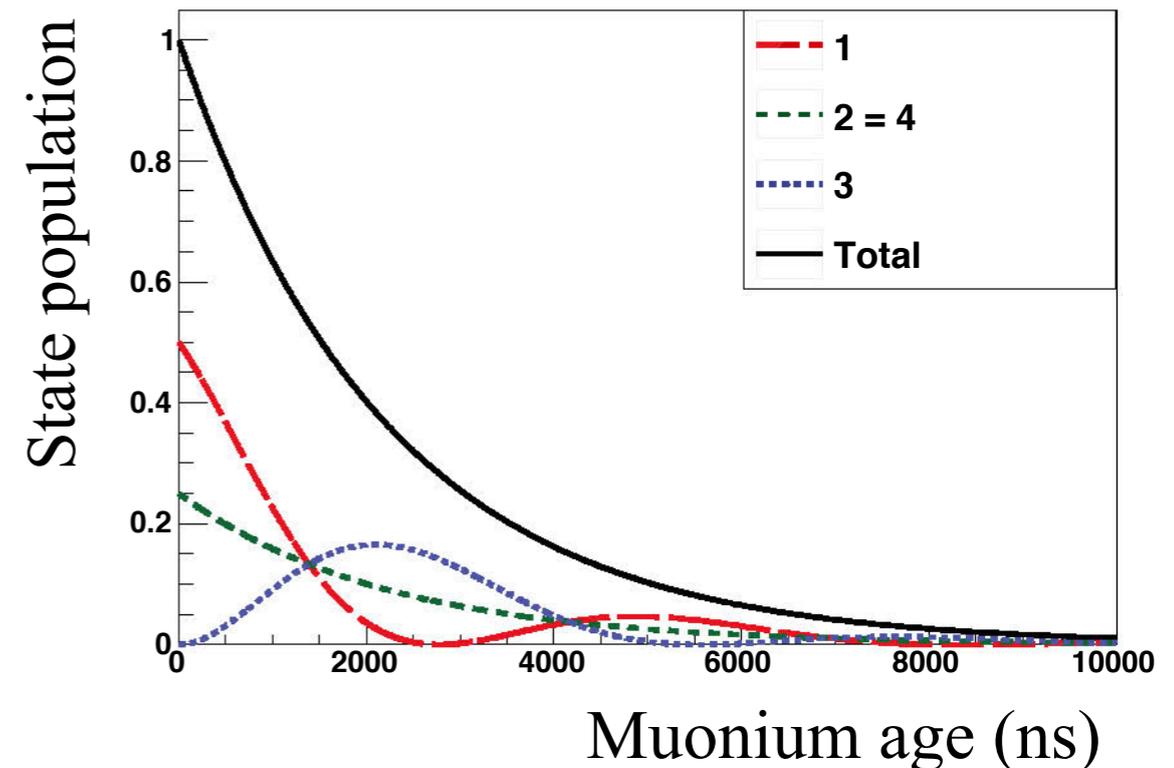
# Parity Violating Muon Decay

9

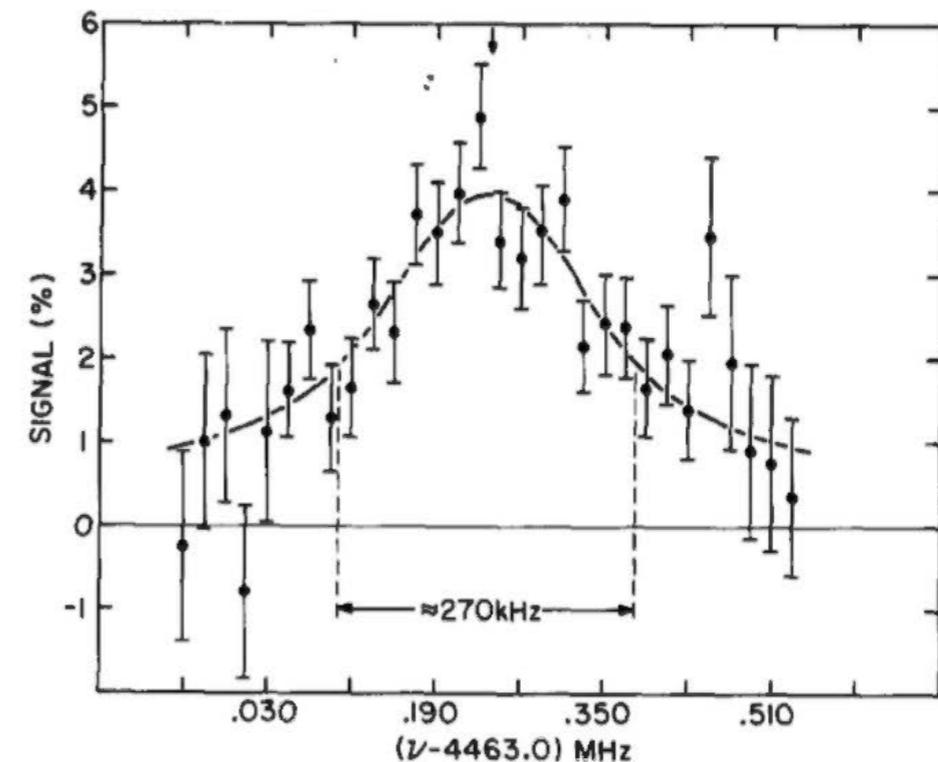
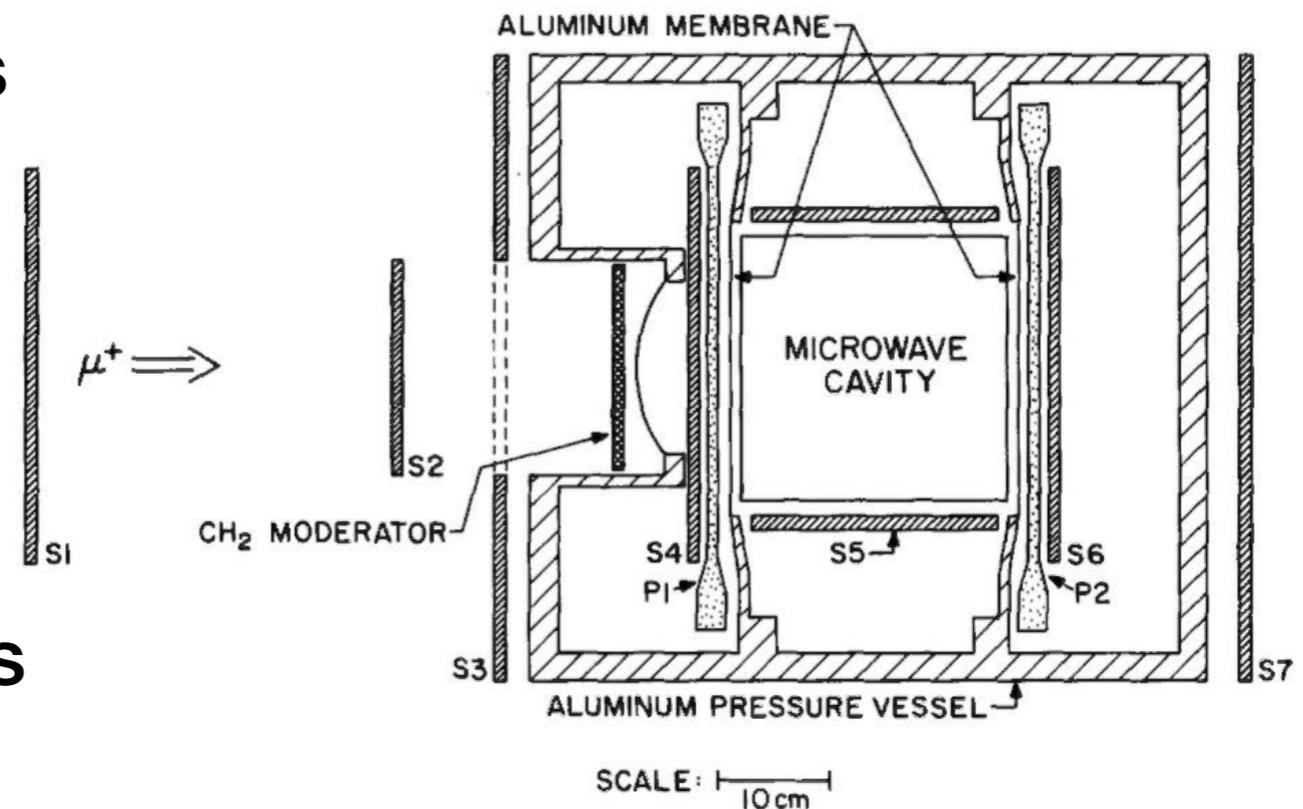
- Muon decay is mediated by the weak boson.
- Muon spin and emission angle of decay positron correlate.
- More positrons are emitted in parallel with the muon spin direction.
- **By measuring the angular asymmetry of the decay positron, the muon spin can be obtained as an ensemble average.**



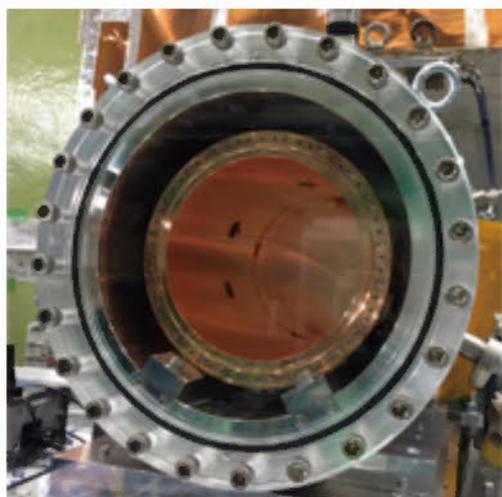
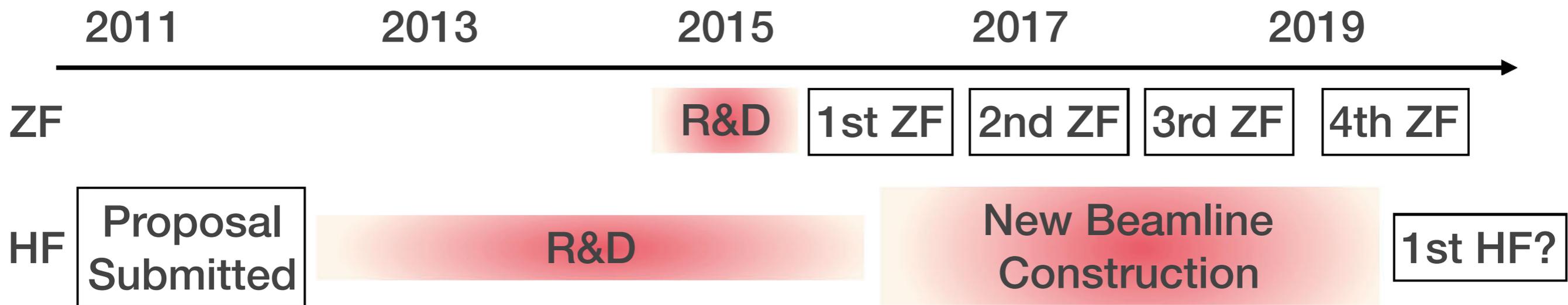
- Muon spin flips after the hyperfine transition induced by the microwave irradiation.
- State population of muonium changes as a function of muonium age.
- Muon spin polarization oscillates with the Rabi frequency.
- Typical oscillation frequency is on the order of MHz.
- **A time-window of several  $\mu\text{s}$  is required to observe the signal.**



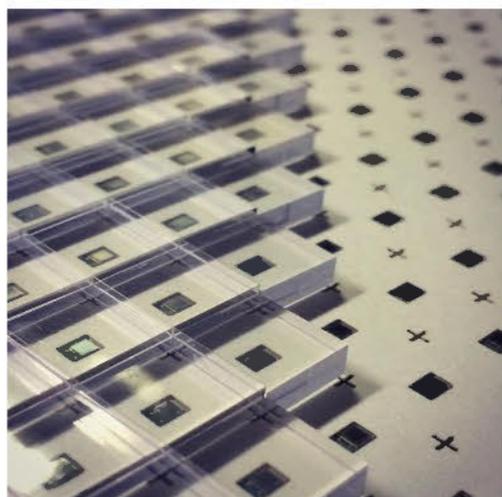
- An experiment at Los Alamos National Laboratory.
- A continuous muon beam irradiates the krypton gas target.
- The muon beam intensity was limited to avoid muon pileup.
- Statistical uncertainty was 1.4 kHz with 600 hours of measurement.
- No quantitative estimation of the systematic uncertainties.
- **High-intensity pulsed muon beam is a key for improvement.**



- Muonium Spectroscopy Experiment Using Microwave
  - Zero field: Demonstration at existing beamline.
  - High field: Highest precision experiment at dedicated beamline under construction.



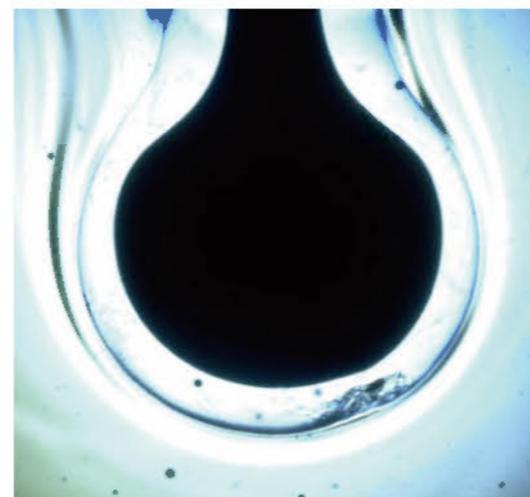
RF Cavity



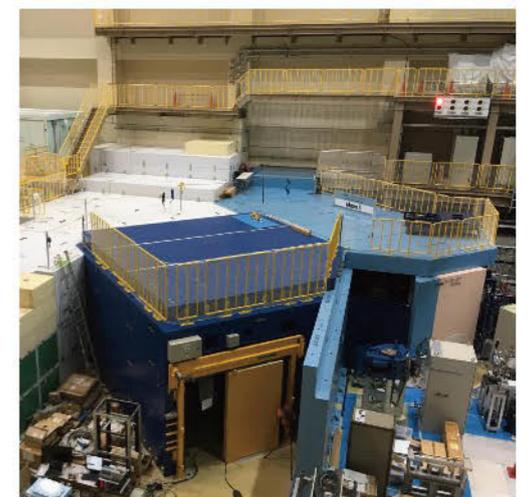
Detector



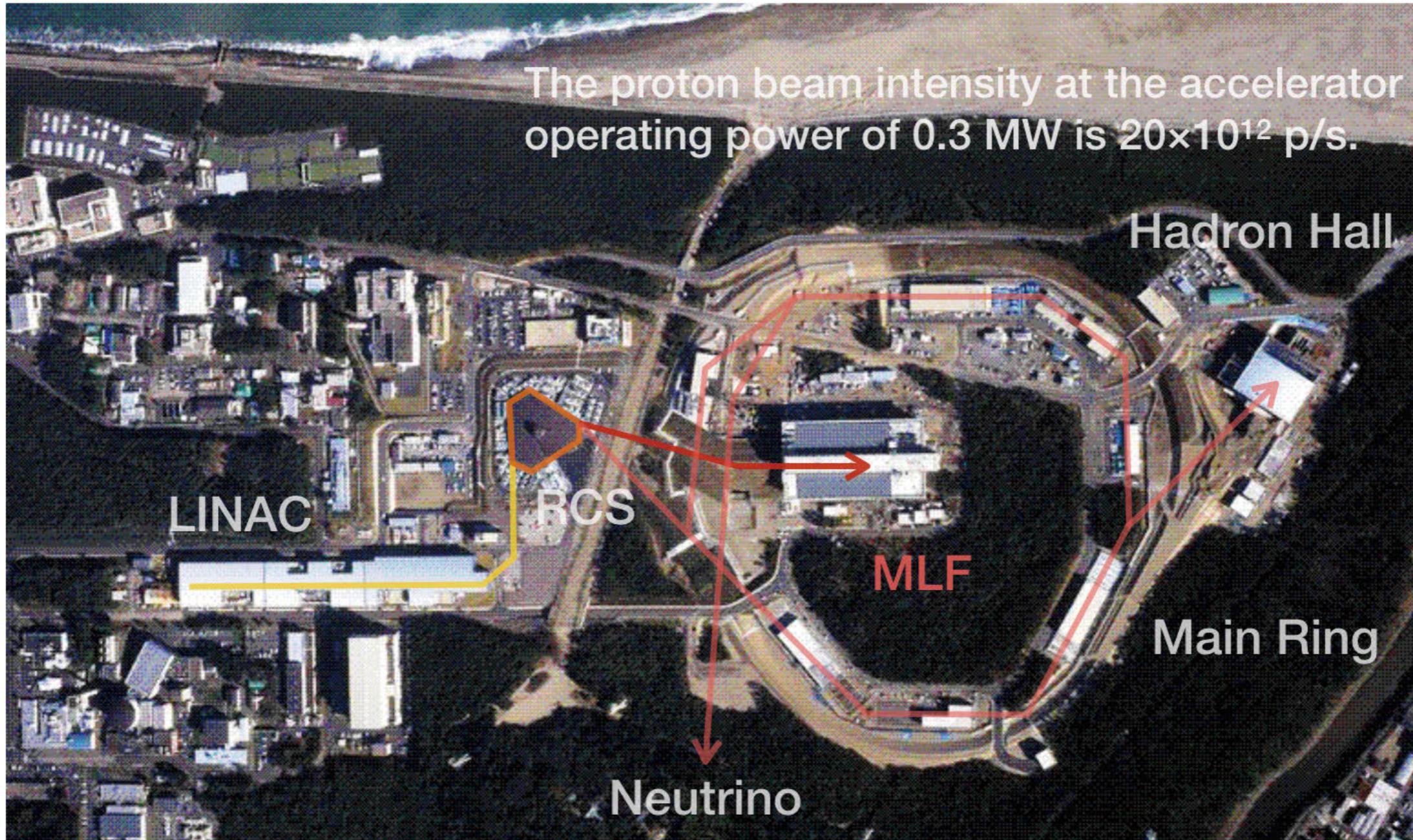
Magnet



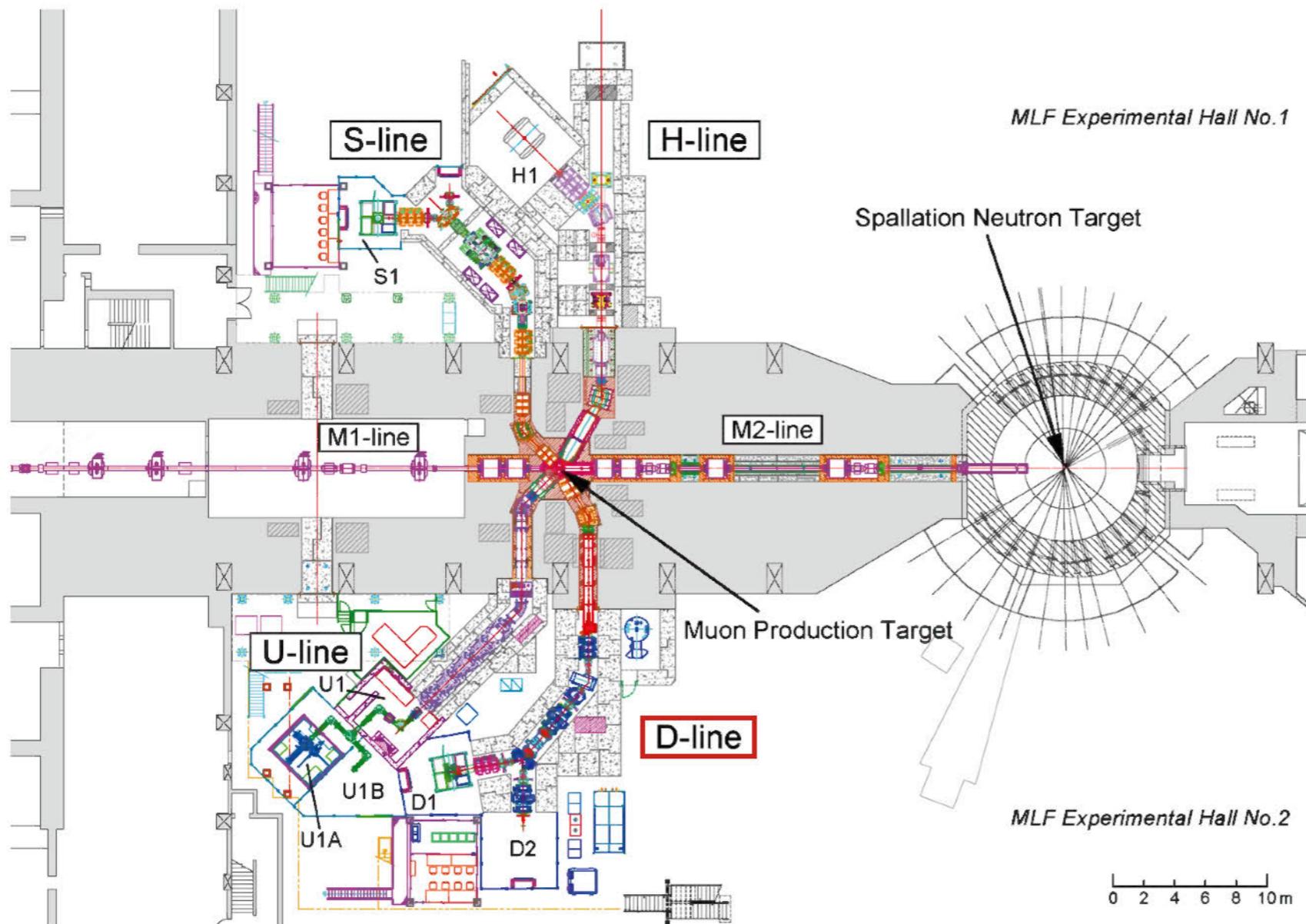
Field probe



New beamline



- **World's highest intensity proton driver.**
- **400 MeV LINAC, 3 GeV RCS, and 30 GeV Main Ring.**
- **Pulsed proton beam with 25 Hz repetition.**



- **Material and Life Science Facility**

- Graphite target for muon production

- Four muon beam lines deliver intense  $\mu^+$  and  $\mu^-$  beams.

- $3 \times 10^6 \mu^+/\text{s}$  at D-Line (0.3 MW).

- D ( $\mu^+$ ,  $\mu^-$ )

- Muon spin rotation

- Muonic X-rays

- General purpose

- U ( $\mu^+$ )

- Ultra slow muon

- S ( $\mu^+$ )

- Muon spin rotation

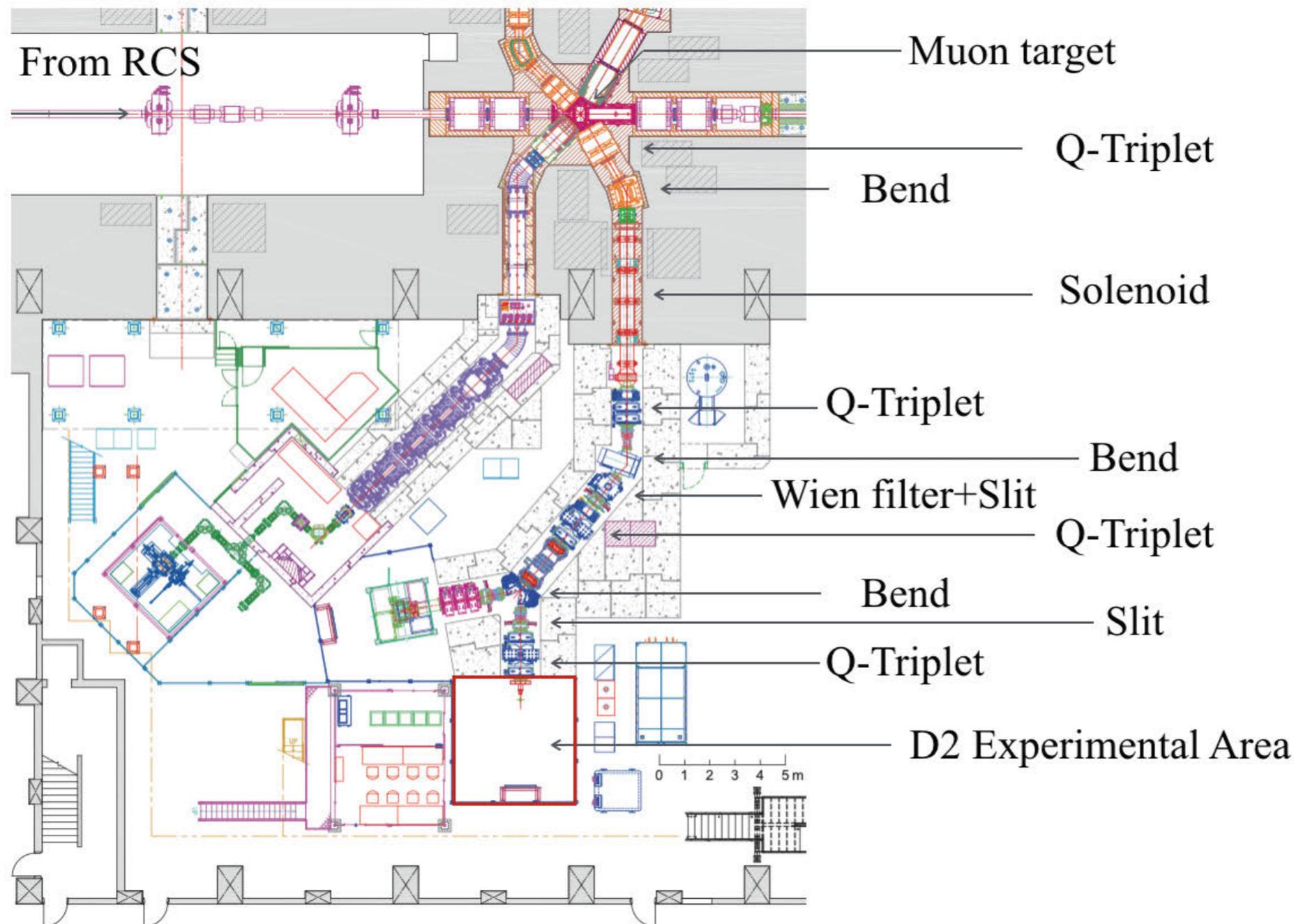
- H ( $\mu^+$ ,  $\mu^-$ )

- Muonium HFS

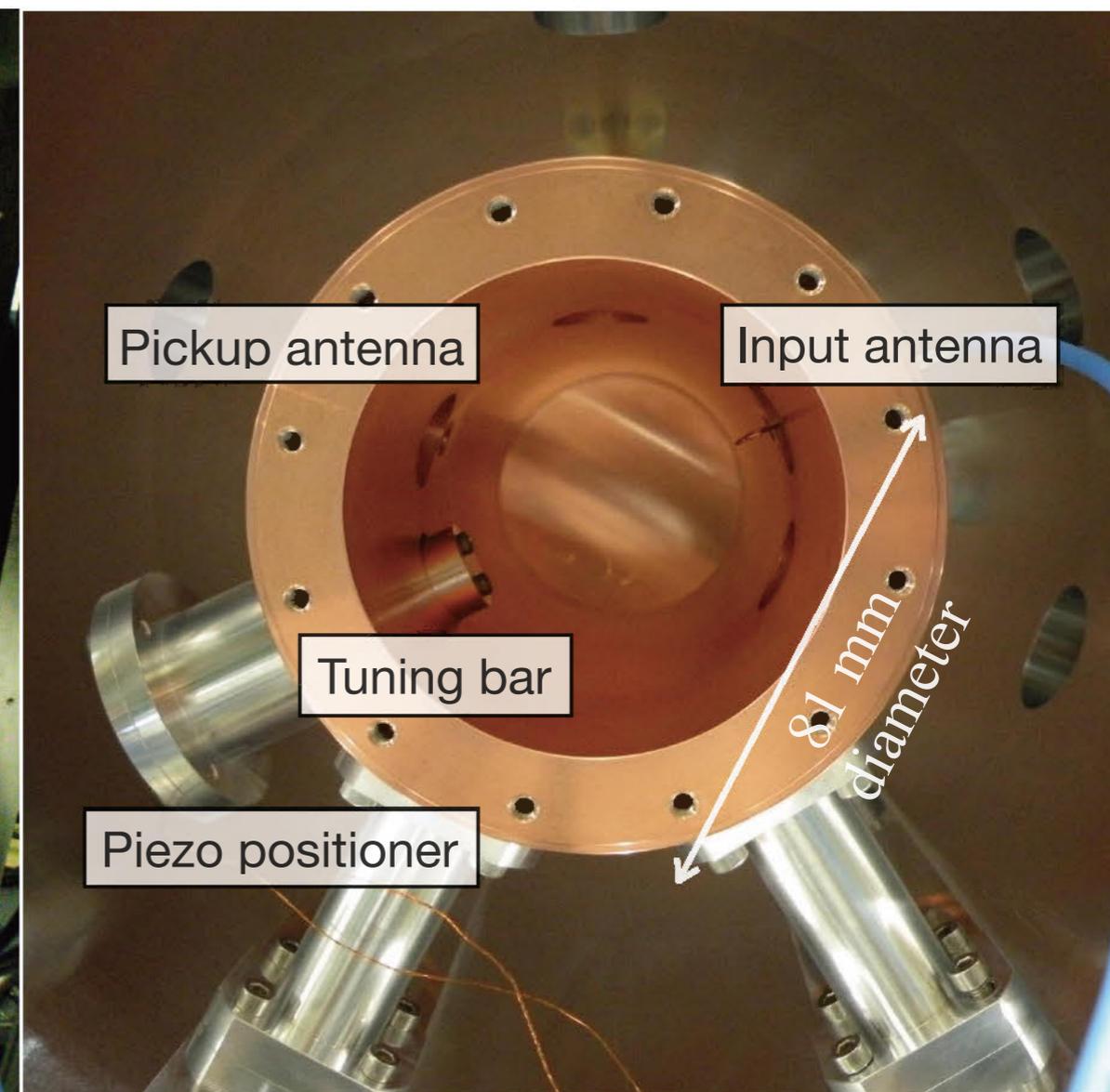
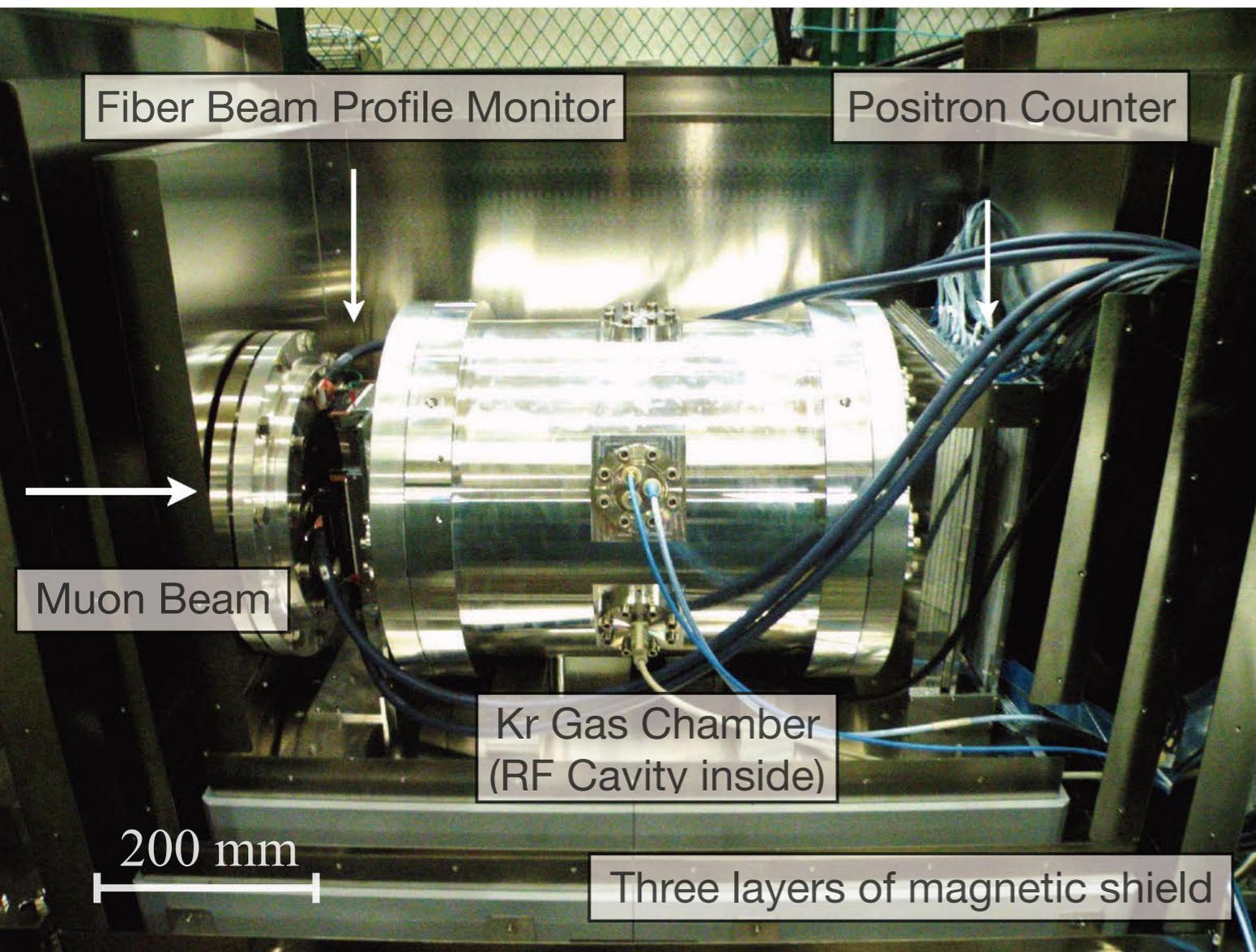
- Muon-to-electron conversion

- Muon g-2

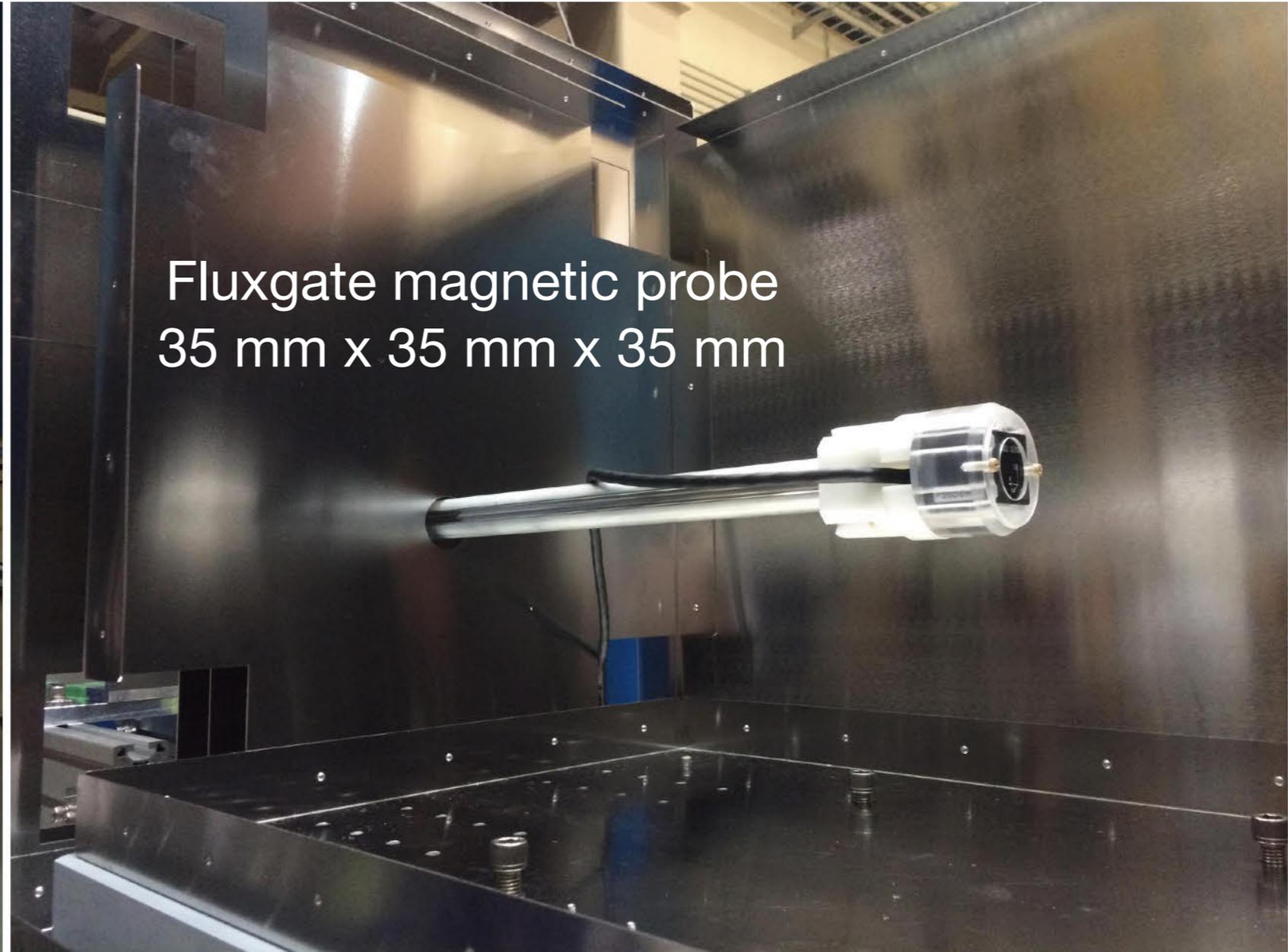
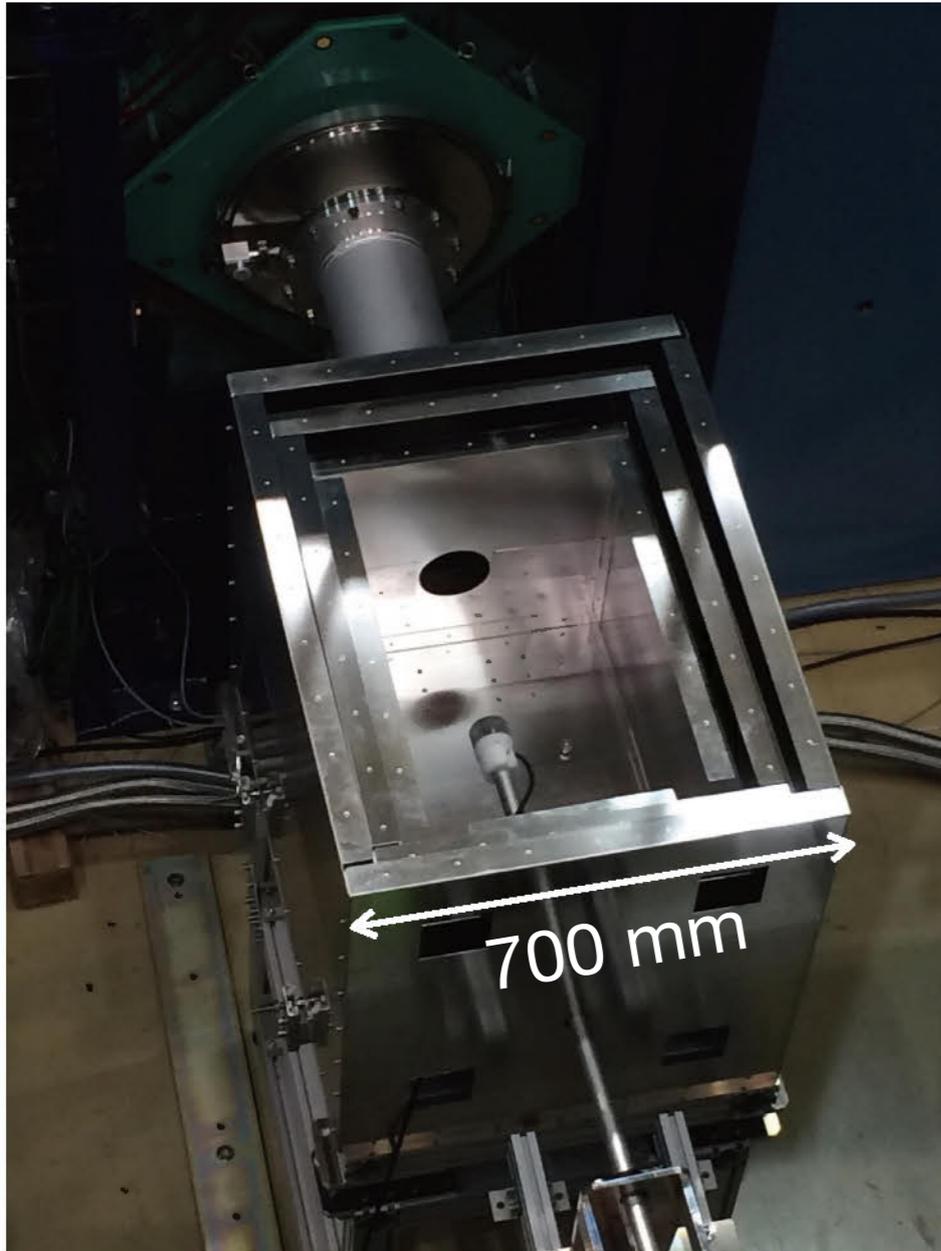
- Muon transmission microscope



- **Multi-purpose beam line for intense, polarized muon beam.**
- **Positive muons produced on the target surface with the momentum of 27.4 MeV/c are transported.**

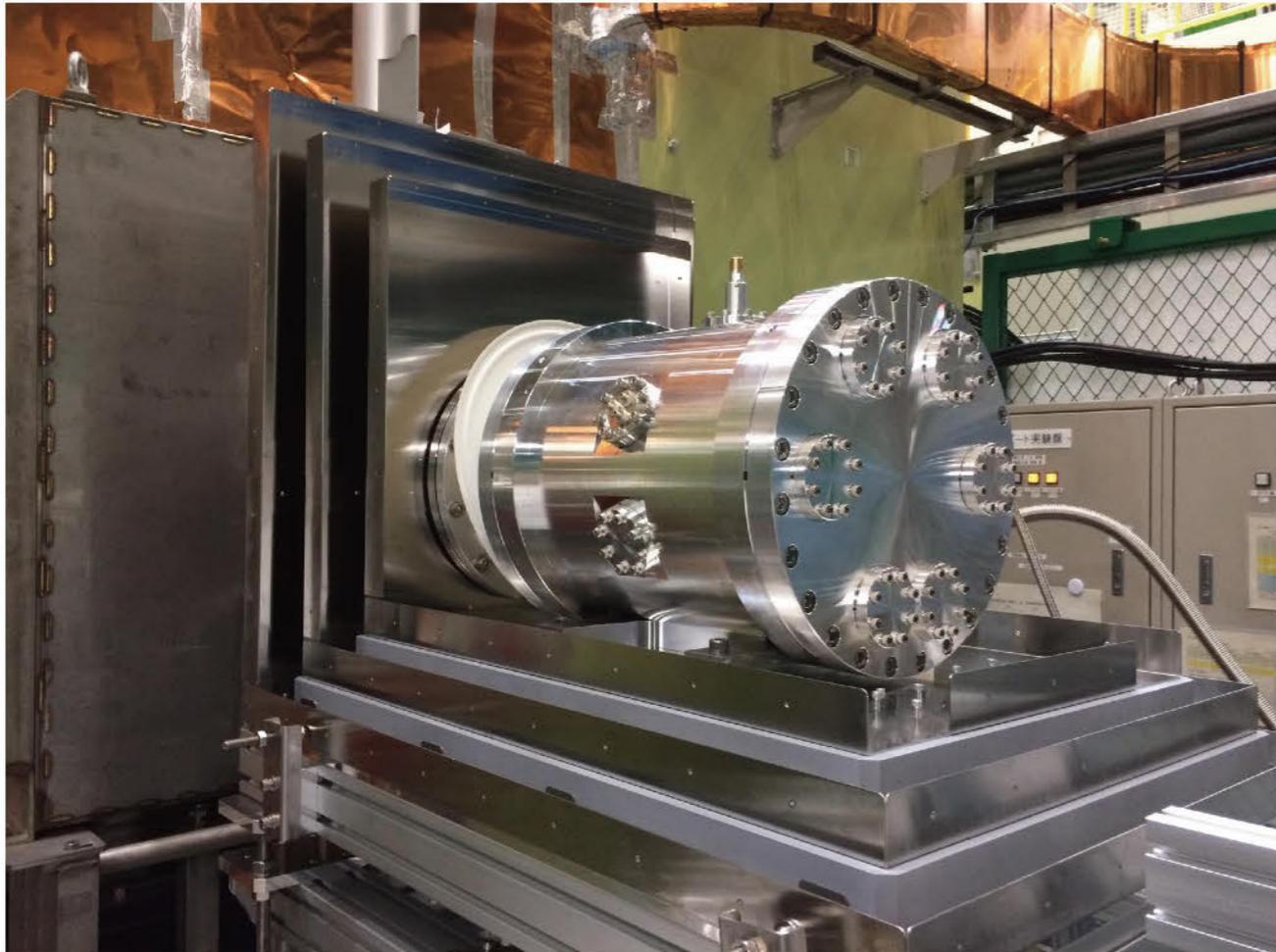


- The apparatus is installed at MLF MUSE D2 experimental area.
- The entire setup was surrounded by a magnetic shield to shield the leakage magnetic field from the beam line and surrounding equipments.

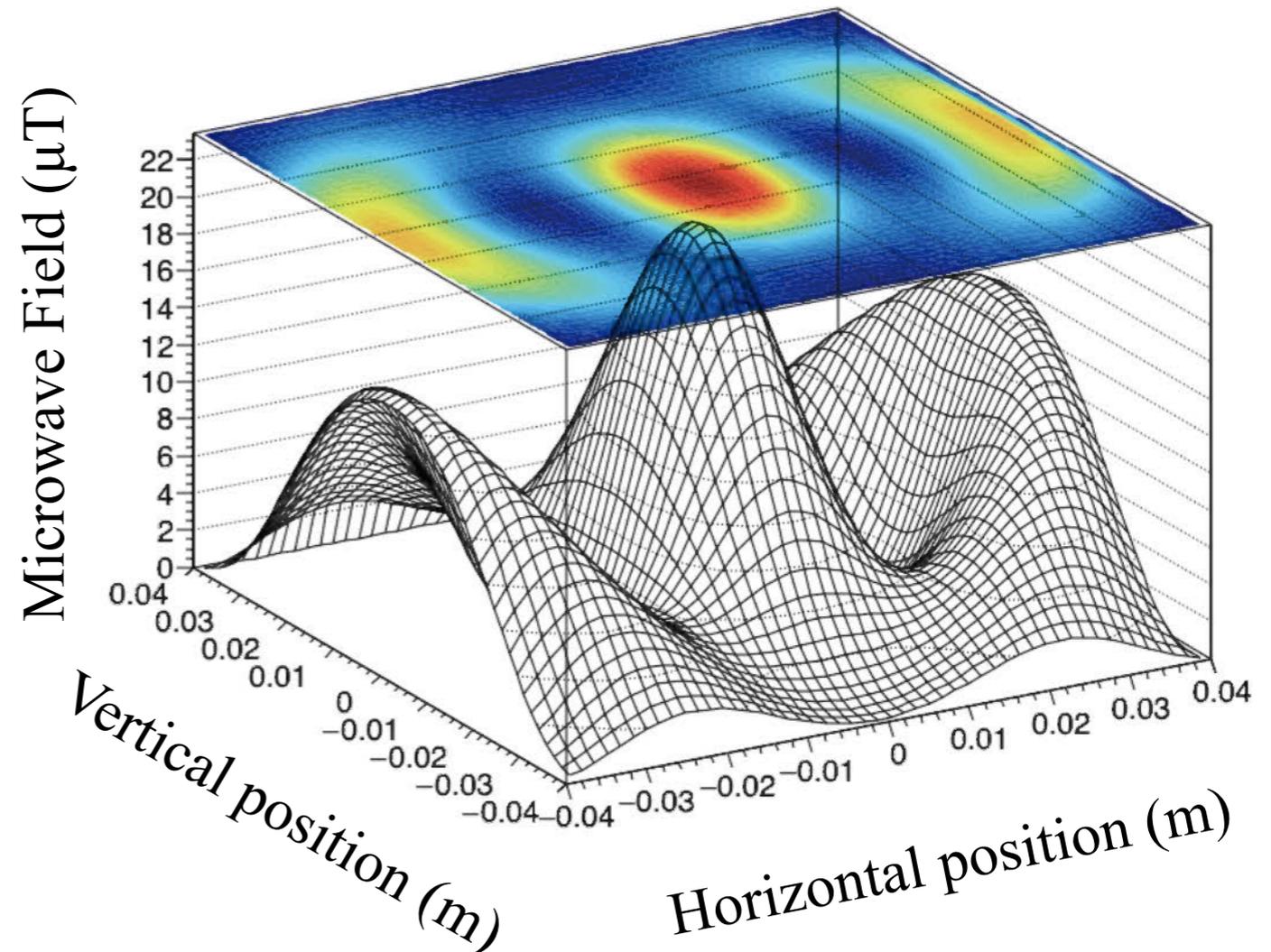


Fluxgate magnetic probe  
35 mm x 35 mm x 35 mm

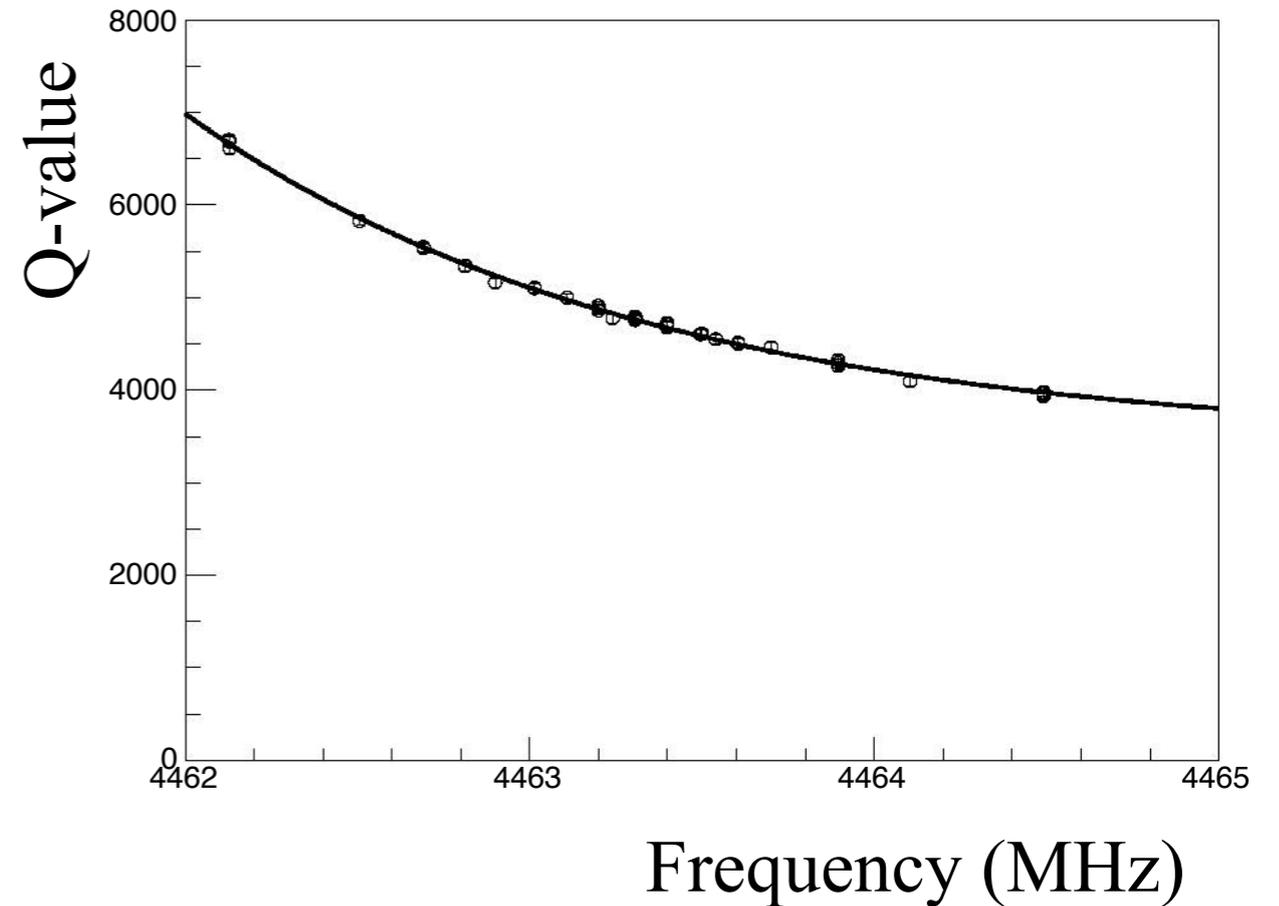
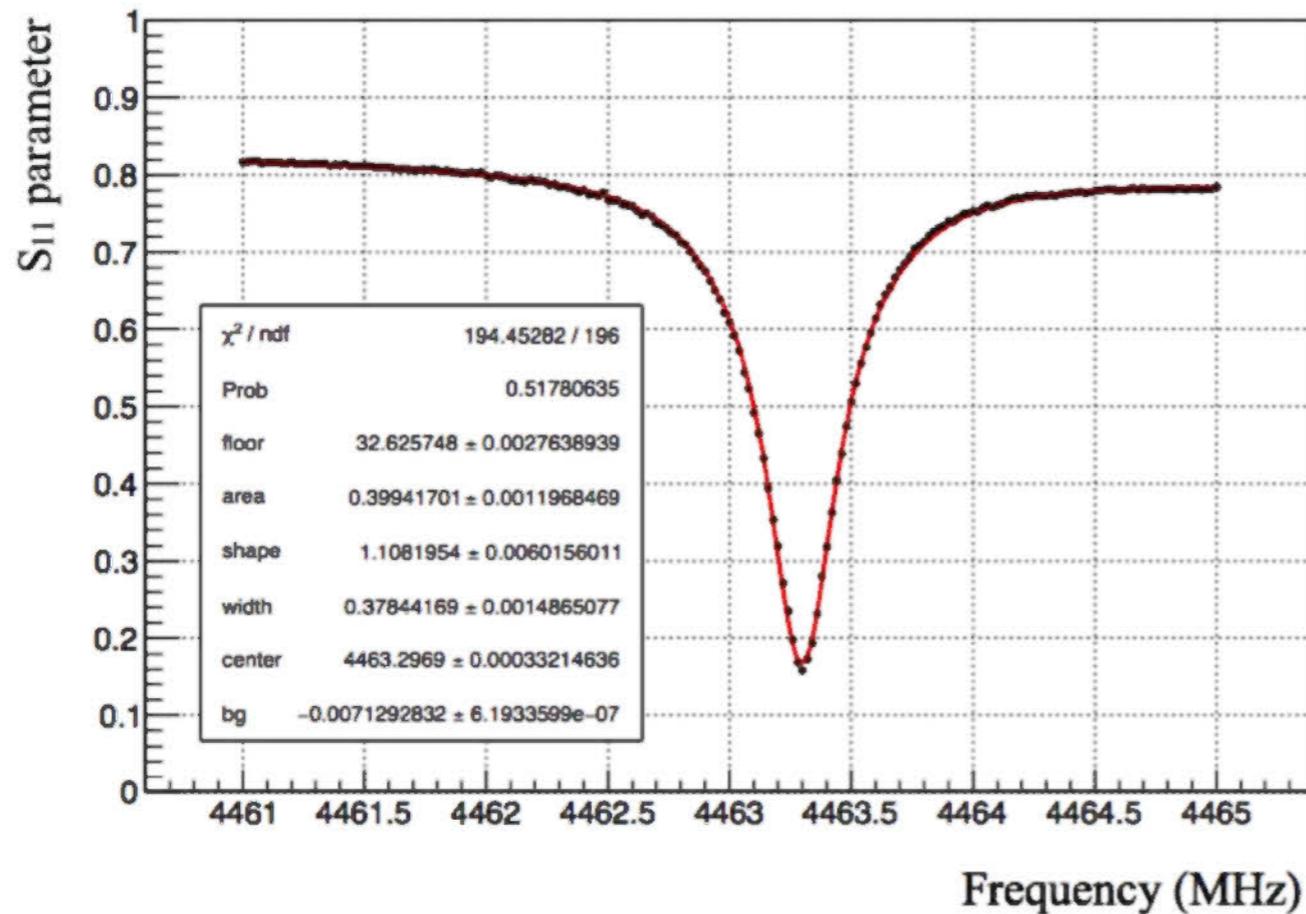
- A three-layer box-type magnetic shield made of an alloy of iron and nickel.
- A magnetic field inside the shield is measured by a fluxgate magneto probe.



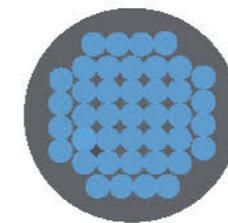
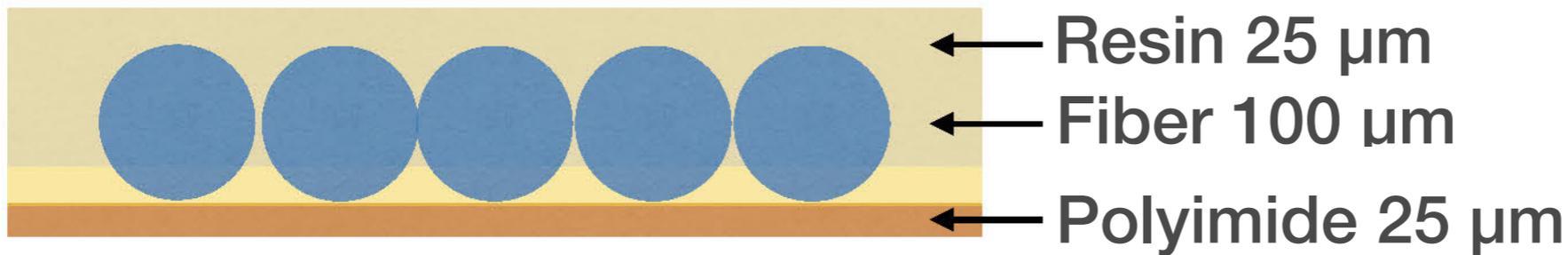
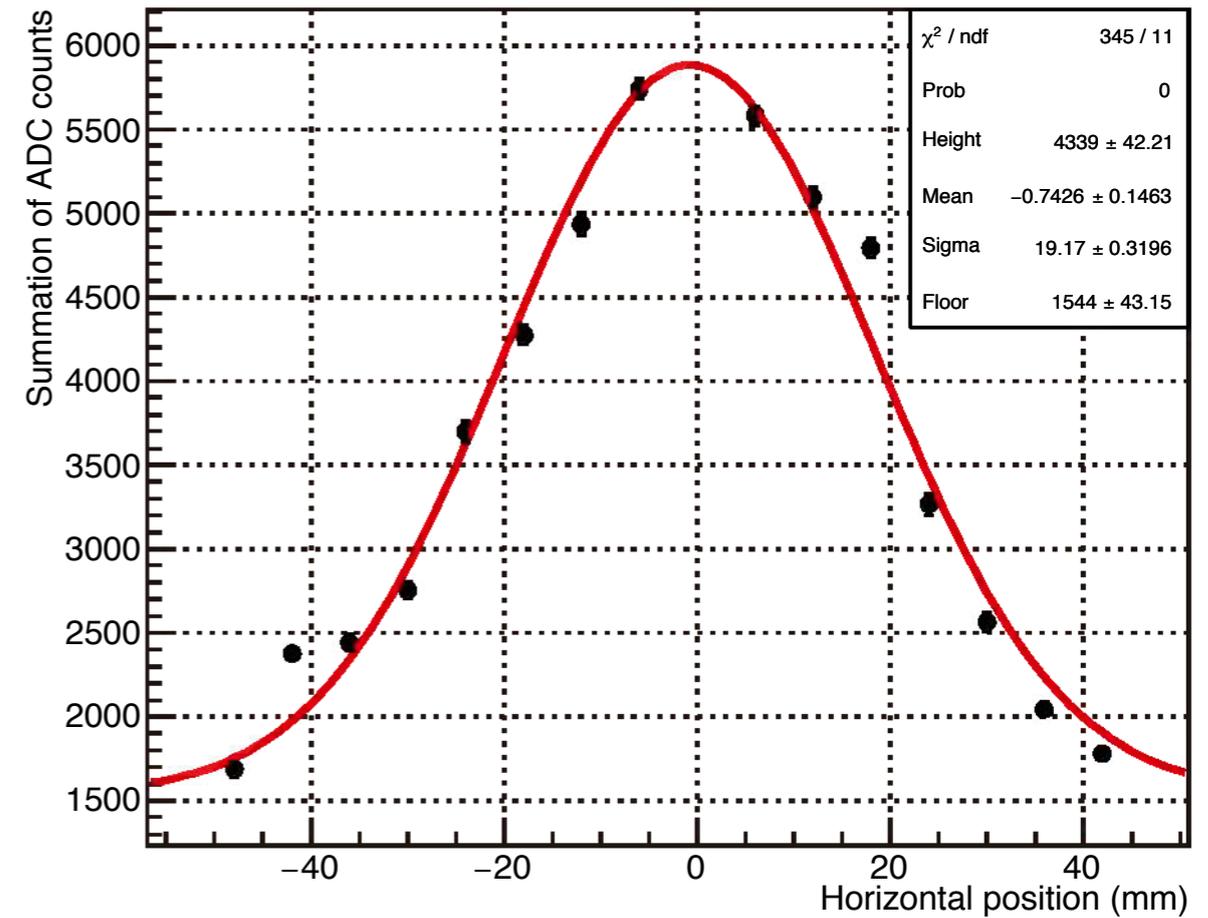
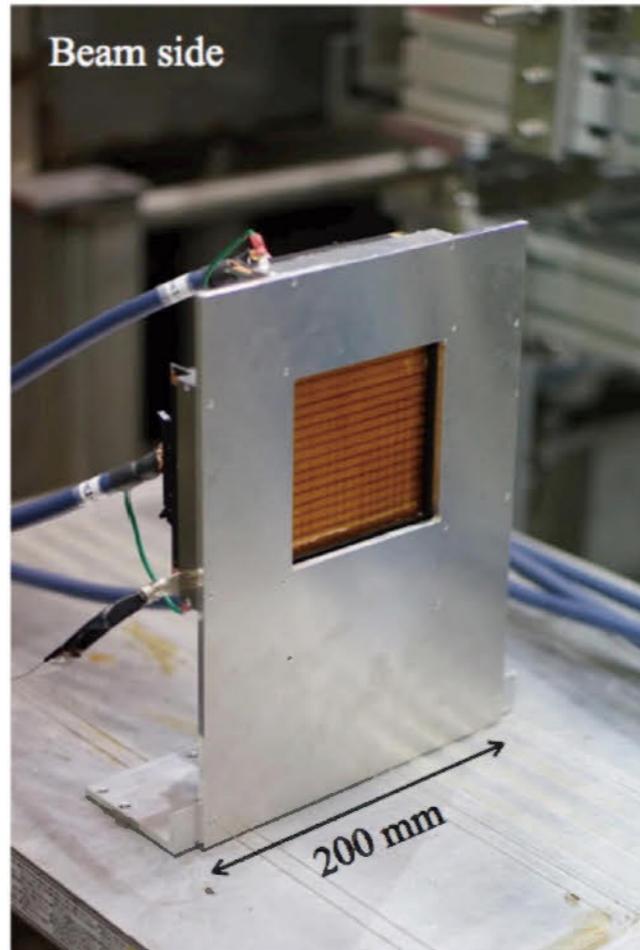
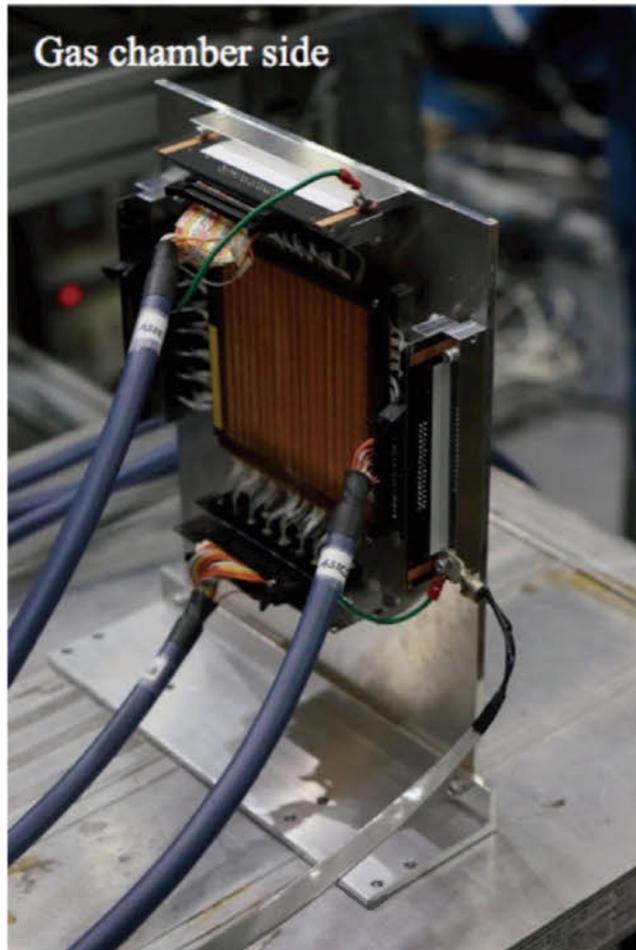
- A cylindrical vessel made of aluminum with a longitudinal length of 450 mm and an inner diameter of 280 mm.
- Target gas is krypton to optimize the energy threshold of muonium formation via electron capture.
- Gas handling panel is developed to fill the chamber with gas, evacuate the chamber, and monitor impurities using a quadrupole mass spectrometer.



- A cylindrical resonator made of oxygen-free copper with a diameter of 81 mm and an axial length of 230 mm.
- TM<sub>110</sub> mode is excited with typical power of 1 W.
- Loop antennas are placed to input and monitor the microwave.
- Resonance frequency is tuned by a piezoelectric positioner.



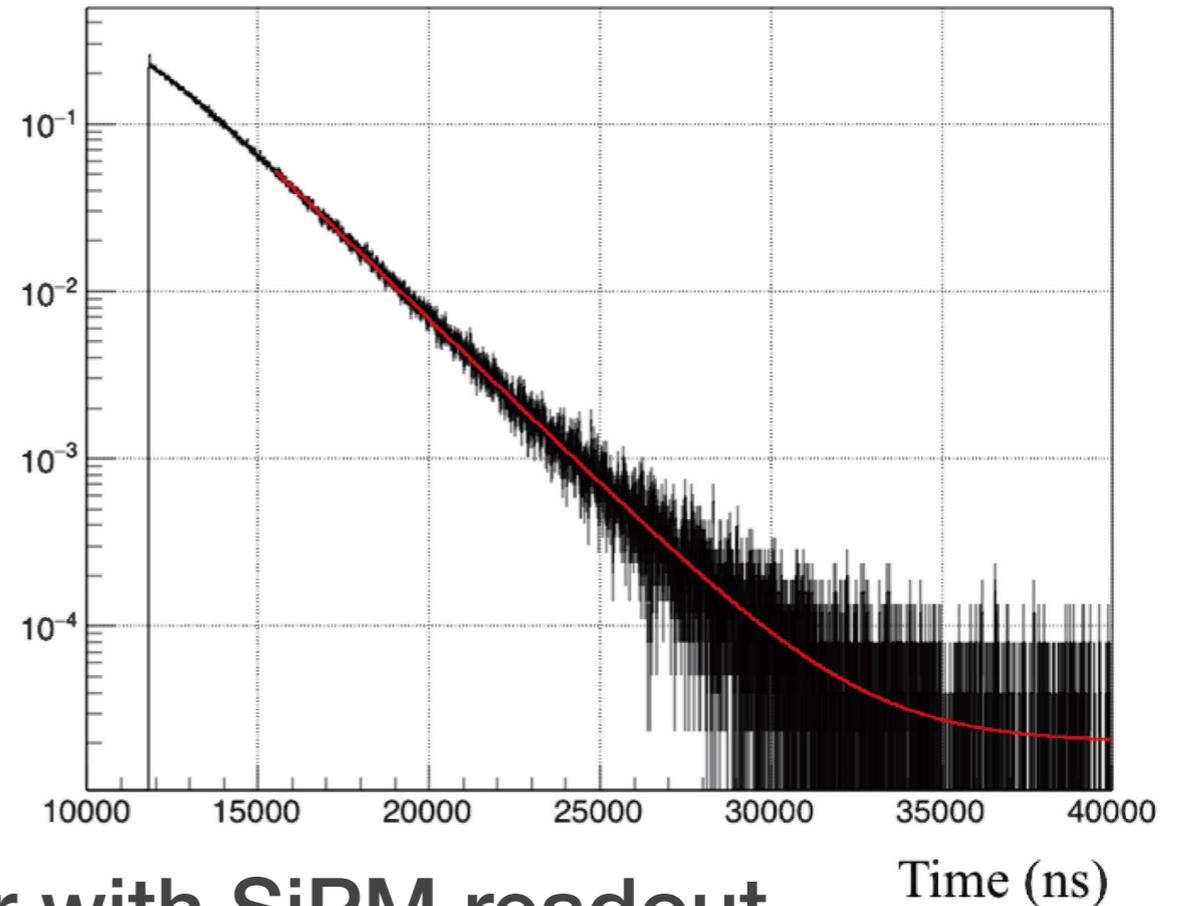
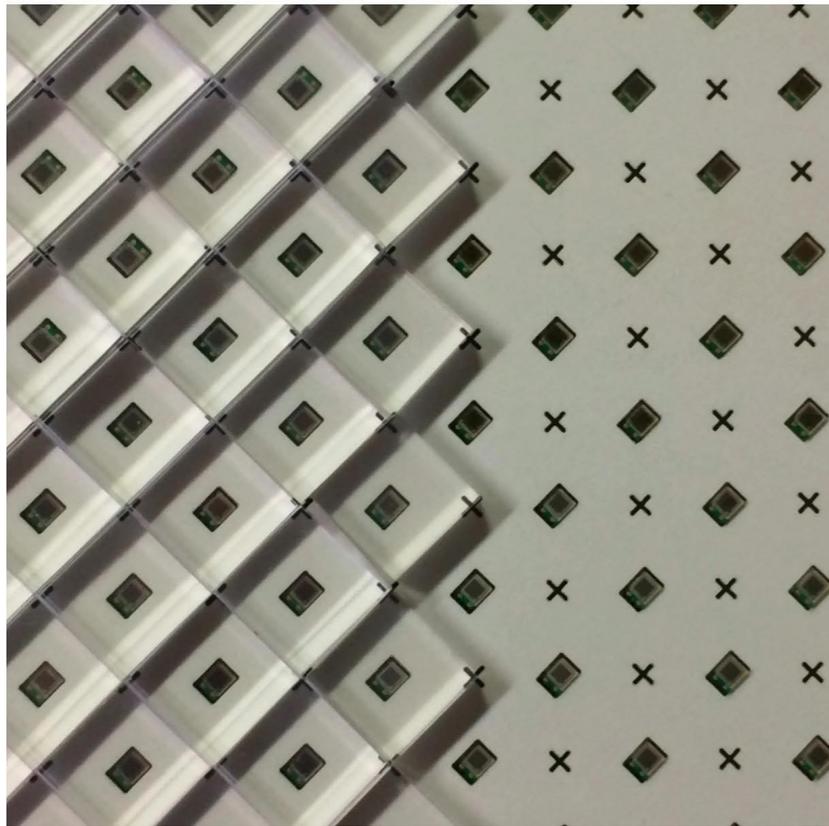
- Resonance of the cavity is observed using a vector network analyzer.
- $S_{11}$  parameter (reflection) is measured to evaluate the quality value.
- Resonance curve is fitted with a pseudo-Voigt function.
- Typical quality value is 5,000.
- Frequency dependence of the quality value is necessary to analyze a resonance lineshape.



40 fibers are bundled for a ch. and connected to MPPC

Fiber array layer structure

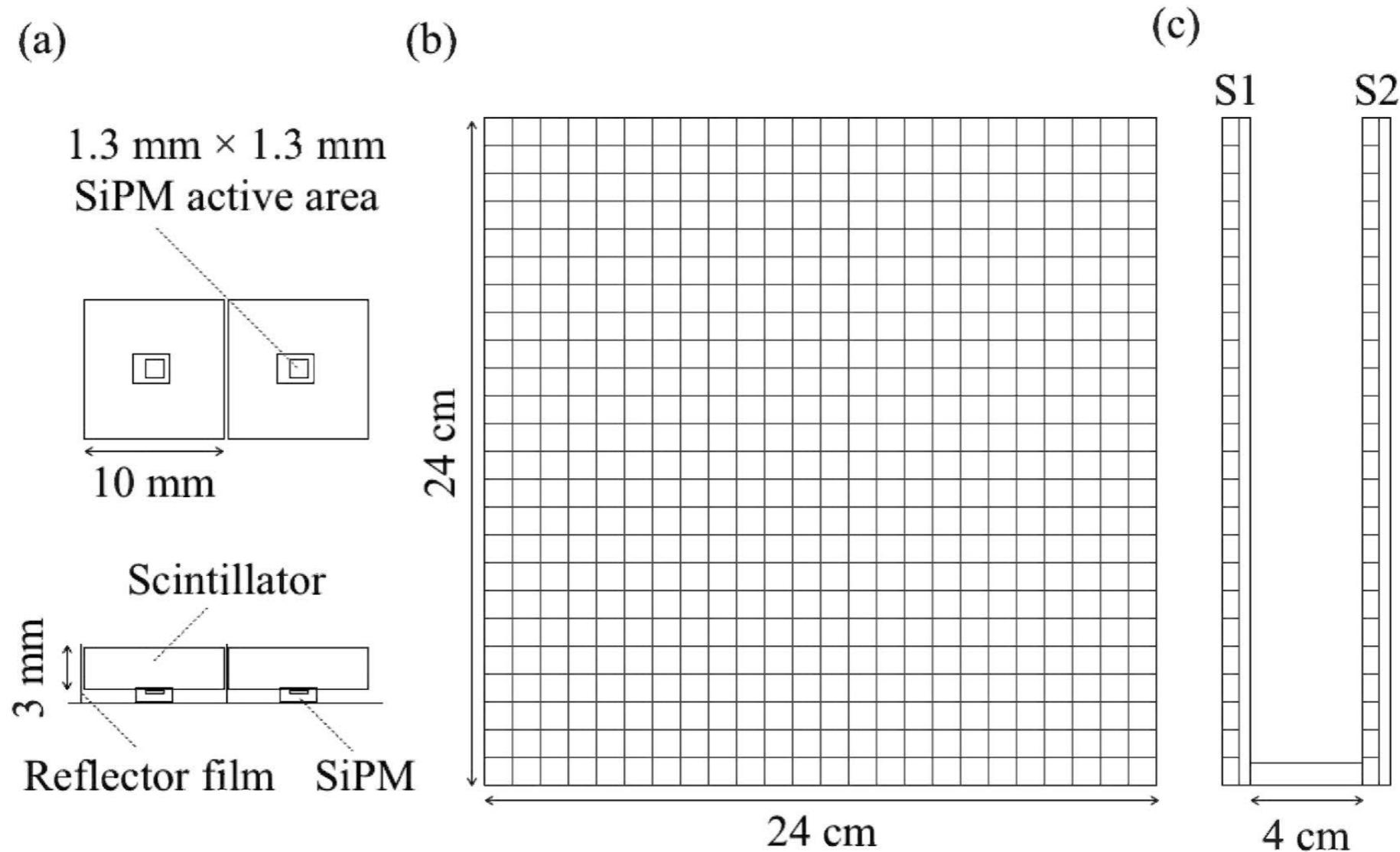
- Cross-configured fiber hodoscope with SiPM readout.
- Online measurement of the beam profile and relative intensity.



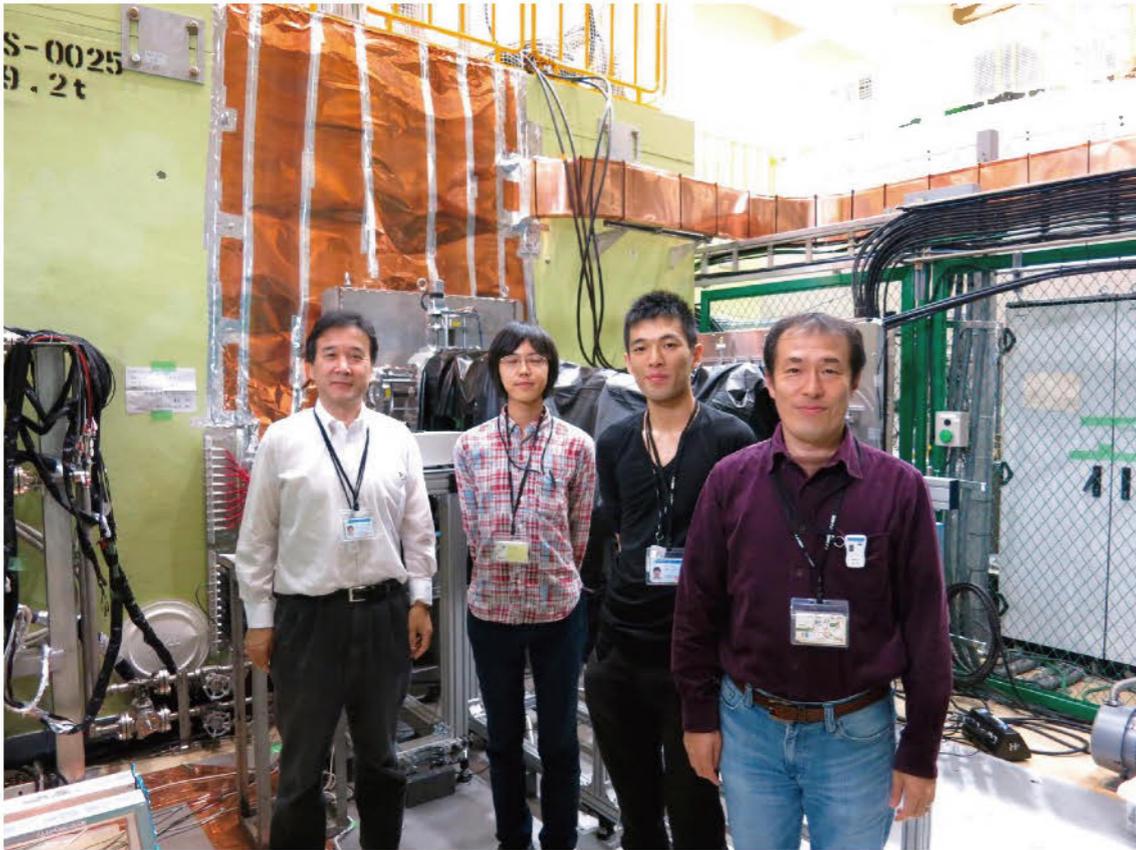
- Segmented scintillation counter with SiPM readout.
  - Unit cell has the dimension of 1 cm x 1 cm x 3 mmt.
  - Reflector film is inserted between scintillators.
  - 240 mm x 240 mm area, 1152 ch. in total.
- Amplifier, shaper, and discriminator are implemented in ASIC.
- FPGA-based multi-hit TDC.

S. Kanda, PoS(PhotoDet16)039 (2016).

K. M. Kojima, S. Kanda et al., J. Phys. Conf. Ser. 551 (2014) 012063.



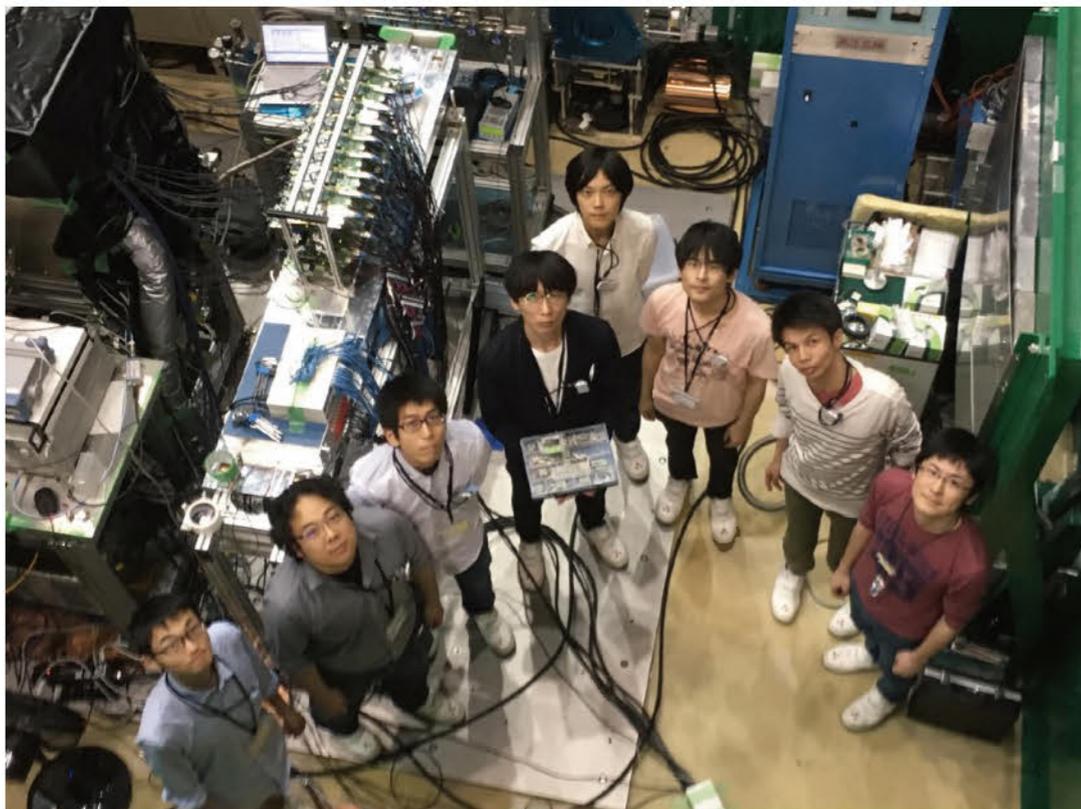
- Two layers of the detector are placed at the interval of 40 mm.
- The decay positron time spectra were analyzed by the coincidence method using two layers of the positron counter.
- Plural simultaneous hits on each layer were merged into a hit cluster to avoid positron over-counting.



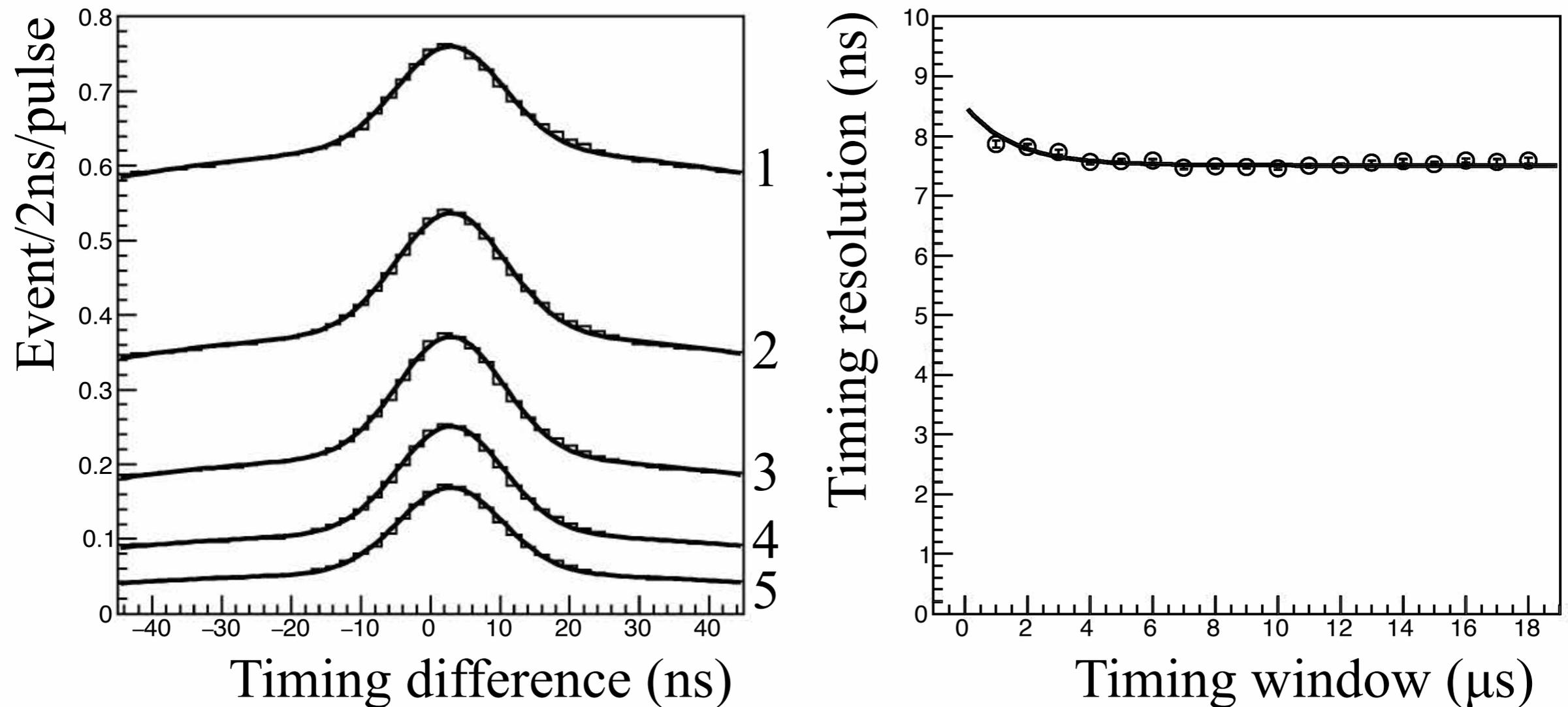
- First trial in 2014.
- No resonance was observed.
- Small signal, severe background.
- No beam delivery in 2015 due to the trouble with the mercury target.



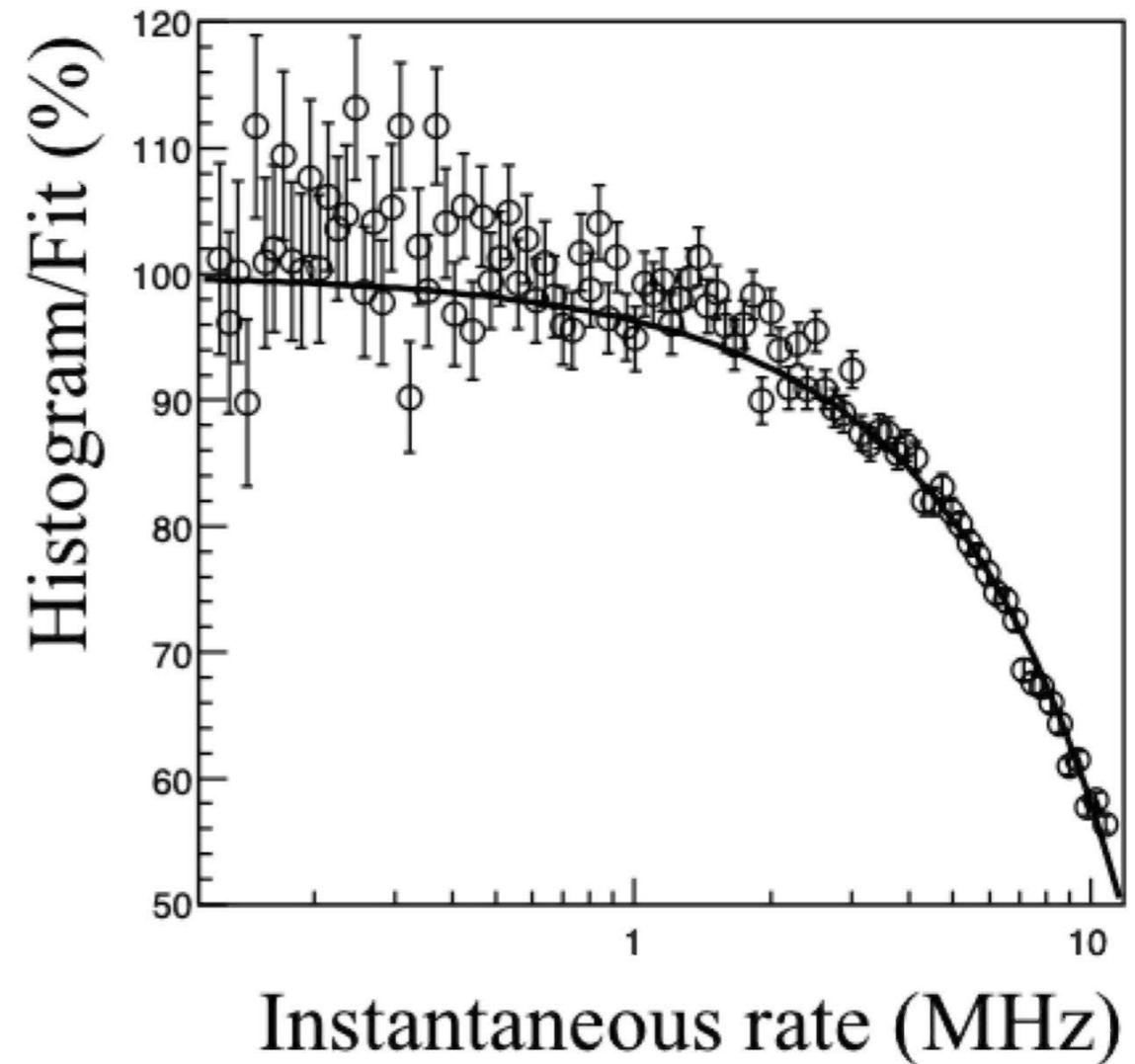
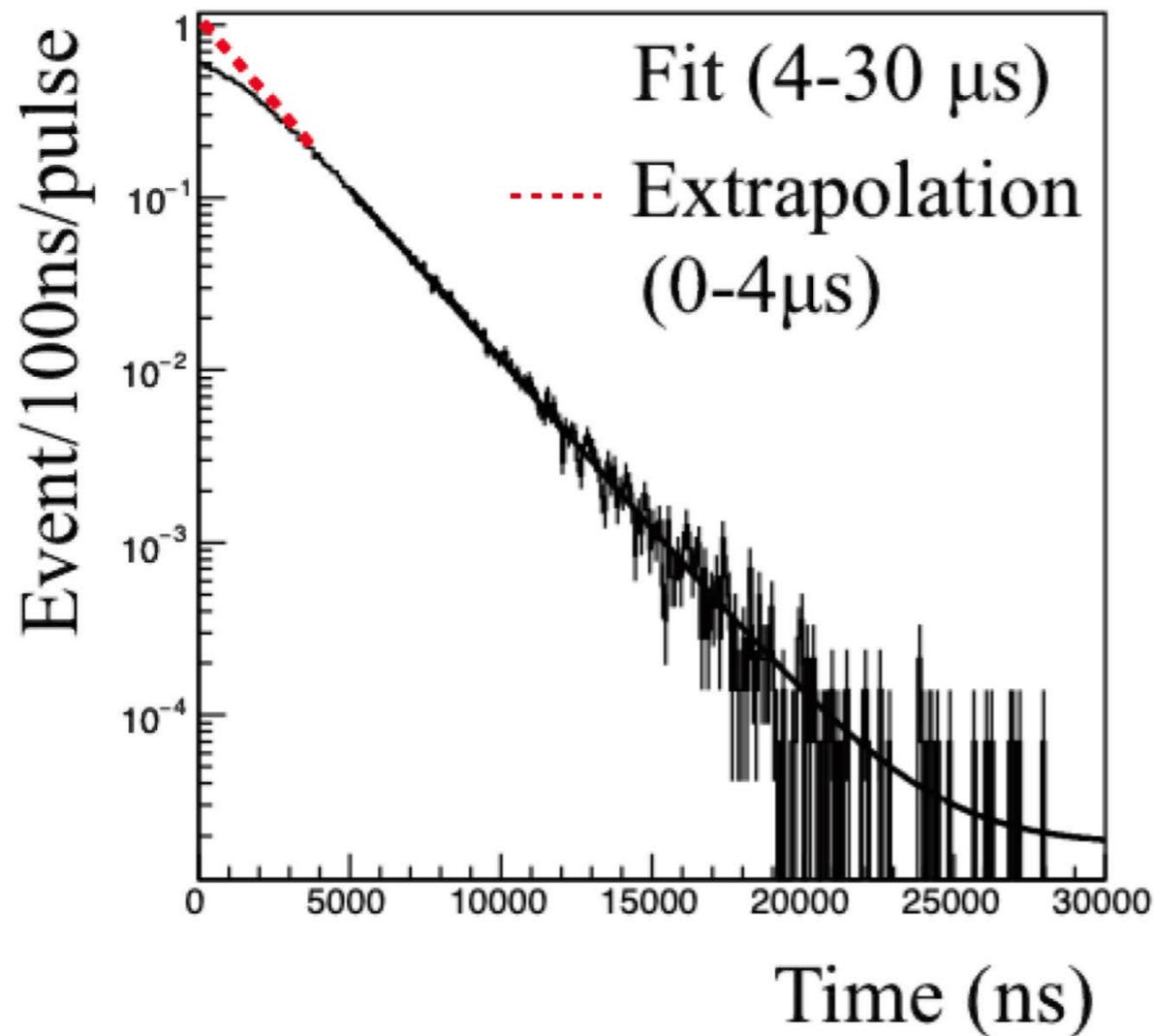
- Second trial in 2016.
- Improvements in the microwave system and suppression of beam-derived background events.
- First observation of the muonium HFS resonance with a pulsed muon beam.



- Third experiment in 2017.
- Improvements in background suppression.
- Microwave power dependence was studied.
- Power optimization.
- Today's main topic.
- Forth experiment in 2018.
- Improvements in stability and controls of measurement environment.
- Gas pressure dependence was studied.
- Fifth experiment in 2019.
- Test of mixture gas target.

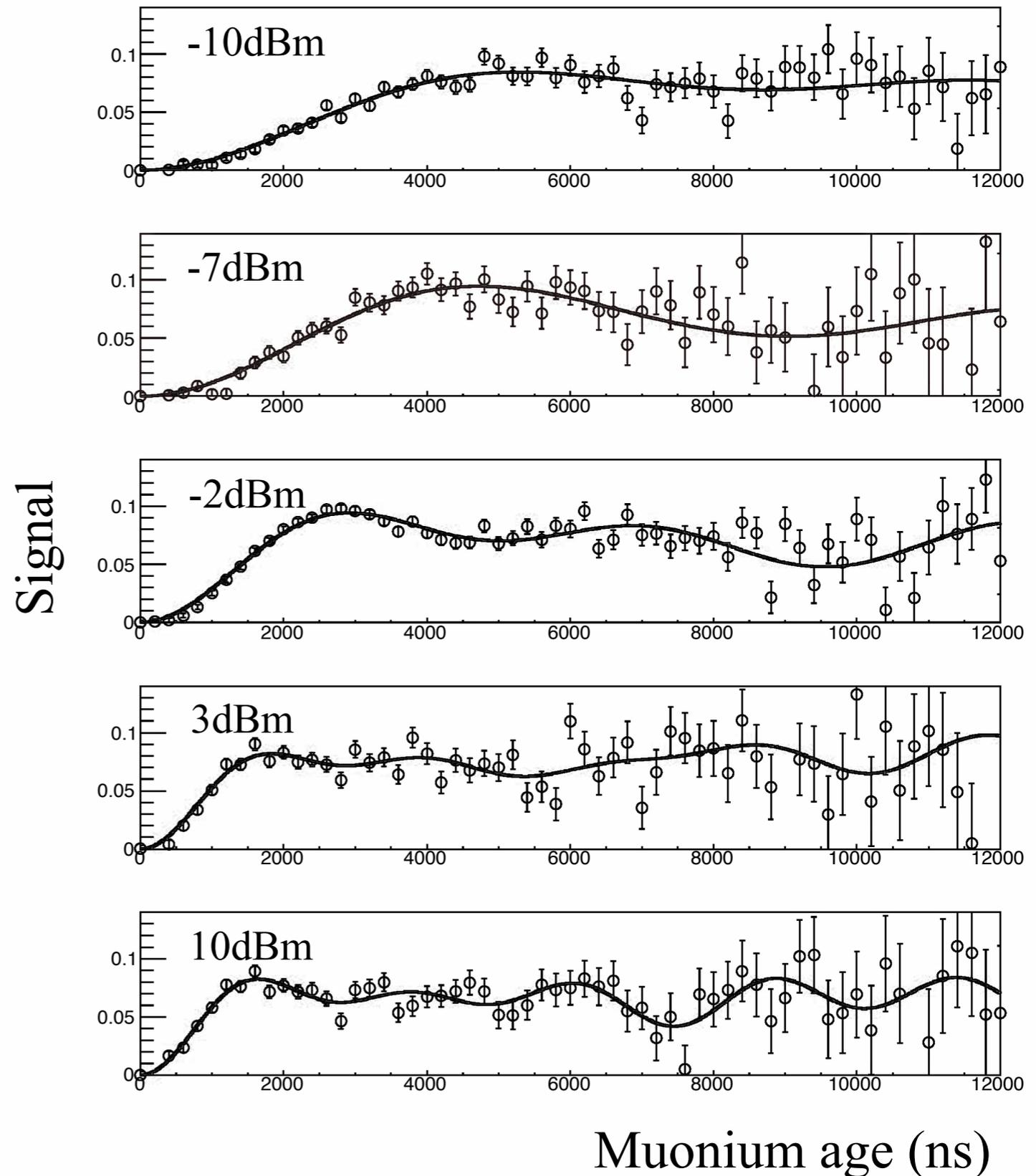


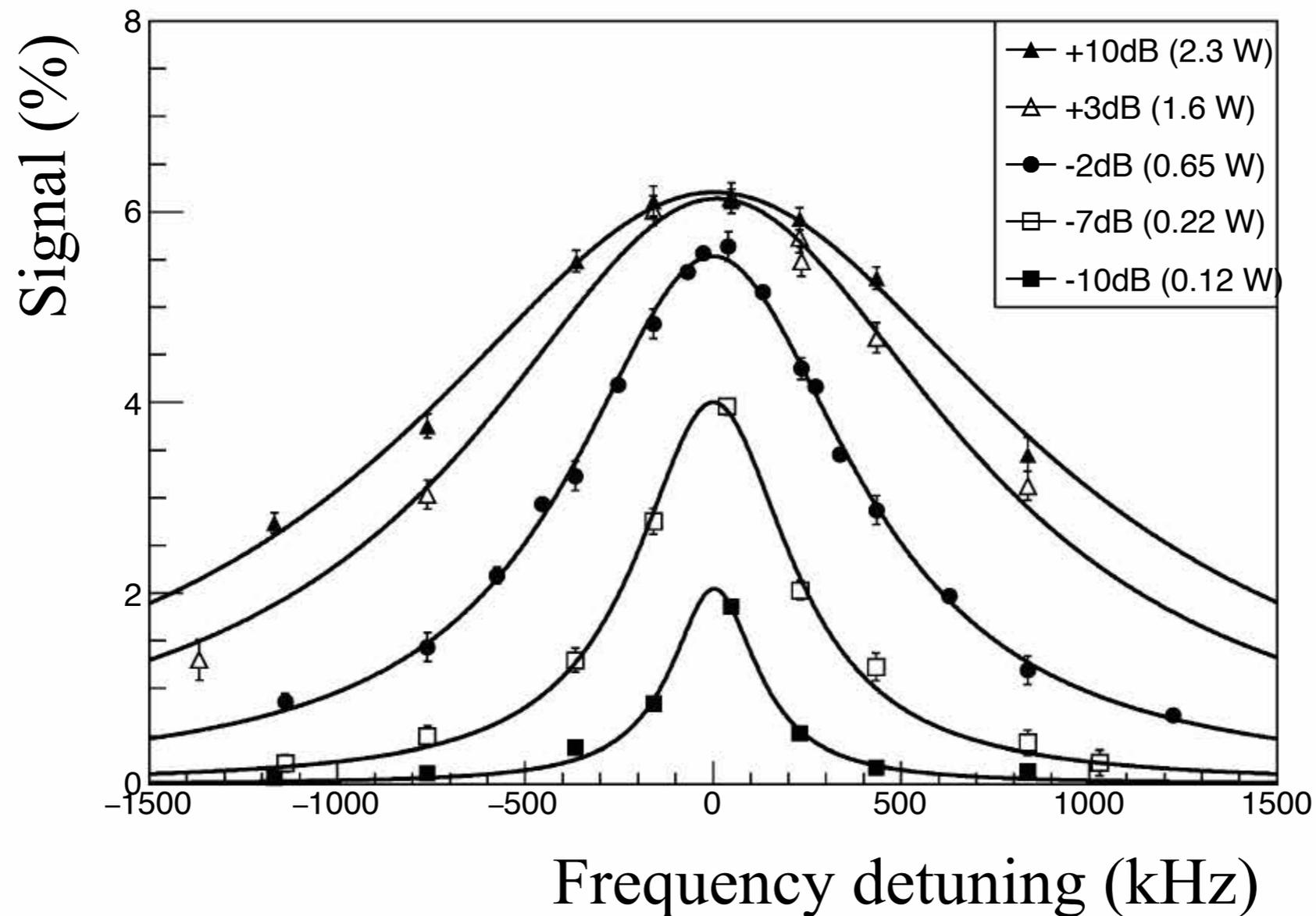
- (left) Time difference between hits on the near and far layers of the detector. The numbers on the right indicate the starting point of the 500 ns timing windows in  $\mu\text{s}$ . Curves are Gaussian on background.
- (right) Timing resolution of the detector, which is defined as  $1\sigma$  of the Gaussian component of the timing difference.



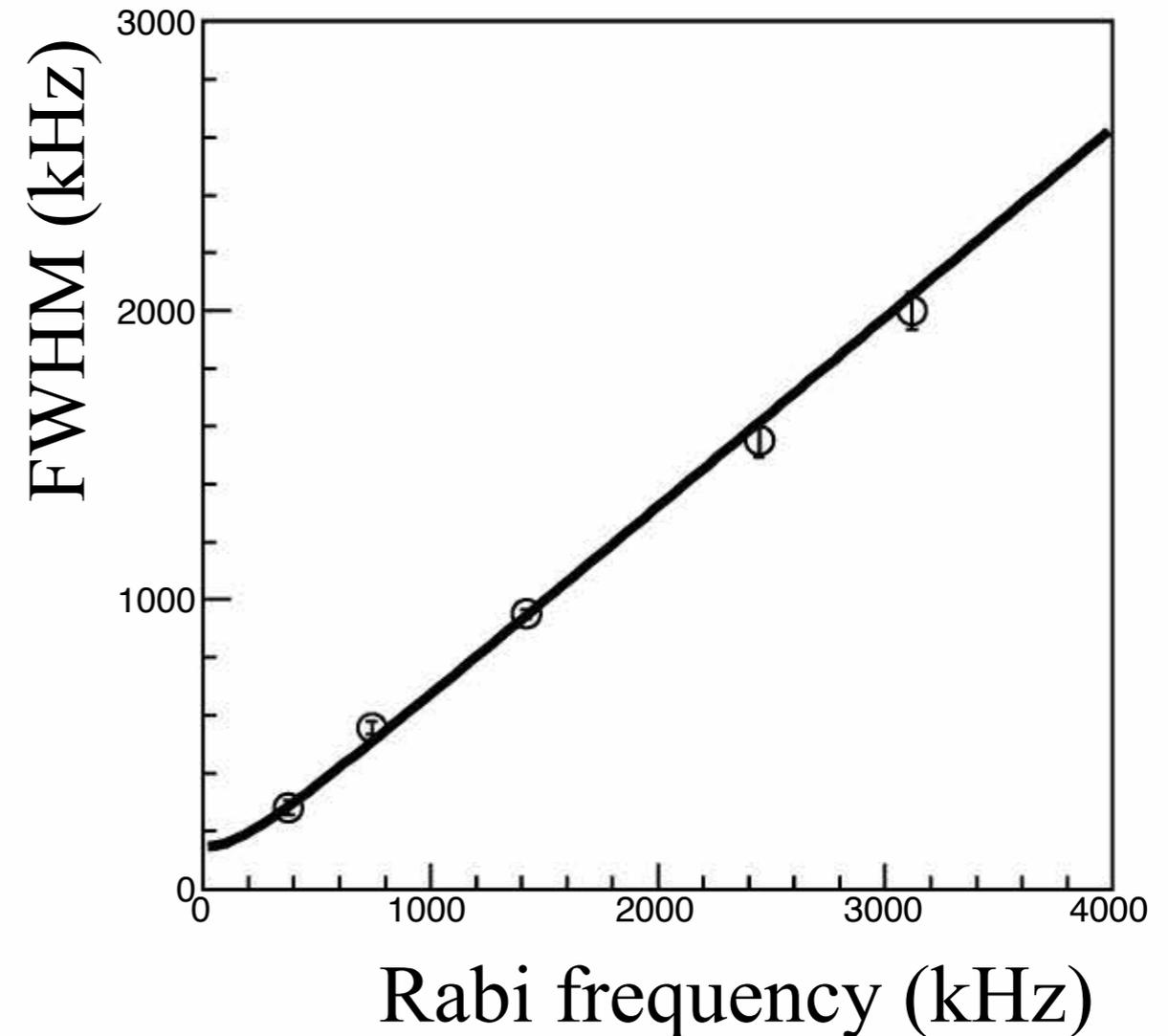
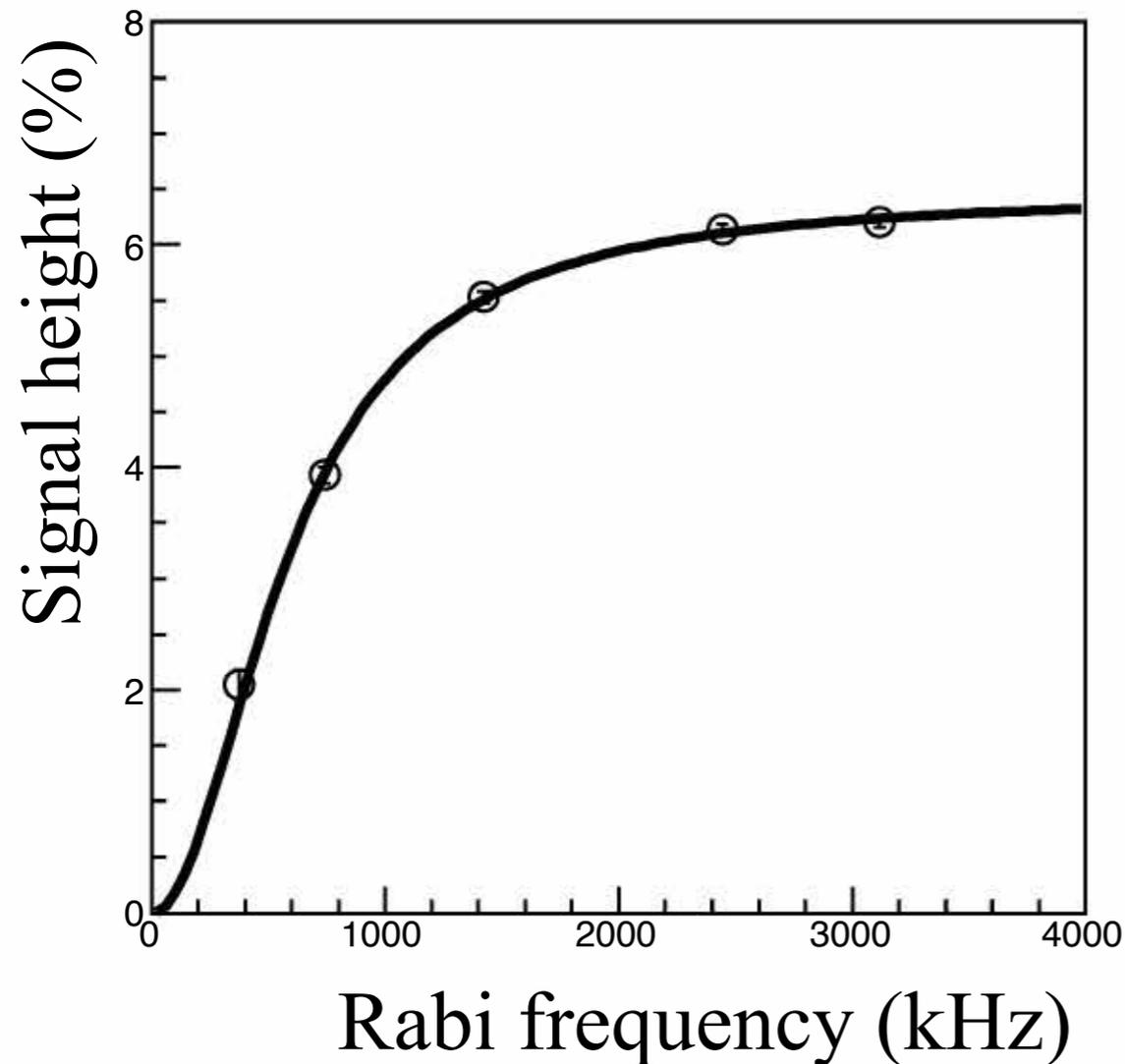
- (left) Time spectrum of decay positrons without microwave irradiation. The black solid curve corresponds to the fitting result with an exponential function on a constant background. The red dashed line indicates an extrapolation of the fitting function.
- (right) Pileup count loss as a function of the instantaneous event rate. The black curve indicates the fitting result with a model function.

- The Rabi oscillation was observed by taking the ratio of positron time spectra with and without microwave.
- Higher microwave power make the oscillation faster.
- Curves are theoretical expression the signal.
- The time integral of the oscillation gives the signal strength at a certain frequency of the resonance curve.

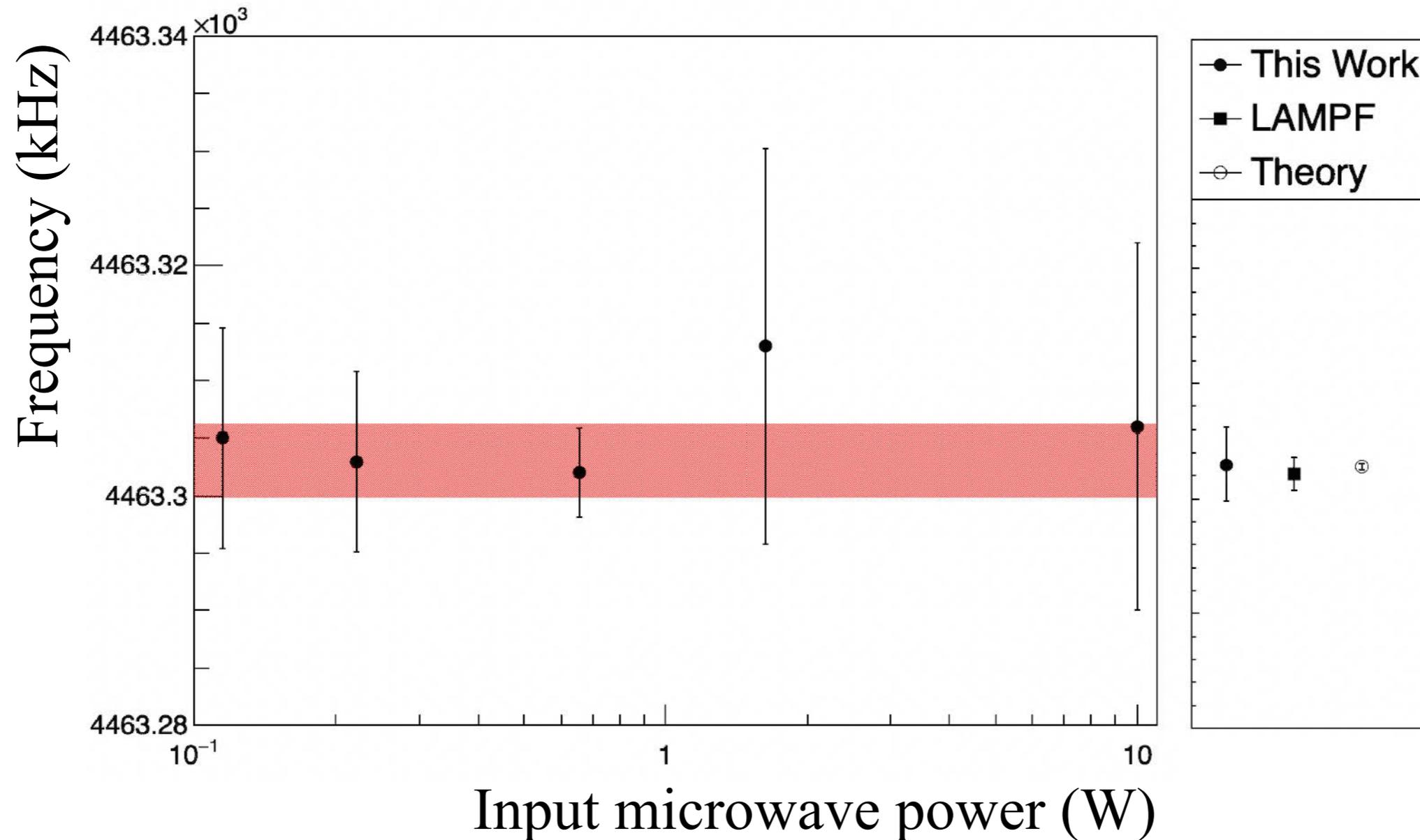




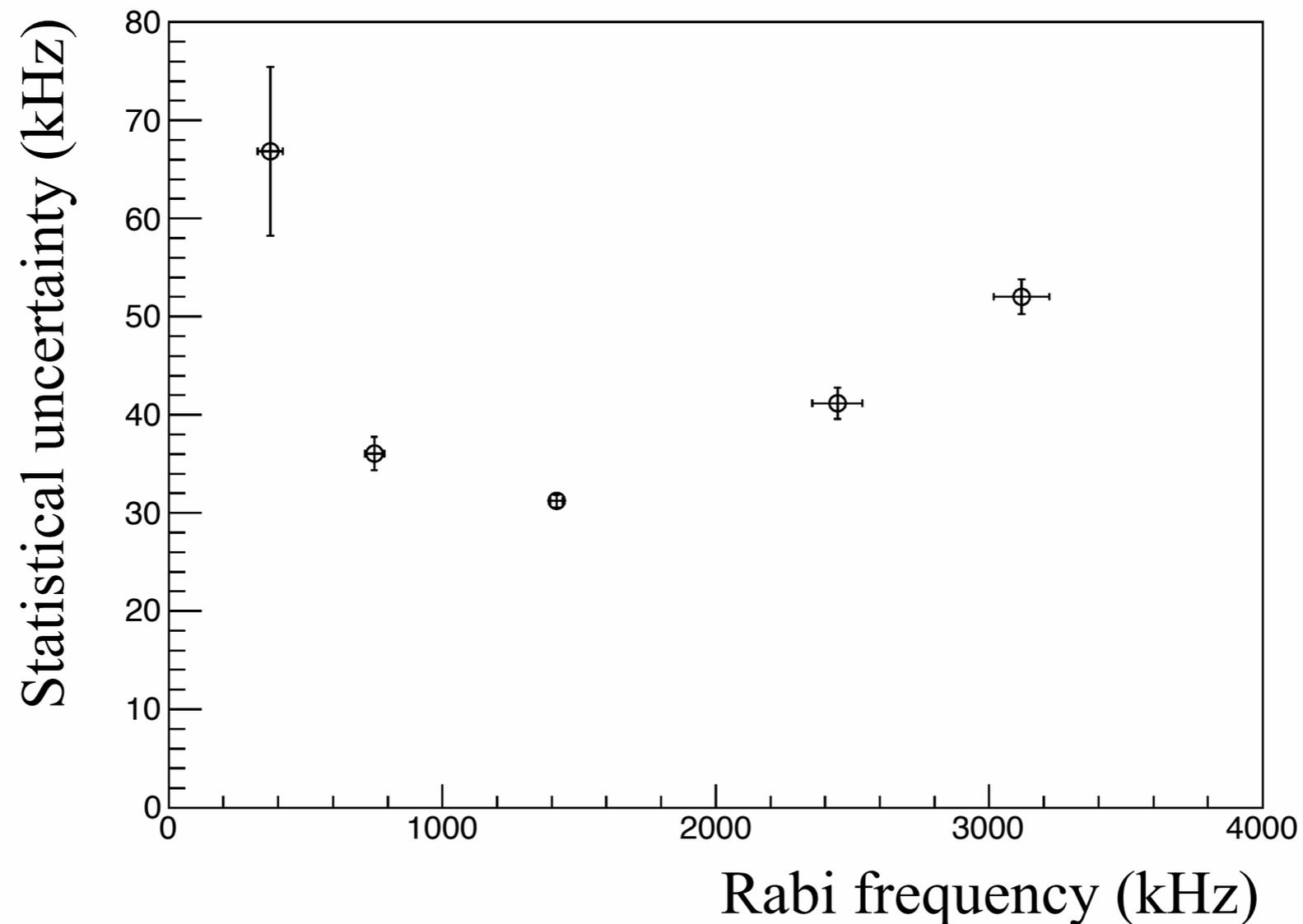
- Resonance lineshapes of the muonium HFS transition.
- The vertical axis shows the time integral of spin flip signal.
- The solid curves indicate to the fitting results with a Lorentz function.



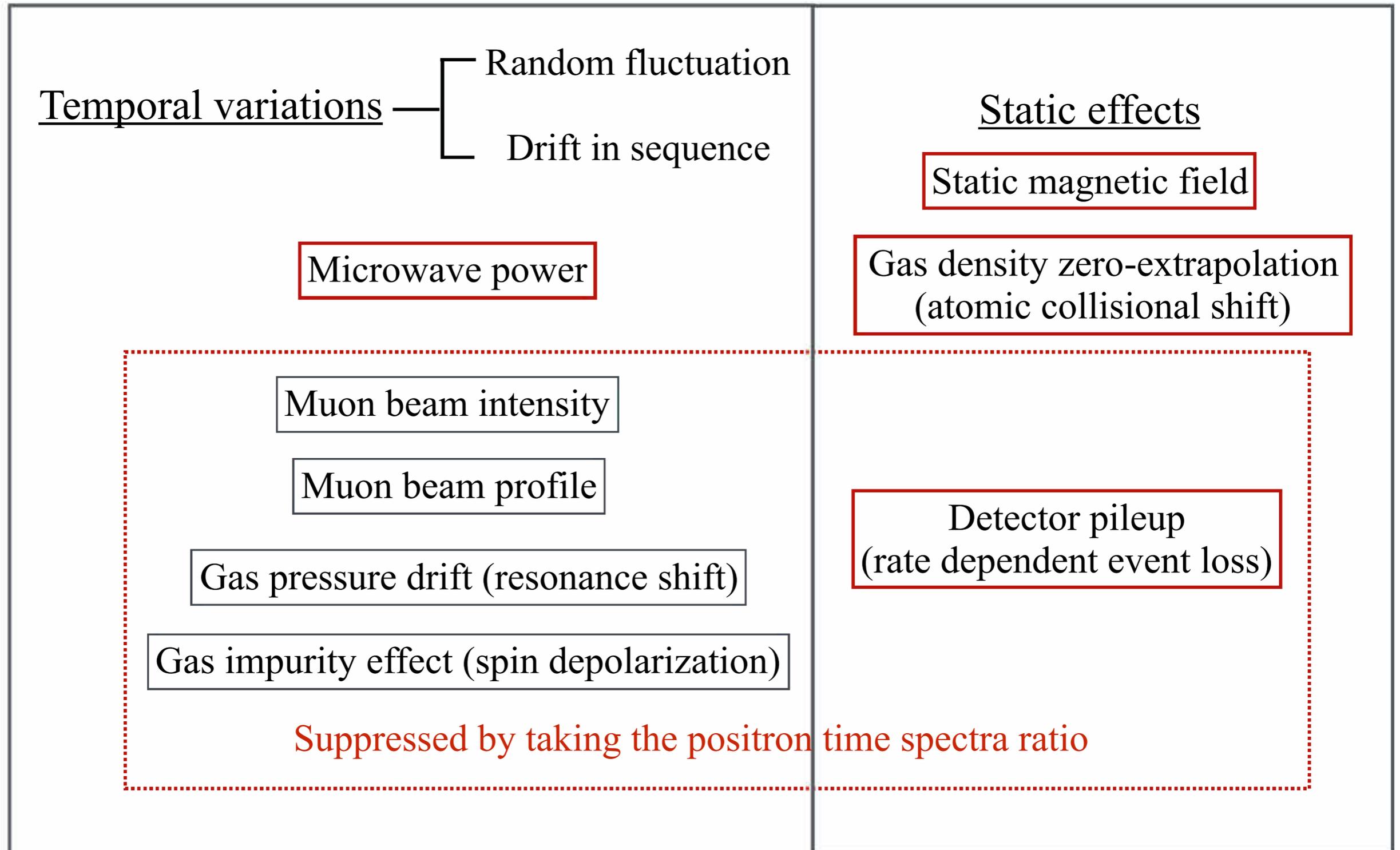
- Microwave power dependence of the resonance line- shape parameters: (a) the signal height; (b) the full-width- half- maximum (FWHM). Solid lines represent theoretical expressions.



- (a) The muonium hyperfine transition frequency for several input microwave power settings. The red band shows the weighted average. (b) Comparison with the precursor experimental result and the theoretical prediction.



- Microwave power optimization. The ordinate indicates the statistical uncertainty in determining the resonance frequency. The positron statistics of  $4 \times 10^8$  for each data point is assumed.



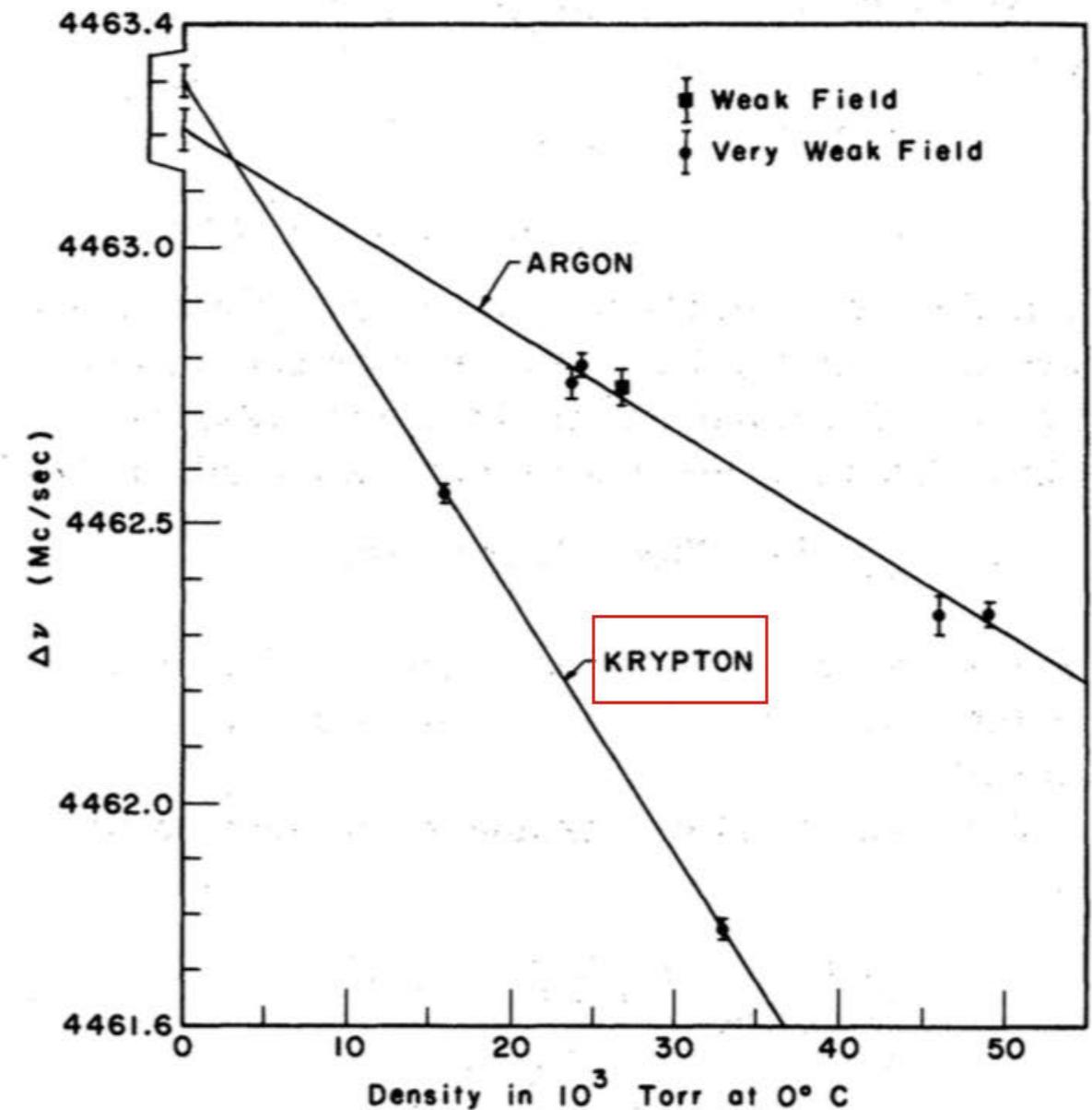
- Resonance frequency shifts due to collisions with atoms in the target.
- The shift depends of the target pressure.

$$\Delta\nu(p) = (1 + ap + bp^2)\Delta\nu(p=0)$$

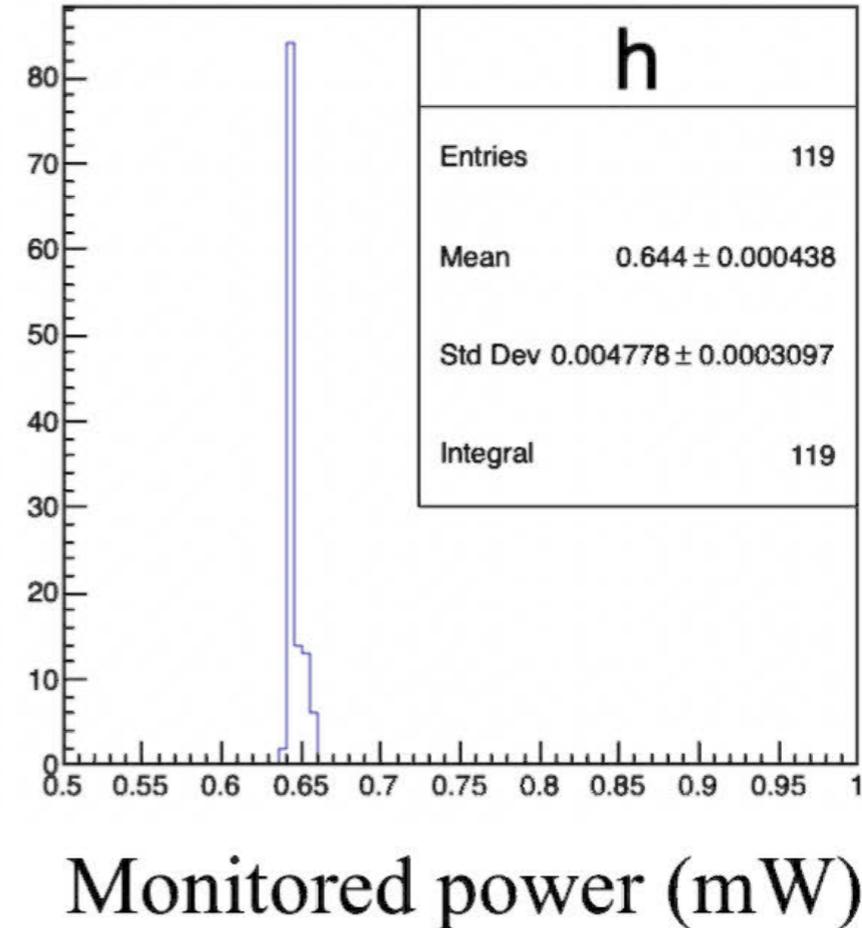
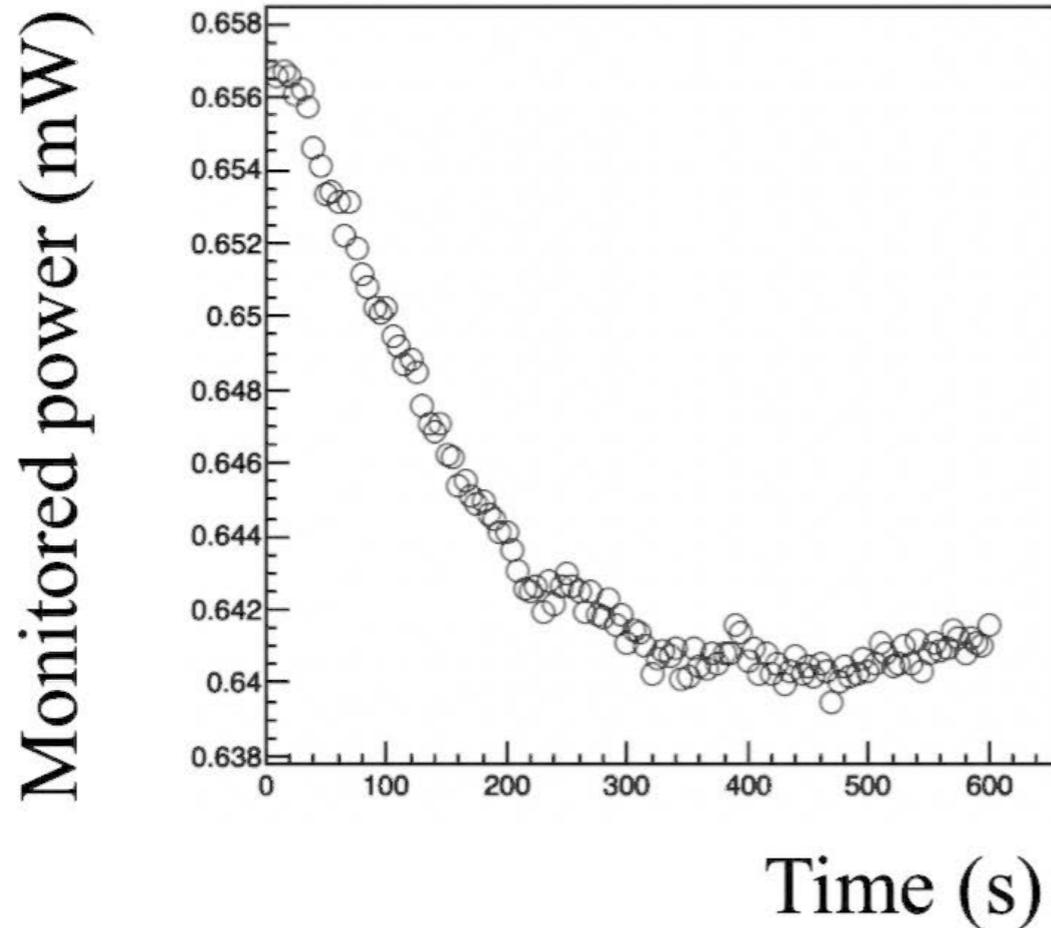
$$a = 7.996(8) \times 10^{-6}/\text{bar} \text{ (two-body)}$$

$$b = 5.5(1.1) \times 10^{-9}/\text{bar}^2 \text{ (three-body)}$$

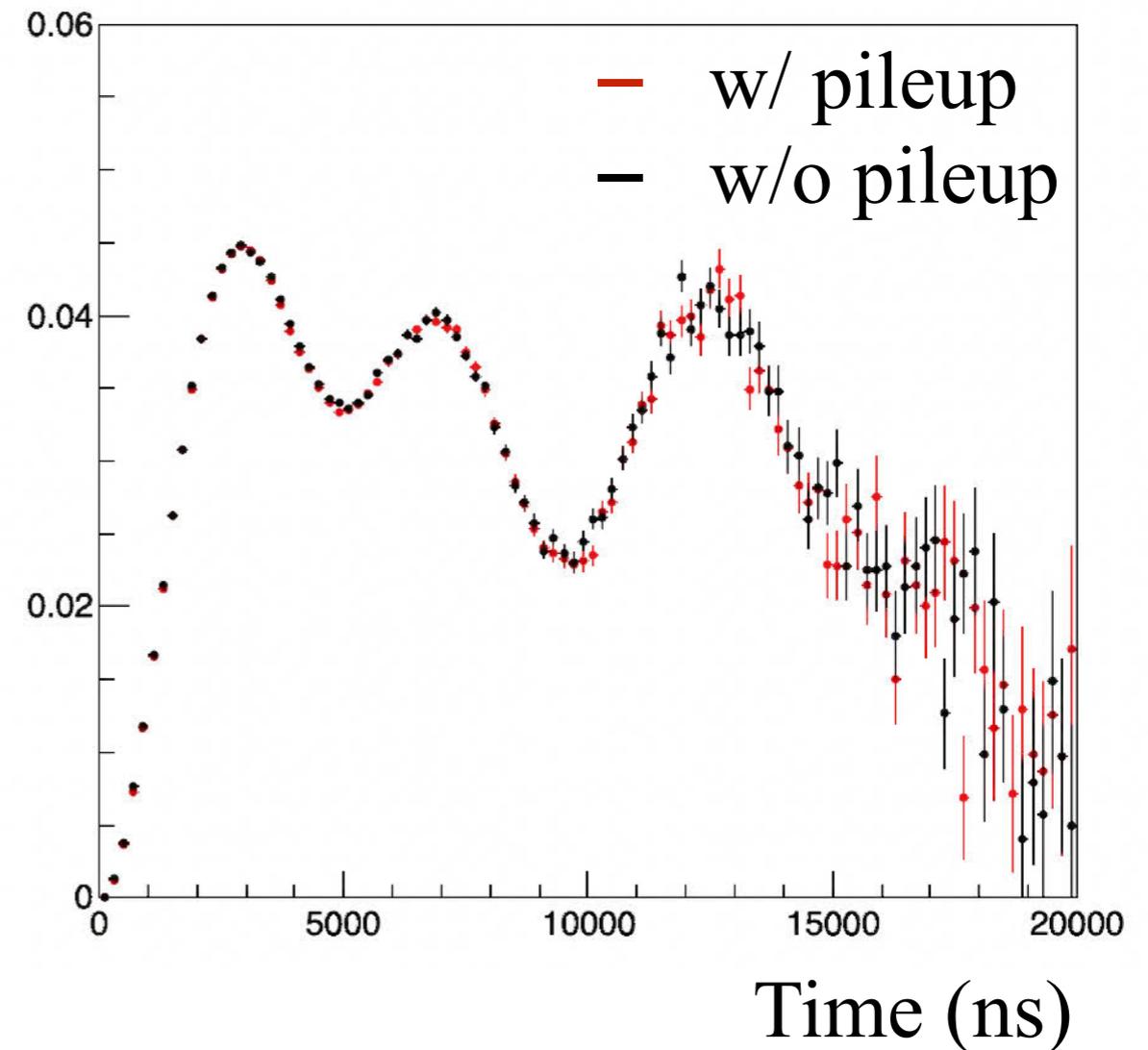
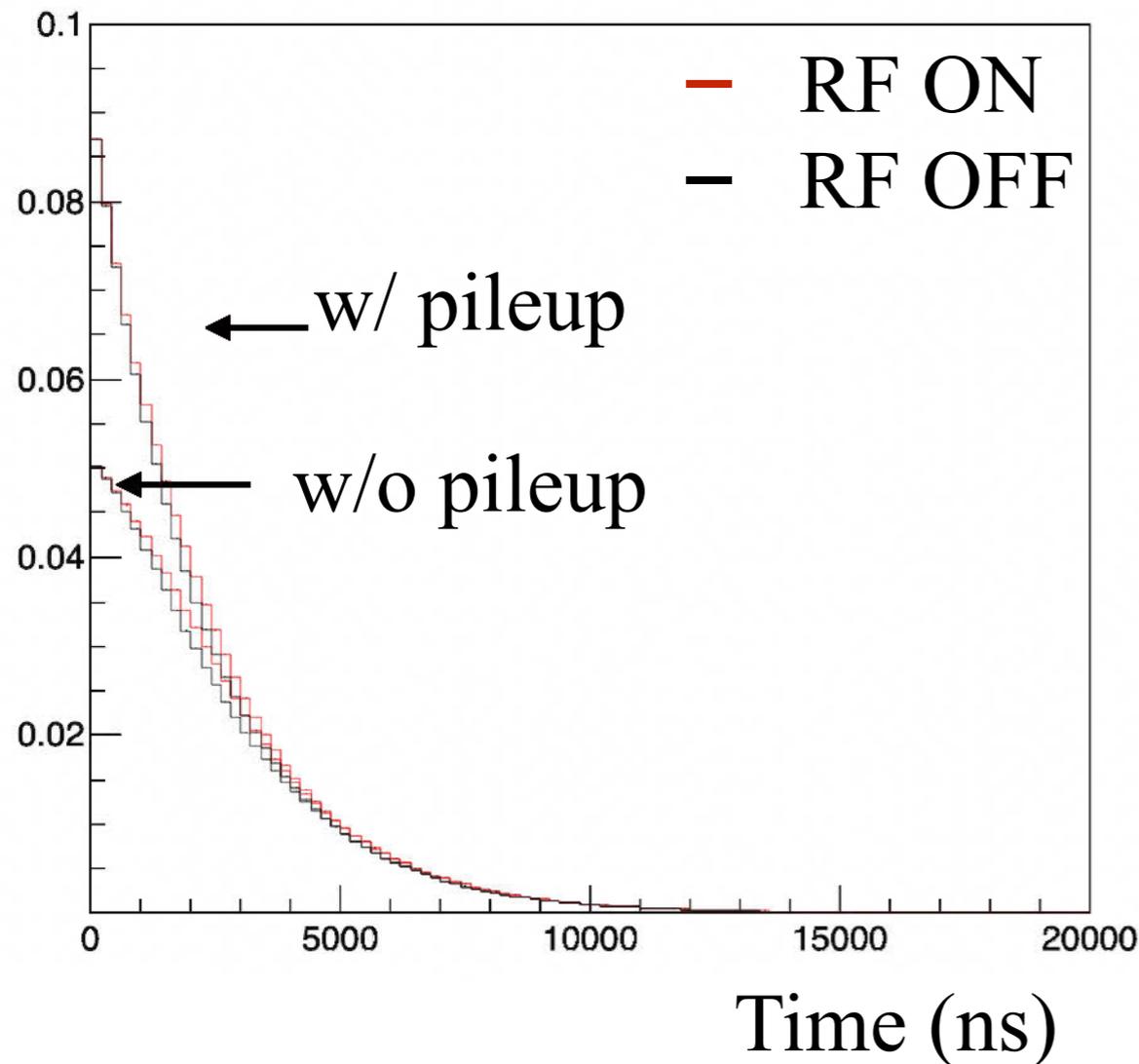
- For the case of krypton with the pressure of 1.0 atm, the resonance shifts -33 kHz.
- Systematic uncertainty due to pressure gauge accuracy is 46 Hz.



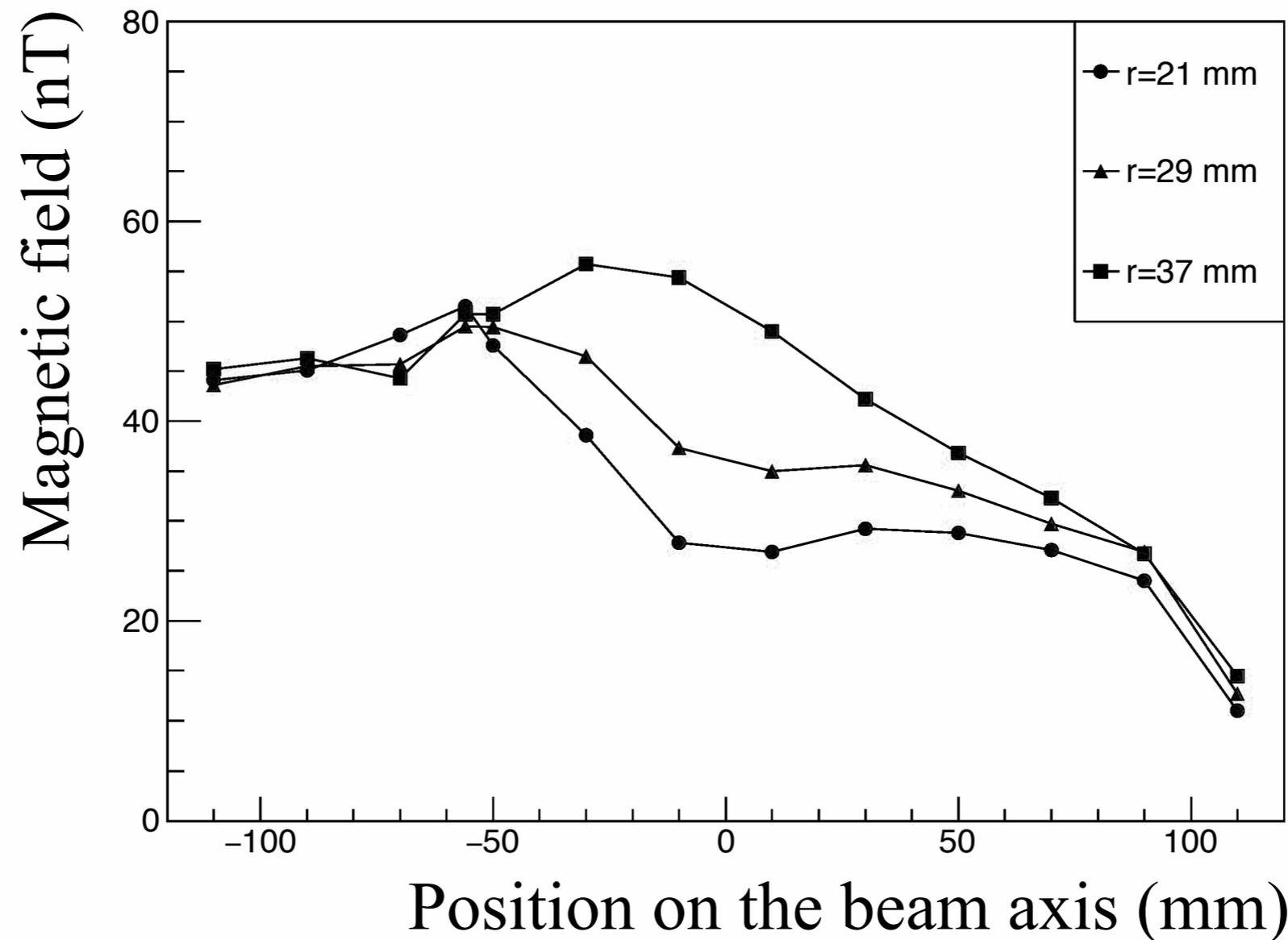
Values from D. E. Casperson *et al.*, Phys. Lett. B, 59, 4 (1975).  
Figure from P. A. Thompson *et al.*, PRL, 22, 5, (1969).



- Power correction was performed using the average value at each frequency data point.
- The systematic error is evaluated using the error in determining the average value.
- The uncertainty is estimated to be 37 Hz.



- Decay positron time spectra are simulated with and without microwave transition considering the pileup count loss.
- The time dependent signal is calculated with and without the pileup effect.
- Systematic uncertainty is estimated to be 19 Hz.

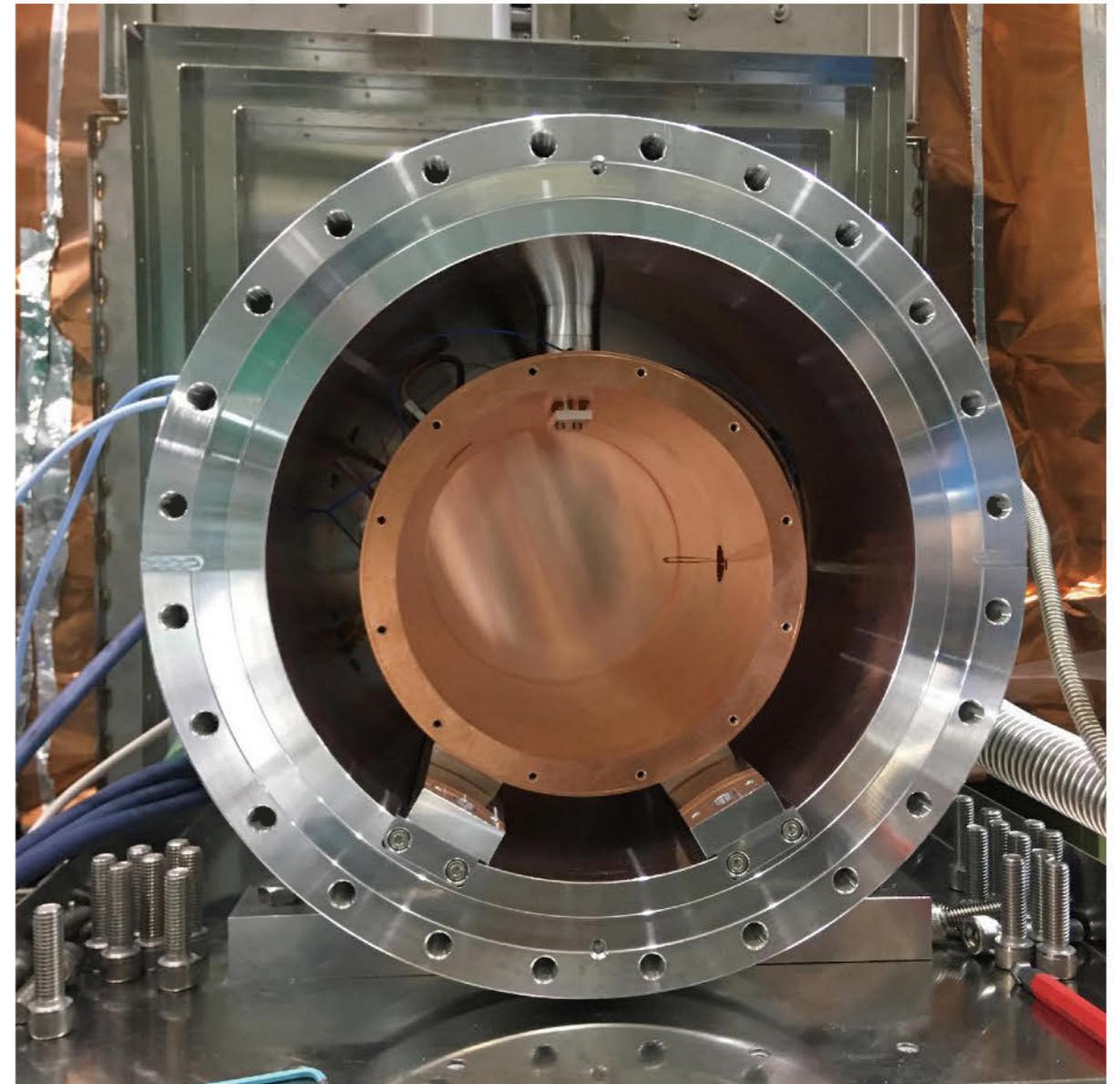


- Magnetic field strength on the longitudinal axis was measured by using the fluxgate probe.
- The Zeeman shifts of muonium levels are calculated.
- Systematic uncertainty is negligibly small.

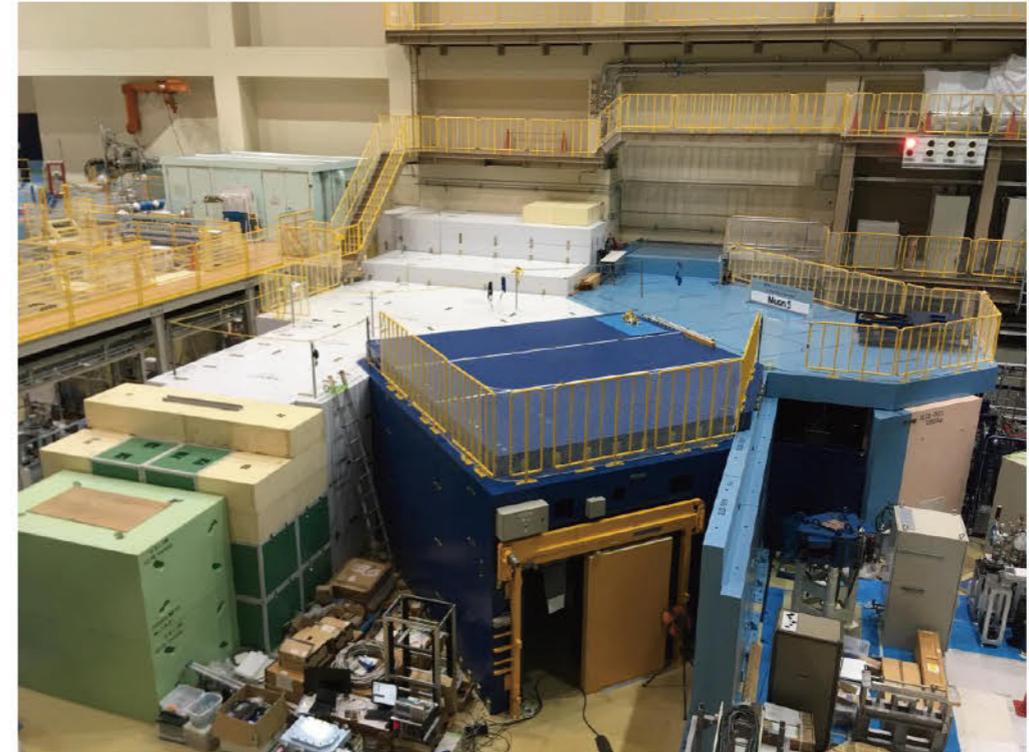
| Source                   | Contribution (Hz) |
|--------------------------|-------------------|
| Atomic collisional shift | 46                |
| Microwave power          | 37                |
| Detector pileup          | 19                |
| Static magnetic field    | 0                 |
| Gas pressure fluctuation | 6                 |
| Gas impurity buildup     | 12                |
| Muon beam intensity      | 0                 |
| Muon beam profile        | 0                 |
| Total                    | 63                |

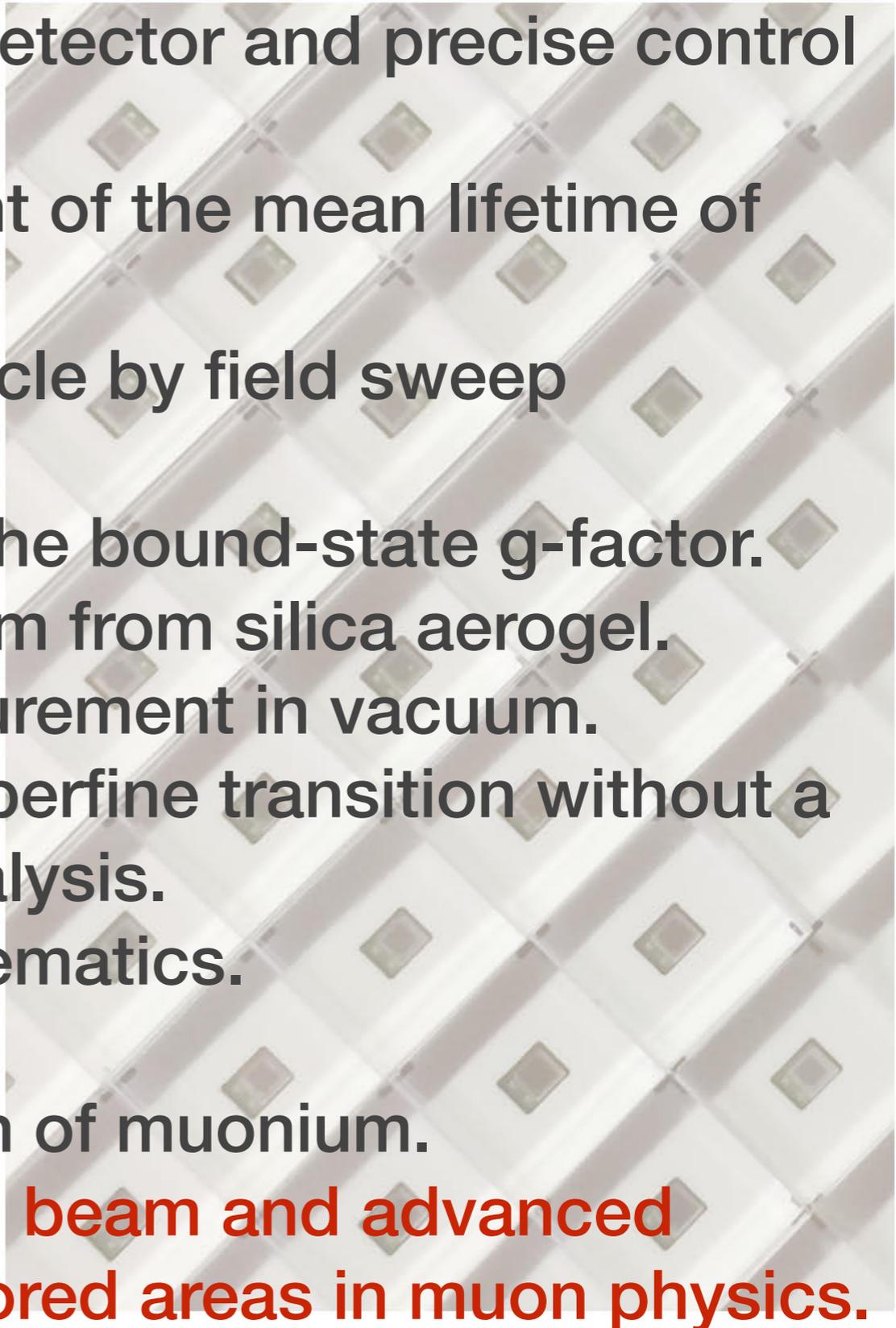
- Currently, the systematic uncertainties are sufficiently small compared to the statistical uncertainty.
- However, after completion of the new beam line, statistical precision is significantly improved and it is more important to understand and suppress systematic uncertainties.
- A pressure gauge with better accuracy, microwave power switching, water cooling of the cavity will help improvements.

- Larger microwave cavity which resonates in TM220 mode
  - Y.Ueno (RIKEN).
- Microwave switching
  - T. Tanaka (U. Tokyo)
- Krypton-Helium mixture gas target to offset the atomic collisional shift
  - S.Seo (U. Tokyo).
- Detailed analysis of the time-dependent spin flip signal.
  - S.Nishimura (U. Tokyo)



- The new beam line at the 1st experimental hall in MLF with larger solid angle is under construction.
- The superconducting magnet for high-field measurement was delivered at J-PARC.
- Passive shimming method for field uniformity was developed.
- NMR probe for field measurement is under development.



- Improvement in the positron detector and precise control of a magnetic field enable:
    - High precision measurement of the mean lifetime of muonium.
    - A search for axion-like particle by field sweep measurement.
    - Precision measurement of the bound-state g-factor.
  - Muonium emission into vacuum from silica aerogel.
    - Atomic-collision-free measurement in vacuum.
  - Direct measurement of the hyperfine transition without a microwave field using FFT analysis.
    - Free from field-related systematics.
  - Laser cooling of muonium.
    - Bose-Einstein condensation of muonium.
  - **Highest intensity pulsed muon beam and advanced instrumentation unveil unexplored areas in muon physics.**
- 

- The ground-state hyperfine splitting in muonium atom is an ideal observable to test bound-state QED theory.
- MuSEUM collaboration aims to improve a measurement precision of the muonium HFS by a factor of ten.
- The high-intensity pulsed muon beam and high-rate capable positron detector enable us to realize a new experiment at J-PARC.
- The first observation of muonium HFS resonance was reported after the experiment in 2016. The principle is proofed.
- Improved experiment was conducted in 2017 after upgrades and the result is comparable to the precursor experimental one.
- Further improvements are under study.