

Magnetic field confined by elasticity in neutron star crust

- YK & Suzuki(2020) arXiv:2004.08006 MNRAS(2020)
- YK, Kisaka & Fujisawa(2021) arXiv:2011.03239 MNRAS(2021)
- YK, Kisaka & Fujisawa(2021) arXiv:2106.14337 MNRAS(2021) in press
- +a

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～中性子星の観測と理論～研究活性化ワークショップ 2021
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Magnetic field confined in crust

Working hypothesis

-> CCO(Magnetar)

-> Formation ? Matsumoto-san's talk

Outer part ($\Delta r/R < 0.1$) of NS

- Core $\rho = 1.4 \times 10^{14} \text{ g/cm}^3$ nuclear density
Pressure $p = 3 \times 10^{33} \text{ erg/cm}^3$
Max. B $B < 2 \times 10^{17} \text{ G} \sim \text{Mev/fm}^3$
 - Outer BC $\rho = 4 \times 10^{11} \text{ g/cm}^3$ neutron drip / Solid crust
Pressure $p = 5 \times 10^{29} \text{ erg/cm}^3$
Max. B $B < 3 \times 10^{15} \text{ G (enough)} \sim \text{kev/fm}^3$
Shear $\mu = 3 \times 10^{27} \rightarrow 10^{30} \text{ erg/cm}^3$
- > Deformation of NS "Mountain"

Core vs crustal field

Core physics may appear via crust Crust first! Negative -> Core next!

- Complicated treatment for core physics

Many species(e,n,p,...) superfluid/superconductor +....

- Simple and possible observational consequence for crust case

Ion lattice + e -> B-evolution by Hall drift & Joule loss

Deformation of crust, “mountain on NS “

-> GW?

Beyond limit Fracture or not? -> Observational events?

-> Quake? Models of FRB?

-> Plastic flow? Effect on magnetic field evolution (YK+20,21)

+...

Outline of my talk

1. (Introduction) Braking strain of solid
2. Astrophysical consequences
 - > A mountain on NS by **elastic** or **magnetic** forces
 - > GW? (Brief review) -----> Ito-san's talk
 - Deformation induced by magnetic evolution (YK+,21a)
3. **Equilibrium of magnetized crust (NEW) (YK+,21b)**
 - **Observation consequences beyond threshold**
 - > Quake? Models of FRB?
 - > Plastic flow?
 - Some discussion in literature by speculation, but is omitted here

Breaking strain of solid crystals

- Material property in high density regime

Crust of NS ($\rho \approx 10^{11} - 10^{14} \text{ g/cm}^3$) outer part $< 1\text{ km}$ thick

pressure ($p \approx 10^{29} - 10^{33} \text{ erg/cm}^3$)

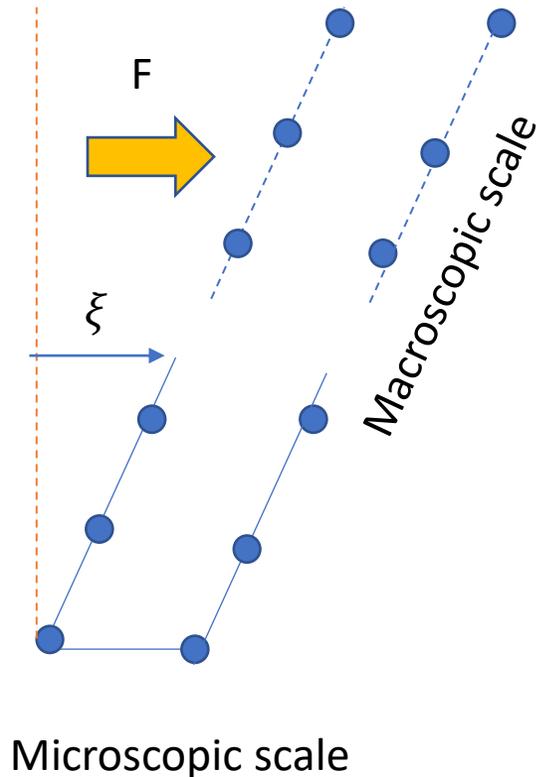
$$\sigma_{ij} = \nabla_i \xi_j + \nabla_j \xi_i - \frac{2}{3} g_{ij} (\nabla \cdot \xi), \quad h_i = \nabla_j (\mu \sigma_i^j)$$

shear modulus $\mu \sim \rho$

Limit of strain

σ_c

- Experimental approach is impossible (at present)
- Theory (Analytic $\sigma_c \sim 0.01$ / Numerical $\sigma_c \sim 0.1$)
- Beyond it, plastic flow or quake
- Possibly it is checked in astrophysical object



Breaking strain

“Introduction to Solid State Physics”

Charles Kittell

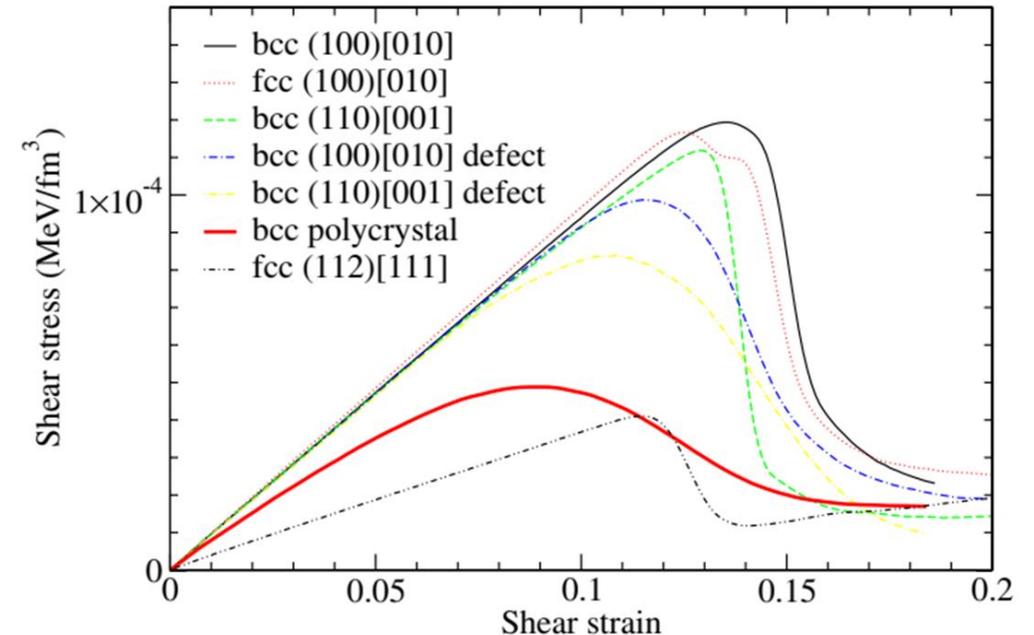
$$\sigma_c = 10^{-2} \text{ to } 10^{-5} \\ 1/6 \text{ for an ideal case}$$

表 18・1 ずれ弾性率と弾性限度の比較 (Mott による)

	ずれ弾性率 G [dyn/cm ²]	弾性限度 σ_c [dyn/cm ²]	$\frac{G}{\sigma_c}$
Sn (単結晶)	1.9×10^{11}	1.3×10^7	15 000
Ag (単結晶)	2.8×10^{11}	6×10^6	45 000
Al (単結晶)	2.5×10^{11}	4×10^6	60 000
Al (純粋, 多結晶)	2.5×10^{11}	2.6×10^8	900
Al (市販圧延品)	$\sim 2.5 \times 10^{11}$	9.9×10^8	250
ジュラルミン	$\sim 2.5 \times 10^{11}$	3.6×10^9	70
Fe (軟, 多結晶)	7.7×10^{11}	1.5×10^9	500
炭素鋼 (熱処理)	$\sim 8 \times 10^{11}$	6.5×10^9	120
ニッケル・クロム鋼	$\sim 8 \times 10^{11}$	1.2×10^{10}	65

Molecular dynamics simulation
Horowitz + (09)PRL102,191102

$$\sigma_c = 0.1$$



$$10^{29} \text{ erg/cm}^3 \sim 10^{-4} p$$

Mountain on NS

e.g. Haskell +(06) Johnson-McDaniel +(13)

- Pressure in crust $\rho = 10^{29} - 10^{33} \text{ erg/cm}^3$

Ellipticity estimated by pressure or energy ratio

- Mountain supported by **elastic** force

-----> Ito-san's talk

$$\varepsilon \approx \mu \sigma_c (4\pi R^2 \Delta R) / (GM^2/R) \approx 10^{-6} \text{ for } \mu = 10^{30} \text{ erg/cm}^3$$

Detailed calculations in literature, e.g, elastic mountain 10^{-8}

(Gittens+ 21)

- Mountain supported by **magnetic** force

$$\varepsilon \approx B^2 R^3 / (GM^2/R) \approx 10^{-6} \text{ for } B = 10^{14} \text{ G}$$

Ellipticity(deformation) of NS will be observed by GW

Magnetic deformation of NS revisited

Suvorov+(16) proposed large deformation $\sim 10^{-6}$ induced by Hall evolution in normal radio pulsars $B \sim 10^{12}$ G

Current upper limit of ellipticity $< 3 \times 10^{-5}$ (-> next slide)

Table 5. Comparison of our models with observational limits of seven selected pulsars, for which $h_0 \lesssim 4h_0^{\text{sd}}$ (Aasi et al. 2014). $|B_s|$ and τ_c were computed from data given in Aasi et al. (2014). The fourth column shows the observational upper limits on ϵ from LIGO and Virgo (Aasi et al. 2014), the fifth column shows ϵ as predicted by the AL model, and the sixth column shows ϵ as predicted by the BL model.

Suvorov+ MNRAS(16)

Pulsar	$ B_s $ (10^{12} G)	τ_c (yr)	Obs. limit $ \epsilon $	AL-predicted $ \epsilon $	BL-predicted $ \epsilon $
<u>J0534+2200 (Crab)</u>	1.84	1.3×10^4	8.6×10^{-5}	9.6×10^{-10}	2.9×10^{-7}
<u>J0537-6910</u>	1.42	4.9×10^3	1.2×10^{-4}	5.8×10^{-8}	3.2×10^{-8}
<u>J1813-1246</u>	1.43	4.3×10^4	3.5×10^{-4}	9.6×10^{-10}	2.9×10^{-7}
J1833-1034	5.52	4.8×10^3	5.7×10^{-3}	5.8×10^{-8}	3.2×10^{-8}
J1913+1011	0.54	1.7×10^5	2.2×10^{-4}	1.3×10^{-8}	2.6×10^{-6}
J0835-4510 (Vela)	5.26	1.1×10^4	6.0×10^{-4}	9.6×10^{-10}	2.9×10^{-7}
<u>J1952+3252</u>	0.74	1.1×10^5	3.0×10^{-4}	1.3×10^{-8}	2.6×10^{-6}

Two models

References

Upper limit from GW

- Crab J0534+2200

Ellipticity $<1 \times 10^{-5}$ LIGO/Virgo O2

- B1951+32/J1952+3252

Ellipticity $<2 \times 10^{-5}$ LIGO/Virgo O2

B. P. Abbott et al 2019 ApJ 879 10

- PSR0537 – 6910

Ellipticity $<3 \times 10^{-5}$ LIGO/Virgo O3(2020 Oct)

Arxiv2012.12926

Our theoretical results

Right figure shows evolution of ellipticity for three models

Unfortunately for GW community, our quadrupole deformation is $\sim 10^{-8}$

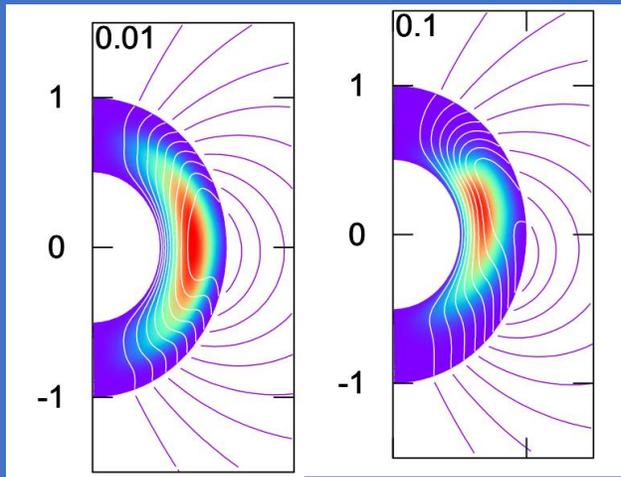
mass moment

- Deformation in shape and that in quadrupole moment are different

Shape deformation in crust is $\sim 3 \times 10^{-5}$

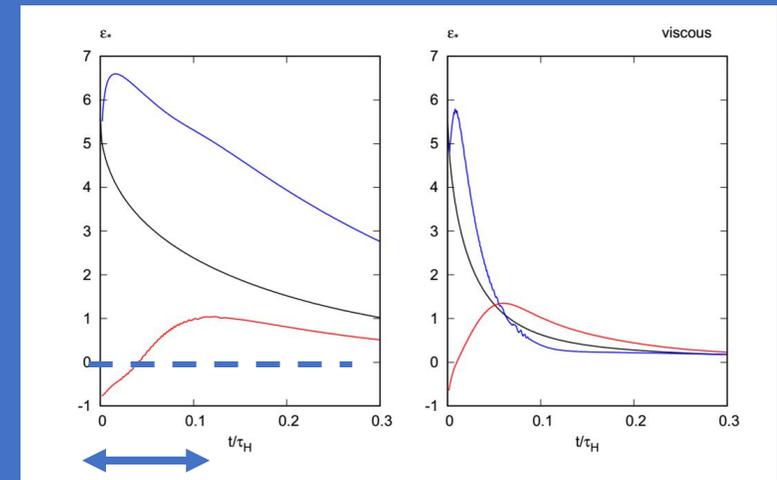
- Different model/treatment

=> Large deformation by core field



Evolution of magnetic field

Evolution of deformation



Evolution of magnetic deformation in neutron star crust
Kojima, Kisaka & Fujisawa, MNRAS(21)

10^5 yr

Is elastic force effective for mag. confinement?

New

In crust $\rho = 4 \times 10^{11} - 1.4 \times 10^{14} \text{ g/cm}^3$

$$p = 5 \times 10^{29} - 3 \times 10^{33} \text{ erg/cm}^3 \sim B^2$$

$$\text{Shear } \underline{\mu = 3 \times 10^{27} - 10^{30} \text{ erg/cm}^3} \quad \underline{\mu/p \sim 10^{-2} - 10^{-3}}$$

Elastic force is too small

✓ Magnetic energy is

$$E_B \sim 10^{46} (B_{14})^2 \text{ erg} \quad (\text{which depends on geometry})$$

enough to activities

✓ Elastic energy is

$$E_{els} \sim 10^{44} \bar{\mu} (\sigma_c / 0.1)^2 \text{ erg}$$

➤ Solution for stability of magnetized neutron star

Mixed poloidal-toroidal fields are stable in conjecture, but toroidal-dominated models are not yet calculated.

See also Bera + (20)



Unimportant!

No !

A pitfall in argument

Magneto-elastic equilibrium of a neutron-star crust

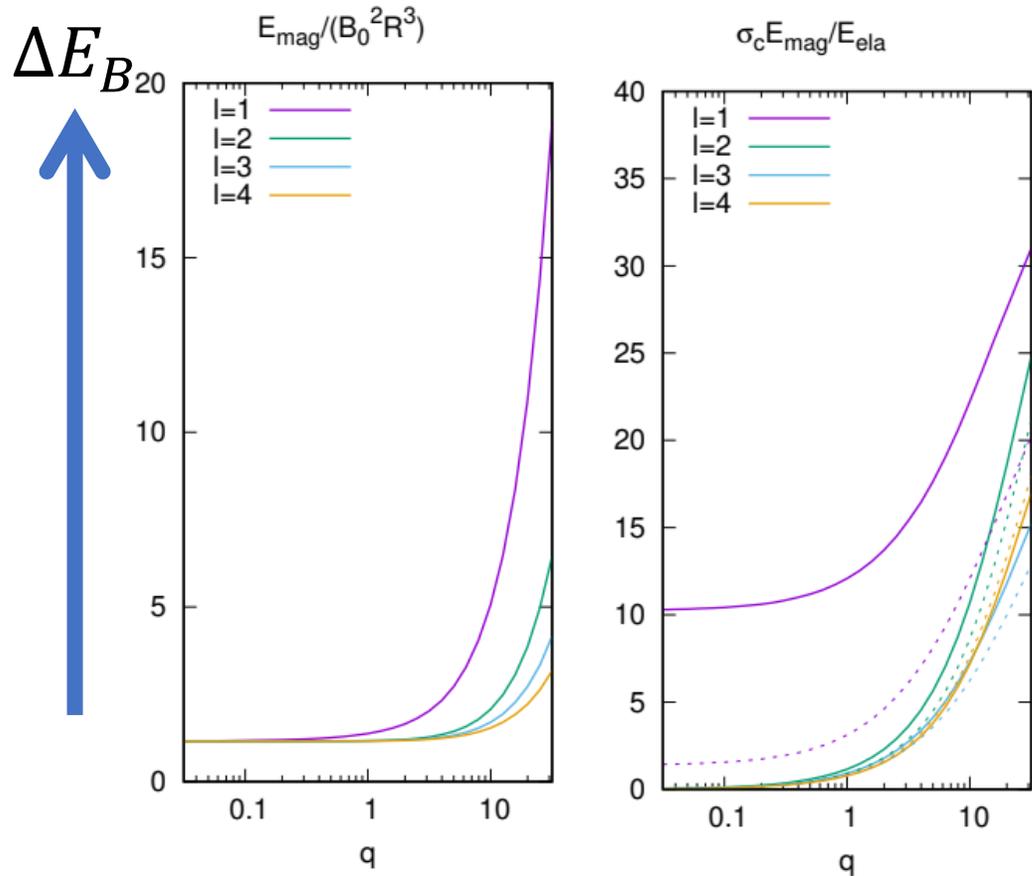
- Static equilibrium determined by
- $\nabla(T_{grav} + T_{fluid} + T_{mag} + T_{ela}) = 0$
- $|T_{grav}| \sim |T_{fluid}| \gg |T_{mag}|, |T_{ela}|$
- **Vector decomposition of the Lorentz force**
Irrotational + **Solenoidal** vector
 $(f = -\nabla F + \nabla \times A)$
- **Solenoidal** part is balanced with elastic force
- Irrotational part is balanced with dominant forces

Magnetic energy

Energy ratio

Results

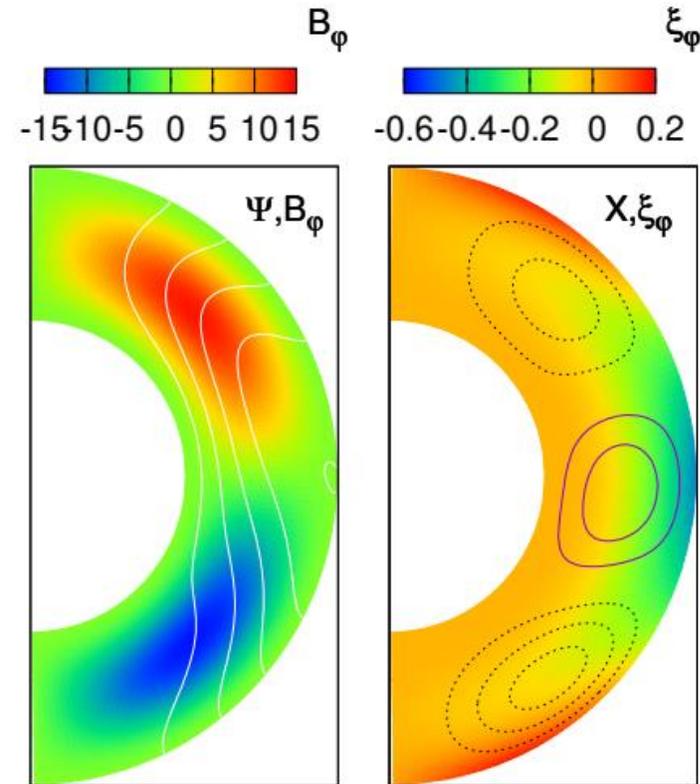
YK+ MN(21b)



Strong magnetic force

Strong elastic force

A large amount of mag. E stored



Magnetic field and elastic displacement

Crust depth enlarged by a factor 5

Summary

- Implication of magneto-elastic equilibrium
 - > Effective /stable confinement of magnetic field in crust
 - Large toroidal magnetic field is also allowed
 - Mixed toroidal-poloidal field may be realized in NS case.
 - => Core field for magnetars?
 - (-> never justify magnetar burst+ evolution models)
 - Braking strain $\sigma_c \sim 0.1$? by theory How do we check?
- We may discuss it in astrophysical objects/events, after lots of theoretical studies, e.g, effect of magnetic field.