Strategy for understanding neutron star matter from HIHR/K1.1 physics

K. Miwa (Tohoku Univ.) on behalf of the HIHR/K1.1 task force



Contents

- Hyperon puzzle in neutron stars
- Strategy for understanding neutron star matter from HIHR/K1.1 physics
- Ap scattering experiment at K1.1 beam line

From Quark to Neutron star

Bound system interacting by strong interaction with completely different scale.



Neutron star

Neutron matter



Strange hadronic matter ?



These system should be understood with the same framework based on the microscopic picture with strong interaction.

Inner core in neutron star

Extremely high-density nuclear matter where baryon-baryon interaction plays essential roles

- ✓ to judge the appearance of strange particles
- \checkmark to understand the mechanism to support the massive neutron star.

Strategy of microscopic understanding of NS (neutron matter)



Constraints on EOS for neutron star from nuclear physics

G.F. Burgio et al., arXiv:2105.03747



S₀ and L₀ values for symmetry energy are constrained from nuclear experiments

Constraints on EOS for neutron star from nuclear physics

G.F. Burgio et al., arXiv:2105.03747



Rather consistent results are obtained up to $\sim 2\rho_0$ including chiral EFT calculation

Constraints on EOS from astrophysical observations



Gravitational waves Tidal deformability of neutron star merger

NICER mission

Simultaneous measurement of both mass and radius

Heavy mass neutron star



What is the key mechanism to make such M-R relation? To answer this question, microscopic approach based on nuclear physics is necessary, with strangeness degree of freedom

G.F. Burgio et al., arXiv:2105.03747

Hyperon puzzle in neutron star

Strange Hadronic Matter in neutron star? Hyperon's appearance is reasonable scenario from the attractive potential in nuclei.



How can we reconcile ?

Softening of EOS w/ hyperon appearance



<u>3 Baryon Force (3BF):</u>

Significant repulsive contribution at high density



We have to derive the nature of the density dependence of ΛN interaction considering the ΛNN force.

Strategy of microscopic understanding of NS with strangeness

Realistic Baryon-Baryon (BB) interaction

Information of all BB interaction ΛN , ΣN , ΞN , $\Lambda \Lambda$

 $SU(3)_{\rm f}$ meson exchange BB models

SU(3)_f chiral EFT force Lattice NLO at present

Lattice BB force

<u>3 Baryon Force with hyperon (3BF)</u>

Phenomenological 3BF TFA + MP repulsion in Nijmegen

Strength of ΛN - ΣN coupling and effect in nuclear medium

 $\Lambda N\mathchar`-\Sigma N$ coupling in free space or a few-body system

BB interaction in medium/nuclear matter

Ab initio calculation No Core Shell Model

hypernuclei

BHF calculation

3BF in chiral EFT for YNN

Hyperon's single particle potential in neutron star matter

EOS with hyperon

This strategy is limited by the insufficient hypernuclear data

Quite scarce YN scattering data

→ Difficult to constrain BB interaction models. More plenty of YN scattering data.

Limited hypernuclear data

Precise hypernuclear spectroscopic data for medium and heavy Λ hypernucleus to determine BE.

What have we provided from J-PARC experiments so far ?

Λ single particle potential at nuclear density



 $U_{\Lambda}(\rho_0)$ ~-30 MeV

Attractive potential at normal nuclear density This is a start point to consider Λ appearance in neutron star matter

High density symmetric nuclear matter 2 + 3 BF2 BF 2 $\rho/
ho_0$

 U_Λ can not be reproduced only with ΛN interaction

This potential value is essential constraint on the ΛNN 3BF at nuclear density



 $U_{\Sigma}(\rho_0) = +30 \pm 10 \text{ MeV} \text{ (real part)}$

Repulsive potential for $\Sigma\text{-nucleus}$ potential from $\Sigma^{\text{--}5}\text{He}$ system

 $U_{\Xi}(\rho_0)$ will be determined by J-PARC experiment Emulsion (E07) + Ξ hypernuclear spectroscopy (E70)

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Toward realistic YN interaction(Experimental side)



Progress on theories

BB interaction Hypernuclear structure





Chiral EFT

Underlying chiral symmetry in QCD

Power counting feature to improve calculation systematically by going to higher order

Multi baryon force appear naturally and automatically in a consistent implementation of the framework



Source of attractive ΛN interaction $\Lambda N\text{-}\Sigma N$ coupling



ΛN-ΣN coupling can be suppressed in nuclear medium due to the Pauli blocking at the intermediate N state
Pauli blocking

w/ ΛN interaction with large ΛN - ΣN coupling + ΛNN three-repulsive force (LECs for ΛNN are adjusted)



Theoretical calculation of Λ binding energy for wide mass range

Comparison between experimental data and theoretical calculation of Λ hypernuclear mass (Λ binding energy) for wide mass range is important to extract the Λ NN force.

Energy spectra of ${}^{13}_{\Lambda}$ C, ${}^{16}_{\Lambda}$ O, ${}^{28}_{\Lambda}$ Si, ${}^{51}_{\Lambda}$ V, ${}^{89}_{\Lambda}$ Y, Λ hypernuclei with chiral YN potential 139 La, 208 Pb with ESC16 model J. Haidenbauer, I. Vidana, Eur. Phys. J. A (2020) 56:55 M.M. Nagels et al. Phys. Rev. C99, 044003 (2019) 0 **NLO19** -5 S Qualitatively good agreement with data 30 Binding energy [MeV] 05 over a wide range of mass number -10p >-15 ₩ □[<] -20 d -25 w/ ΛN int. w/o ΛNN int. 10 g w/ ΛN int. w/ ΛNN int. -30 0.15 0.00 0.05 0.10 0.20 0.05 0.1 0.15 0.2 0.25 0.3 0.35 0 Over bound w/ only ΛN int. $A^{-2/3}$ Sizable cutoff dependence ΛNN repulsive interaction is necessary to \rightarrow would be reduced at higher order and after explain Λ hypernuclear bounding energy inclusion of 3BF

Possibility of density dependence of ΛN int. from excited states



 k_{F} from ADA ${}^{40}{}_{\Lambda}K(s): k_{F} = 1.26 \text{ fm}^{-1}$ ${}^{40}{}_{\Lambda}K(p): k_{F} = 1.15 \text{ fm}^{-1}$ ${}^{40}{}_{\Lambda}K(d): k_{F} = 1.02 \text{ fm}^{-1}$

Possibility to check the density dependent ΛN interaction from each Λ orbital state



Example: ${}^{40}_{\Lambda}$ K



<u>B</u>_A values reliable?

Hotchi et al.; PRC 64 (2001) 044302 250 g_{Λ} f_{Λ} ⁸⁹Y(π^+ ,K⁺) reaction 200 $\Delta E = 1.64 \text{ MeV} (FWHM)$ Counts / 0.25MeV \mathbf{d}_{Λ} 89 49 39Λ p_Λ 50 \mathbf{S}_{Λ} 2 20 0 **-**30 -25 -20 -15 -10 5 10 -5 D B_{Λ} (MeV) Λ 's $\hbar \omega \sim 4$ MeV for ${}^{208}_{\Lambda}$ Pb

Spacing of hole states (n orbit) \gtrsim 300 keV

=> a few 100 keV (FWHM) resolution is necessary to decompose the hole states





 $^{208}{\rm Pb}(\pi^+,{\rm K}^+)^{208}_{\Lambda}{\rm Pb},\,{\rm p}_{\pi}=1.06~{\rm GeV/c}$





HIHR/K1.1 strategy



HIHR/K1.1 physics (microscopic understanding of NS matter) <u>Hyperon puzzle in neutron star</u>

Softening of EOS w/ hyperon appearance How can massive NS be supported ?

Realistic ΛN interaction based on Λp scattering experiment

 \rightarrow K1.1 beam line : maximum Λ production cross section $d\sigma/d\Omega$ and Spin observables (100 times better statistics)

Understanding of 3Baryon Force (3BF) including hyperon is essential





Essential input to establish realistic ΛN interaction

Investigation of ΛNN 3BF by ultra high-resolution Λ hypernuclear spectroscopy
 →HIHR beam line

<u>New generation spectroscopy with momentum-dispersion</u> <u>matching method</u>

High-resolution Λ hypernuclear spectroscopy at HIHR



Strategy of microscopic understanding of NS with strangeness

Realistic Baryon-Baryon (BB) interaction



Short summary of HIHR/K1.1 physics

• We propose our strategy of hypernuclear physics based on the outputs from the HIHR/K1.1 beam lines

We are going to provide

- ✓ high-statistics YN scattering data at K1.1 beam line
- \checkmark Ultra precise Λ hypernuclear spectroscopic data at HIHR beam line

By collaborating with theorist, we are going to establish "realistic YN interaction" Extract the strength of YNN 3BF phenomenologically (or determine LEC's for 3BF) by comparing reliable many-body calculations based on the realistic YN interaction and high precision hypernulear data.

- We expect the high-density matter can be microscopically examined in NN sector and also with YN/YY sectors.
- These strategies are essential for the understanding of baryonic matter in neutron star from the microscopic picture based on nuclear physics

New proposal for J-PARC experiment

Measurement of the differential cross section and spin observables of the $\Lambda\,p$ scattering with a polarized Λ beam

ΛN interaction and its uncertainty

Comparison between experimental data and theoretical calculation of Λ hypernuclear mass (Λ binding energy) for wide mass range is important to extract the Λ NN force.

Energy spectra of ${}^{13}{}_{\Lambda}$ C, ${}^{16}{}_{\Lambda}$ O, ${}^{28}{}_{\Lambda}$ Si, ${}^{51}{}_{\Lambda}$ V, ${}^{89}{}_{\Lambda}$ Y, ${}^{139}{}_{\Lambda}$ La, ${}^{208}{}_{\Lambda}$ Pb with ESC16 model



 Λ NN repulsive interaction is necessary to explain Λ hypernuclear bounding energy

Choice of two-body ΛN interaction has large effect of theoretical calculation

M. Isaka et al., Phys. Rev. C 95, 044308 (2017)

Calculation with only two-body interaction of different ESC versions



ESC12 can explain $-B_{\Lambda}$ w/o Λ NN repulsive int. ESC14 need Λ NN repulsive int.

ΛN interaction and its uncertainty

M. Isaka et al., Phys. Rev. C 95, 044308 (2017)

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Large uncertainty of P-wave potential

Calculation with only two-body interaction of different ESC versions

ESC12 : more repulsive in P-wave potential

 \rightarrow Experimental Λ binding energy can be reproduced with only two-body Λ N interaction

ESC14 : moderately repulsive in P-wave potential

 \rightarrow Λ binding energy with only two-body Λ N interaction is too large.

 ΛNN repulsive force is necessary to reproduce experimental binding energy

Difference in P-wave contribution should appear as different angular distribution in $d\sigma/d\Omega$ of Λp scattering

AN interaction and its uncertainty

Paper of chiral EFT J. Haidenbauer et al. Eur. Phys. J. A. (2020) 56:91

allow one to pin down the interaction in the S = -1 sector. It should be emphasized that the aspects discussed above apply only to the interaction in the S waves. Since there are practically no data for differential observables, it is impossible to fix the YN contact terms in the P-waves. In this case, imple-

Total cross section data exist for Λp scattering. But no differential cross section data.

Total cross section might be similar for each model.

However, there are large uncertainty in the P-wave interaction.





Such difference clearly appears in differential observables ! We should measure these differential cross sections and spin observables to constrain theoretical model significantly.

Conventional representation of elastic scattering
Scattering amplitude in
$$\frac{1}{2} + \frac{1}{2} \rightarrow \frac{1}{2} + \frac{1}{2}$$
 scattering : $\rightarrow 4 \times 4$ matrix
 $\rightarrow 6$ components from the restriction of parity conservation and time-reversal invariance
spin-independent spin-spin symmetric LS (Λ S=0) anti-symmetric LS (Λ S=1)
Tensor
T matrix
 $M = V_c + V_\sigma(s_a \cdot s_b) + V_{SLS}(s_a + s_b) \cdot L + V_{ALS}(s_a - s_b) \cdot L + V_T([s_a \otimes s_b]^{(2)} \cdot Y_2(\hat{r})),$
Scalar amplitude
 $U_a \equiv \langle k_f | V_a | k_i \rangle = \langle k_f | V_a | k_i \rangle$
 $M = v_c + l_{\sigma}(s_a \cdot s_b) + V_{SLS}(s_a + s_b) \cdot L + V_{ALS}(s_a - s_b) \cdot L + V_T([s_a \otimes s_b]^{(2)} \cdot Y_2(\hat{r})),$
Scalar amplitude
 $U_a \equiv \langle k_f | V_a | k_i \rangle = \langle k_f | V_{ALS} L_a | k_i \rangle, S_{ALS} \equiv \langle k_f | V_{SLS} L_a | k_i \rangle = T_{2} = \frac{1}{2} \langle k_f | V_T Y_{2j-1} | k_i \rangle$
 $M = are going to measure following observables.
Differential cross section
 $(\frac{d\sigma}{d\Omega}) = \frac{1}{4} \text{Tr}(MM^1) = |U_a|^2 + \frac{3}{16} |U_\beta|^2 + \frac{1}{2} (|S_{SLS}|^2 + |S_{ALS}|^2) + \frac{1}{4} |T_1|^2 + \frac{1}{2} (|T_2|^2 + |T_3|^2).$
Analyzing power
(Polarization)
 $D_y^u = \frac{1}{\sigma(\theta)} \text{Re} \left\{ \frac{1}{2\sqrt{3}} (U_0 + \frac{1}{\sqrt{3}} U_1)^* U_1 + \frac{1}{2} (U_0 - \frac{1}{\sqrt{3}} U_1)^* (\frac{1}{\sqrt{6}} T_1 + T_3) - S_1^* S_2 + \frac{1}{2} |S_3|^2 - \frac{1}{\sqrt{6}} T_1^* (\frac{1}{\sqrt{6}} T_1 - T_3) - \frac{1}{2} |T_2|^2 \right\}.$
Number of observables is still limited to determine each component senarately.$

Number of observables is still limited to determine each component separately. But measurements of many observables contribute to impose constraints on YN theoretical models.

Conventional representation of elastic scattering Analyzing power Depolarization (D_v^y)

Left/Right asymmetry of Ap scattering

1 × y /

Change the spin polarization after the Λp scattering



We are going to measure following observables.

Differential cross section

Analyzing power (Polarization)

Depolarization

$$\left(\frac{d\sigma}{d\Omega}\right) = \frac{1}{4} \operatorname{Tr}(MM^{\dagger}) = |U_{\alpha}|^{2} + \frac{3}{16} |U_{\beta}|^{2} + \frac{1}{2} (|S_{SLS}|^{2} + |S_{ALS}|^{2}) + \frac{1}{4} |T_{1}|^{2} + \frac{1}{2} (|T_{2}|^{2} + |T_{3}|^{2}).$$

$$A_{y}(Y) = -\frac{1}{\sqrt{2}\sigma(\theta)} \operatorname{Im} \left\{ (U_{\alpha} + \frac{1}{4}U_{\beta})^{*}S_{SLS} + (U_{\alpha} - \frac{1}{4}U_{\beta})^{*}S_{ALS} - \frac{1}{2}T_{\alpha}^{*}(-S_{ALS} + S_{SLS}) \right\},$$

$$D_{y}^{y} = \frac{1}{\sigma(\theta)} \operatorname{Re} \left\{ \frac{1}{2\sqrt{3}} \left(U_{0} + \frac{1}{\sqrt{3}}U_{1} \right)^{*} U_{1} + \frac{1}{2} \left(U_{0} - \frac{1}{\sqrt{3}}U_{1} \right)^{*} \left(\frac{1}{\sqrt{6}}T_{1} + T_{3} \right) - S_{1}^{*}S_{2} + \frac{1}{2} |S_{3}|^{2} - \frac{1}{\sqrt{6}}T_{1}^{*} \left(\frac{1}{\sqrt{6}}T_{1} - T_{3} \right) - \frac{1}{2} |T_{2}|^{2} \right\}.$$

Number of observables is still limited to determine each component separately. But many observables contribute to impose constraints on YN theoretical models.

Proposal for an experiment at the 50-GeV PS

Measurement of the differential cross section and spin observables of the Λp scattering with a polarized Λ beam

New project: Λp scattering experiment at K1.1 beam line with polarized Λ beam

K. Miwa(spokesperson), S. H. Hayakawa, Y. Ishikawa, K. Itabashi, K. Kamada, T. Kitaoka, T. Morino, S. Nagao, S. N. Nakamura, F. Oura, T. Sakao, H. Tamura, H. Umetsu, S. Wada *Tohoku University, Japan*

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> T. Nanamura Kyoto University, Japan

K. Shirotori Research Center for Nuclear Physics (RCNP), Osaka University, Japan

> P. Evtoukhovitch Joint Institute for Nuclear Research (JINR), Russia

Z. Tsamalaidze Joint Institute for Nuclear Research (JINR), Russia Georgian Technical University (GTU), Tbilisi, Georgia
K1.1 beam line at extended hadron hall



Pmax = 1.2 GeV/c Two stages of Elec. Separator

Advantage of Λ production at 1.05 GeV/c

<u>Maximum Λ production cross section</u>

<u>High spin polarization of Λ for Λ production plane</u>



Ap scattering experiment at K1.1 beam line

 Λ beam identification

Tagged by $\pi^-p \rightarrow K^0\Lambda$ reaction at p=1.05 GeV/c



<u>Ap scattering identification</u>

Detected by CATCH



Ap scattering experiment at K1.1 beam line

Λ beam identification



Momentum-tagged Λ beam 4000 3500 0.4 < p(GeV/c) < 0.83000 2500 2000 1500 1000 500 8.4 0.8 0.85 0.45 0.75 Momentum (GeV/c)

Goal of experiment

30 days production : 50 M Λ beam \rightarrow (d σ /d Ω)₀, A_y(Λ)

+30 days production : 100 M Λ beam in total $\rightarrow D^{y}_{y} + (d\sigma/d\Omega, A_{y}(\Lambda) \text{ w/ improved statistics})$

Feasibility study in E40 (Σp scattering)

By using by-product data in E40, we already checked feasibility of Ap scattering



Feasibility study in E40 byproduct data (Ap scattering)



Ap scattering with polarized Λ beam



- ✓ Selection of model (chiral EFT, Nijmegen, Julich) is possible
- ✓ Selection in each model
 - ✓ NSC97f, ESC16 OK
 - ✓ ChiralEFT13, 19 OK



cost

Simulated results w/ 100M Λ

Simulated results w/ 100M Λ



In the middle momentum range $(0.5 \sim 0.7 \text{ GeV/c})$, 10% level accuracy can be achieved.

We believe that these new scattering data becomes important constraint to determine spin-dependent AN interaction

Summary of Λp scattering experiment

- BB interactions are essential ingredient to understand NS matter Both theoretical progresses and experimental progress ! Now is a time to make effort to establish the realistic BB interaction in collaboration with theory and experiment.
- Experimental method of YN scattering is established Systematic measurement of Σp scattering was performed at J-PARC Ap scattering experiment is also feasible.
- There are large uncertainties in the P- and higher waves in ΛN interaction Measurement of differential observables (d σ /d Ω , spin observables) are important
- New project to measure dσ/dΩ and spin observables of Λp scattering at K1.1 beam line Λp scattering with ~100% polarized Λ beam
 pΛ: 0.4 ~ 0.8 GeV/c (Threshold region of ΣN channel)
 30-days beam time : 50 M Λ beam
 → dσ/dΩ and A_y(Λ) measurements
 +30-days beam time : total 100 M Λ beam
 → D_{yy} and precice dσ/dΩ and A_y(Λ)

Thank you for your attention

We need closer collaboration with theorists and experimentalists in the world.

Please give us your comments and suggestion in any time.

Chiral EFT

Underlying chiral symmetry in QCD

Power counting feature to improve calculation systematically by going to higher order

Multi baryon force appear naturally and automatically in a consistent implementation of the framework



Chiral EFT toward NNLO

Two body YN force

Three body YNN force



These 3BF should be considered for calculation of neutron star matter in future

In near future, chiral EFT can be extended to

- NNLO for two body YN w/ accurate YN scattering data •
- NLO for three-body YNN w/ Λd scattering or few-body hypernuclear binding energies ٠

U_{Λ} dependence for M-R relation

Y. Yamamoto et al. PRC90, 045805 (2014)



Density dependence of ANN 3BF is roughly quadratic ($\propto \rho^2$).

ANN strength at normal nuclear density (ρ_0) becomes 9~16 times more effect at 3~4 ρ_0 region.

Precise determination of binding energy is very important.

ΛN interaction and its uncertainty

Isaka et al. pointed out that choice of two-body AN interaction has large effect of theoretical calculation





M. Isaka et al., Phys. Rev. C 95, 044308 (2017)

Large P-wave uncertainty of Potential

ESC12 can explain $-B_{\Lambda}$ w/o Λ NN repulsive int. ESC14 need Λ NN repulsive int. Uncertainty in the theoretical microscopic manybody calculation based on realistic NN interactions

Pure neutron matter

Symmetric nuclear matter



EOS with Quark-Hadron Crossover





Cross over $5.5n_0 > Quark phase$

Repulsive force in vector channel is necessary for supporting heavy neutron star

A hyperon appear at 0.42 fm^{-3} in Togashi EOS

Interpolation region from baryon to quark

Hyperon interaction is necessary to get better description at this region



This calculation depends on two-body ΛN interaction very much.

 \leftarrow Owing to Large uncertainty in P-wave ΛN potential

Baryon-baryon interaction model's picture

There exists various theoretical models based on its own fundamental degree of freedom.



Nijmegen model

Widely used interaction in hypernuclear physics

Nijmegen potential + G-matrix calculation is powerful method to study hypernuclei and neutron star



Yamamoto et al. Phys. Rev. C 90 (2014) 045805





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Underlying chiral symmetry in QCD

Power counting feature to improve calculation systematically by going to higher order

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Recent experimental progress at J-PARC



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Recent experimental progress at J-PARC

K. Miwa et al. arXiv:2104.13608



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Neutron star and YN interaction

Two-body YN scattering is essential to understand the internal structure of neutron star.

- Interaction at short range
- Basic information to derive 3 body force from hypernuclear structure

Hypernuclear physics based on Realistic YN interaction





Λ binding energy and ΛNN interaction

Energy spectra of ${}^{13}{}_{\Lambda}$ C, ${}^{16}{}_{\Lambda}$ O, ${}^{28}{}_{\Lambda}$ Si, ${}^{51}{}_{\Lambda}$ V, ${}^{89}{}_{\Lambda}$ Y, ${}^{139}{}_{\Lambda}$ La, ${}^{208}{}_{\Lambda}$ Pb with ESC16 model

M.M. Nagels et al. Phys. Rev. C99, 044003 (2019)



ESC16 need Λ NN repulsive interaction to explain Λ hypernuclear bounding energy



Sizable cutoff dependence → Would be reduced at higher order and after inclusion of 3BF

Λ binding energy and ΛNN interaction

Energy spectra of ${}^{13}{}_{\Lambda}$ C, ${}^{16}{}_{\Lambda}$ O, ${}^{28}{}_{\Lambda}$ Si, ${}^{51}{}_{\Lambda}$ V, ${}^{89}{}_{\Lambda}$ Y, ${}^{139}{}_{\Lambda}$ La, ${}^{208}{}_{\Lambda}$ Pb with ESC16 model

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Essential measurement to derive ΛNN int.

ESC16 need Λ NN repulsive interaction to explain Λ hypernuclear bounding energy

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M. Isaka et al., Phys. Rev. C 95, 044308 (2017)



ESC12 (similar with NSC97f) can explain $-B_{\Lambda}$ w/o Λ NN repulsive int.

ESC14 (similar with ESC16) need Λ NN repulsive int.

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Energy spectra of ${}^{13}_{\Lambda}$ C, ${}^{16}_{\Lambda}$ O, ${}^{28}_{\Lambda}$ Si, ${}^{51}_{\Lambda}$ V, ${}^{89}_{\Lambda}$ Y, ${}^{139}_{\Lambda}$ La, ${}^{208}_{\Lambda}$ Pb with ESC16 model

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allow one to pin down the interaction in the S = -1 sector. It should be emphasized that the aspects discussed above apply only to the interaction in the S waves. Since there are practically no data for differential observables, it is impossible to fix the YN contact terms in the P-waves. In this case, imple-

Even for the ΛN interaction, there are large uncertainty in the P-wave interaction.

Total cross section might be similar for each model. However, the angular dependence is very different.

We need accurate scattering data in the P-wave and higher regions

As many scattering observables as possible to construct realistic YN interaction



Conventional representation of elastic scattering
Scattering amplitude in
$$\frac{1}{2} + \frac{1}{2} \rightarrow \frac{1}{2} + \frac{1}{2}$$
 scattering : $\rightarrow 4 \times 4$ matrix
 $\rightarrow 6$ components from the restriction of parity conservation and time-reversal invariance
spin-independent spin-spin symmetric LS (Δ S=0) anti-symmetric LS (Δ S=1)
Tensor
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 $M = V_c + V_\sigma(s_a \cdot s_b) + V_{SLS}(s_a + s_b) \cdot L + V_{ALS}(s_a - s_b) \cdot L + V_T([s_a \otimes s_b]^{(2)} \cdot Y_2(\hat{r})),$
Scalar amplitude
 $U_a \equiv \langle k_f | V_a | k_i \rangle = \langle k_f | V_{aLS} \rangle = \langle k_f | V_{ALS} L_1 | k_i \rangle, S_{ALS} \equiv \langle k_f | V_{SLS} L_1 | k_i \rangle = T_2 = \frac{1}{2} \langle k_f | V_T Y_{2j-1} | k_i \rangle$
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Scalar amplitude
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 $M = are going to measure following observables.$
Differential cross section
 $(\frac{d\sigma}{d\Omega}) = \frac{1}{4} Tr(MM^{\dagger}) = |U_a|^2 + \frac{3}{16} |U_\beta|^2 + \frac{1}{2} (|S_{SLS}|^2 + |S_{ALS}|^2) + \frac{1}{4} |T_1|^2 + \frac{1}{2} (|T_2|^2 + |T_3|^2).$
Analyzing power
(Polarization)
 $D_y^2 = \frac{1}{\sigma(\theta)} \operatorname{Re} \left\{ \frac{1}{2\sqrt{3}} \left(U_0 + \frac{1}{\sqrt{3}} U_1 \right)^* U_1 + \frac{1}{2} \left(U_0 - \frac{1}{\sqrt{3}} U_1 \right)^* \left(\frac{1}{\sqrt{6}} T_1 + T_3 \right) - S_1 S_2 + \frac{1}{2} |S_3|^2 - \frac{1}{\sqrt{6}} T_1^* \left(\frac{1}{\sqrt{6}} T_1 - T_3 \right) - \frac{1}{2} |T_2|^2 \right\}.$
Number of observables is still limited to determine each component senarately.

But many observables contribute to impose constraints on YN theoretical models.

Analyzing power



We will measure $d\sigma/d\Omega$ for left and right scatted event separately



Up/Down asymmetry for depolarization measurement

Spin polarization in the final state

$$P_{scat} = \frac{2}{\alpha} \frac{N_U - N_D}{N_U + N_D} = \frac{P + D_y^{\mathcal{Y}} P_{beam}}{1 + P P_{beam}}$$

 P_{beam} : Polarization of beam $P_{beam} = P(\phi = 0) \times \cos \phi$ P: Induced polarization by the unpolarized beam D_y^y : Depolarization

U/D asymmetry of detector is important



Up/Down asymmetry for depolarization measurement


Up/Down asymmetry for depolarization measurement

Spin polarization in the final state

$$P_{scat} = \frac{2}{\alpha} \frac{N_U - N_D}{N_U + N_D} = \frac{P + D_y^{\mathcal{Y}} P_{beam}}{1 + P P_{beam}}$$

 P_{beam} : Polarization of beam $P_{beam} = P(\phi = 0) \times \cos \phi$ P : Induced polarization by the unpolarized beam D_y^y : Depolarization

U/D asymmetry of detector is important



Spin observables in Ap scattering

Simulated results w/ 100M Λ



In the middle momentum range $(0.5 \sim 0.7 \text{ GeV/c})$, 10% level accuracy can be achieved.

Toward construction of realistic YN interaction

Even for YN case

We need help of theorist and experience of NN experimentalist. We are very happy if we make a such collaboration for the phase-shift analysis

Phase-shift analysis

Two-body scattering data



Small relative-momentum region



0.5

cosθ

-0.5

n

0.5

cost

Mass number dependence

Widely used interaction in hypernuclear physics

Nijmegen potential + G-matrix calculation is powerful method to study hypernuclei and neutron star



Yamamoto et al. Phys. Rev. C 90 (2014) 045805

Neutron star matter and YN/YNN interactions



HIHR/K1.1 physics

- $\frac{\text{K1.1 beam line}}{\text{Optimized for S}=-1 \text{ physics}} \begin{array}{l} \text{Suitable to produce} \\ \overrightarrow{\text{polarized }} \Lambda \end{array}$
- Ap scattering experiment w/ polarized Λ beam



Time We need accumulation of many YN scattering observables

High-Intensity and High-Resolution beam line

Dispersion-matching beam line → NO need to track beam particles Potential high-intensity beam at J-PARC can be used





d

-15

-10

 $-B_{\Lambda}$ (MeV)

-5

10

(0.1

Counts

T.Hasegawa et

al., Phys. Rev. C

53 (1996) 1210

190

185

180

M_{HY}-M_A (MeV)

170

0.1

3000

2000

1000

-30

S

-25

-20

for wide-mass range

Density dependence of ΛN interaction (ANN interaction)

Calculate U_{Λ} at high density region Untangle hyperon puzzle in neutron stat



Chiral EFT

Underlying chiral symmetry in QCD

Power counting feature to improve calculation systematically by going to higher order

Multi baryon force appear naturally and automatically in a consistent implementation of the framework



Λ binding energy and ΛNN interaction

Energy spectra of ${}^{13}{}_{\Lambda}$ C, ${}^{16}{}_{\Lambda}$ O, ${}^{28}{}_{\Lambda}$ Si, ${}^{51}{}_{\Lambda}$ V, ${}^{89}{}_{\Lambda}$ Y, ${}^{139}{}_{\Lambda}$ La, ${}^{208}{}_{\Lambda}$ Pb with ESC16 model

M.M. Nagels et al. Phys. Rev. C99, 044003 (2019)



ESC16 need Λ NN repulsive interaction to explain Λ hypernuclear bounding energy



Sizable cutoff dependence → Would be reduced at higher order and after inclusion of 3BF

Toward construction of realistic YN interaction

Even for YN case

We need help of theorist and experience of NN experimentalist. We are very happy if we make a such collaboration for the phase-shift analysis

Phase-shift analysis

Two-body scattering data





Single particle potential at nuclear density



Attractive potential at normal nuclear density

 $U_{\Xi}(\rho_0)$ will be determined by J-PARC experiment Emulsion (E07) + Ξ hypernuclear spectroscopy (E70)

Other experiment

LN interaction

LN interaction can be modified easily due to the LN-SN coupling

LN interaction in free space should be determined.