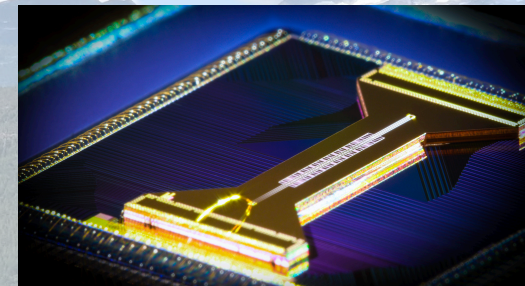
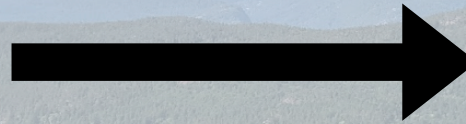
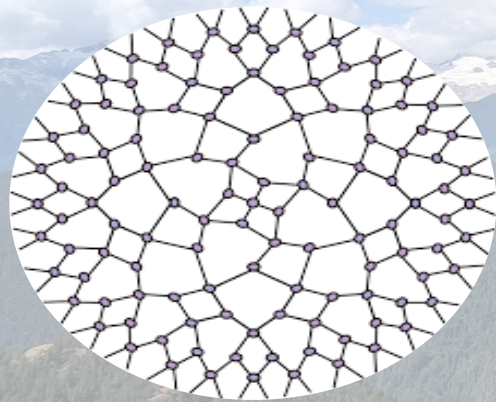


# Simulating highly-entangled matter with quantum tensor networks

Andrew C. Potter (UT Austin -> UBC)



U.S. DEPARTMENT OF  
**ENERGY**



Alfred P. Sloan  
FOUNDATION

**Honeywell**

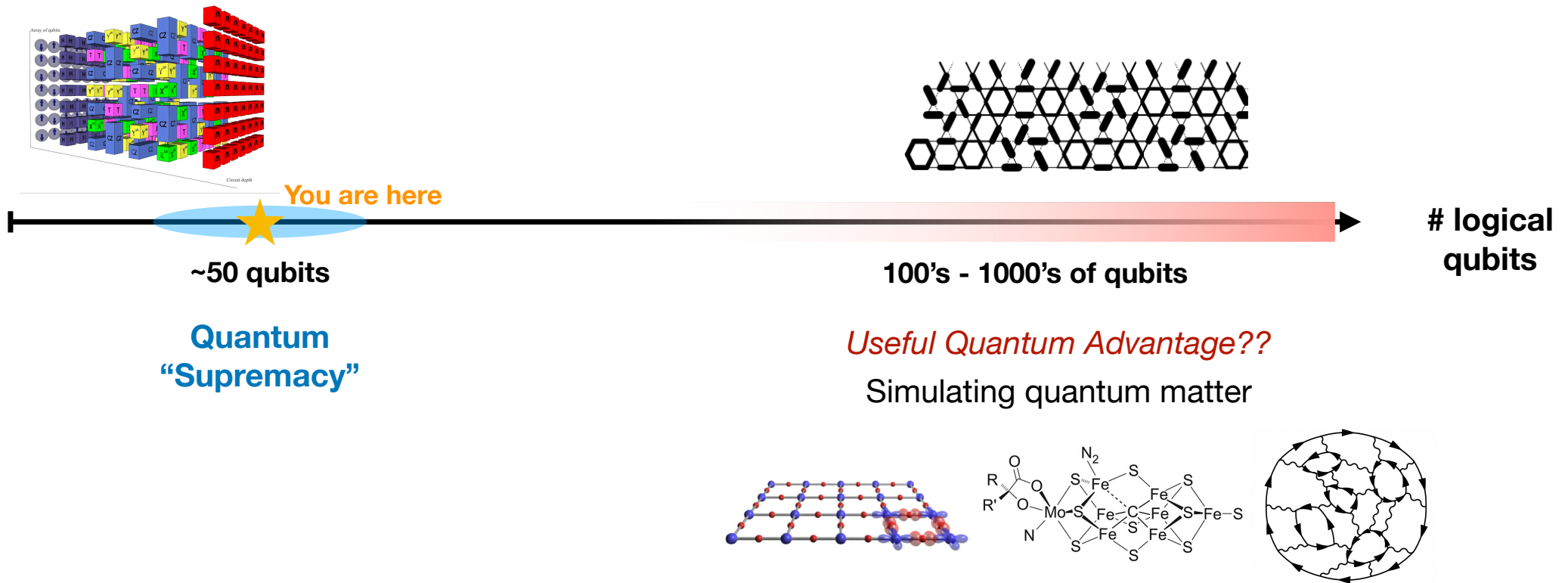


THE UNIVERSITY OF  
**TEXAS**  
— AT AUSTIN —



Stewart Blusson  
**Quantum Matter Institute**  
THE UNIVERSITY OF BRITISH COLUMBIA

# Practical quantum computation?

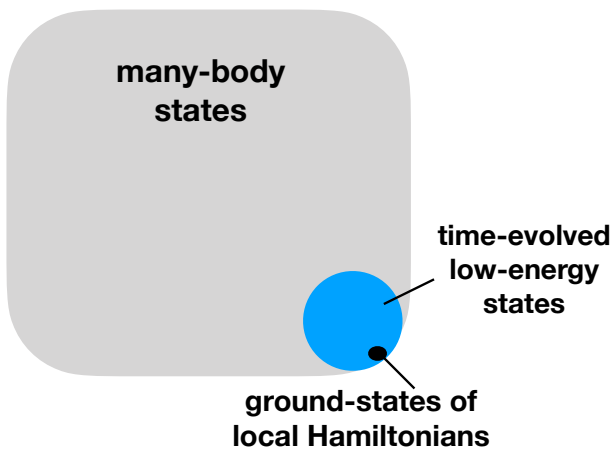


**Goal:** Port successful classical strategies (**tensor-network representations**) to the quantum realm

**Strategy:** Focus quantum resources on classically-hard aspects of computation



# Tensor Network States (TNS)

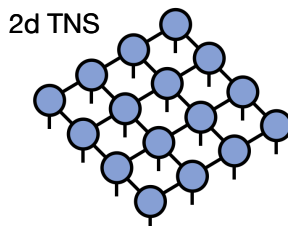
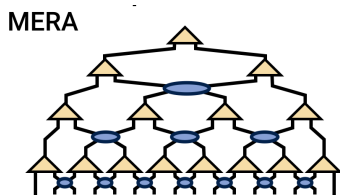
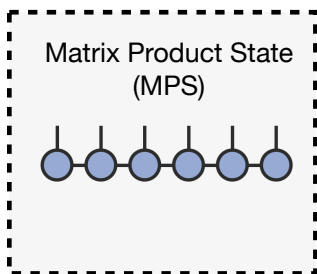


$$\langle s_1, s_2, s_3 | \Psi \rangle = \text{Diagram} = \sum_{ijk=1}^{\chi} A_{ik}^{s_1} B_{ij}^{s_2} C_{kj}^{s_3}$$

Physical DOF

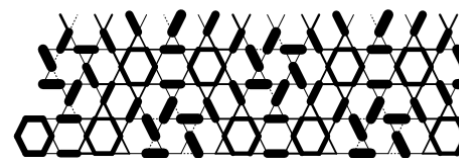
Hidden/"Bond" DOF

"Bond-dimension"



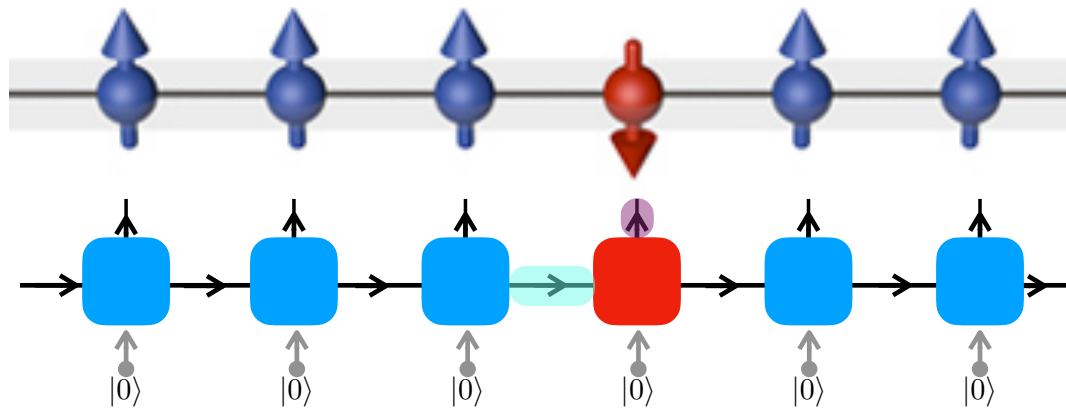
Higher-d MPS:  $\chi \sim e^{L^{d-1}} L^p$  (gapless)

Can simulate 100's-1000's of spins!



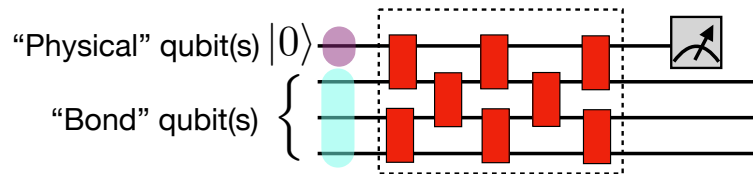
e.g. Yan, Huse, White, Science '10

# Quantum Matrix Product States (qMPS)



$$\chi = 2^{\#\text{bond-qubits}}$$

$$\# \text{ hardware qubits} = \log_2 \chi + L^{d-1}$$

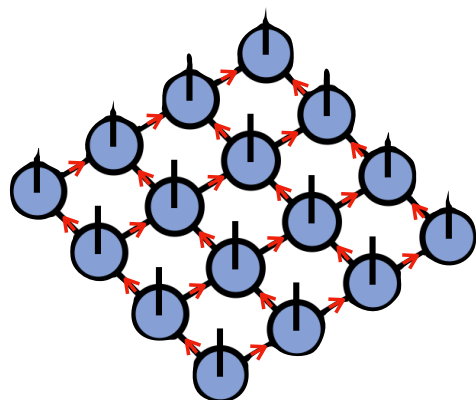


**“Holographic” simulation:**  
 Simulate  $d$ -dimensions w/  $(d-1)$   
 dimensions’ worth of qubits

*Schoen, Hammerer, Wolf, Cirac PRA ‘07*  
*Foss-Feig, Hayes, ... ACP PRR ‘21*  
*Barratt, Dborin, Bal, Stojevic, Green, Pollmann NPJQI ‘21*

# Representational Power of qTNS?

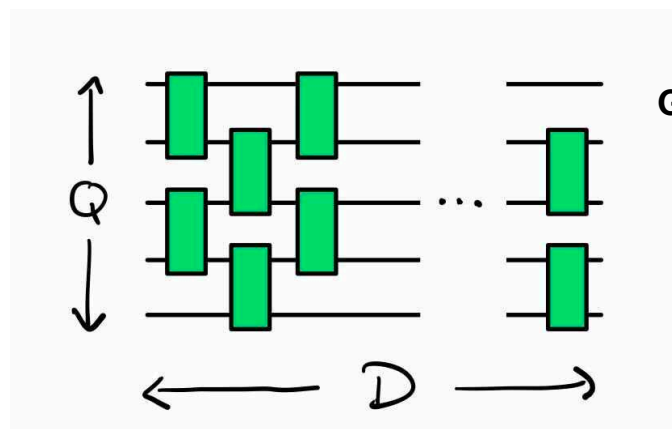
## Restriction 1: Tensors = Isometries



- **1d:** No Restriction (“Canonical Form” )
- **2d, 3d: ??** [Many interesting states have low-bond dimension isoTNS representations]

*Soejima, Siva, Bultinck, Chatterjee, Pollmann, Zaletel PRB '20*

## Restriction 2: Circuit Resources



Generic “Dense” Unitary:

$$D \sim e^Q$$

*(e.g. Solovay, Kitaev)*

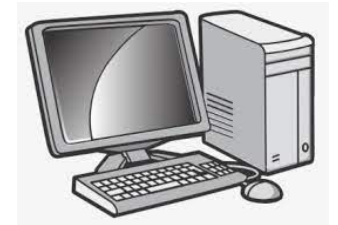
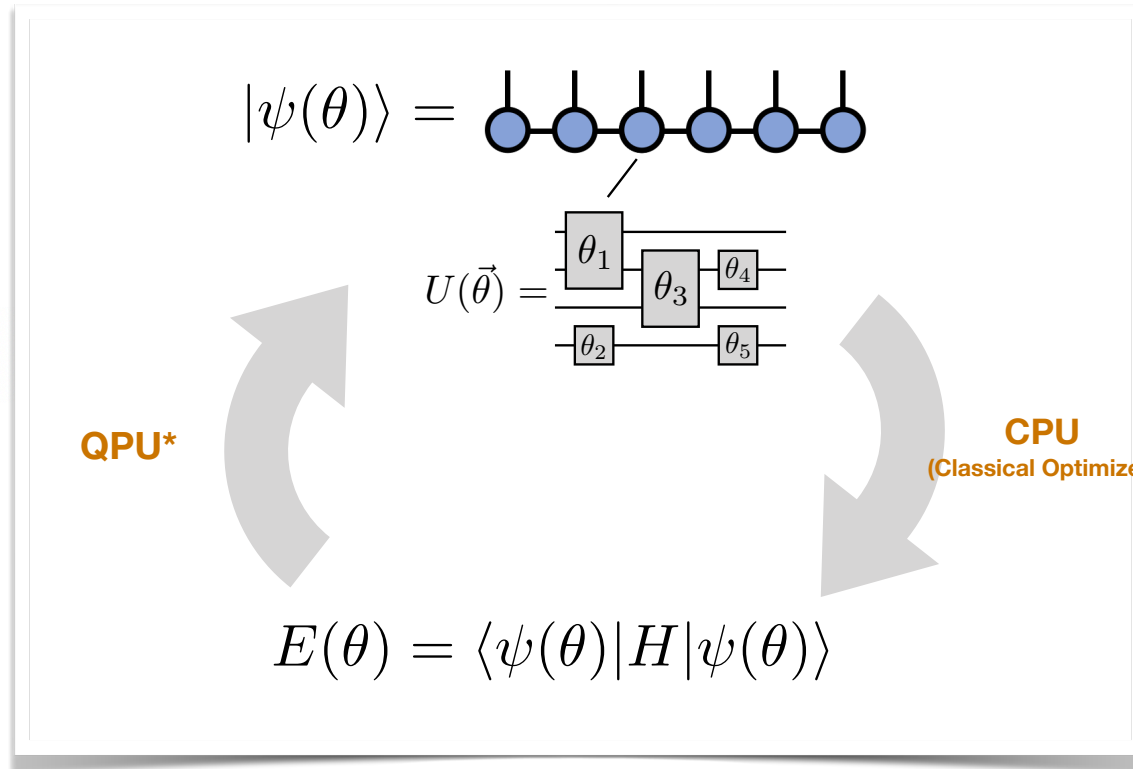
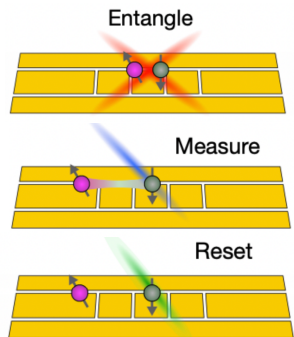
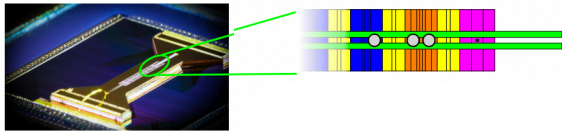
- **Key Question:** What class of states can be captured w/  $D \sim \text{poly}(Q, L)$ ?
- **(Partial) Answer:**
  - Anything continuously connected to non-interacting electrons
  - (Non-chiral) topological orders

*Niu, Zhang, Haghenas, Chan, ACP [to appear]*



# qMPS Ground-State Preparation: “Holographic VQE”

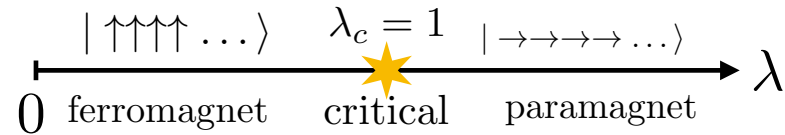
Honeywell H1  
Trapped-Ion Processor



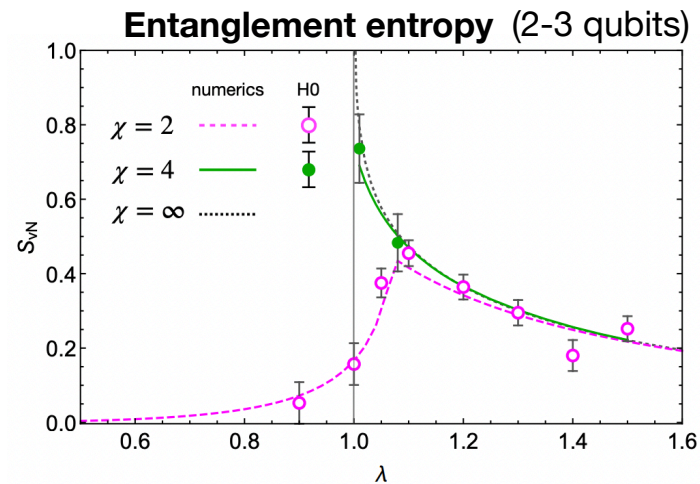
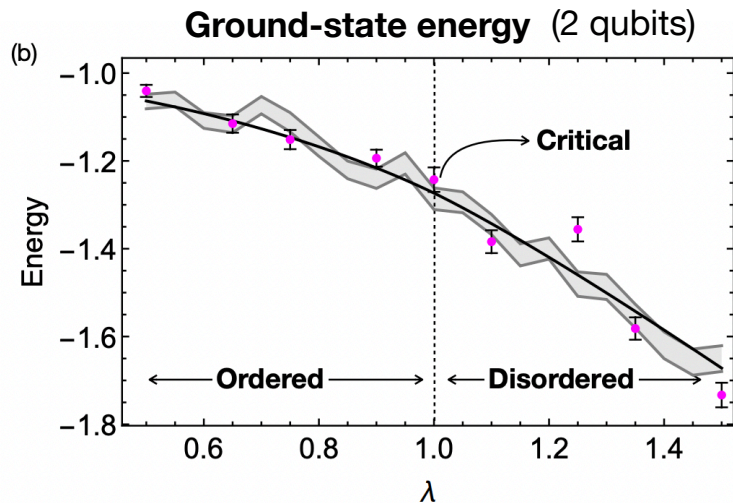
\*Data shown will be classically pre-optimized, then implemented on QPU

# Variational qMPS: Experimental Demo

$$H_{\text{TFIM}} = - \sum_j (\sigma_j^z \sigma_{j+1}^z + \lambda \sigma_i^x)$$

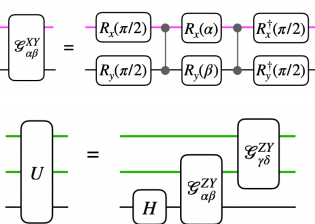


$$\langle \sigma_r \sigma_0 \rangle \sim 1/r^{2\Delta} \quad S \sim \frac{1}{12} \log L \quad \text{Cardy, Calabrese}$$

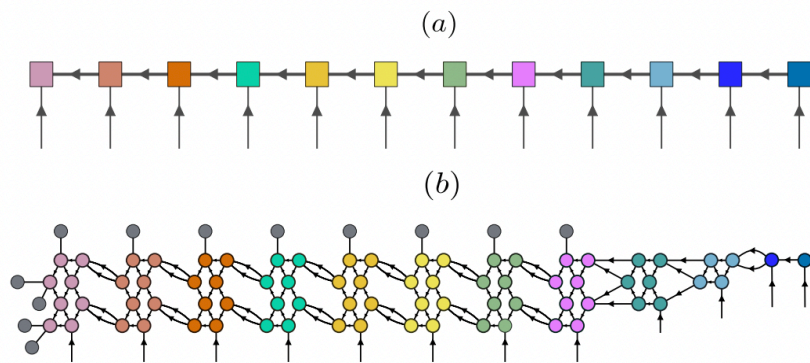


**# Qubits used: 2-3  
Vs. >~ 40 qubits!!**

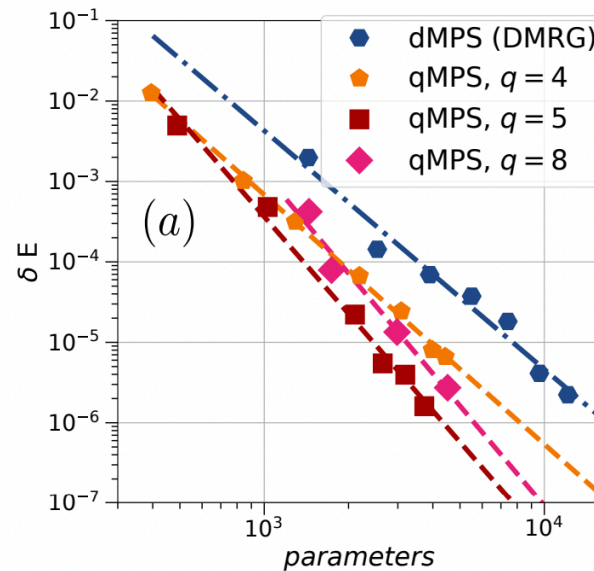
*Foss-Feig, Hayes, et al. [Honeywell + ACP] arXiv:2104.11235*



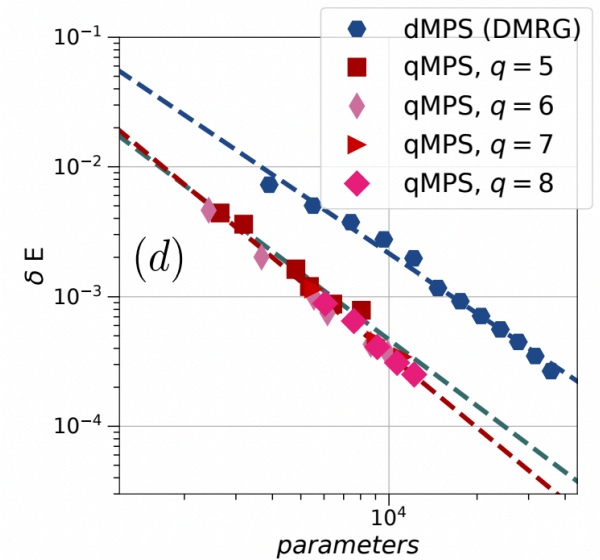
# Variational qMPS: Systematics



Heisenberg spin-chain



Fermi-Hubbard Chain  
(partial filling = metal)



## Takeaways:

- Local circuits can achieve similar performance to DMRG
- Extrapolation: qMPS can be more expressive than dense-MPS

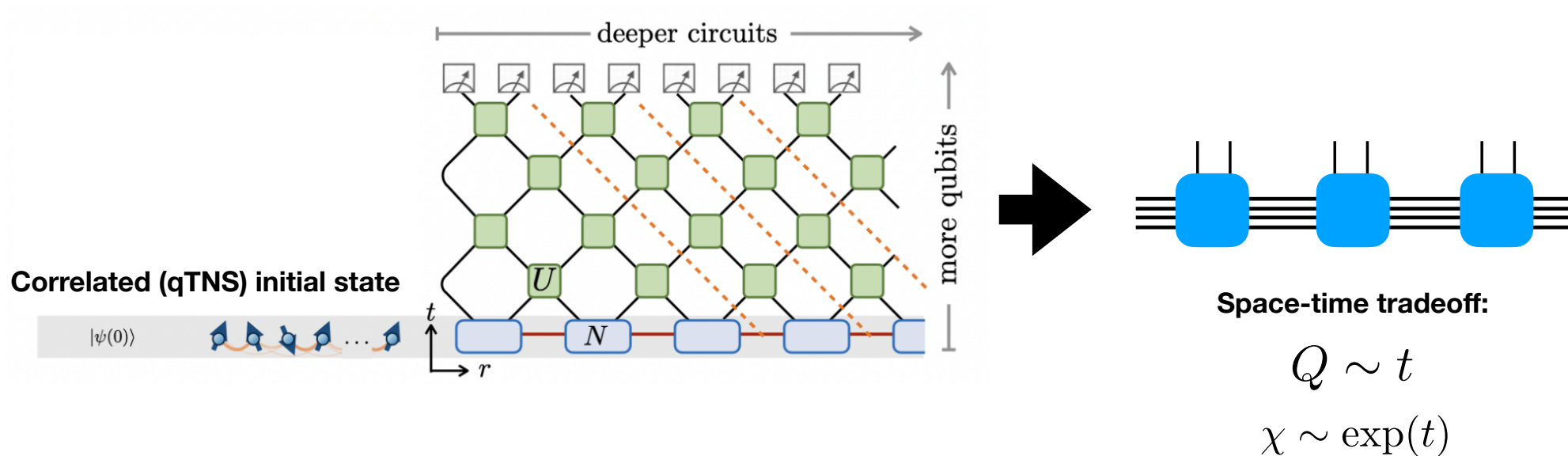
*Haghenas, Gray, ACP, G Chan arXiv:2107.01307*



# Non-equilibrium dynamics

Transport, Scattering, Thermalization, Localization, Quantum-Chaos,  
Scrambling, Dynamical critical phenomena & phases...

Classical simulations (TEBD):  $\chi \sim \exp(\text{time})$  **Hard even in 1d!!**



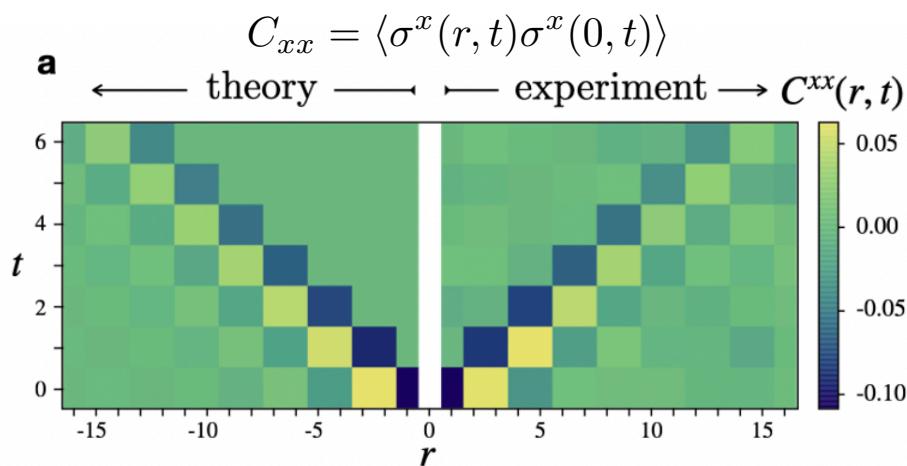
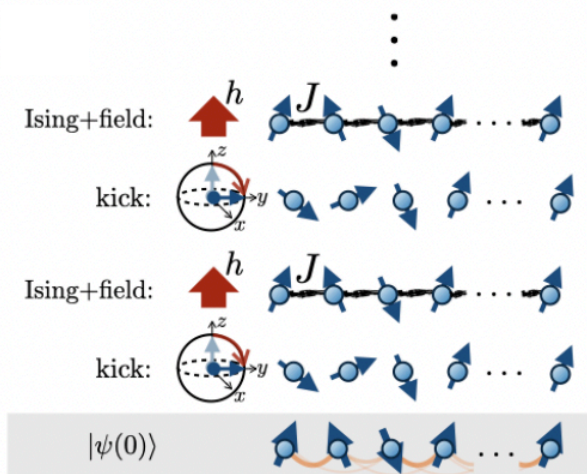
*Foss-Feig, Hayes, ... ACP PRR '21*

# Benchmark: Chaotic Floquet Dynamics

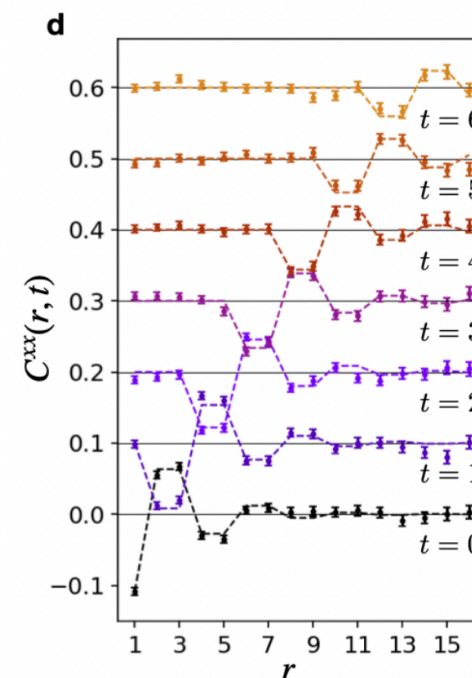
**Self-Dual Kicked Ising Model:** Non-integrable, chaotic, but solvable w/ certain MPS initial states  
*(quantum algorithm does not exploit solvability)*

$$H(t) = \sum_i \left( \frac{\pi}{4} \sigma_i^z \sigma_{i+1}^z + h \sigma_i^z + \frac{\pi}{4} \sum_{n \in \mathbb{Z}} \delta(t - n) \sigma_i^x \right)$$

*Bertini, Kos, Prosen*



- Quantitatively Accurate Simulations of  $L=32$  spins w/  $\leq 8$  qubits
- Essentially no post-processing or error-mitigation



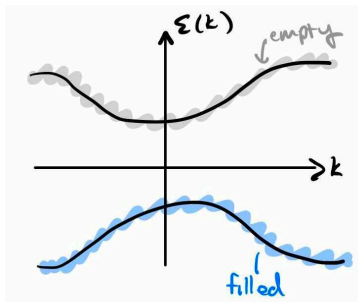
*Chertkov, ... Foss-Feig (Honeywell + ACP) arXiv: 2105.09324*

# Simulating Electrons

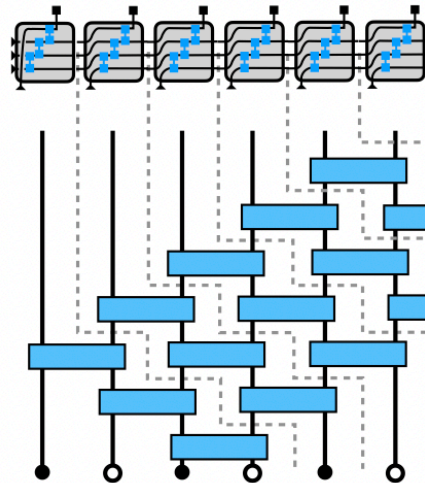
## Gaussian (Mean-field) States

$$G_{ij} = \langle c_i^\dagger c_j \rangle = W \begin{pmatrix} \mathbb{1} & 0 \\ 0 & 0 \end{pmatrix} W^\dagger$$

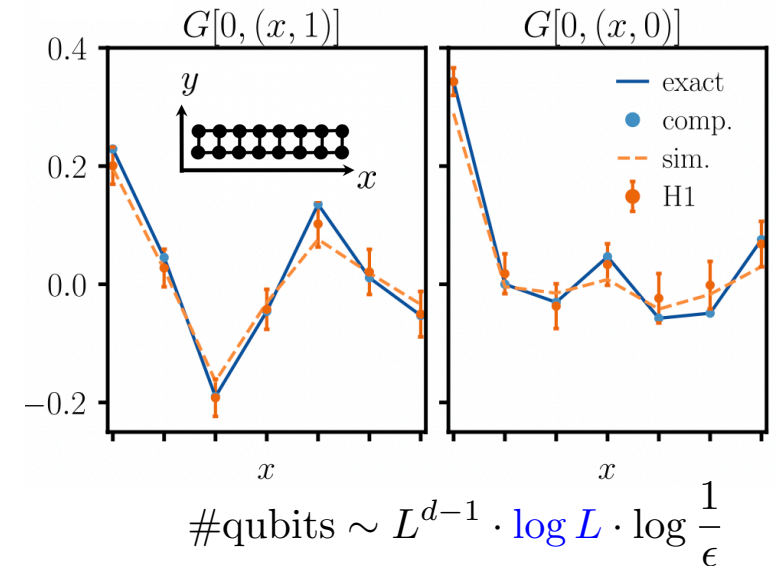
occupied  
empty



## GMPS Compression: *Fishman, White*



## Experimental demo



- Approximate localized basis for filled states
- Construct circuit that rotates into that basis
- Convert to qMPS

- Compare to Google's "Hartree-Fock demo"  $\#Q \sim L^{2d}$  *Arute et al. Science '20*
- + **adiabatic time evolution** => **efficient qMPS for large-class of states** (metals, insulators, superconductors, magnets, some spin-liquids, etc...)



# Outlook

- **Future directions:**

- Other tensor network architectures (2d, 3d, qMERA, etc...)
- Mixed/Thermal states, open-system dynamics (qMPDO)
- Alternative hardware platforms (circuit QED)



**Yuxuan Zhang**  
(UT Austin)



**Daoheng Niu**  
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**Reza Haghenas**  
(Caltech)



**Johannes Hauschild**  
(UC Berkeley)



**Sajant Anand**  
(UC Berkeley)



**Mike Foss-Feig**  
(Honeywell)



**Dave Hayes**  
(Honeywell)



**Mike Zaletel**  
(Berkeley)



**Garnet Chan**  
(Caltech)



**Shyam Shankar**  
(UT Austin)

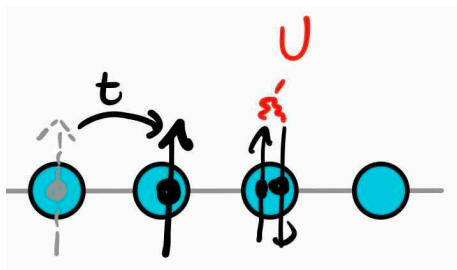


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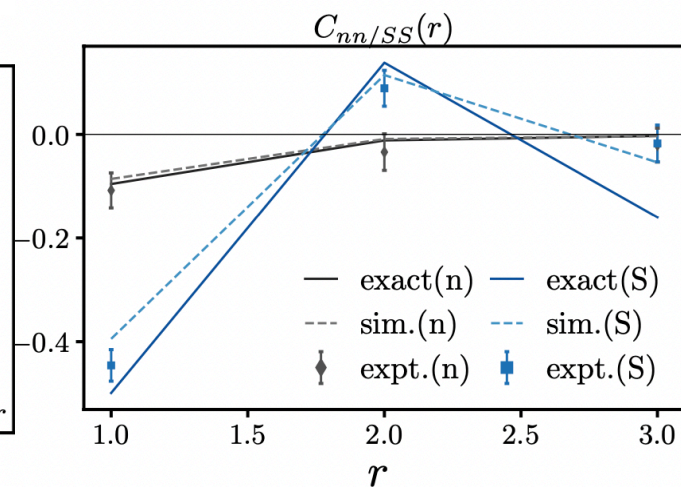
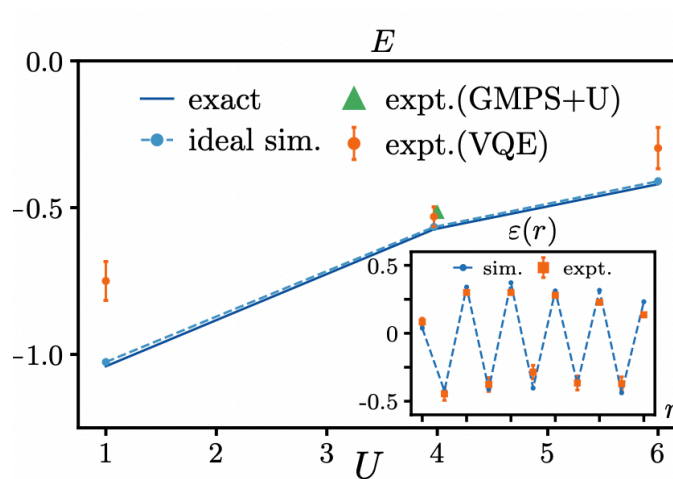
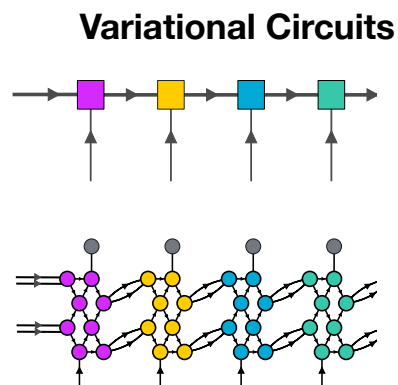
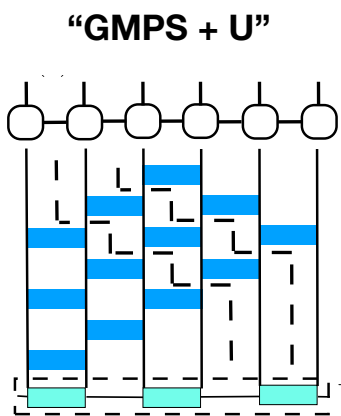


Alfred P. Sloan  
FOUNDATION

# Simulating Correlated Electrons



$$H_{\text{FH}} = - \sum_{x,\sigma} \left( c_{x+1,\sigma}^\dagger c_{x,\sigma} + \text{h.c.} + \mu n_x \right) + \sum_x \frac{U}{2} n_x (n_x - 1)$$



$$C_{SS}(r) = \langle S_0^z S_r^z \rangle - \langle S_0^z \rangle \langle S_r^z \rangle$$

$$C_{nn}(\mathbf{r}, \mathbf{r}') = \langle n_{\mathbf{r}} n_{\mathbf{r}'} \rangle - \langle n_{\mathbf{r}} \rangle \langle n_{\mathbf{r}'} \rangle$$

Niu, Heghanas, Zhang, Chan, ACP (to appear)