

# Single transverse spin asymmetry of very forward neutral pion

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# Outline

- *Introduction*
- *Born amplitudes*
- *Single transverse spin asymmetry*
- *Results*
- *Summary*

# Introduction

# Introduction

- Single transverse spin asymmetry(SSAs) is one of the interests in high energy reactions, but it has not been well understood yet.
- The significant  $A_N$  for pion production was reported for the first time in 1976 [1].
- The SSAs for various **very**-forward productions( $\pi, n, \gamma$ ) was measured in 2007 [2] and can not be explained by pQCD calculations(TMD, twist-3, etc).
- Sizable  $p_T$  and  $x_F$  distributions of  $A_N$  for very forward neutron with  $\sqrt{s} = 62, 210$  and  $500$  GeV was also measured by PHENIX Collaboration [3].
- Recent observations indicate that the diffractive process might account for a large fraction of  $A_N$  in the forward direction.

SSAs :

$$A_N = \frac{d\sigma^{\uparrow} - d\sigma^{\downarrow}}{d\sigma^{\uparrow} + d\sigma^{\downarrow}}$$

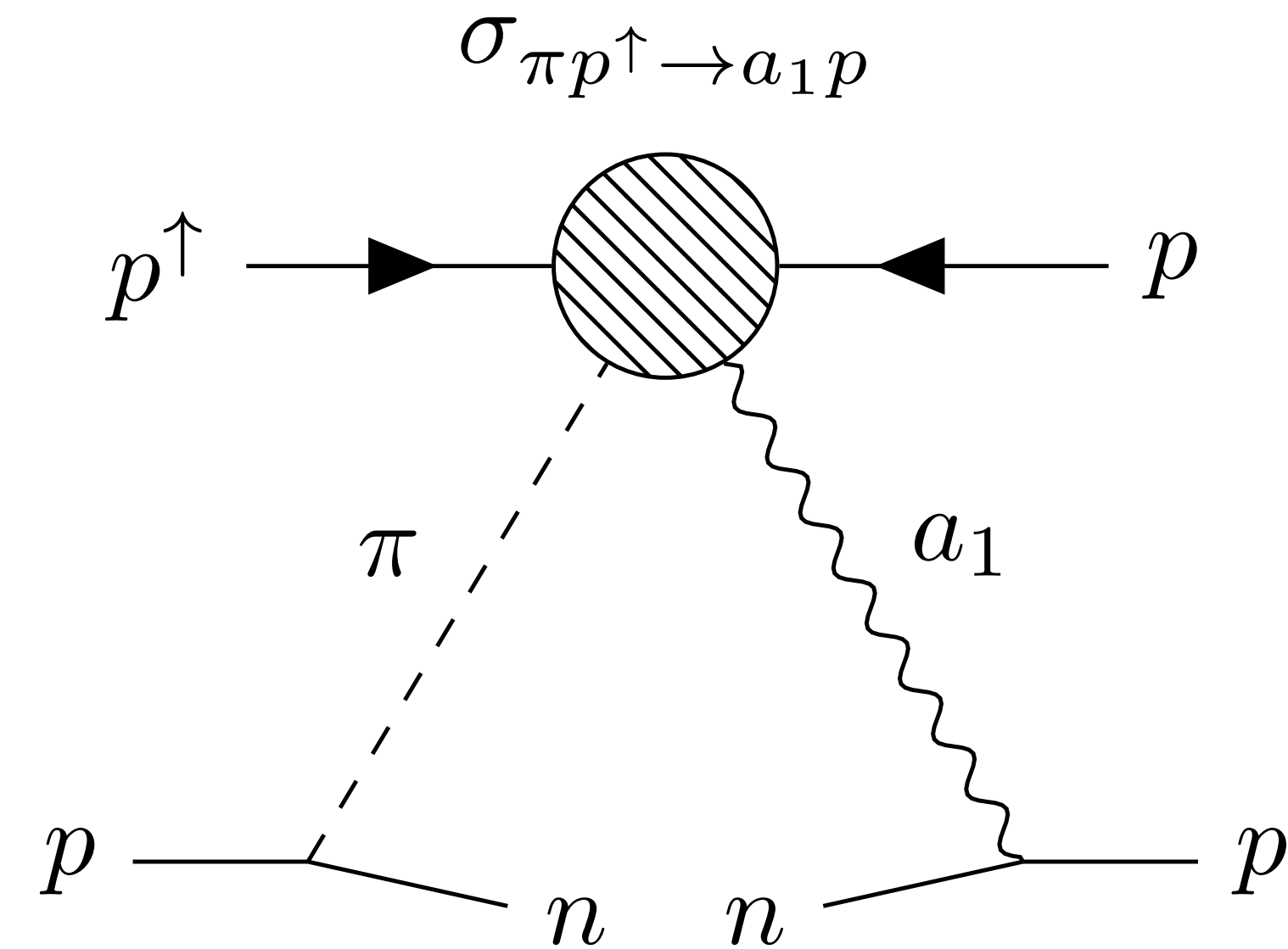
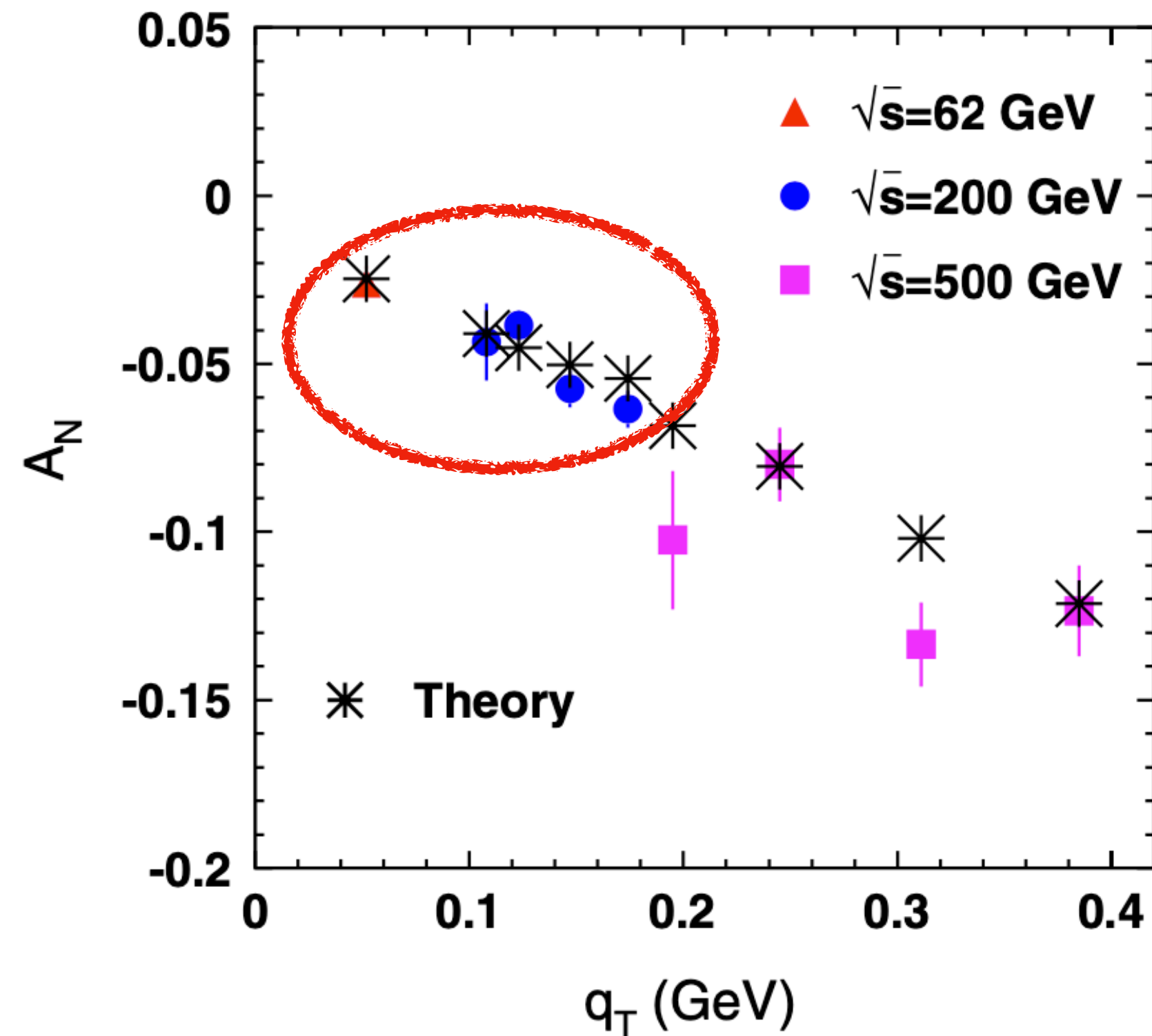
Spin direction of the polarized proton beam

[1] R. Klem et al. Phys. Rev. Lett. 36 (1976) 929

[2] Y. Fukao et al., Phys. Lett. B650 (2007) 325

[3] K. Tanida et al.(PHENIX Collaboration), J.Phys.Conf.Ser.295(2011)

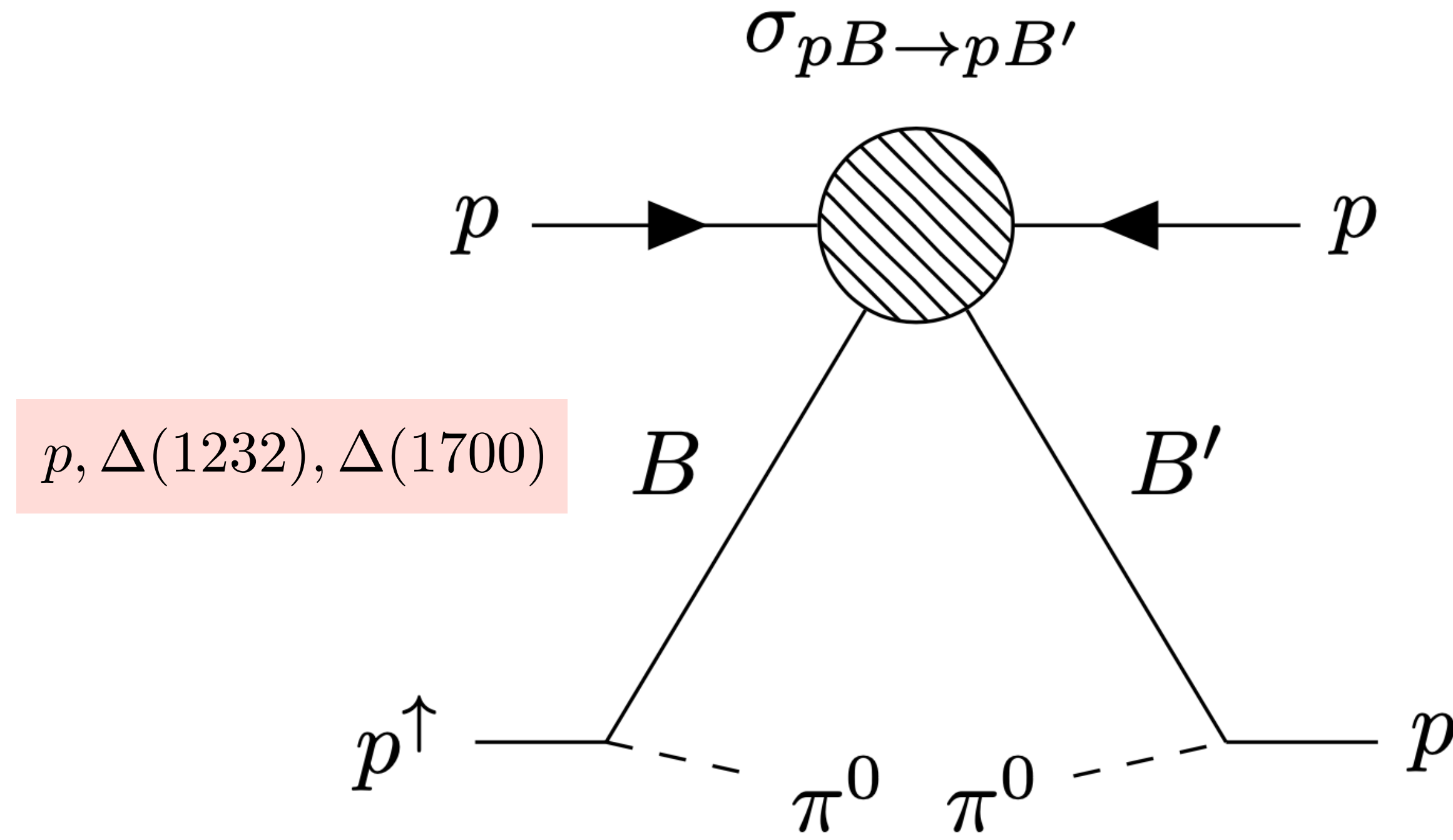
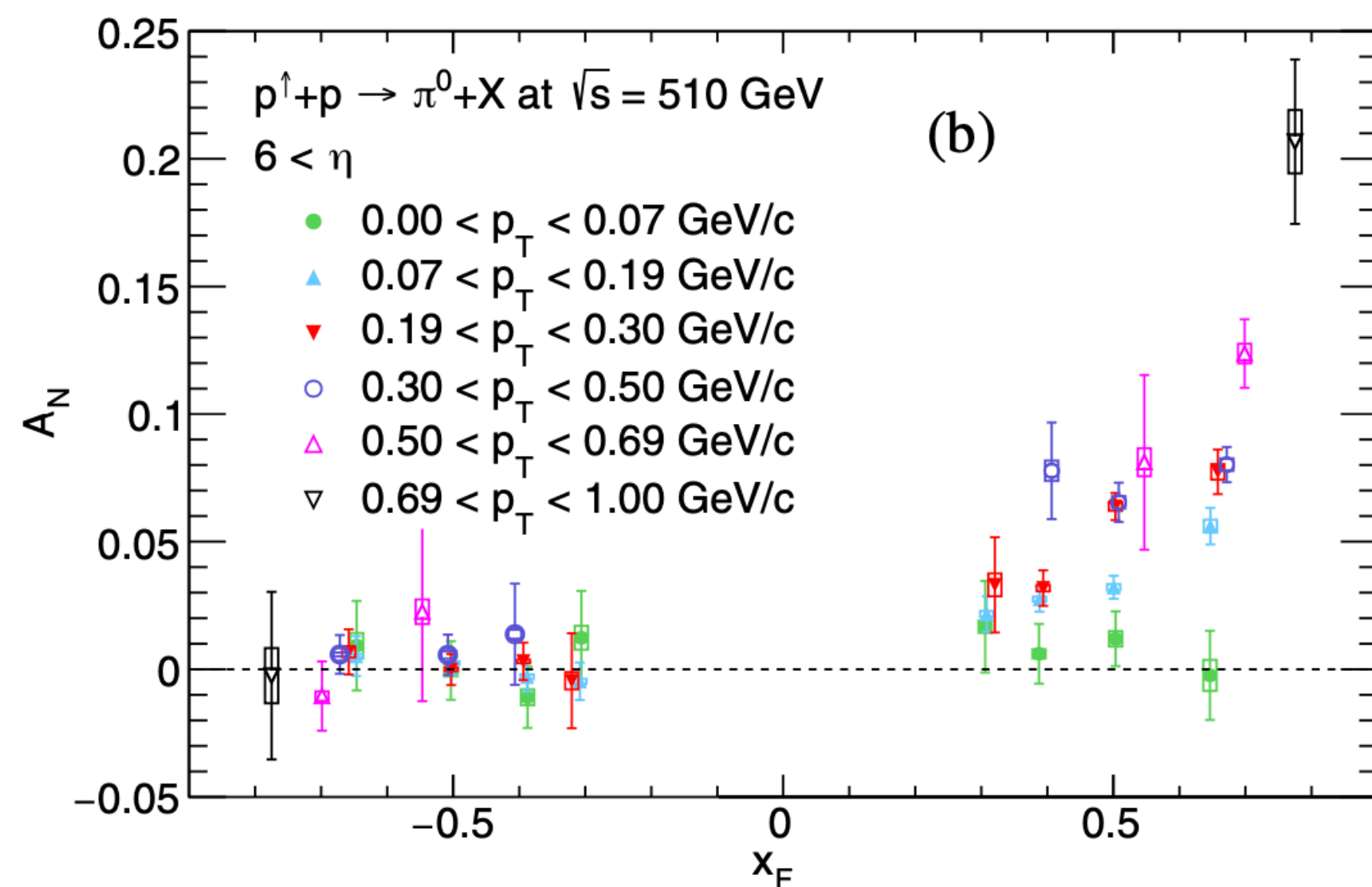
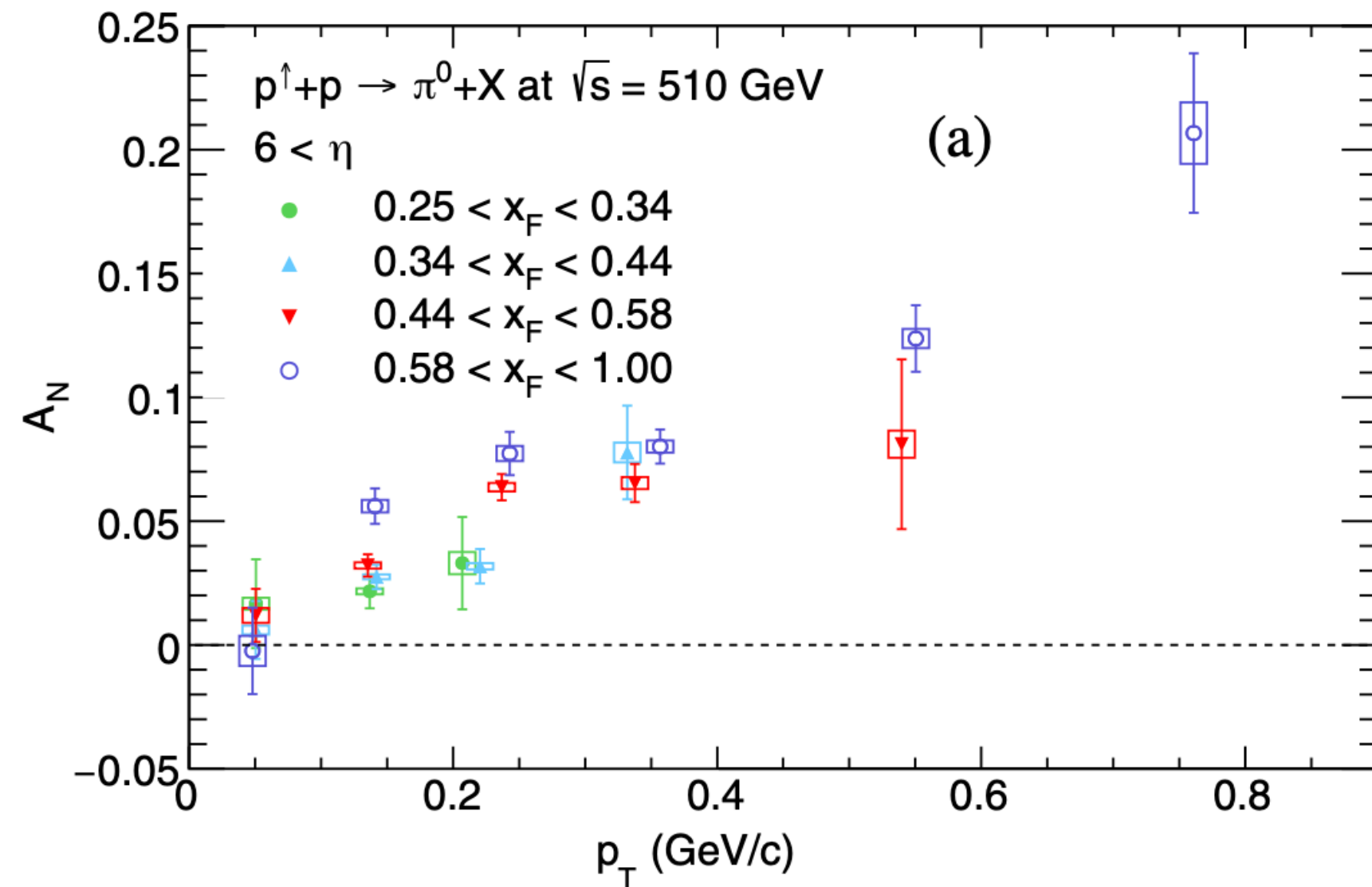
- A theoretical approach to the SSAs for the very forward neutron production was successfully achieved by OBE exchange based on Regge theory [4].



- The contribution of the  $\pi$ - $a_1$  interference for  $A_N$  matches the experiment results very well.



- Recently the SSAs of very forward neutral pion was measured in RHIC experiment [5].

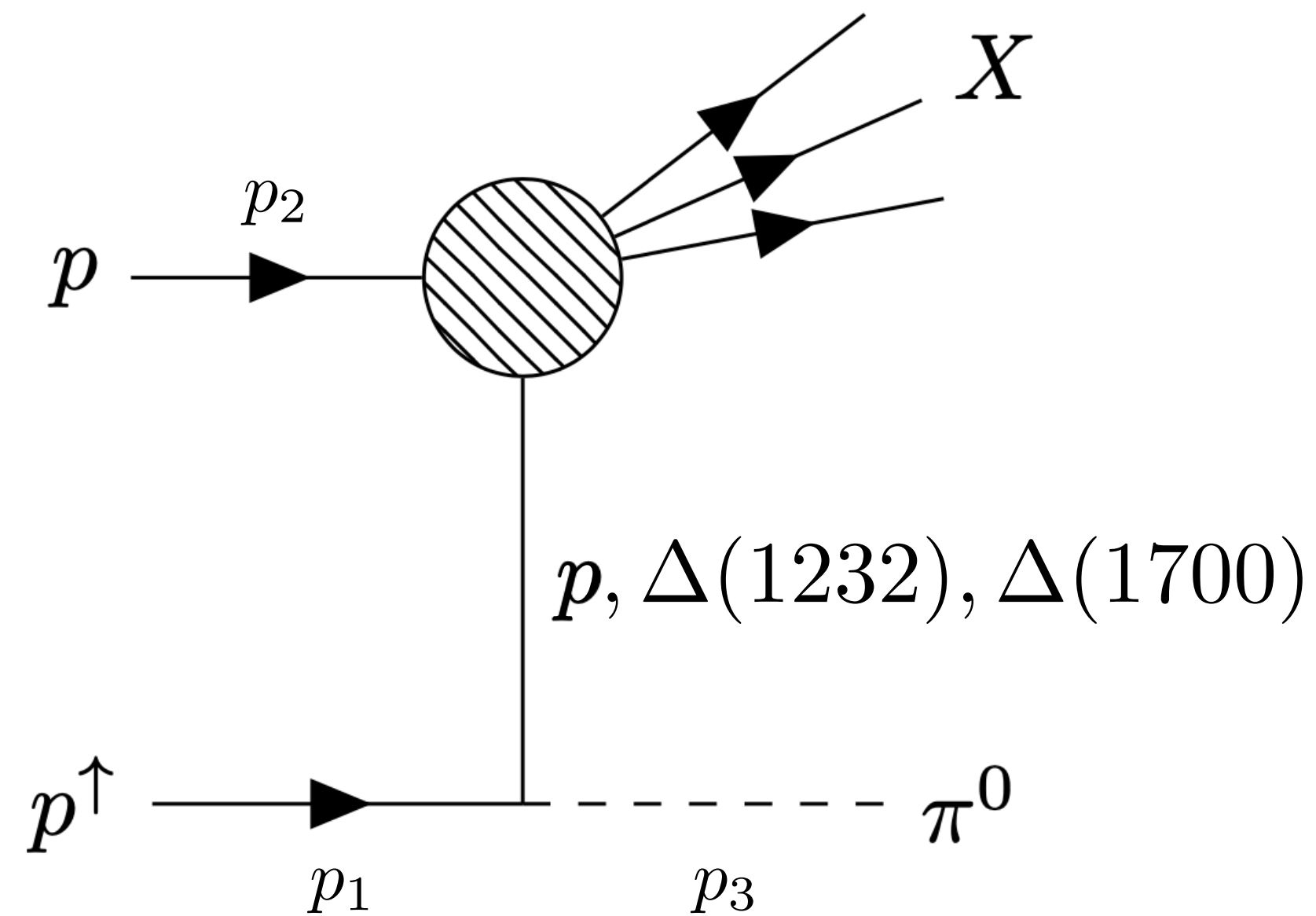


- We take into account the interferences between the proton and  $\Delta(1700)$ .
- The  $p$ - $\Delta(1700)$ - $\pi$  triple-Regge process plays an essential role for production of TSSA.

# Born amplitudes

# Born approximation

- Born diagram for the  $p + p^\uparrow \rightarrow \pi^0 + X$  :



- Kinematics for single diffractive(SD) process

$$p_1 = (E_1, 0, 0, p_z), \quad p_2 = (E_2, 0, 0, -p_z), \quad p_3 = (E_3, \mathbf{p}_T, p'_z)$$

$$M_X^2 \equiv (p_1 + p_2 - p_3)^2$$

$$x_F \simeq 1 - \frac{M_X^2}{s}, \quad t \simeq (1 - x_F)m_N^2 - \frac{\mathbf{p}_T^2}{x_F}$$

The SD processes can be described in terms of  $s$ ,  $x_F$ , and  $\mathbf{p}_T^2$ .

- Effective Lagrangians

$$\mathcal{L}_{NN\pi} = -g_{\pi NN} \bar{\psi} \gamma_\mu \gamma_5 \boldsymbol{\tau} \cdot \partial^\mu \boldsymbol{\pi} \psi$$

$$\mathcal{L}_{N\Delta^*\pi} = ig_{\pi N\Delta^*} \bar{\psi}_{\Delta^*}^\mu (g_{\mu\nu} + a\gamma_\mu \gamma_\nu) \gamma_5 \mathbf{T} \cdot \partial^\nu \boldsymbol{\pi} \psi$$

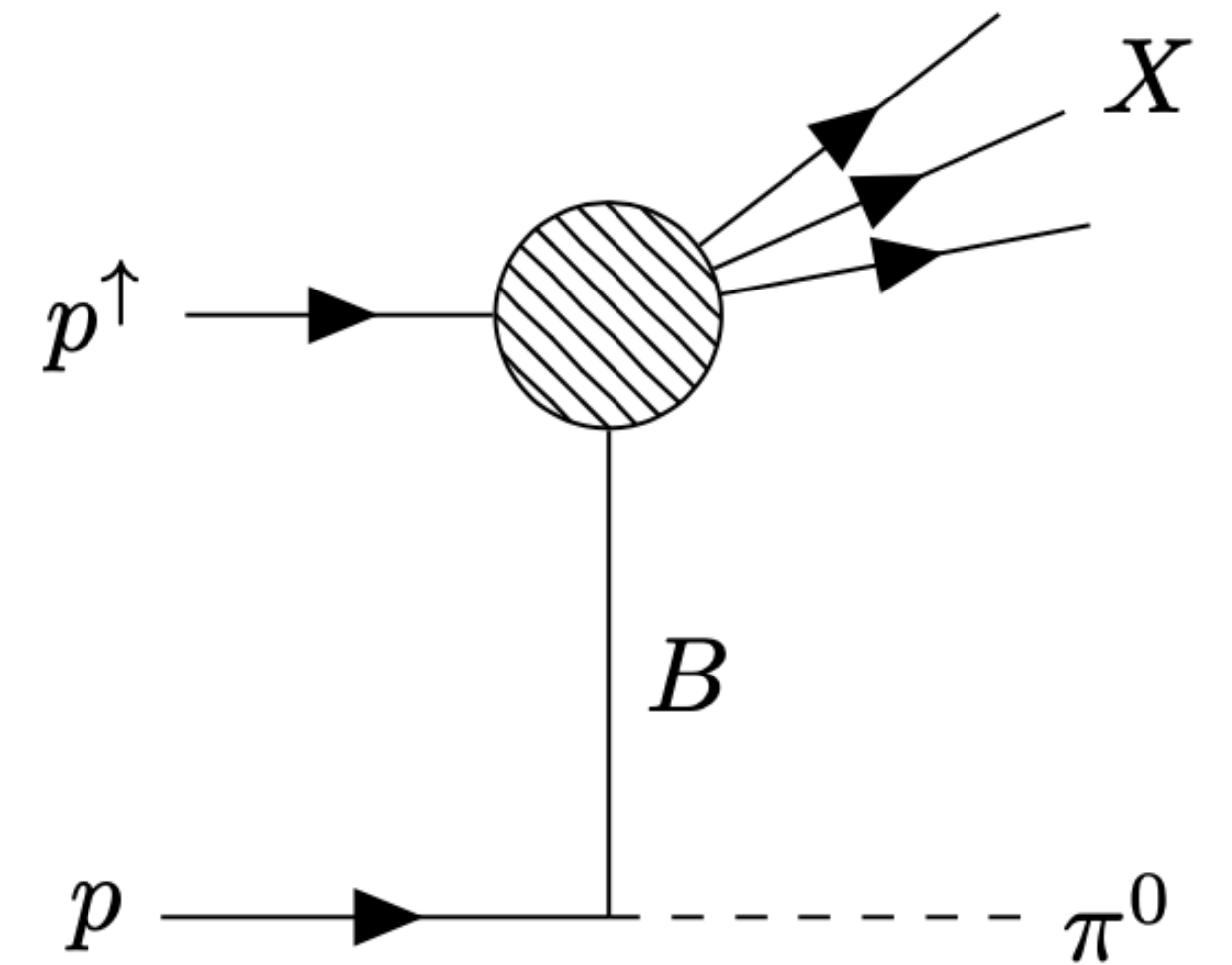


# Born amplitudes

$$A_{p \rightarrow \pi^0}^N(s, s') = g_{NN\pi} \beta_N(s, s'; p_T) \phi_N(p_T, x_F),$$

$$A_{p \rightarrow \pi^0}^\Delta(s, s') = g_{N\Delta\pi} \beta_\Delta(s, s'; p_T) \phi_\Delta(p_T, x_F),$$

$$A_{p \rightarrow \pi^0}^{\Delta^*}(s, s') = g_{N\Delta^*\pi} \beta_{\Delta^*}(s, s'; p_T) \phi_{\Delta^*}(p_T, x_F).$$



**(Reggeized) Baryon amplitude :**

$$\phi_B = \frac{\alpha'_B}{2} (1 \pm \exp\{-i\pi(\alpha_B(t) - J_B)\}) \Gamma(J_B - \alpha_B(t)) (1 - x_F)^{-\alpha_B(t) + J_B} A_{pB \rightarrow X} \quad [6]$$

provides the correct Regge poles.

$$\frac{s}{M_X^2} \approx (1 - x_F)^{-1}$$

**Regge trajectories**

$$\alpha_p(t) = -0.30 + 0.96t,$$

$$\alpha_\Delta(t) = 0.16 + 0.89t,$$

$$\alpha_{\Delta^*}(t) = -1.30 + 0.96t$$

# Single transverse spin asymmetry

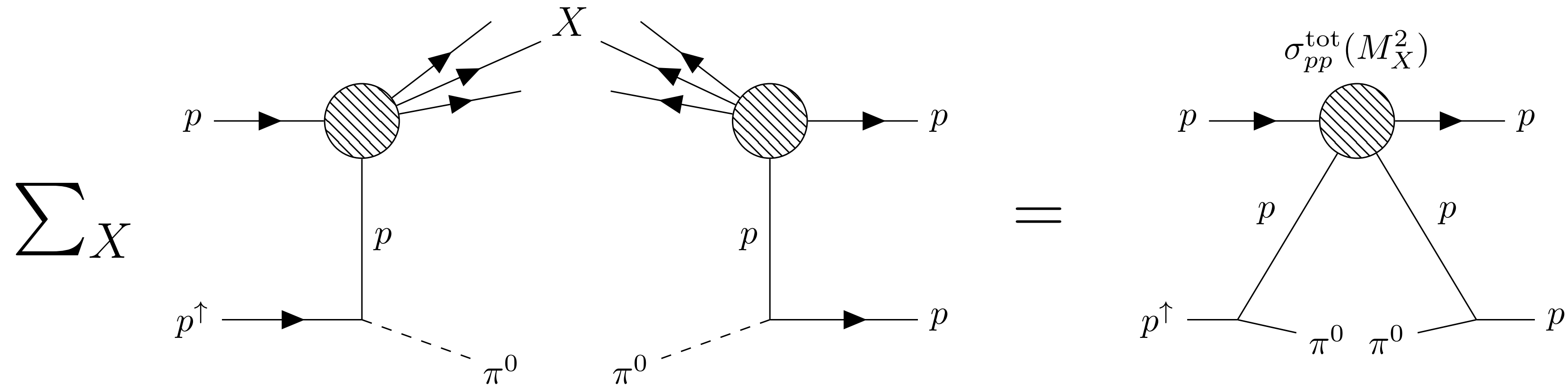
# Single transverse spin asymmetry(SSAs)

- SSAs

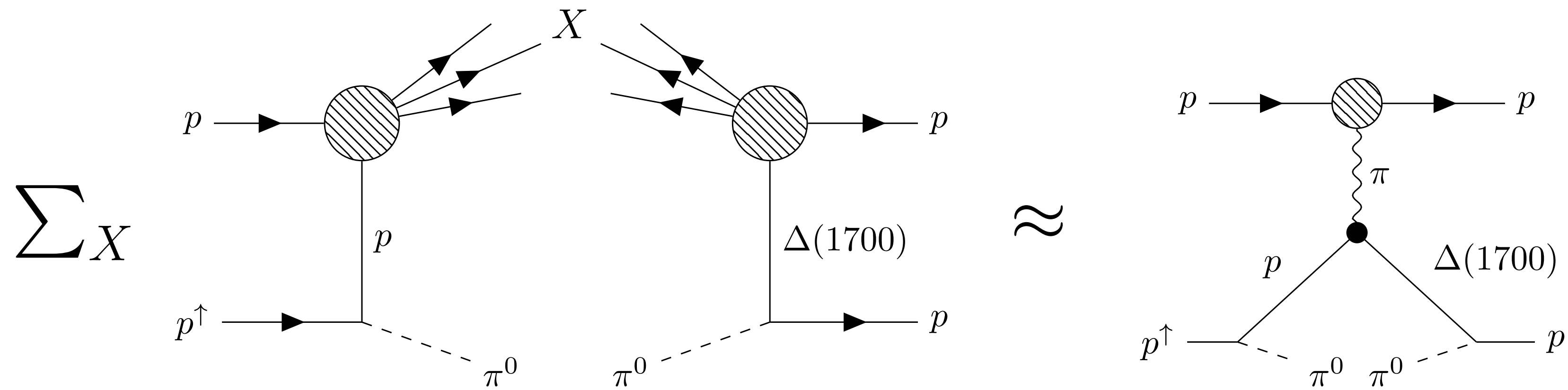
$$A_N = \frac{i(A_{p \rightarrow \pi^0}^{+*} A_{p \rightarrow \pi^0}^- - A_{p \rightarrow \pi^0}^{-*} A_{p \rightarrow \pi^0}^+)}{|A_{p \rightarrow \pi^0}^+|^2 + |A_{p \rightarrow \pi^0}^-|^2}$$

- In order to avoid producing nonzero  $A_N$  at  $p_T=0$ , we exclude the interference between **natural** and **unnatural** parity states: Natural =  $\{p, \Delta(1700)\}$
- The inclusive part of the interference terms can be approximated as the triple-Regge process.
- The inclusive part of the denominator is normalized in terms of the  $pp \rightarrow pp$  and  $p\Delta \rightarrow p\Delta$  differential cross sections.

Normalization of the  $pp \rightarrow X$  amplitude :  $\sum_X |A_{pp \rightarrow X}(M_X^2)|^2 = M_X^2 \sigma_{pp}^{\text{tot}}(M_X^2)$

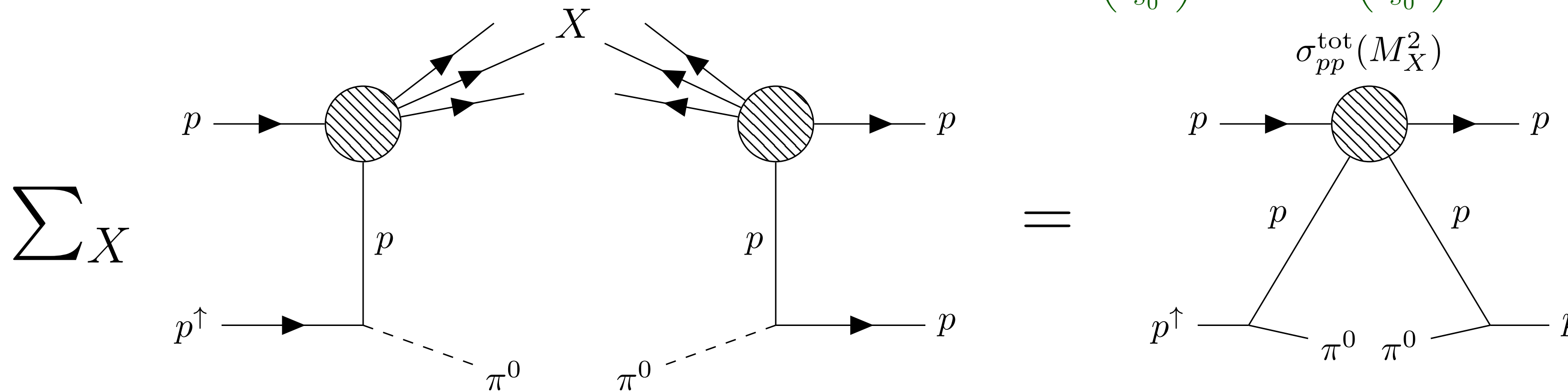


The triple-Regge diagram:



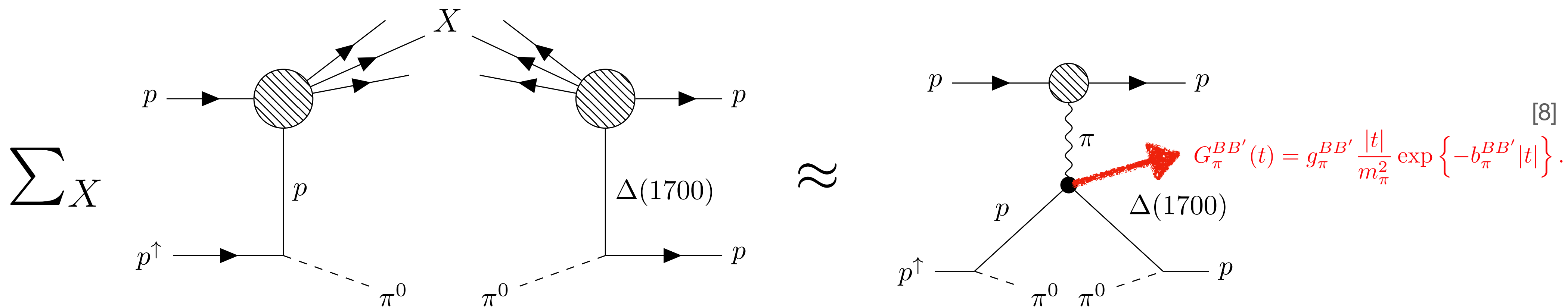
Normalization of the  $pp \rightarrow X$  amplitude :  $\sum_X |A_{pp \rightarrow X}(M_X^2)|^2 = M_X^2 \sigma_{pp}^{\text{tot}}(M_X^2)$  [7]

$$\sigma_{pp}^{\text{tot}}(M_X^2) = 42.6 \left(\frac{M_X^2}{s_0}\right)^{-0.46} - 33.4 \left(\frac{M_X^2}{s_0}\right)^{-0.545} + 35.5 + 0.307 \log^2 \left(\frac{M_X^2}{29.1}\right) \text{ [mb]}$$



The triple-Regge diagram:

$p\Delta$  scattering :  $C_{\Delta\Delta} \equiv \frac{\sigma_{p\Delta}^{\text{tot}}(M_X^2)}{\sigma_{pp}^{\text{tot}}(M_X^2)}$

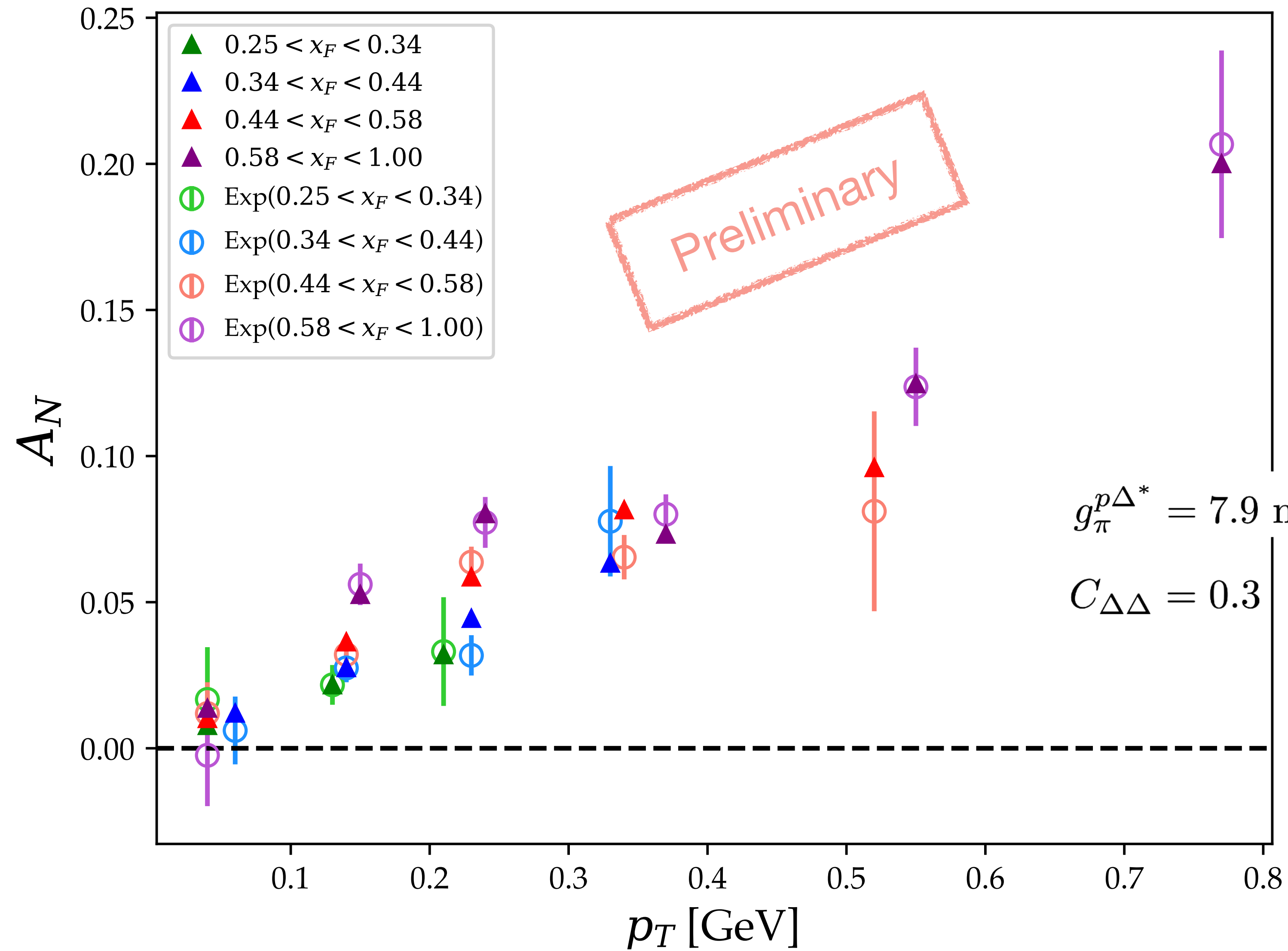




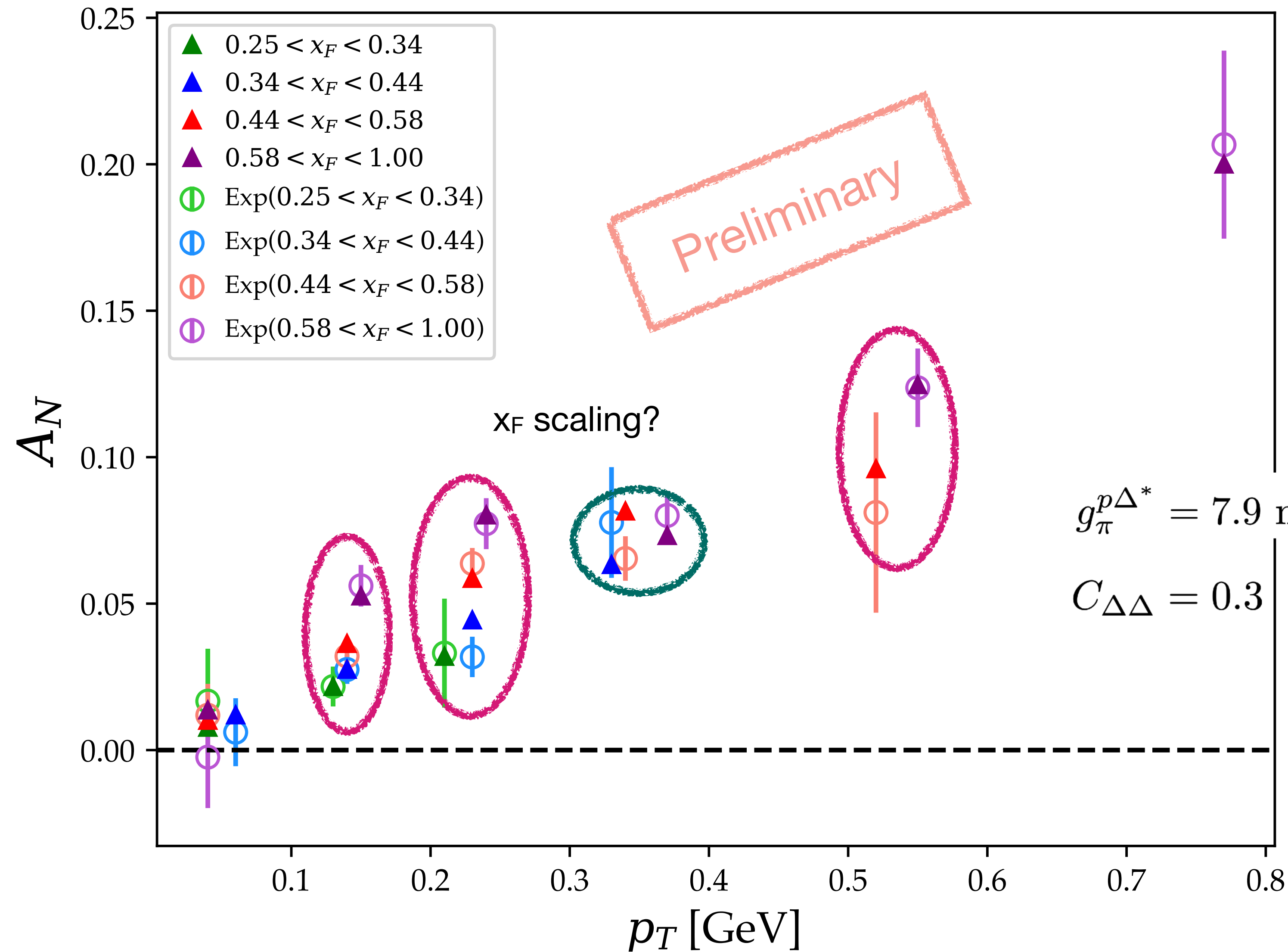


# Results

# SSAs vs $p_T$



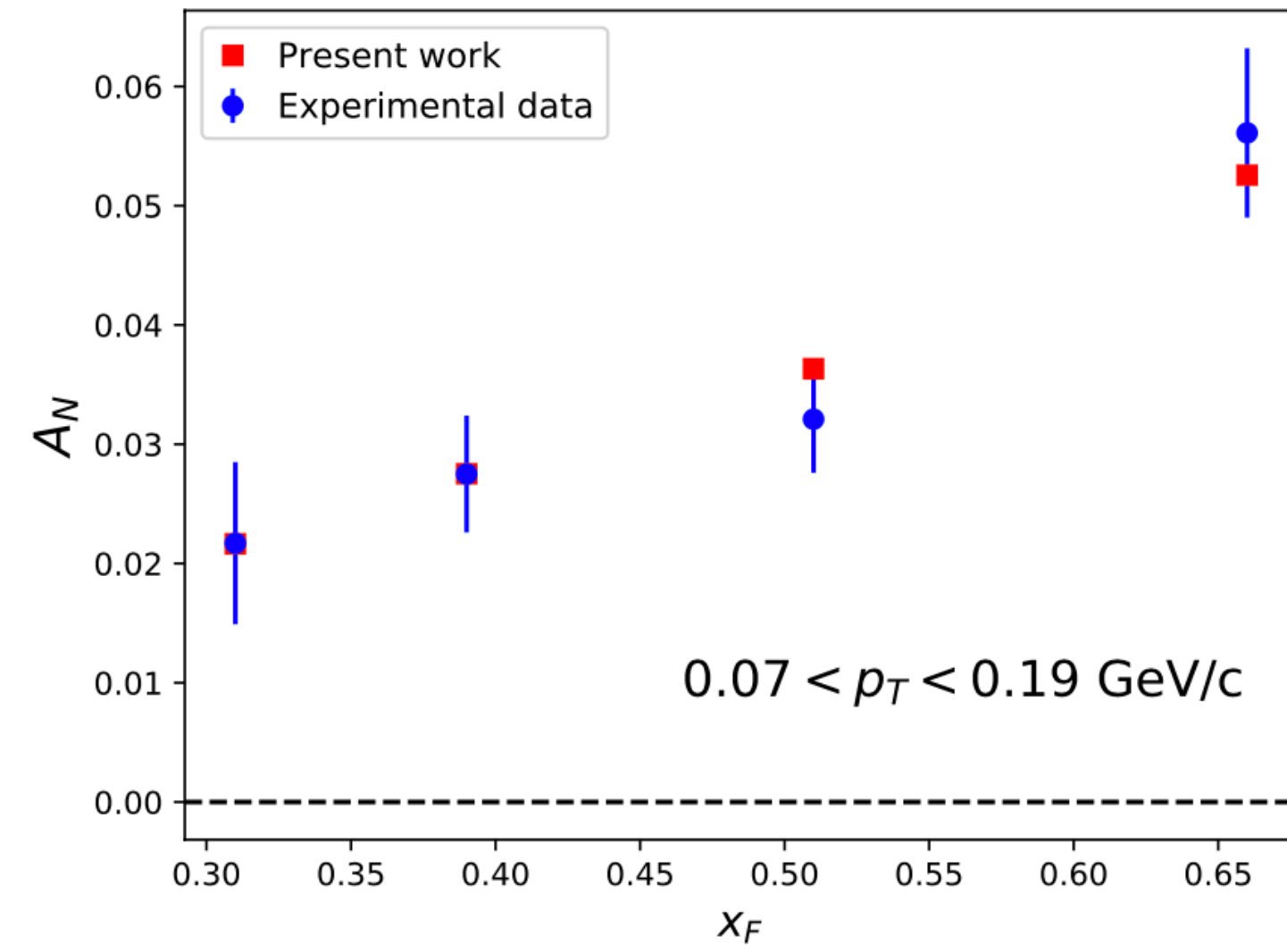
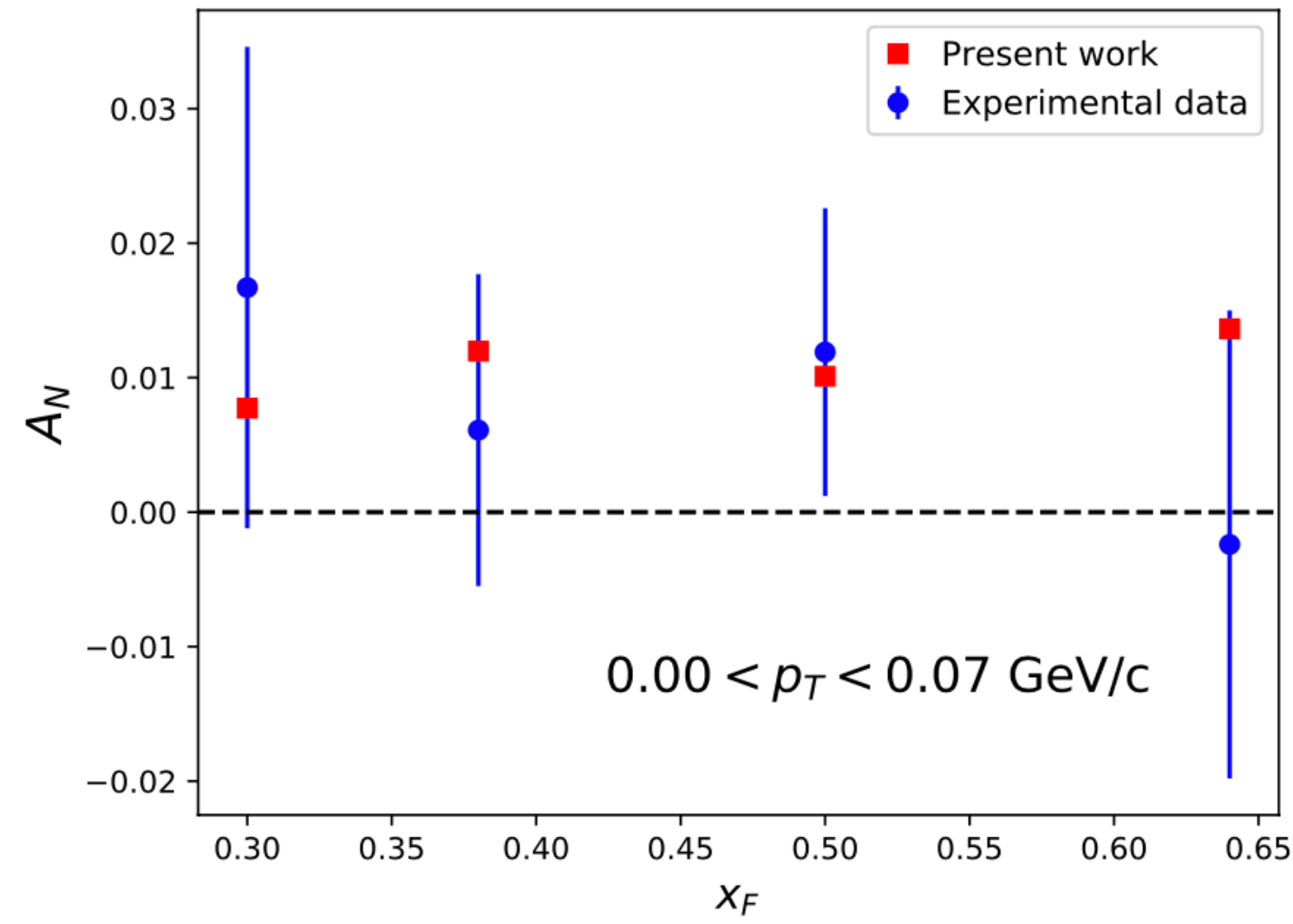
# SSAs vs $p_T$



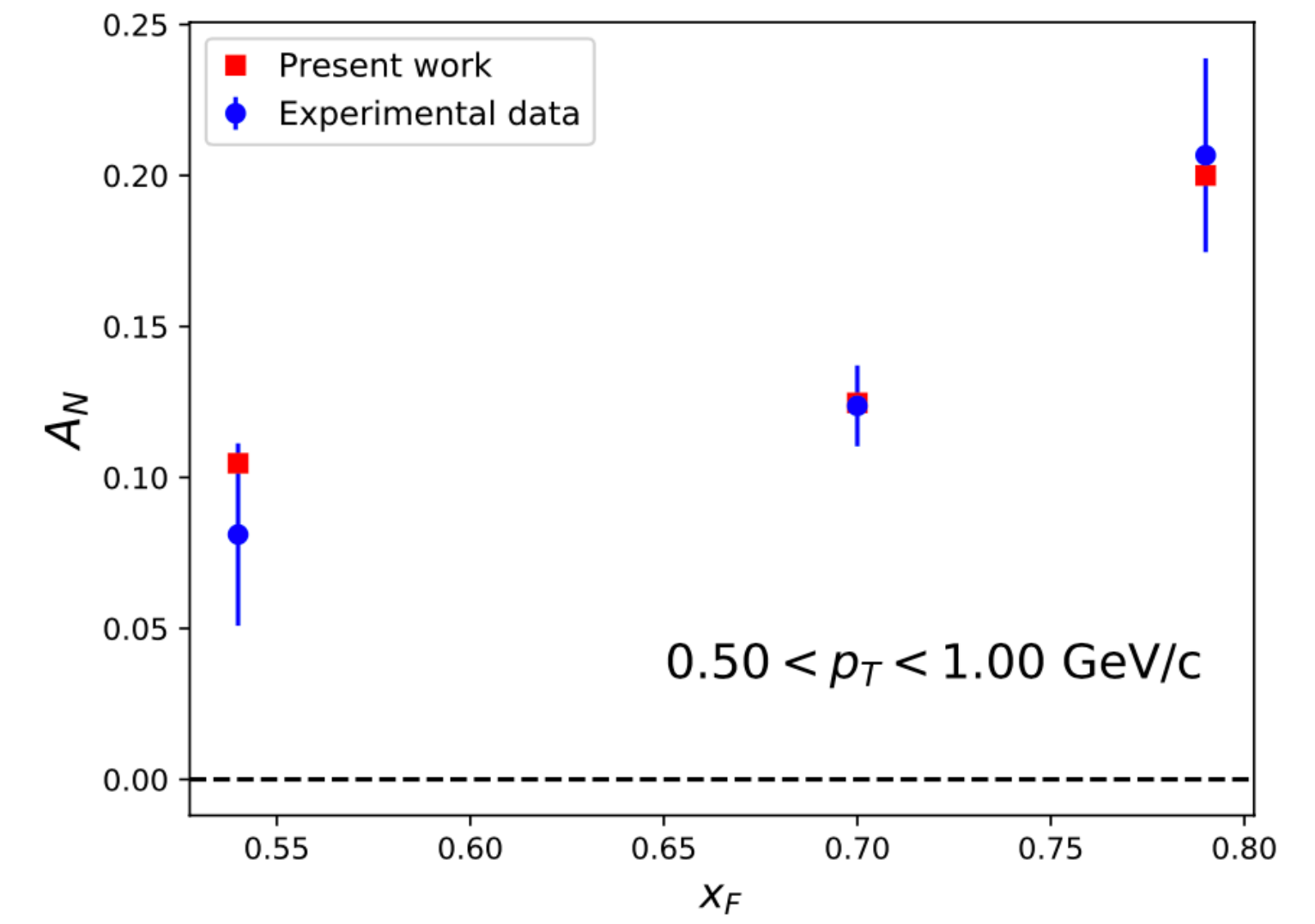
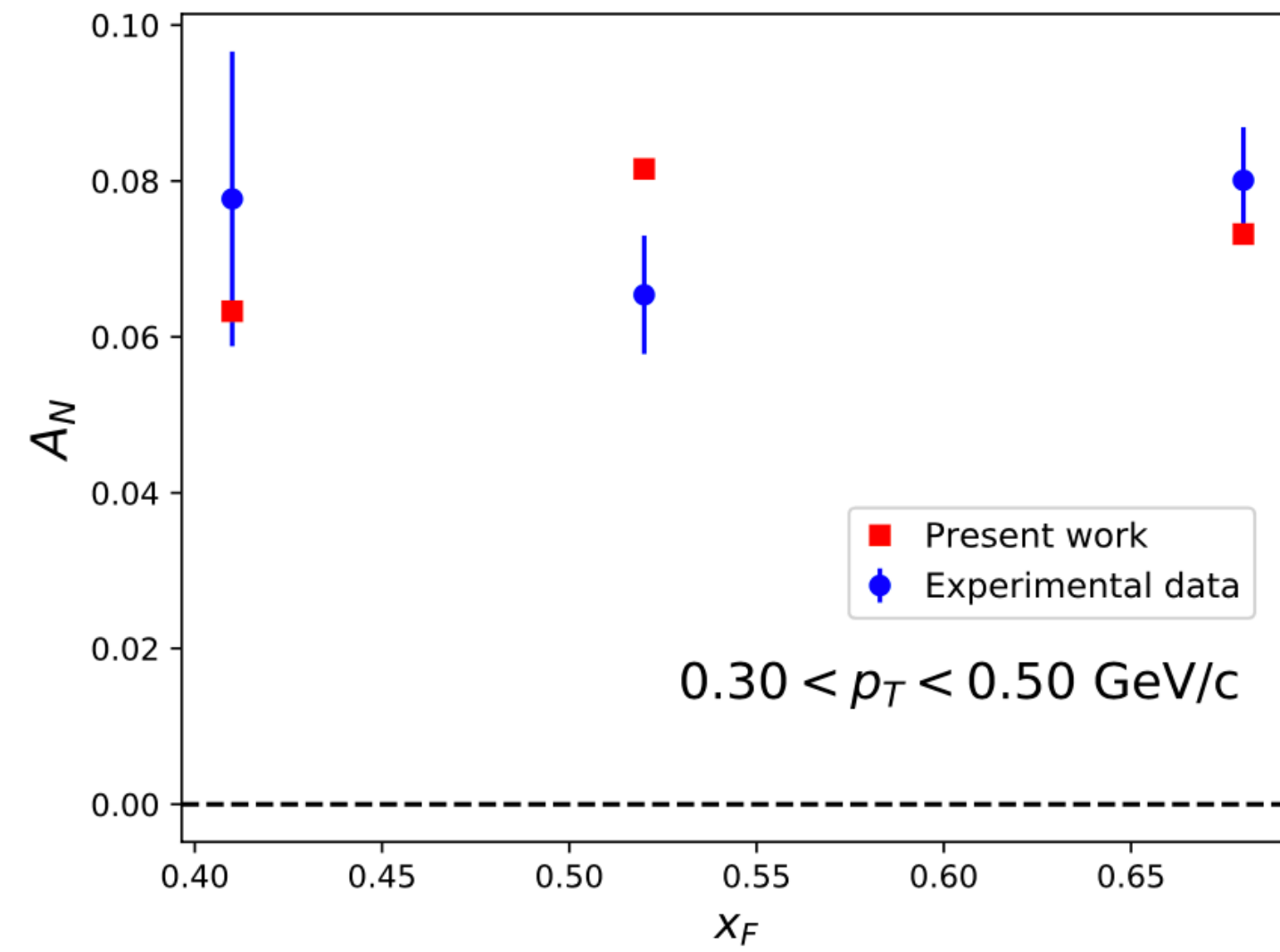
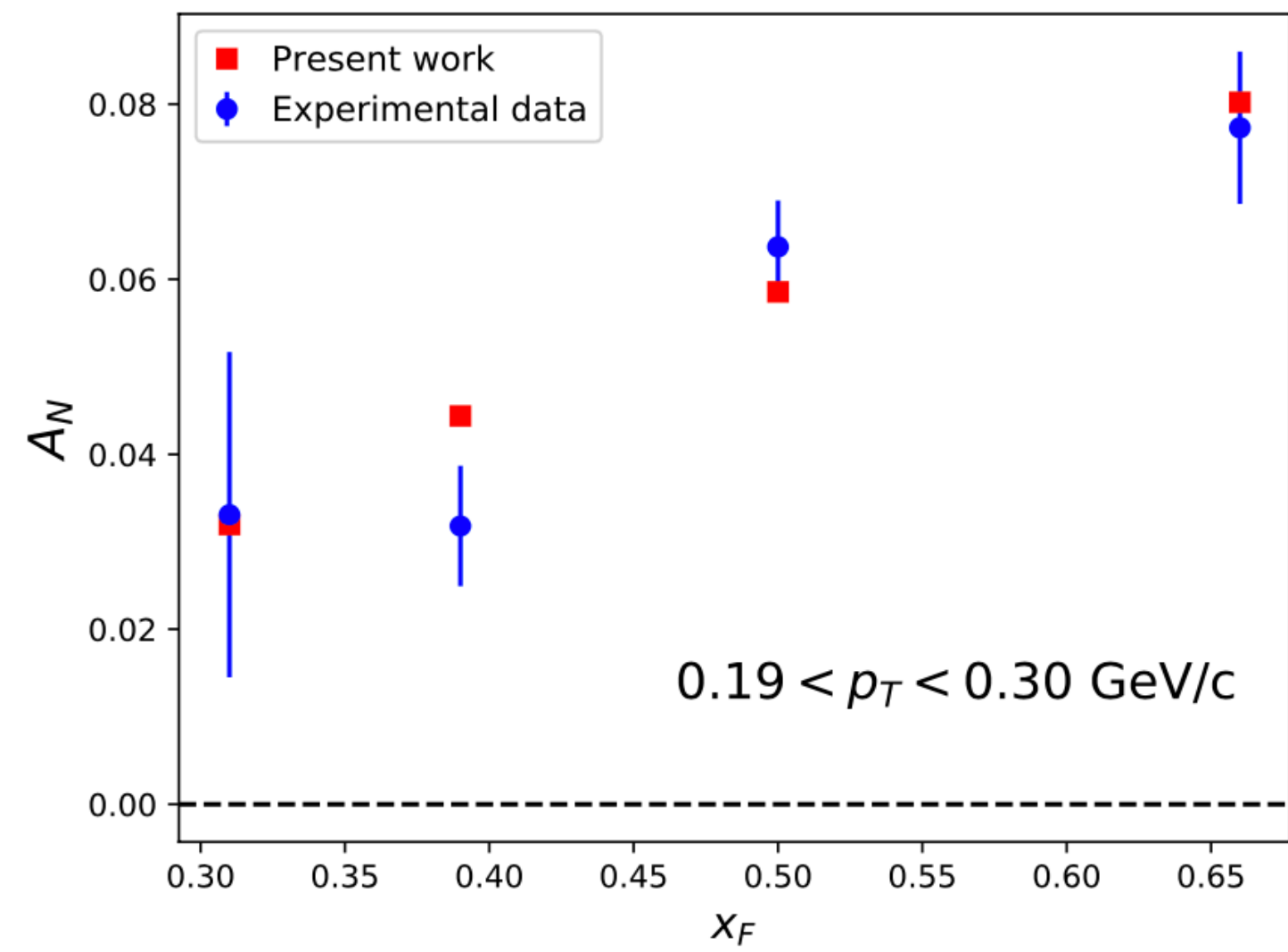
$$g_{\pi}^{p\Delta^*} = 7.9 \text{ mb TeV}^2, \quad b_{\pi}^{p\Delta^*} = 4 \text{ GeV}^{-2}$$

$$C_{\Delta\Delta} = 0.3$$

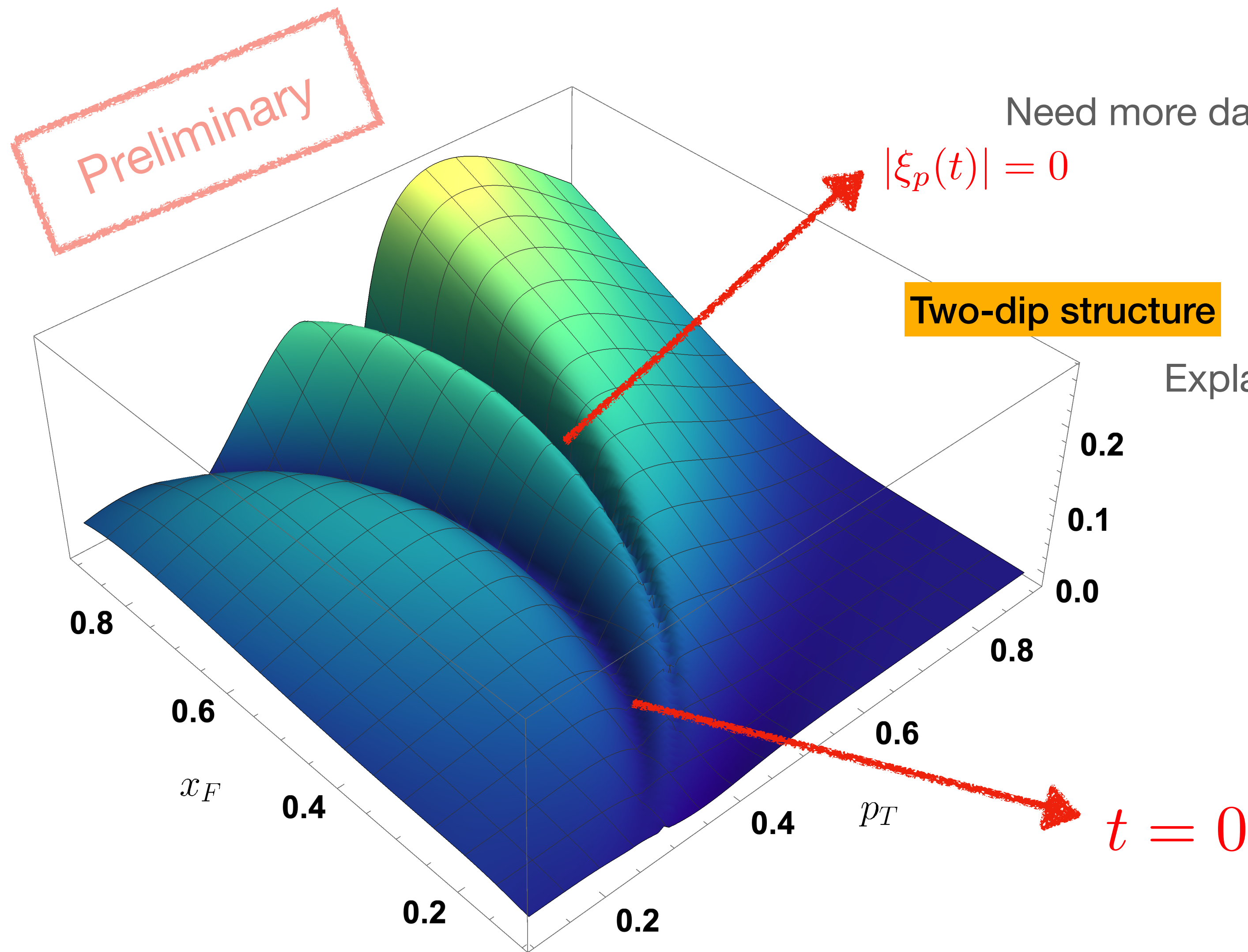
# SSAs vs $X_F$



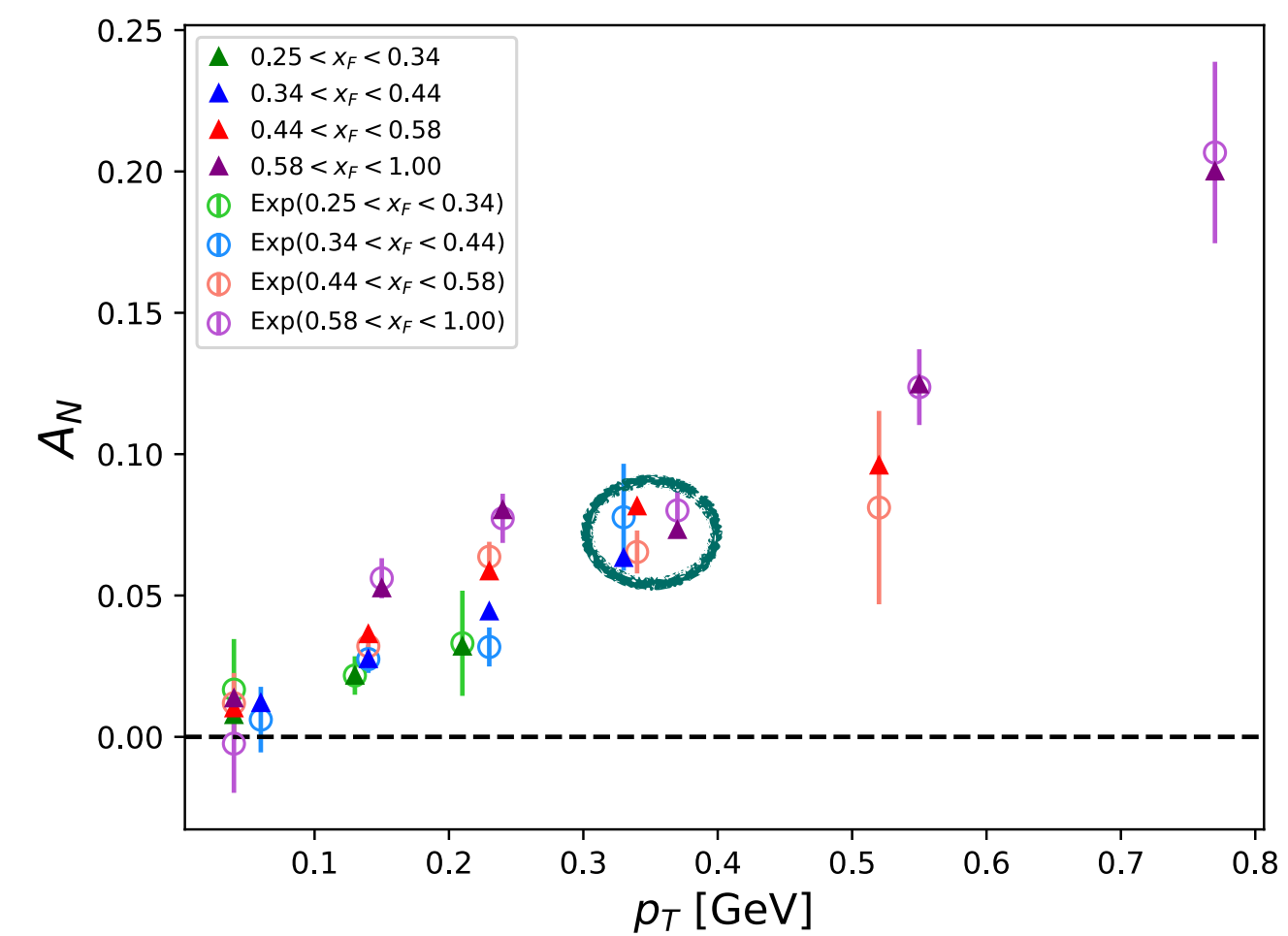
Preliminary



# 3D plot



Explain the vanishing of  $x_F$  scaling effect at a specific  $p_T$  range.



# Summary



# Summary

- We investigated the SSAs for forward neutral pion through the Reggeon exchange processes.
- The interference between Reggeon exchanges are approximated as the triple-Regge processes.
- The  $p\Delta$  total cross sections and the triple-Regge coupling are parametrized due to the lack of experimental data.
- Our results match the RHICf data of both transverse momentum and  $x_F$  distribution quite well.
- We found that the very forward neutral pion  $A_N$  is produced through the diffractive processes.

**Thank you**