

5次元Cassiniパラメータを用いた 核分裂の動力学的研究

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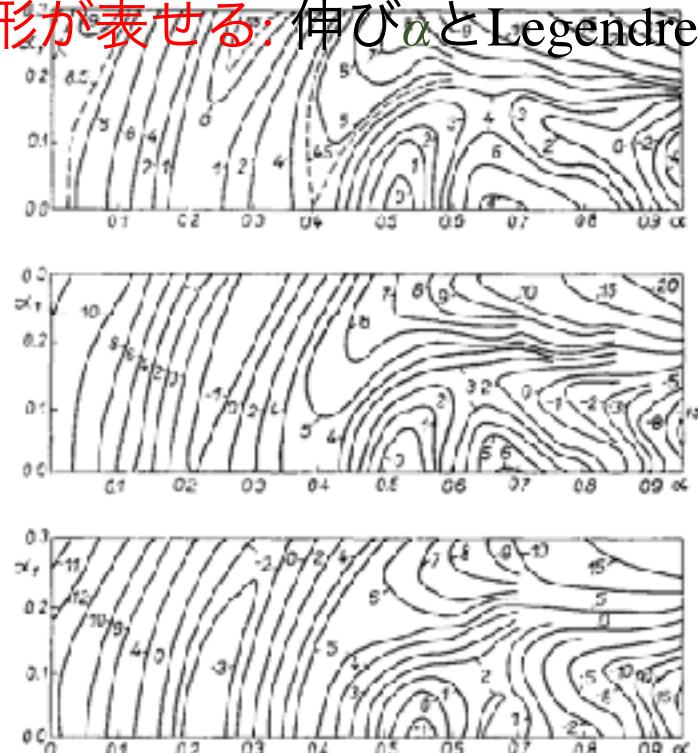
研究背景

5次元Cassiniパラメータを用いた核分裂のLangevin計算

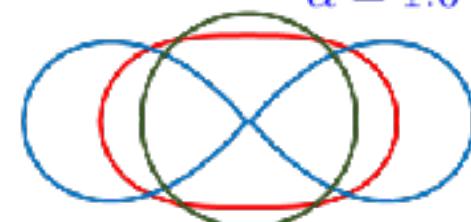
- ・超重核領域における変形核形状の描写
- ・多次元Langevin方程式への適用

Cassiniパラメータの特徴

- ・多様な変形が表せる: 伸び α と Legendre 展開係数 α_n



$$\alpha = 0.0 \quad \text{green circle}$$
$$\alpha = 0.5 \quad \text{red circle}$$
$$\alpha = 1.0 \quad \text{blue circle}$$



Cassini橙円($R = \text{一定}$)

$$R(x) = 1 + \sum_{n=1}^{\infty} \alpha_n P_n(x)$$

V. V. Pashkevich, Nucl. Phys. A 477, 1(1971)

Fig. 8. The same as in fig. 7 for ^{236}U at the top, ^{246}Pu in the middle and ^{252}Cf at the bottom.

研究背景

Cassini shape parameterization

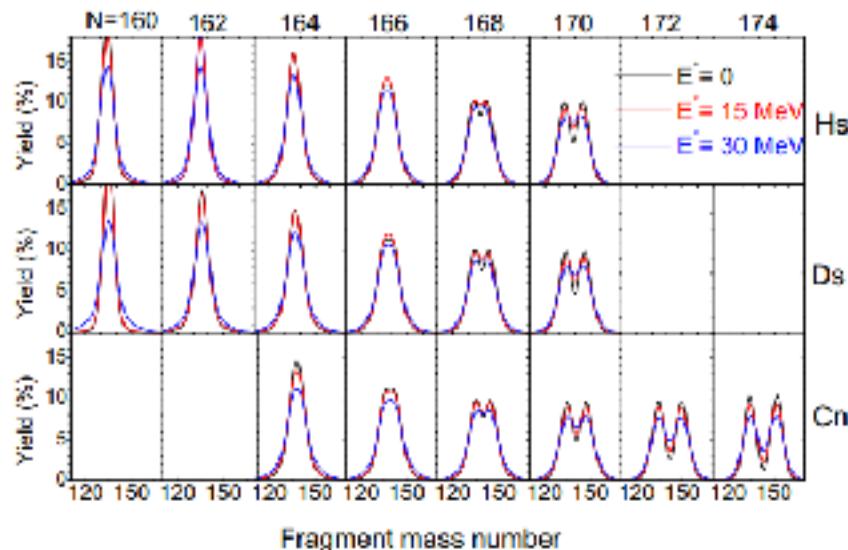


FIG. 3. The calculated fragment mass distributions for isotopes of Hs, Ds, and Cn for which spontaneous fission has been detected. Three values of the excitation energy E^* have been considered. $T_{\text{coll}} = 2 \text{ MeV}$, $E_d = 40 \text{ MeV}$.

$\alpha_1, \alpha_3, \alpha_4, \alpha_6$ が重要 (超重核)

→ 5次元 $\{\alpha, \alpha_1, \alpha_3, \alpha_4, \alpha_6\}$ Langevin方程式
系統的な核分裂計算

Static approach (not Langevin)

- pre-scission shape
- potential energy of deformation

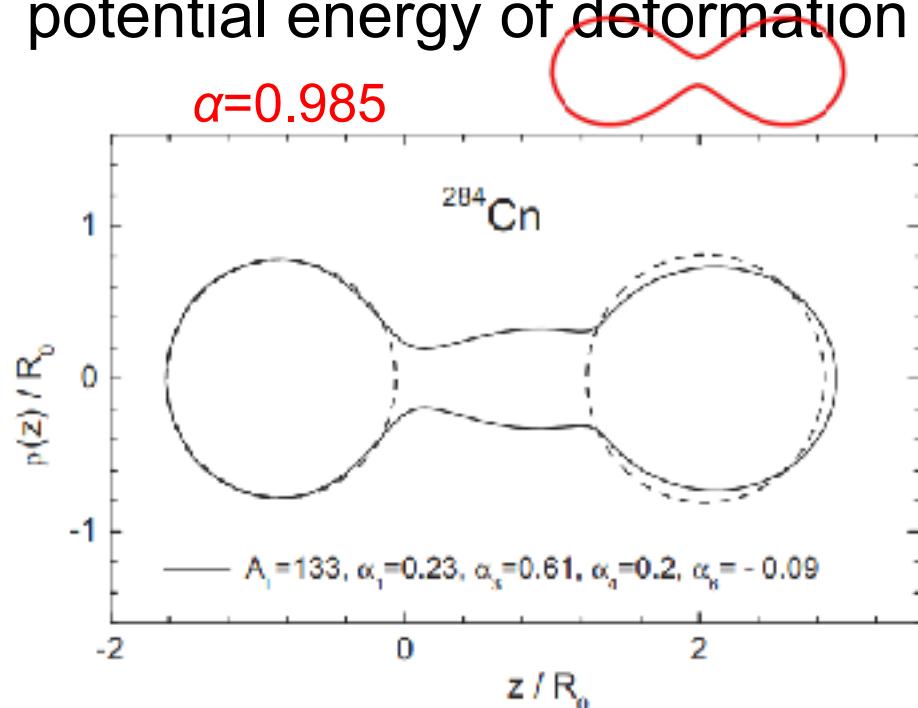
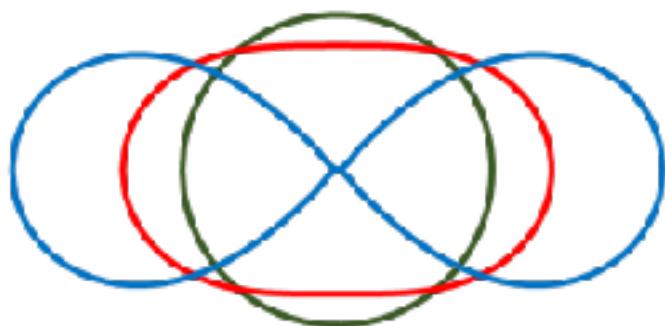


FIG. 2. The just-before-scission shape of ^{284}Cn for $A_L = 133$

N. Carjan, F. A. Ivanyuk, Yu. Ts. Oganessian, Phys. Rev. C **999**, 064606 (2019).

Cassini Shape Parameterization



Cassini ovals ($R = \text{const.}$)

$$\alpha = 0.0$$

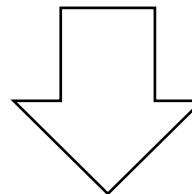
$$\alpha = 0.5$$

$$\alpha = 1.0$$

$$\alpha = \frac{z_R^2 + z_L^2 - 2\rho_{\text{neck}}^2}{z_R^2 + z_L^2 + 2\rho_{\text{neck}}^2}$$

α : elongation

Pashkevich, Nucl. Phys. A
169 (1971) 275.

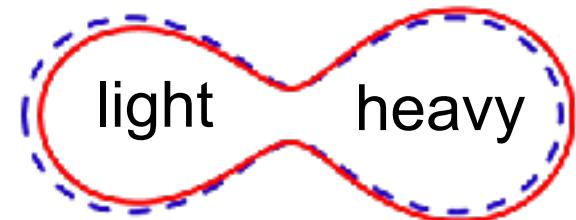


$$R(x) = 1 + \sum_{n=1}^{\infty} \alpha_n P_n(x)$$

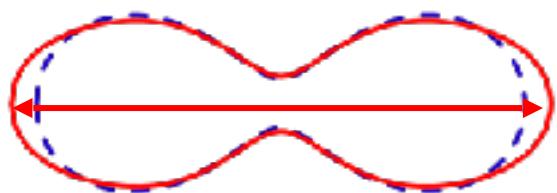
Legendre expansion of R
for additional parameters

n	odd: asymmetric deformation	α_1	α_3
	even: symmetric deformation	α_2	α_4

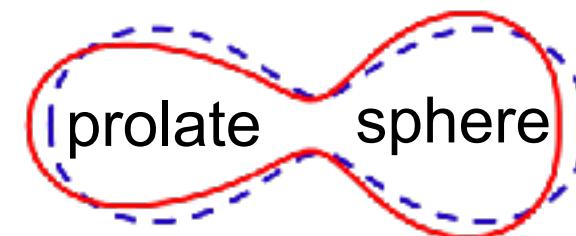
Cassini Shape Parameterization



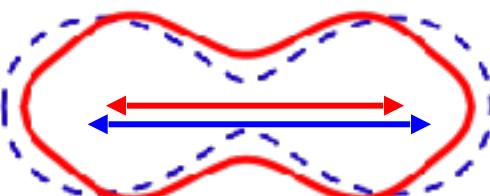
α_1 mass asymmetry
FMN



α_4 quadrupole deformation
of fragments
GS, 1stB, 2ndM



α_3 shape asymmetry
double magic ^{132}Sn



α_6 octupole deformation
of fragments
super-long (short) shape

necessary parameters

additional parameters

Langevin方程式

$$\{q_i\} = \{\alpha, \alpha_n\}$$

$$\frac{dq_i}{dt} = m_{ij}^{-1} p_j$$

$$\frac{dp_i}{dt} = -\gamma_{ij} m_{jk}^{-1} p_k - \frac{\partial V}{\partial q_i} - \frac{1}{2} \frac{\partial m_{jk}^{-1}}{\partial q_i} p_j p_k + g_{ij} R_j(t)$$

$$V(\mathbf{q}, T) = V_{\text{LDM}}(\mathbf{q}) + V_{\text{micro}}(\mathbf{q}) \exp(-aT^2/E_D) \quad E_D = 25 \text{ MeV}$$

変形ポテンシャル macroscopic-microscopic approach

V. M. Strutinsky, Nucl. Phys. A 95 (1967) 420

$m_{ij}(\mathbf{q})$ 慣性テンソル Werner-Wheeler method K. T. R. Davies et al., Phys. Rev. C 13, 2385 (1976)

$\gamma_{ij}(\mathbf{q})$ 摩擦テンソル Completed wall-and-window formula

J. Randrup and W. J. Swiatecki, Nucl. Phys. A 429, 105 (1984)

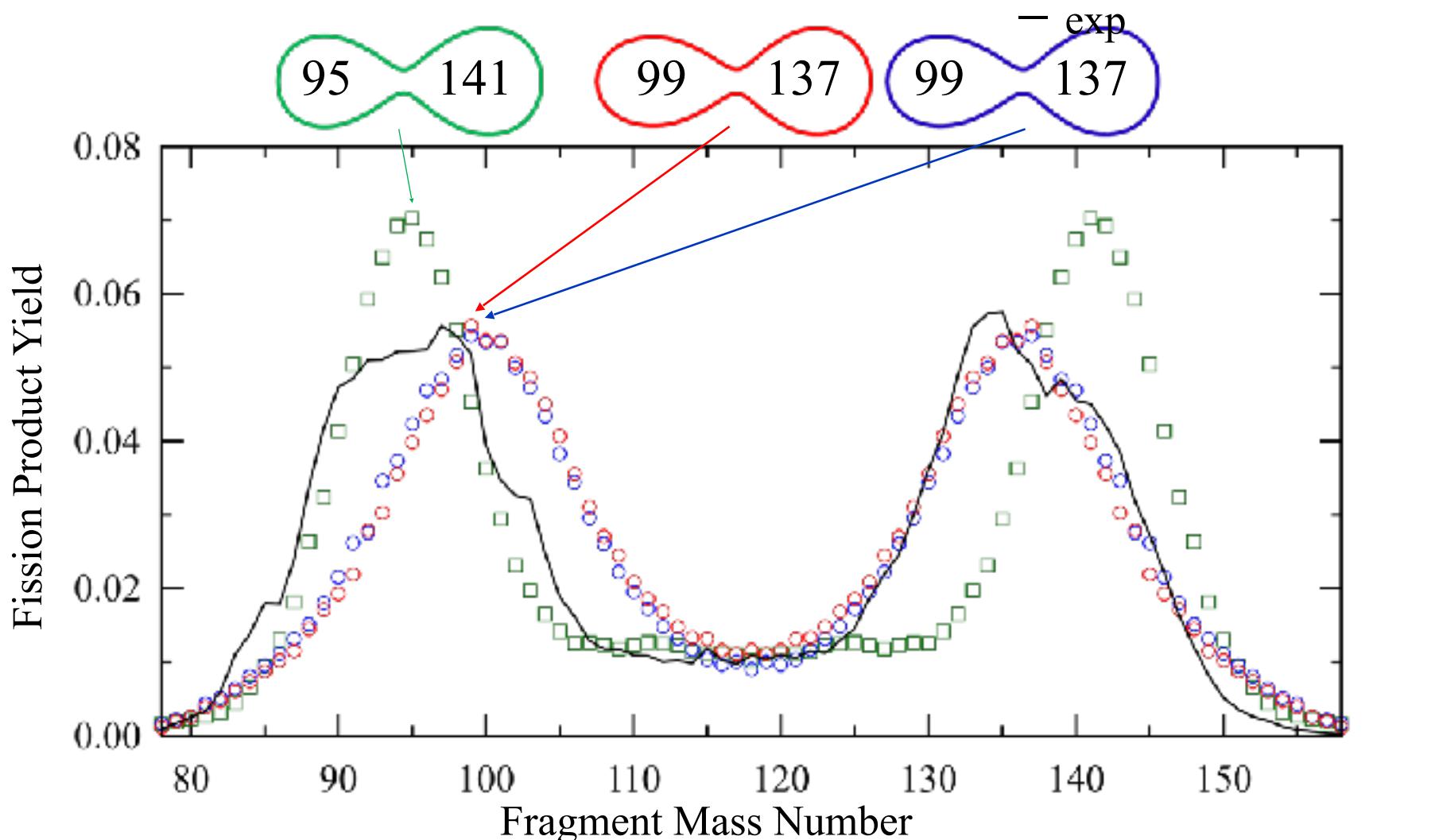
3D $\{\alpha, \alpha_1, \alpha_4\} \cdot 4D \begin{cases} \{\alpha, \alpha_1, \alpha_3, \alpha_4\} \\ \{\alpha, \alpha_1, \alpha_4, \alpha_6\} \end{cases} \cdot 5D \{\alpha, \alpha_1, \alpha_3, \alpha_4, \alpha_6\}$ での比較

Result: ^{236}U Fragment Mass Distributions

14 MeV $n + ^{235}\text{U}$

JENDL FPY-2011

K. Shibata et al. (2011)



3D \rightarrow 4D ($3\text{D} + \alpha_3$), 5D ($4\text{D} + \alpha_6$): prolate + sphere

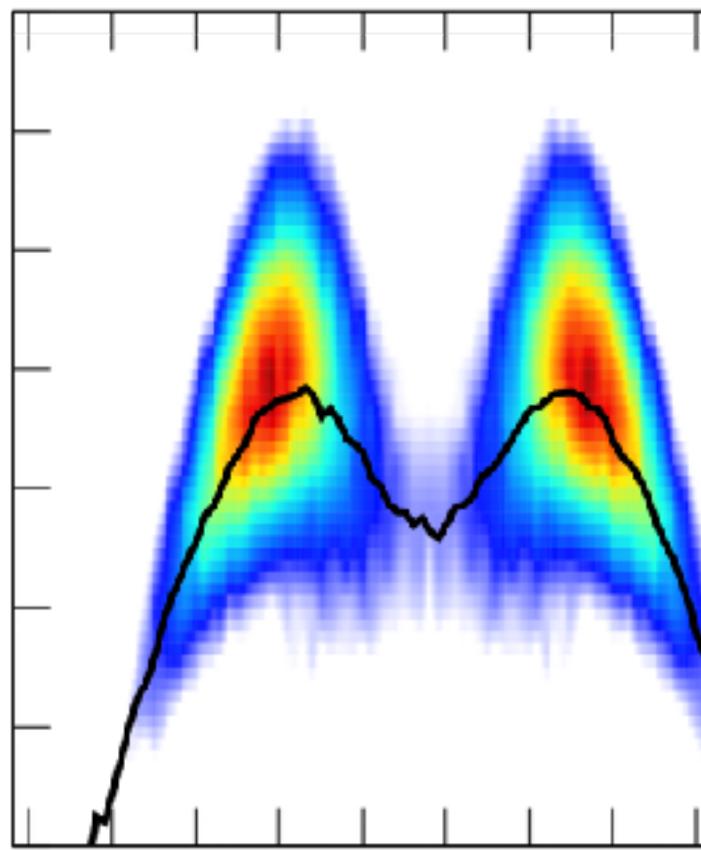
Result: ^{236}U FMN-

TKE [MeV]

4D ($\alpha, \alpha_1, \alpha_3, \alpha_4$)

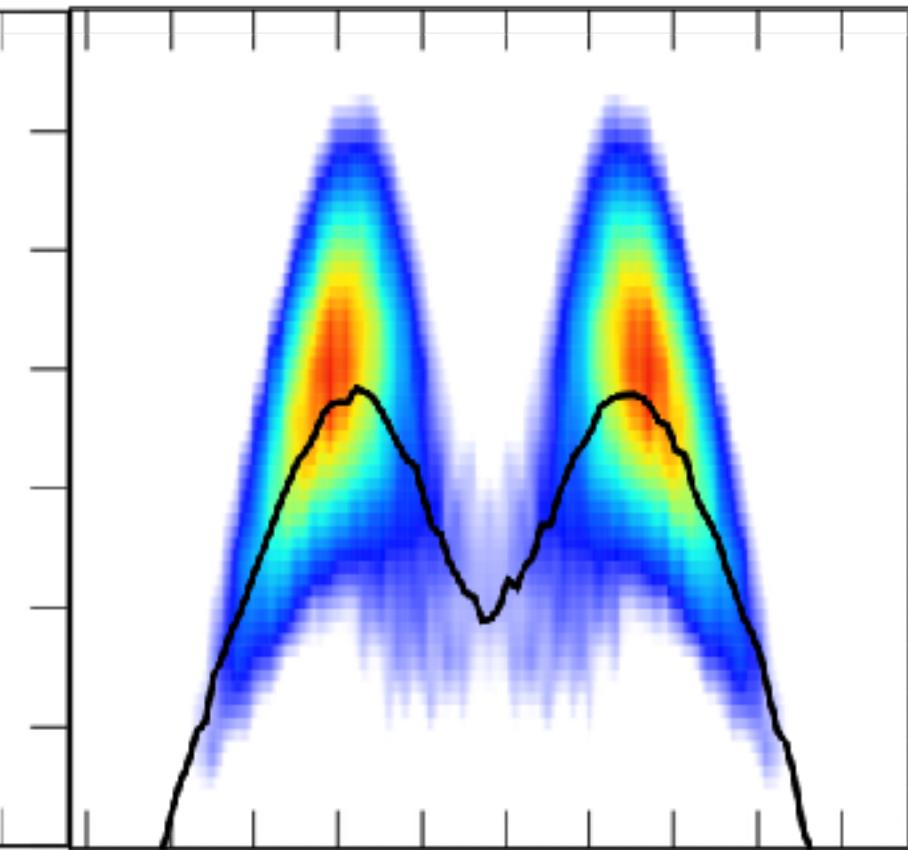
Total Kinetic Energy

195
190
185
180
175
170
165
160



Fragment Mass Number

5D ($\alpha, \alpha_1, \alpha_3, \alpha_4, \alpha_6$)



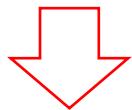
Fragment Mass Number

super-long symmetric fission

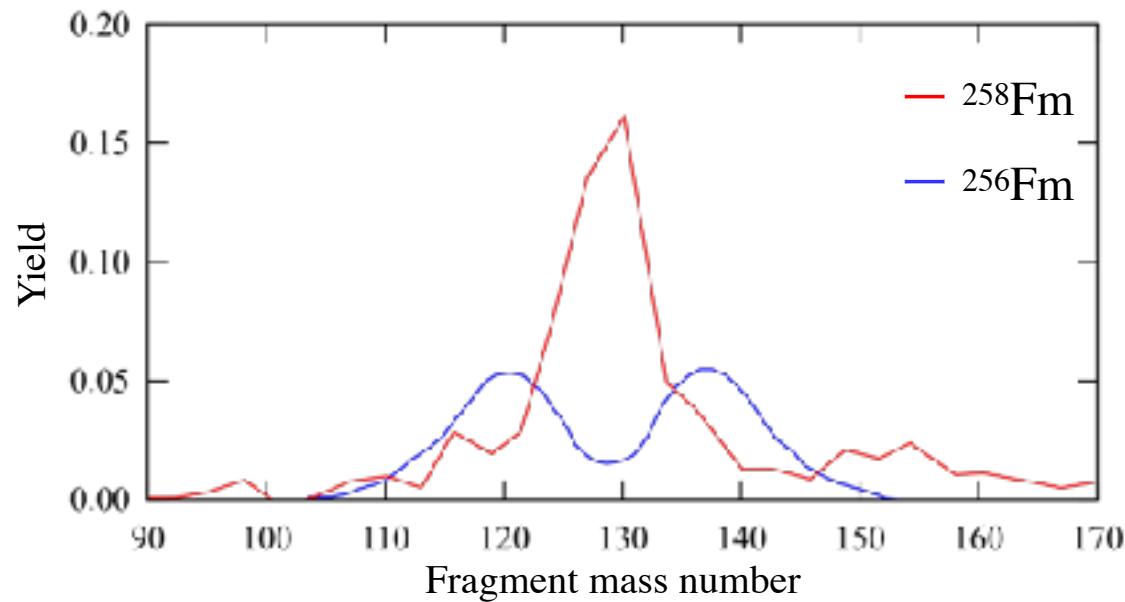
Fm同位体の分裂片質量分布

- ・質量分布が変化

非対称分裂(^{256}Fm)



対称分裂(^{258}Fm ~)



Fm分裂片質量分布 (実験値)

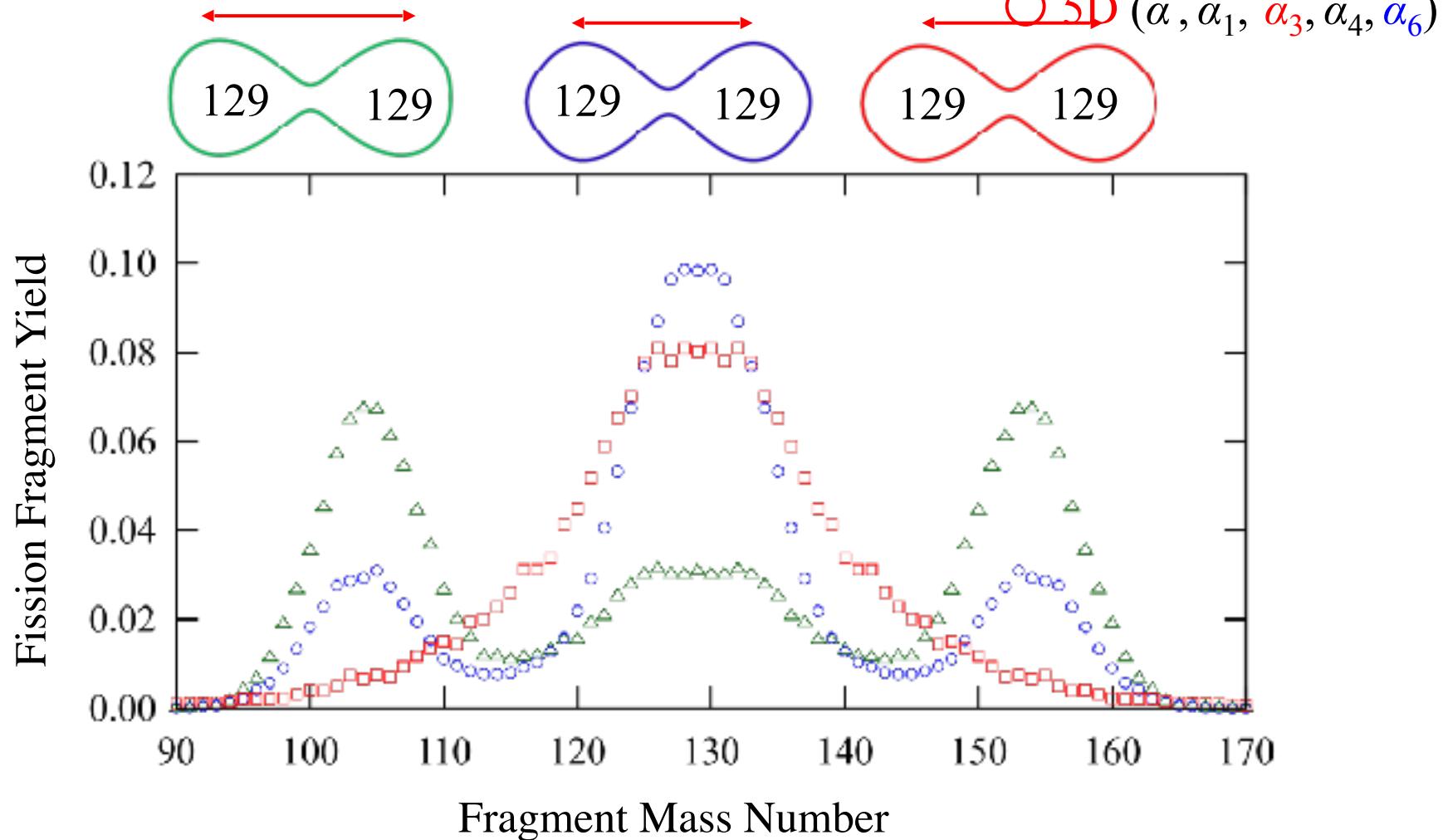
- ・対称分裂を描写できる自由度
- ・Cassiniパラメータの追加

D. C. Hoffman et al., Phys. Rev. C 21, 972–981, (1980).

→

$\alpha, \alpha_1, \alpha_4, \alpha_6 + \alpha_3$

- ・非対称分裂の考慮

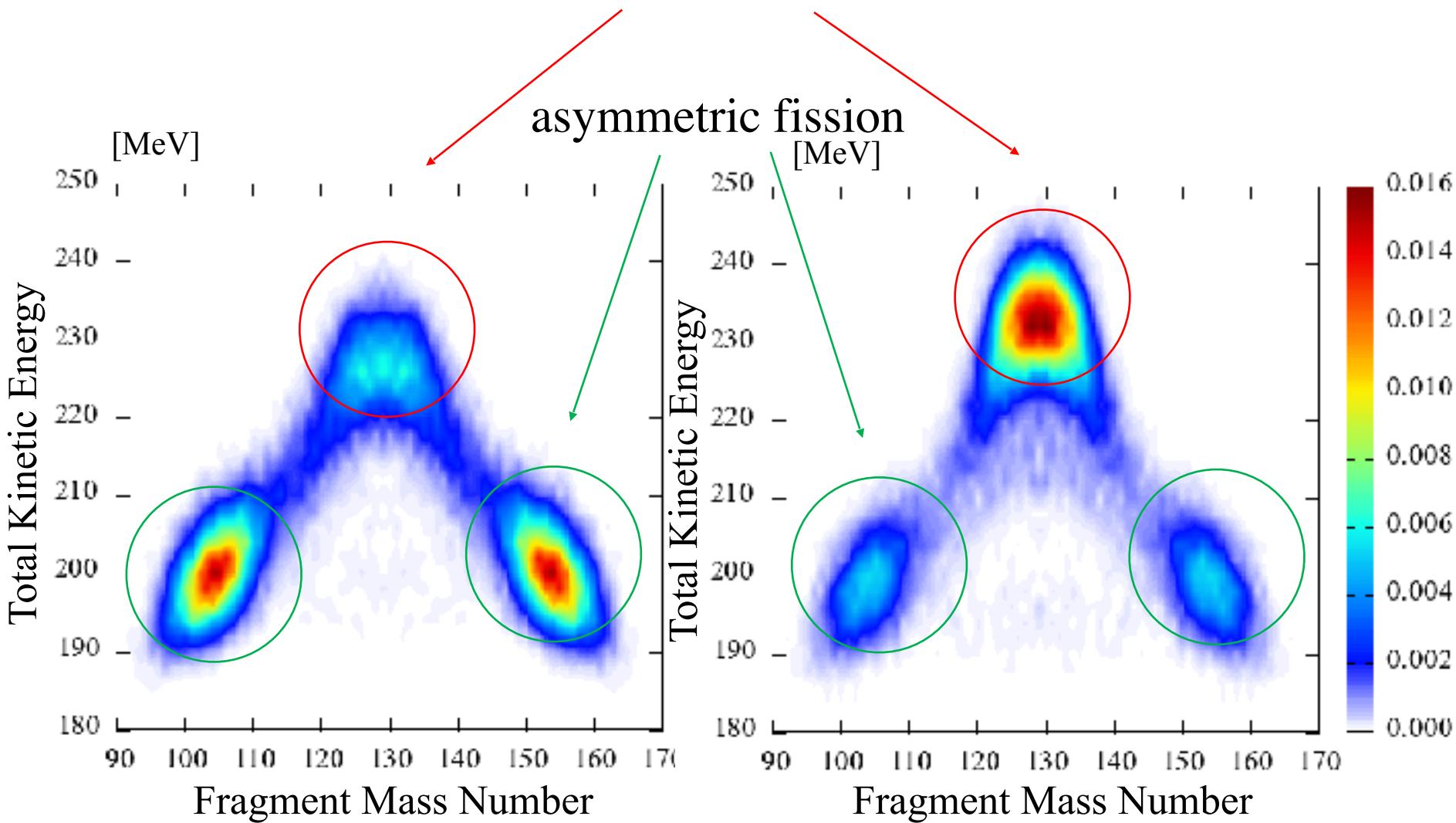
\triangle 3D ($\alpha, \alpha_1, \alpha_4$) \square 4D ($\alpha, \alpha_1, \alpha_4, \alpha_6$) \circlearrowleft 5D ($\alpha, \alpha_1, \alpha_3, \alpha_4, \alpha_6$)

Result: ^{258}Fm TKE-FMN

$^{258}\text{Fm} \{\alpha, \alpha_1, \alpha_4\}$

$^{258}\text{Fm} \{\alpha, \alpha_1, \alpha_4, \alpha_6\}$

super-short symmetric fission

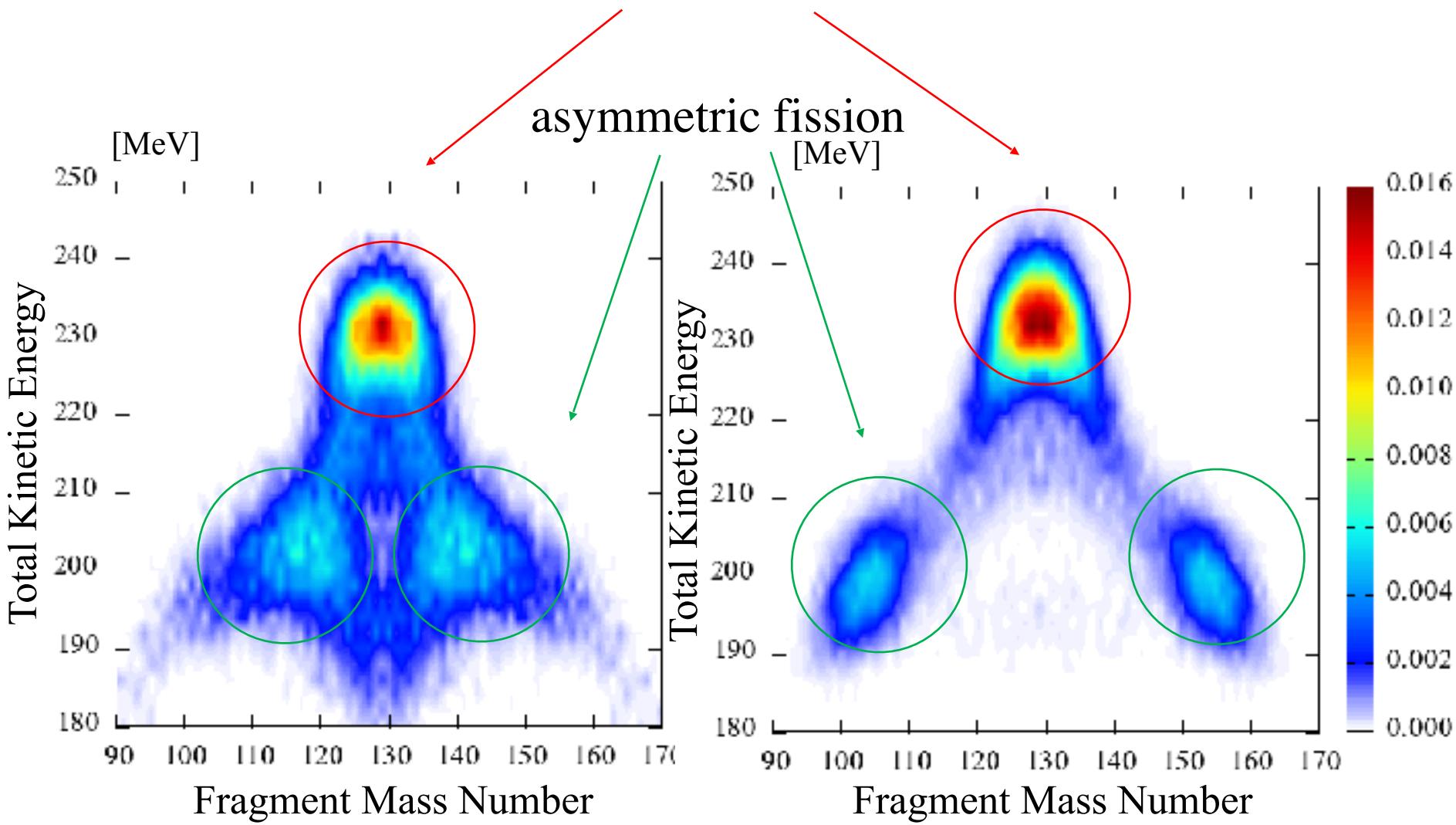


Result: ^{258}Fm TKE-FMN

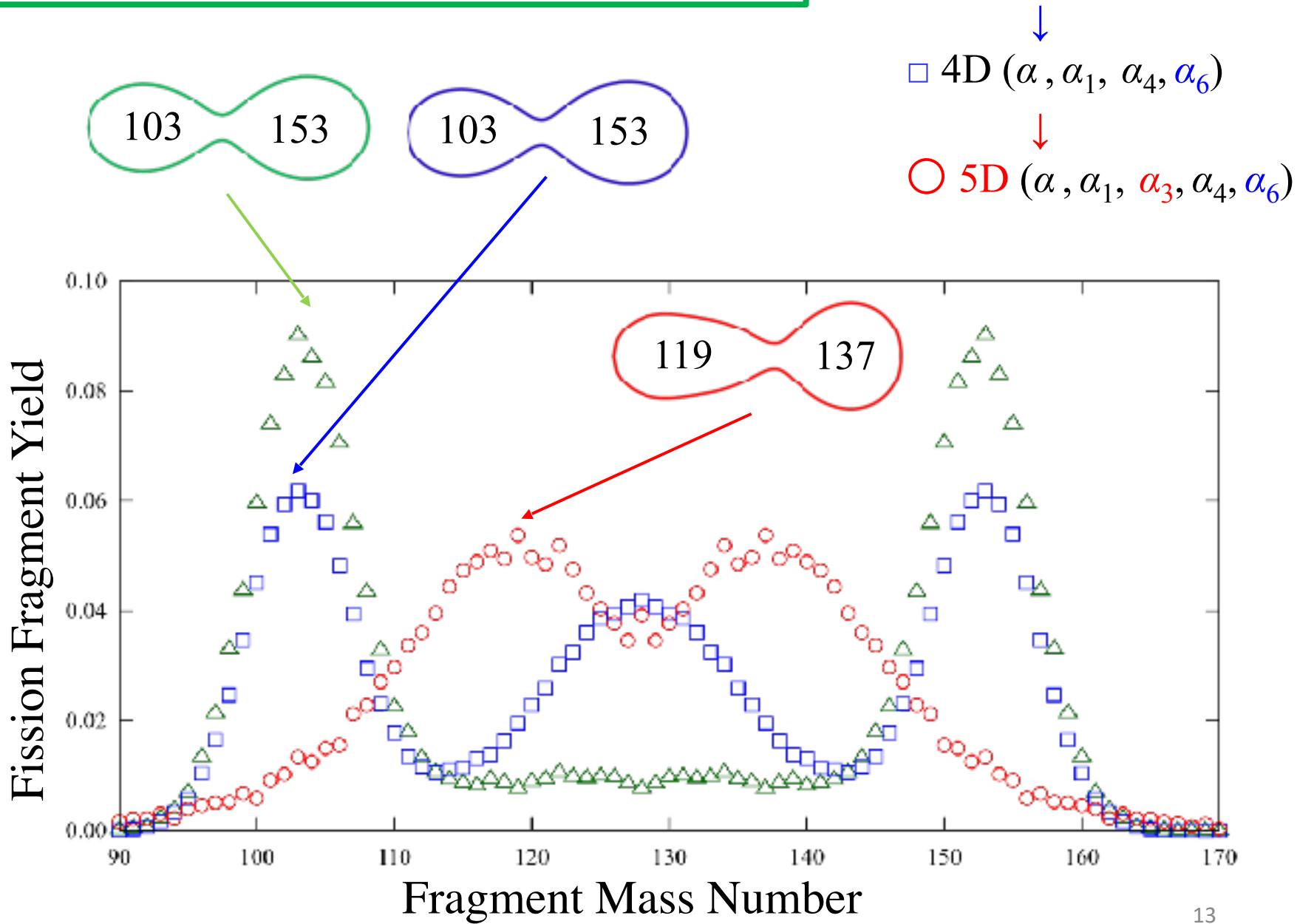
$^{258}\text{Fm} \{\alpha, \alpha_1, \alpha_3, \alpha_4, \alpha_6\}$

$^{258}\text{Fm} \{\alpha, \alpha_1, \alpha_4, \alpha_6\}$

super-short symmetric fission



Result: ^{256}Fm Fragment Mass Distributions

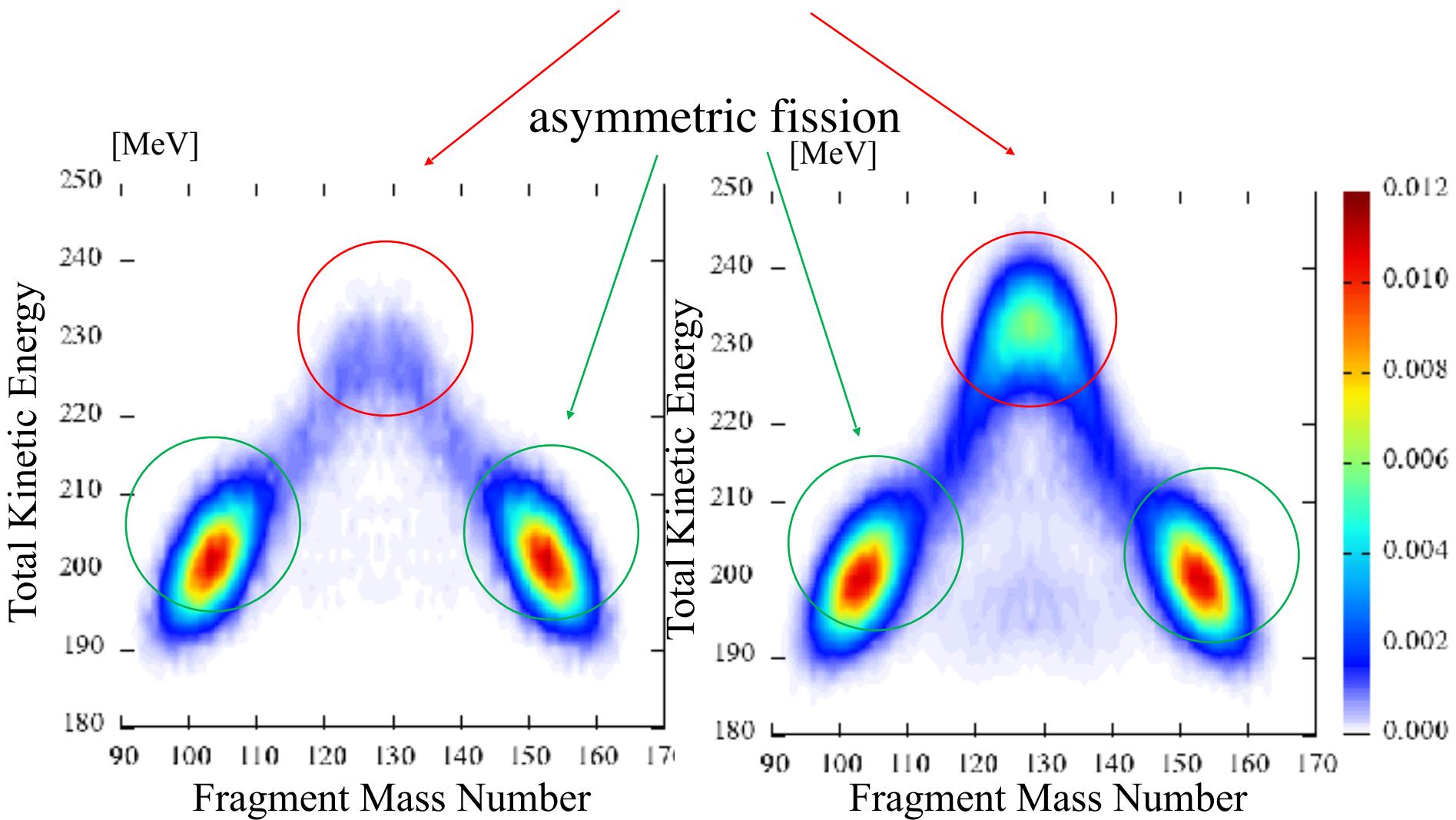


Result: ^{256}Fm TKE-FMN

$^{256}\text{Fm} \{\alpha, \alpha_1, \alpha_4\}$

$^{256}\text{Fm} \{\alpha, \alpha_1, \alpha_4, \alpha_6\}$

super-short symmetric fission

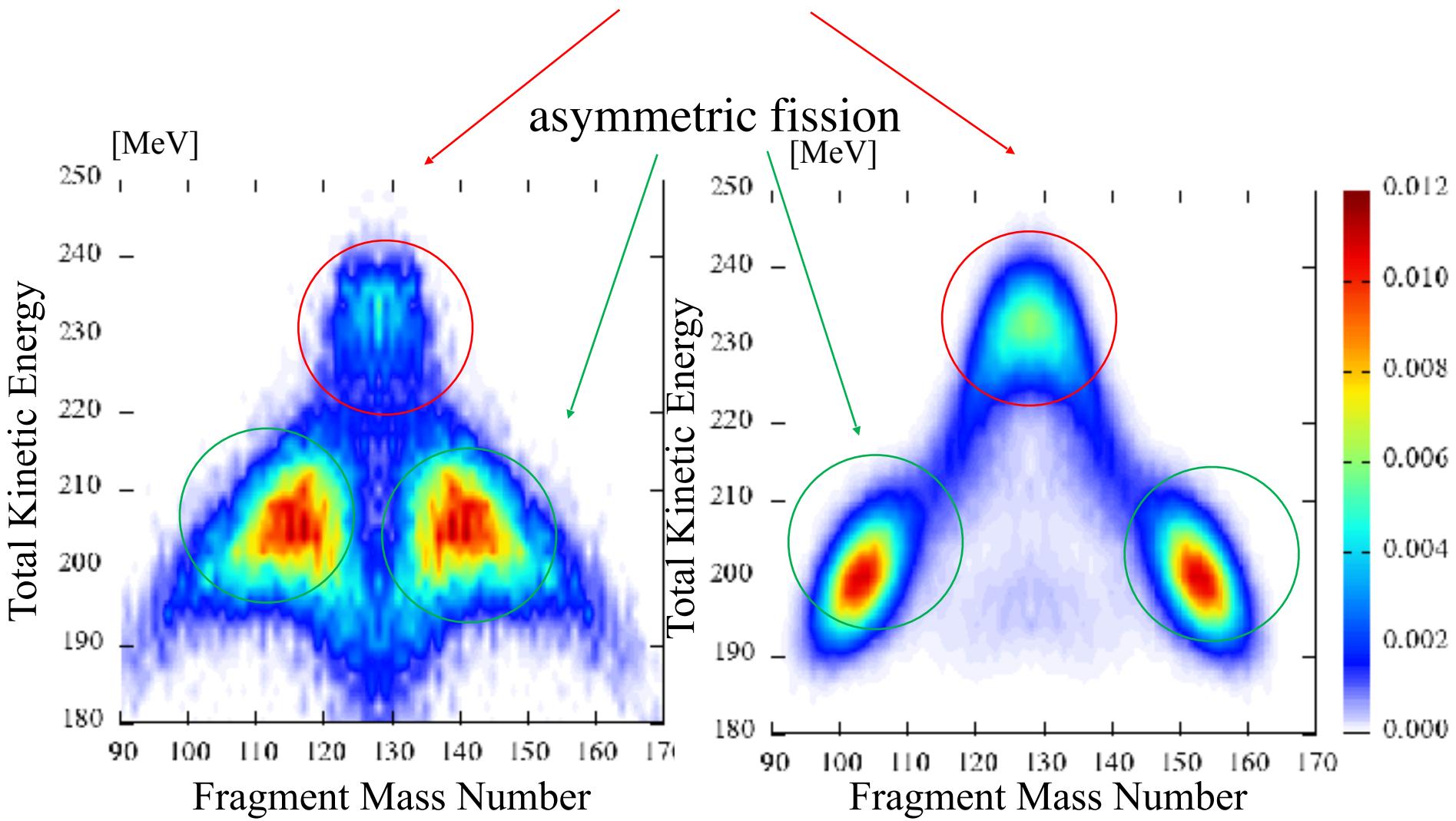


Result: ^{256}Fm TKE-FMN

$^{256}\text{Fm} \{\alpha, \alpha_1, \alpha_3, \alpha_4, \alpha_6\}$

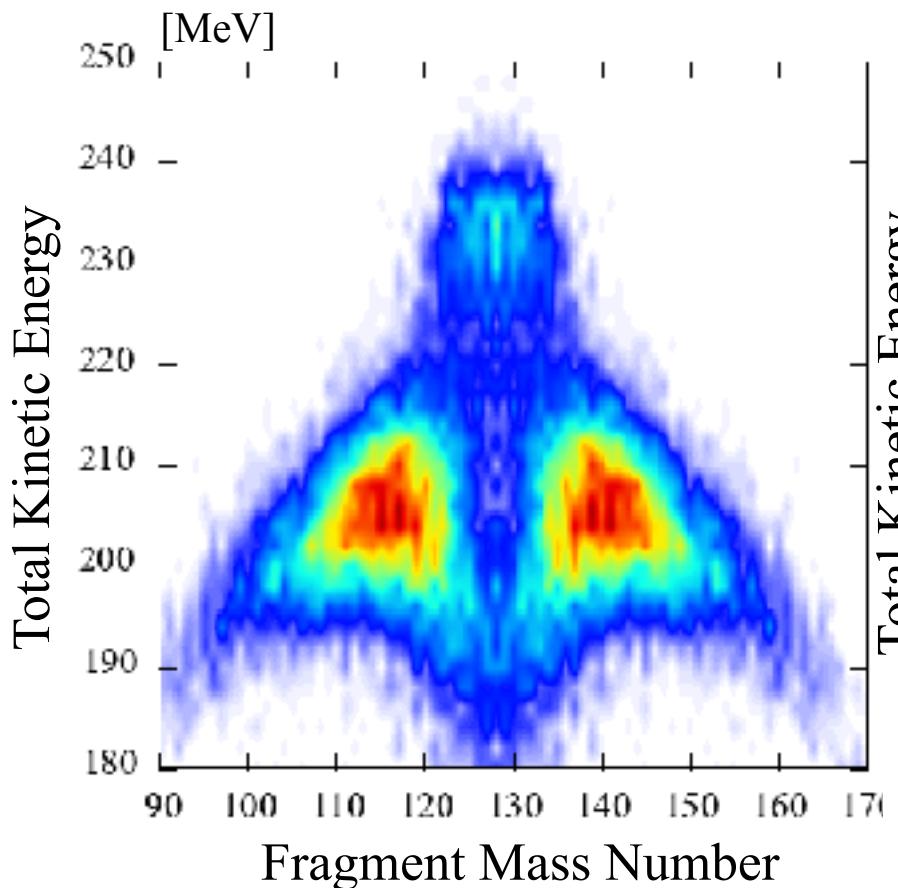
$^{256}\text{Fm} \{\alpha, \alpha_1, \alpha_4, \alpha_6\}$

super-short symmetric fission

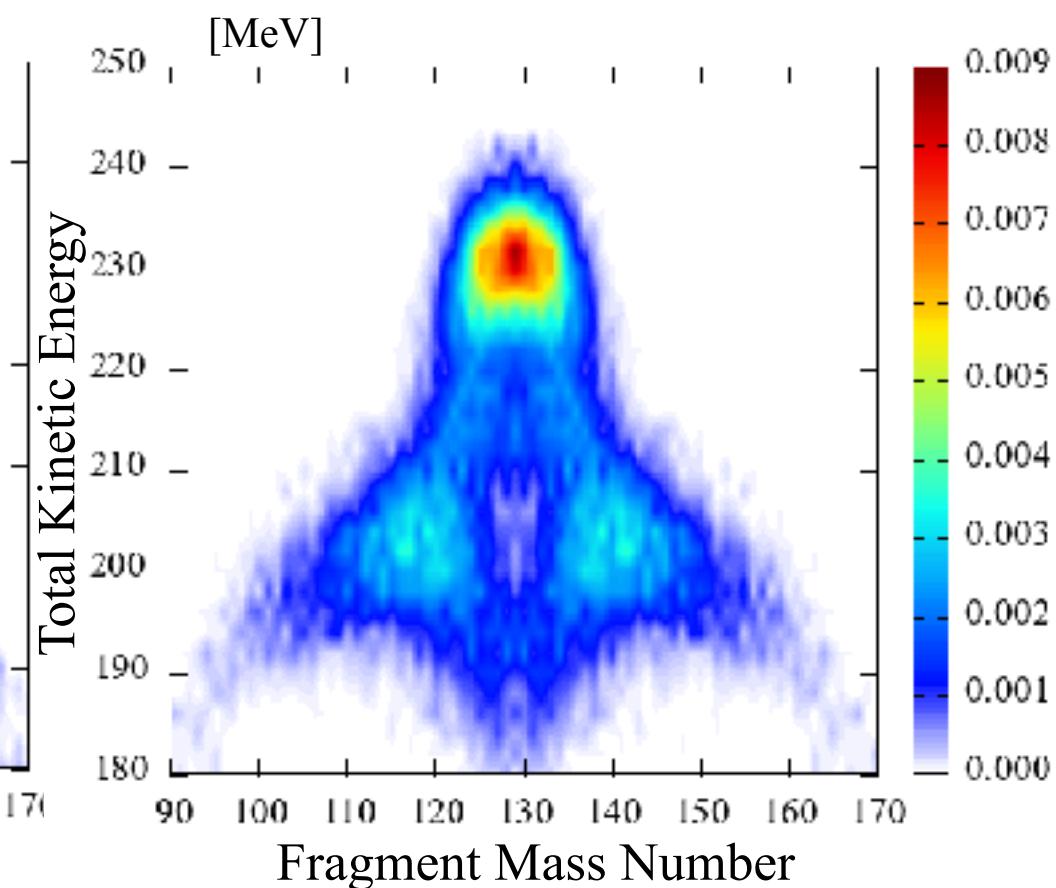


Result: TKE-FMN

$^{256}\text{Fm} \{\alpha, \alpha_1, \alpha_3, \alpha_4, \alpha_6\}$



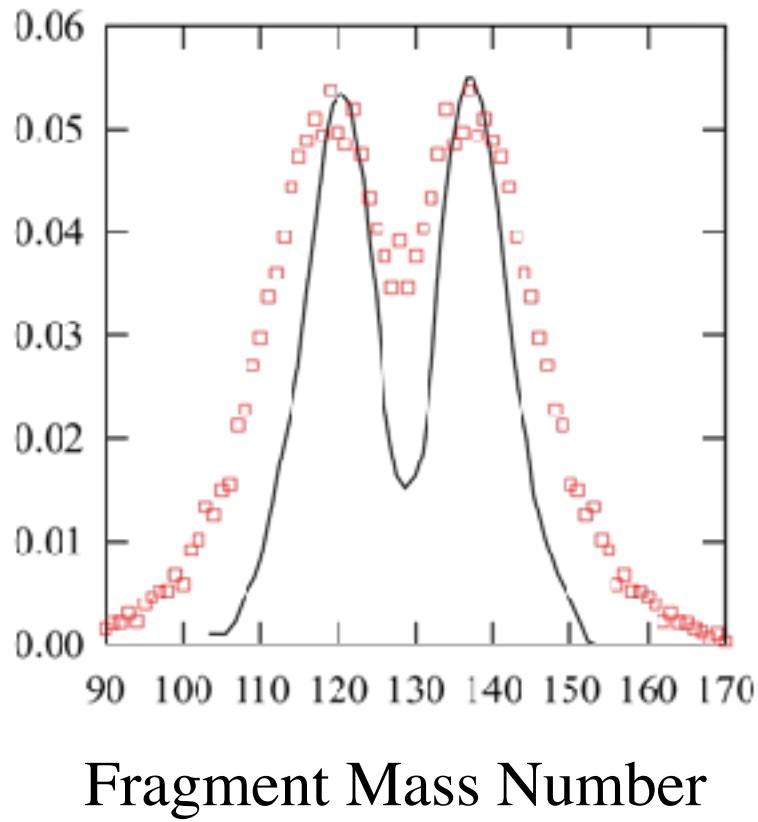
$^{258}\text{Fm} \{\alpha, \alpha_1, \alpha_3, \alpha_4, \alpha_6\}$



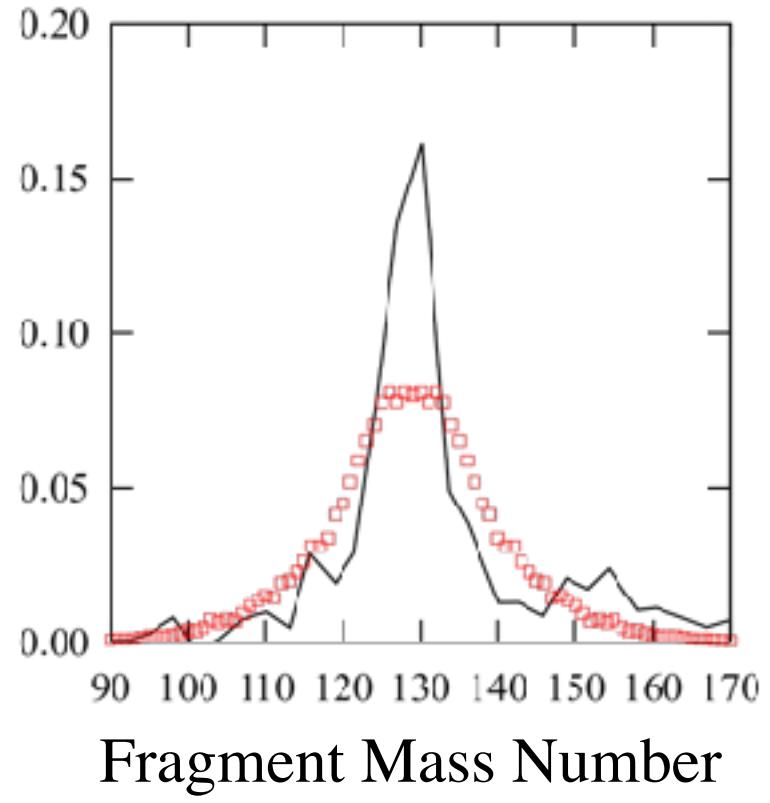
Result: Fm Fragment Mass Distributions

○ 5D ($\alpha, \alpha_1, \alpha_3, \alpha_4, \alpha_6$)
— exp

D. C. Hoffman et al., Phys.
Rev. C 21, 972–981, (1980).



^{256}Fm



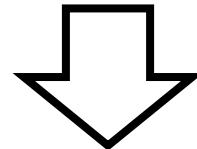
^{258}Fm

まとめ

- Cassiniパラメータを用いた核分裂Langevin計算の5次元化

非対称分裂に特化 $\{\alpha, \alpha_1, \alpha_3, \alpha_4\}$

対称分裂に特化 $\{\alpha, \alpha_1, \alpha_4, \alpha_6\}$



$\{\alpha, \alpha_1, \alpha_3, \alpha_4, \alpha_6\}$

系の分裂様相に関わらず
採用すべきパラメータセット

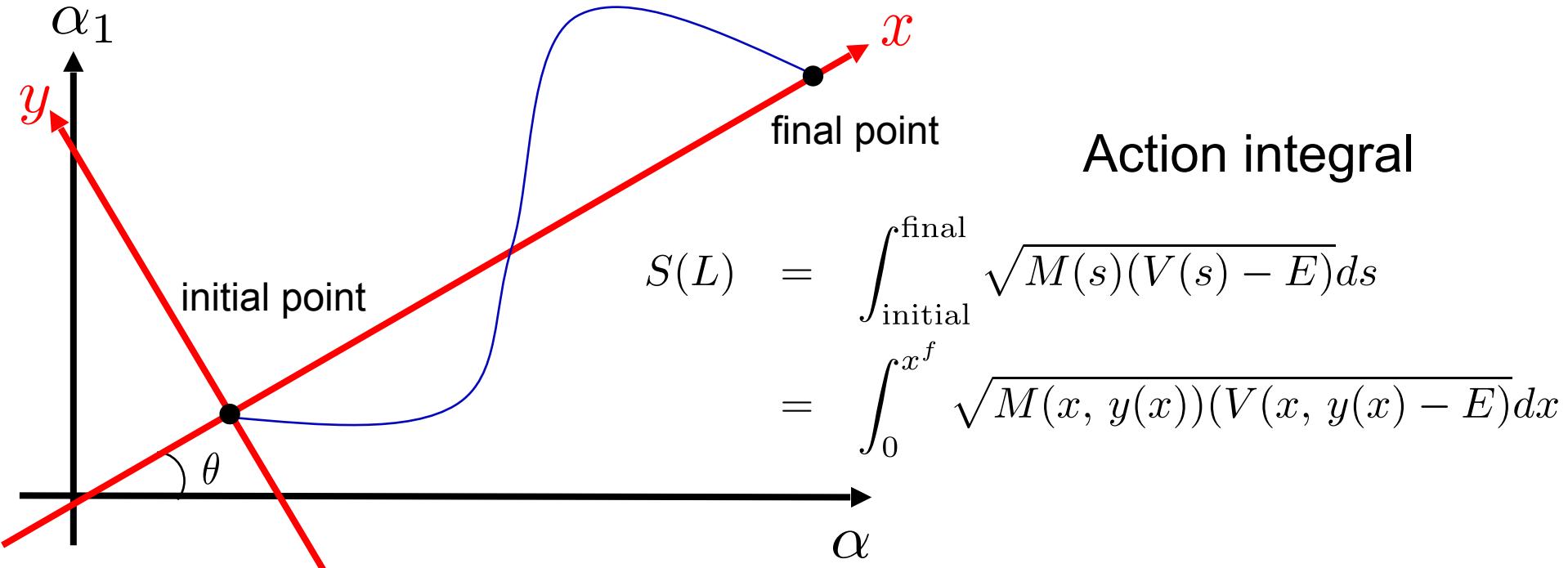
展望

今後: アクチノイド領域での検討

Langevin計算の改良 (particle emission, linear response)

5次元ポテンシャル上fission pathの解析 (次スライド)

Future Work (1): Action Integral



initial point: $(\alpha, \alpha_1) = (\alpha^0, \alpha_1^0)$

$$(x, y) = (0, 0)$$

final point: $(\alpha, \alpha_1) = (\alpha^f, \alpha_1^f)$

$$(x, y) = (x^f, 0)$$

- Formulation of Ritz method

$$y(x) = \sum_k a_k \sin \frac{k\pi x}{x^f}$$

Future Work (1): Action Integral

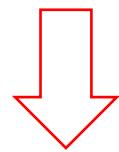
2D → 多次元化

- Multi-dimensional Ritz method

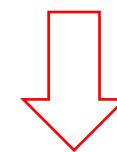
$$\bar{\alpha}_i(\bar{\alpha}) = \sum_{k=1}^{N_i} b_k^i \sin \frac{k\pi\bar{\alpha}}{\bar{\alpha}^f}$$

Action integral

$$\begin{aligned} S(L) &= \int_{\text{initial}}^{\text{final}} \sqrt{M(s)(V(s) - E)} ds \\ &= \int_0^{x^f} \sqrt{M(x, y(x))(V(x, y(x)) - E)} dx \end{aligned}$$



Werner-Wheeler mass
or
Cranking mass



deformation
energy

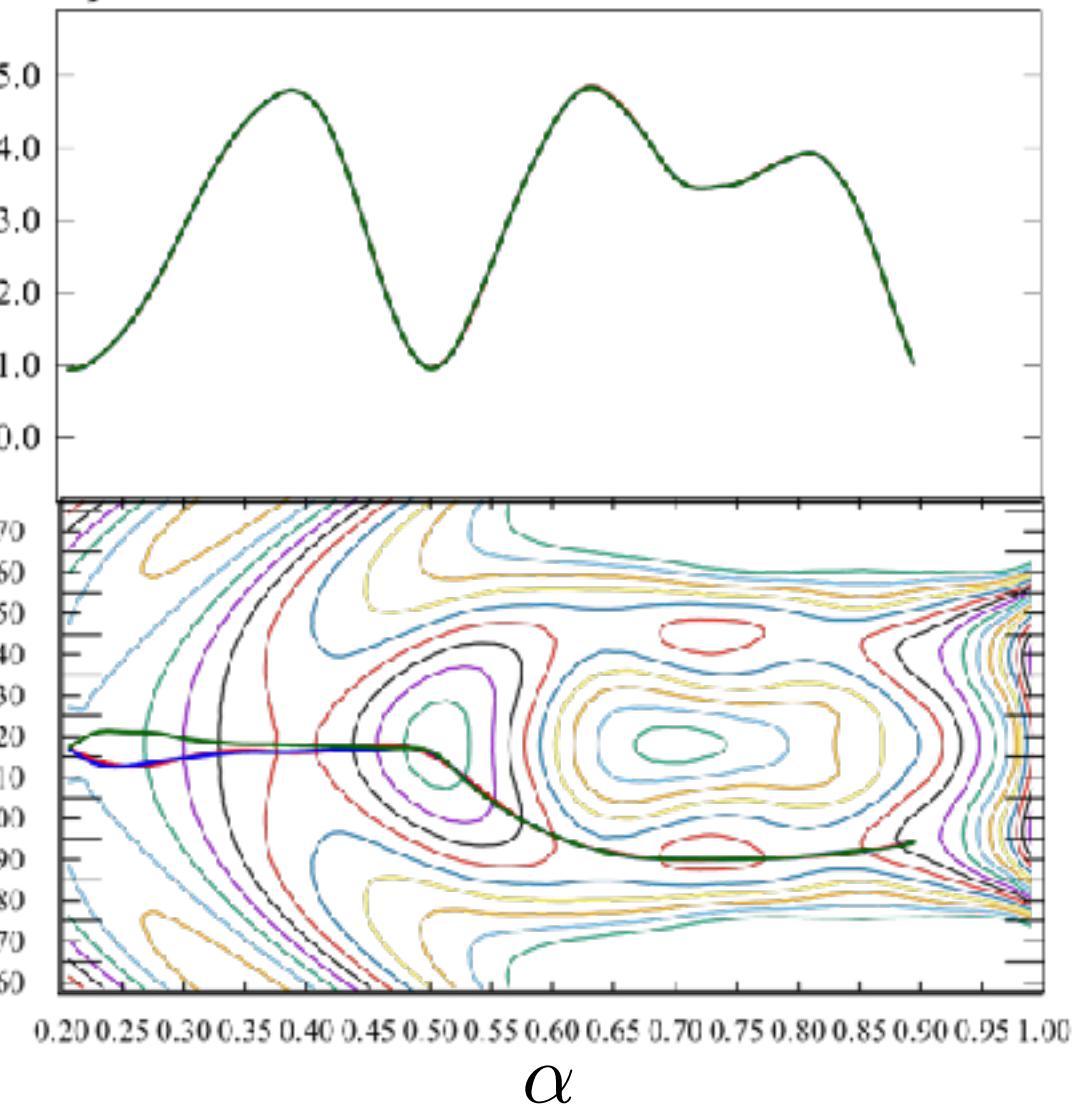
$^{256}\text{Fm}, ^{258}\text{Fm}$ 自発核分裂へのアプローチ

5次元ポテンシャルの解析

Future Work (1): Action Integral

2D $\{\alpha, \alpha_1\}$

^{236}U E
Fragment Mass Number



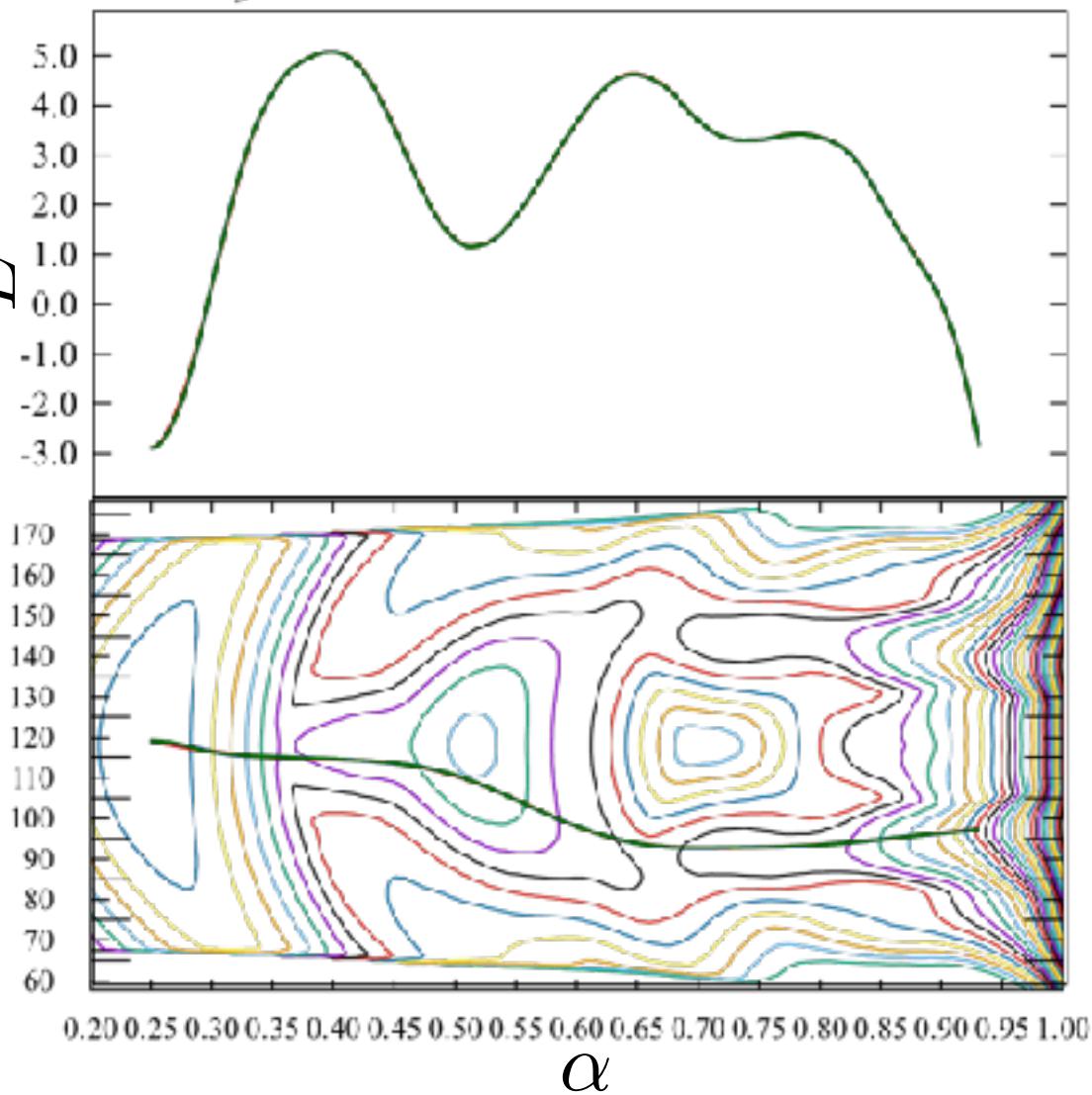
red: $N_i=10$
blue: $N_i=20$
green: $N_i=30$

Future Work (1): Action Integral

3D $\{\alpha, \alpha_1, \alpha_4\}$

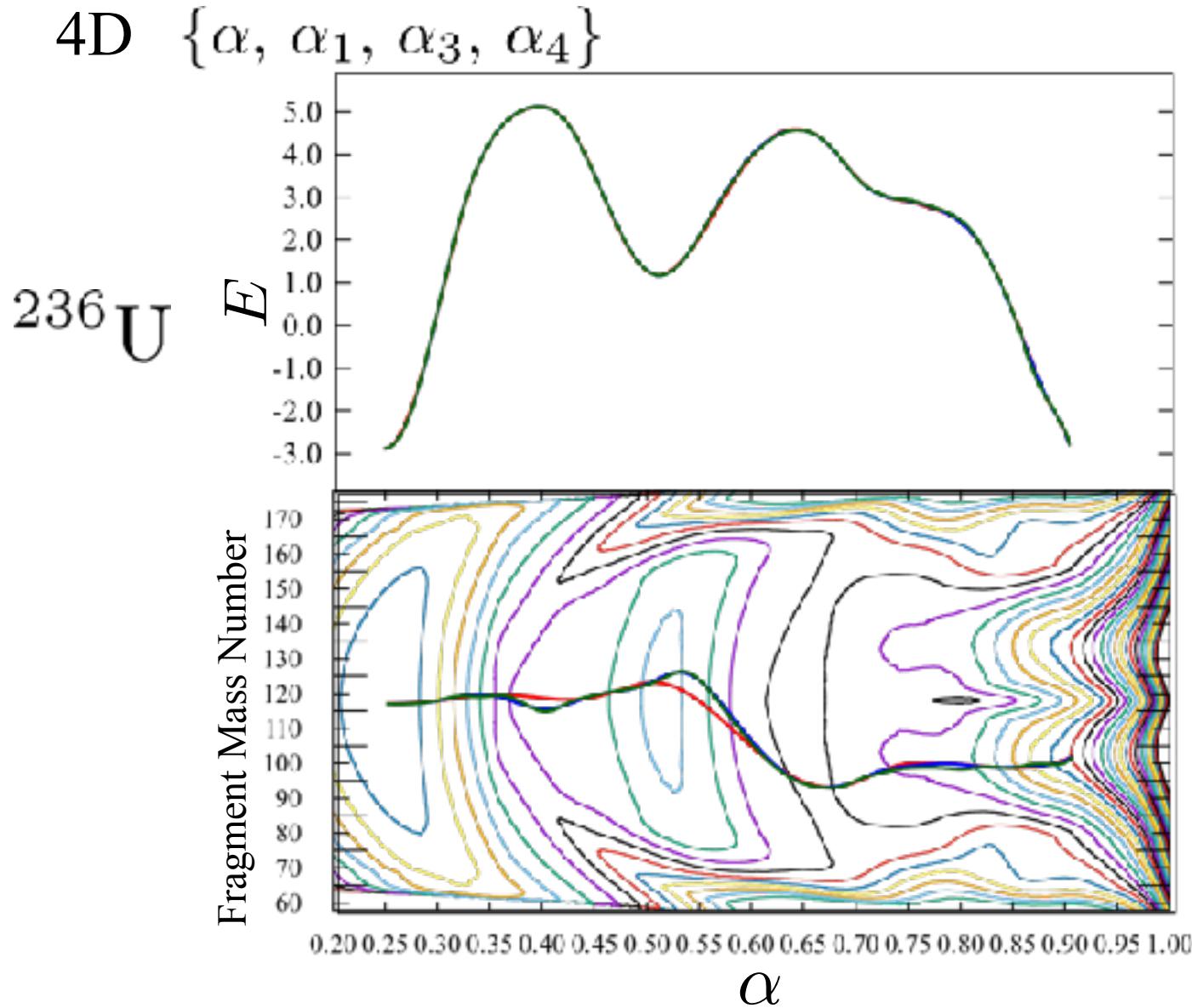
^{236}U

Fragment Mass Number



red: $N_i=10$
blue: $N_i=20$
green: $N_i=30$

Future Work (1): Action Integral



展望

5次元Cassiniパラメータを用いた核分裂研究

$$\{\alpha, \alpha_1, \alpha_3, \alpha_4, \alpha_6\}$$

fission pathの解説

- Langevin計算から得られる軌跡を解析
- Minimum action integral計算

→ 中性子数変化に対する劇的な分裂様相の変化にアプローチ

Thank you for your kind attention!