# '56Ni problem' in Supernova Explosions

# **Ryo Sawada** The University of Tokyo/ JSPS postdoctoral fellow

### Reference :

• "Nucleosynthesis Constraints on the Energy Growth Timescale of a Core-collapse Supernova Explosion"

#### Sawada & Maeda (2019), ApJ, 886, 47 (arXiv.1910.06972)

- "A Consistent Modeling of Neutrino-driven Wind with Accretion Flow onto a Protoneutron Star and its Implications for 56Ni Production"
   Sawada & Suwa (2021), ApJ, 908, 6 (arXiv.2010.05615)
- "Constraints on the Explosion Timescale of Core-collapse Supernovae Based on Systematic Analysis of Light Curves"

#### Saito, Tanaka, Sawada, & Moriya (2022), ApJ 931, 153 (arXiv:2205.00624)

• "Updating the \$^{56}\$Ni Problem in Core-collapse Supernova Explosion"

#### Sawada & Suwa (2023) arXiv:2301.03610

(2) discussion.

(3) future studies.

# What is the core-collapse supernova?





## • <u>Standard scenario : Neutrino-driven model</u>





supernova brightness (hereafter, refer 56Ni as 'nickel').





Figure P30.87 The light curve for Supernova 1987A.

• the amount of 56Ni synthesized can be estimated fairly accurately and easily from the SN light curve alone.

Overview.(1) in Canonical SN.(2) discussion.Role of isotope 56Ni in supernova explosions

 Radioactive decay heat of 56Ni is the energy source of supernova brightness (hereafter, refer 56Ni as 'nickel').





(3) future studies.

# Overview.(1) in Canonical SN.(2) discussion.Role of isotope 56Ni in supernova explosions

• Light-curve calculation: Based on One-zone Arnett model (Arnett 1982)

$$\frac{\partial E_{\text{int}}(t)}{\partial t} = -P \frac{\partial V(t)}{\partial t} + \dot{Q}(t) - L(t)$$
$$\frac{\partial E_{\text{kin}}(t)}{\partial t} = P \frac{\partial V(t)}{\partial t}.$$

Assuming that the internal energy is dominated by radiation,

$$P\frac{\partial V(t)}{\partial t} = \frac{\epsilon}{3} \cdot 4\pi R_{\rm ej}^2 v_{\rm ej}(t) = \frac{E_{\rm int}(t)}{R_{\rm ej}} v_{\rm ej}(t) = \frac{E_{\rm int}(t)}{t_{\rm dyn}} , \qquad \left[ t_{\rm dyn} = \frac{R_{\rm ej}(t)}{v_{\rm ej}(t)} \right]$$

When the luminosity is described by the spherical diffusion equation,

$$L(t) = -4\pi r^2 \frac{c}{3\kappa\rho} \frac{\partial e}{\partial r} = -4\pi r^2 \frac{c}{3\kappa\rho} \frac{\partial}{\partial r} \left( \frac{E_{\rm int}(t)}{V(t)} \cdot \left( \frac{\pi}{3} \frac{\sin(\pi x)}{x} \right) \right)$$
$$= E_{\rm int} \cdot \frac{t_{\rm dyn}}{t_{\rm diff}^2}, \qquad \left[ t_{\rm diff} = \frac{3}{4\pi} \frac{\kappa M_{\rm ej}}{c v_{\rm ej}} \frac{1}{\xi} \right]$$

(3) future studies.

# Overview.(1) in Canonical SN.(2) discussion.Role of isotope 56Ni in supernova explosions



Arnett model : for Type-I Supernova

$$\frac{\partial E_{\text{int}}(t)}{\partial t} = -\frac{E_{\text{int}}(t)}{t_{\text{dyn}}} + \dot{Q}(t) - E_{\text{int}} \cdot \frac{t_{\text{dyn}}}{t_{\text{diff}}^2}$$
$$\frac{\partial E_{\text{kin}}(t)}{\partial t} = \frac{E_{\text{int}}(t)}{t_{\text{dyn}}}.$$
$$\dot{Q}(t) = f_{\text{dep}} \cdot [M({}^{56}\text{Ni}) \cdot q_{\text{Ni}}(t)] + Q_{\text{heat}}$$



(3) future studies.





 $M(^{56}Ni) \propto \begin{cases} 1. explosion energy. (How much extends the 56Ni synthesis region?) \\ 2. compactness. (How much mass is contained in the 56Ni synthesis region?) \end{cases}$ 



### Observational



### Theoretical



 $\therefore M({}^{56}\text{Ni}) \propto \begin{bmatrix} 1. \text{ explosion energy.} (How much extends the 56Ni synthesis region?) \\ 2. \text{ compactness.} (How much mass is contained in the 56Ni synthesis region?) \end{bmatrix}$ 



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## • <u>Standard scenario : Neutrino-driven model</u>



\*unclear : Can reach to  $10^{51}$  [erg] ? (see, e.g., Janka 2012, and references therein)



- ✓ **The state-of-the-art simulations** employ Multi-D / Full-Boltzmann *v* -transport / GR-effect etc.
- ✓ Most, if not all, of these simulations do not reach sufficient explosion energy.
- ✓ those simulations now suggest a "slower" explosion than previously thought.





Sawada & Maeda (ApJ, 2019) arxiv.1910.06972















# Overview.(1) in Canonical SN.(2) discussion.(3) future studies.[topic 1] Can modern supernova simulations synthesize sufficient 56Ni?

### ※ note : This model is not a fine-turned model to SN1987A



Note: Multi-D effect may enhance M(44Ti) (Nagataki et al. 1998; Maeda & Nomoto 2003).

Overview.(1) in Canonical SN.(2) disucussion(3) future studies.[topic 1] Can modern supernova simulations synthesize sufficient 56Ni?

## • NOTE (1): all models ?

some models in the state-of-the-art simulations have succeeded in producing the typical mass  $0.07 M_{\odot}$  of 56Ni in CCSNe,

## But 'nickel mass problem'(Ni problem)

→ unclear whether we can reproduce sufficient 56Ni amount <u>as a canonical nature.</u>





Overview.(1) in Canonical SN.(2) disucussion(3) future studies.[topic 1] Can modern supernova simulations synthesize sufficient 56Ni?

• NOTE (2): choice of typical value?



# How to estimate the progenitor structure before the explosion



1D hydrodynamics and nucleosynthesis with explosion timescale as a parameter

Overview.(1) in Canonical SN.(2) disucussion(3) future studies.[topic 1] Can modern supernova simulations synthesize sufficient 56Ni?

### • NOTE (2): choice of typical value?



A few tens of percent of observations can be explained by the *non-standard supernova explosions*,

But, the timescales that reproduce less than 50% are difficult to adopt as standard explosion models.

"Constraints on the Explosion Timescale of Core-collapse Supernovae Based on Systematic Analysis of Light Curves" Saito, Tanaka, Sawada, & Moriya (2022), ApJ 931, 153 (arXiv:2205.00624)



• NOTE (3): Opposition to this study (Imasheva + 2022)





Our work 1D hydrodynamics and nucleosynthesis

simulation

with more realistic explosion calculations /with lightbulb approximation (not detailed, but mimics neutrino-driven explosions to some extent)

$$\begin{aligned} \frac{\partial r}{\partial M_r} &= \frac{1}{4\pi r^2 \rho} , \qquad (5) \\ \frac{Dv}{Dt} &= -\frac{GM_r}{r^2} - 4\pi r^2 \frac{\partial P}{\partial M_r} , \qquad (6) \\ \frac{D\epsilon}{Dt} &= -P \frac{D}{Dt} \left(\frac{1}{\rho}\right) + \mathcal{H} - \mathcal{C} , \qquad (7) \end{aligned}$$
$$\mathcal{H} &= 1.544 \times 10^{20} \text{ erg g}^{-1} \text{ s}^{-1} \\ &\times \left(\frac{L_{\nu_e}}{10^{52} \text{ MeV}}\right) \left(\frac{r_{\nu_e}}{100 \text{ km}}\right)^{-2} \left(\frac{T_{\nu_e}}{4.0 \text{ MeV}}\right)^2 , \end{aligned}$$
$$\mathcal{C} &= 1.399 \times 10^{20} \text{ erg g}^{-1} \text{ s}^{-1} \times \left(\frac{T}{2.0 \text{ MeV}}\right)^6 . \end{aligned}$$





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Figure 11. Schematic picture of core-collapse supernova.