

～中性子星の観測と理論～研究活性化ワークショップ 2023 @Kyoto  
6-8 Sept. 2023

# 中性子星の熱的進化とパルサーの回転進化

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Introduction

# Neutron star cooling and rotation

Superfluidity affects Neutron Star Cooling



Superfluidity also affects the rotation of the star



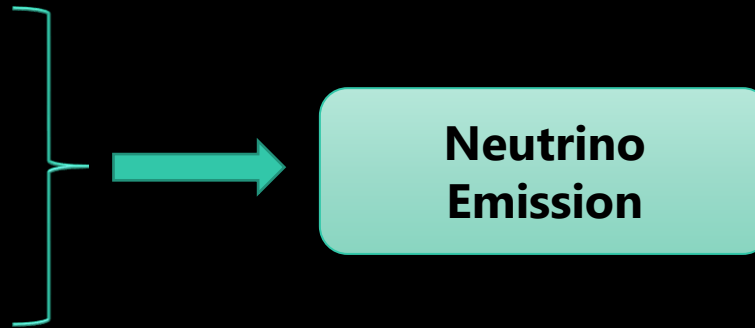
**Concurrent calculations of cooling and rotation of neutron stars allows comparison with pulsar spin-downs.**

Adding rotational effects with cooling models

# Cooling of Neutron Stars

Cooling processes of NS strongly depend on the interior state

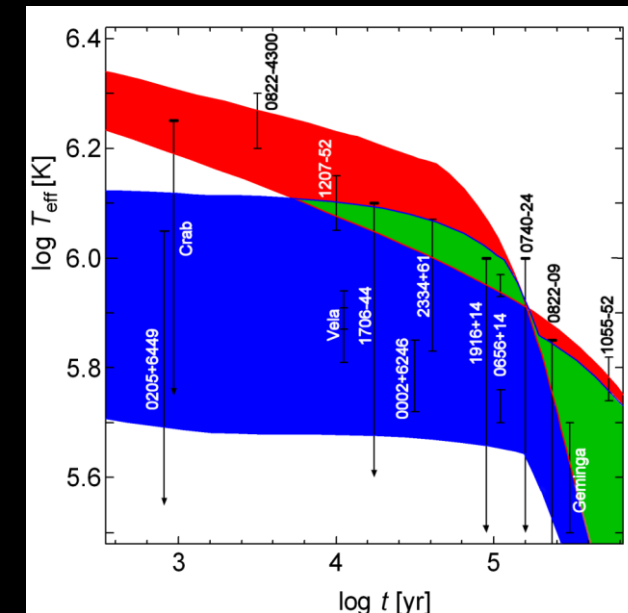
- **Normal nuclear matter**
- $\pi$  condensation
- K condensation
- **Quark matter**
- **Superfluidity**
- etc...



**Exotic phase** appears in high density state, **cools star rapidly**

Central density above the threshold density = Heavy NS

Comparing the calculation results and isolated NS observation  
 $\Rightarrow$  Constraining the high-density state



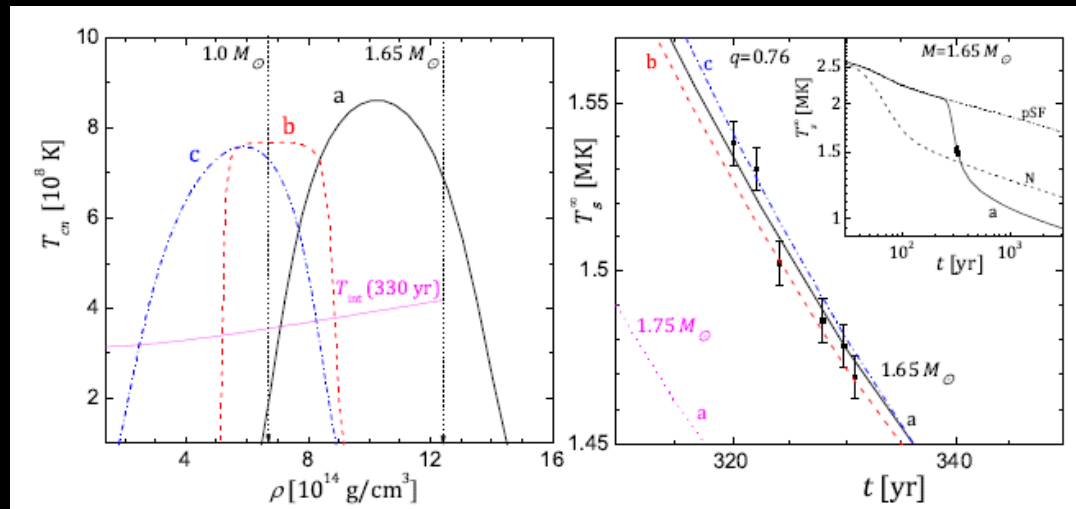
(TN+ 2006)

# Neutrino Emission

- Occurring in all NS (**Standard**)
  - Modified URCA: **Weak**
  - Nucleon superfluidity: **Marginally Strong** (at transition)
- Occurring in heavy NS (**Exotic**)
  - Quark  $\beta$ -decay (Considering **Quark Matter**): **Strong**
  - Direct URCA ( $y_p < 1/9$ ): **Strong**
- Superfluidity, Superconductivity
  - Neutrino emission at transition
  - Superfluid state **suppresses Other Neutrino Emission**  
→ mild the "Strong" emission

# Effects of Superfluidity on cooling

- Superfluidity has 2 effects on cooling
  - Transition from Normal to Super: Neutrino emission (**Accelerate cooling**)
  - After the transition: Suppression other Neutrino emission processes (**Decelerate cooling**)
- Neutron  ${}^3P_2$ 
  - The density dependence of the critical temperature determine the cooling history of neutron stars.



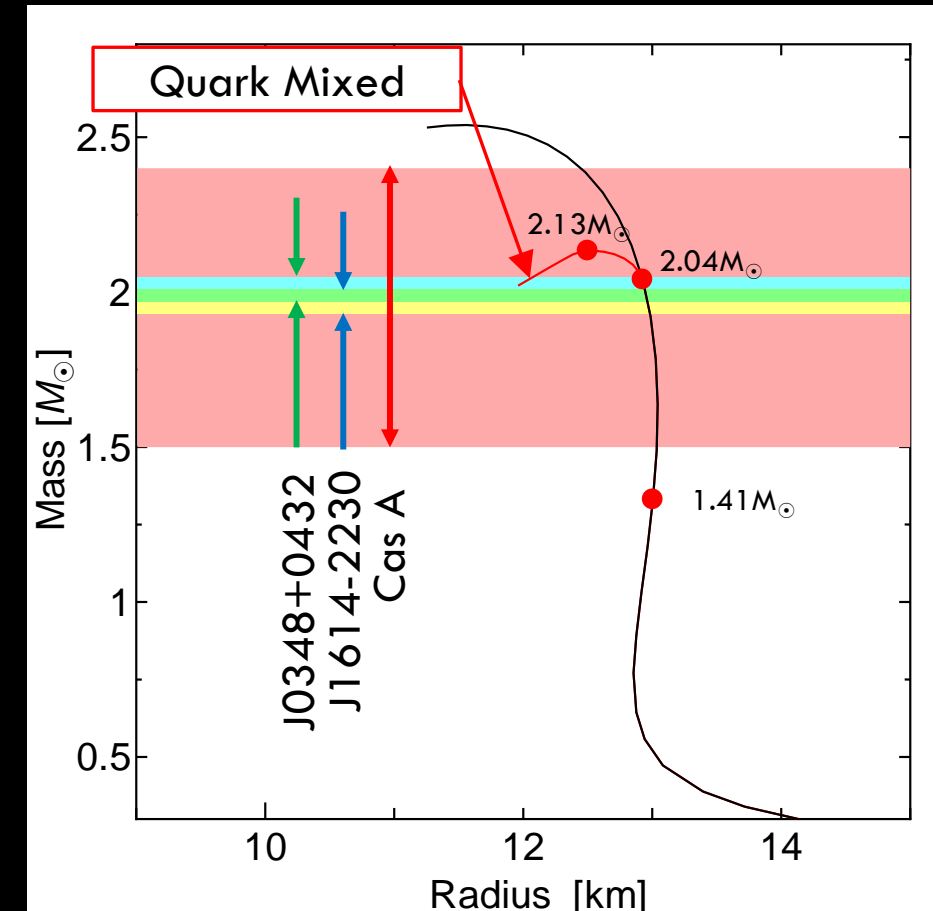
Shternin et al. 2011

# Models

- EoS with Maximum mass above  $2M_{\odot}$ 
  - Brueckner-Hartree-Fock (HM) + Dyson-Schwinger (QM)
  - Mixed phase between HM-QM (Yasutake+ 2016)

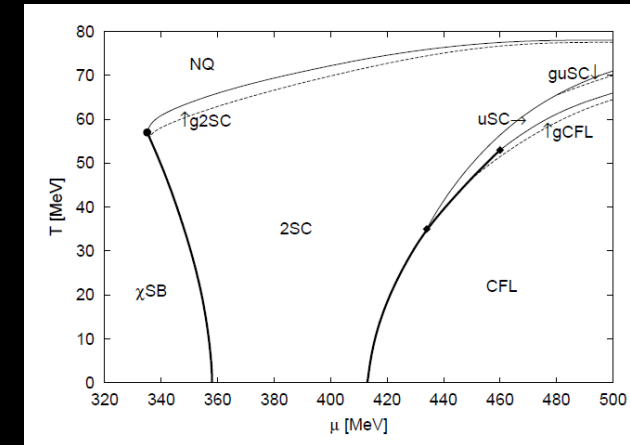
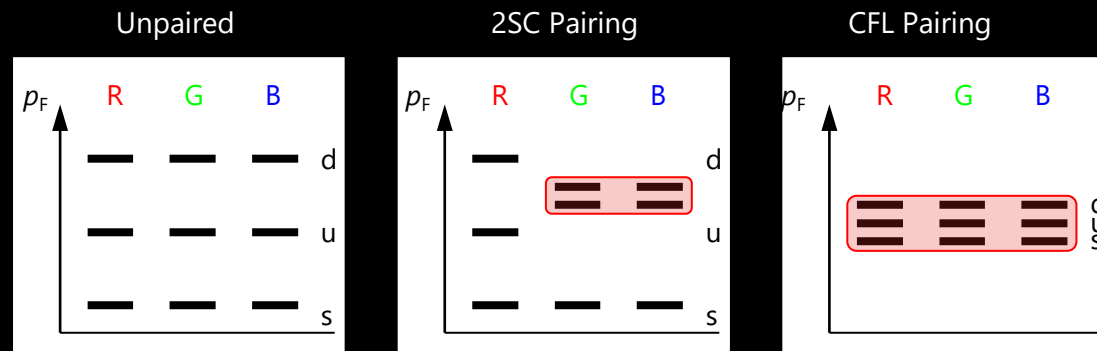
**Even at the maximum mass, the centre is mixed phase**

- Surface composition:  $^{56}\text{Fe}$
- Cooling processes
  - Modified URCA + Bremsstrahlung
  - n-Super( $^1S_0$ ,  $^3P_2$ ), p-Super( $^1S_0$ )
  - **Direct URCA** ( $y_p > 1/9$ )
  - **Quark Cooling** with Colour Superconductivity (CSC)
- Parameters
  - **Masses**
  - **n, d- $^3P_2$  Superfluidity** Critical Temperature
  - **CSC Paring** (CFL / 2SC / 2SC+X)



# Pairing of Quarks

- CSC in quark matter has **Multiple pairings**
  - Degrees of freedom of colour and flavour
- **CFL** (Colour Flavour Locking • **Higher density** X)
  - All colours and flavours can make pairs
  - **All quarks (RGB) in superconducting** → **Suppressing neutrino emission** ✓
- **2SC** (Two Flavour Superconductivity • **Lower density** ✓)
  - 2 of colours/flavours can make pairs
  - **1/3 of normal quark remains** → **Strong neutrino emission** X



Rüster et al. (2006)

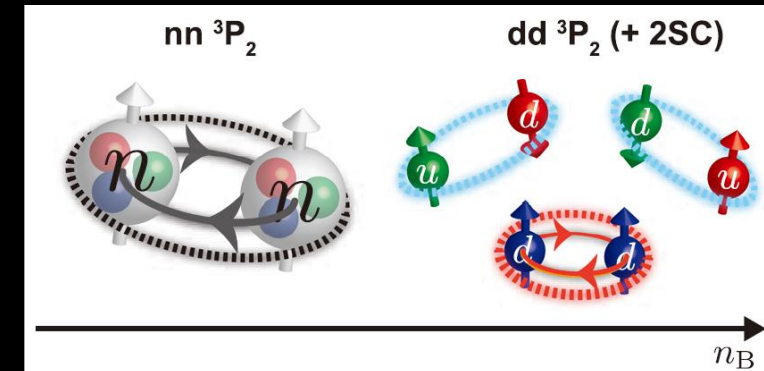
# Quark-Hadron Continuity

Neutron  ${}^3P_2 \rightarrow$  Quark  ${}^3P_2 + 2SC$  Continuous transition (Fujimoto+ PRD 101, 094009, (2020))

- Neutron  ${}^3P_2$  has been continued by d-quarks  $\rightarrow$  Other can make 2SC

- **All quarks can make pairs**

- **Suppressing neutrino emission in 2SC (2SC+X)**



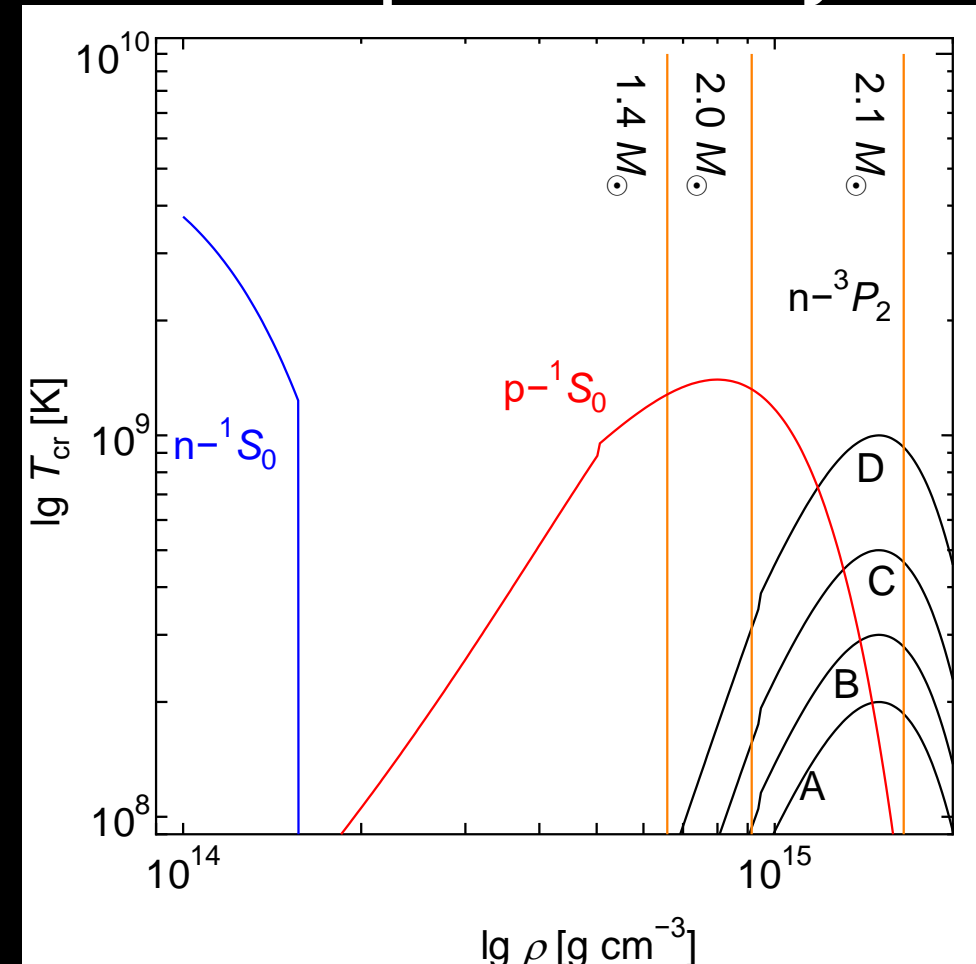
- Assumption: **Critical temperature of Neutron  ${}^3P_2$  is carried by d-quark's  ${}^3P_2$**

- No effects for proton  ${}^1S_0$
- $\Delta$  of 2SC / CFL are few tens of MeV
- No s-quarks



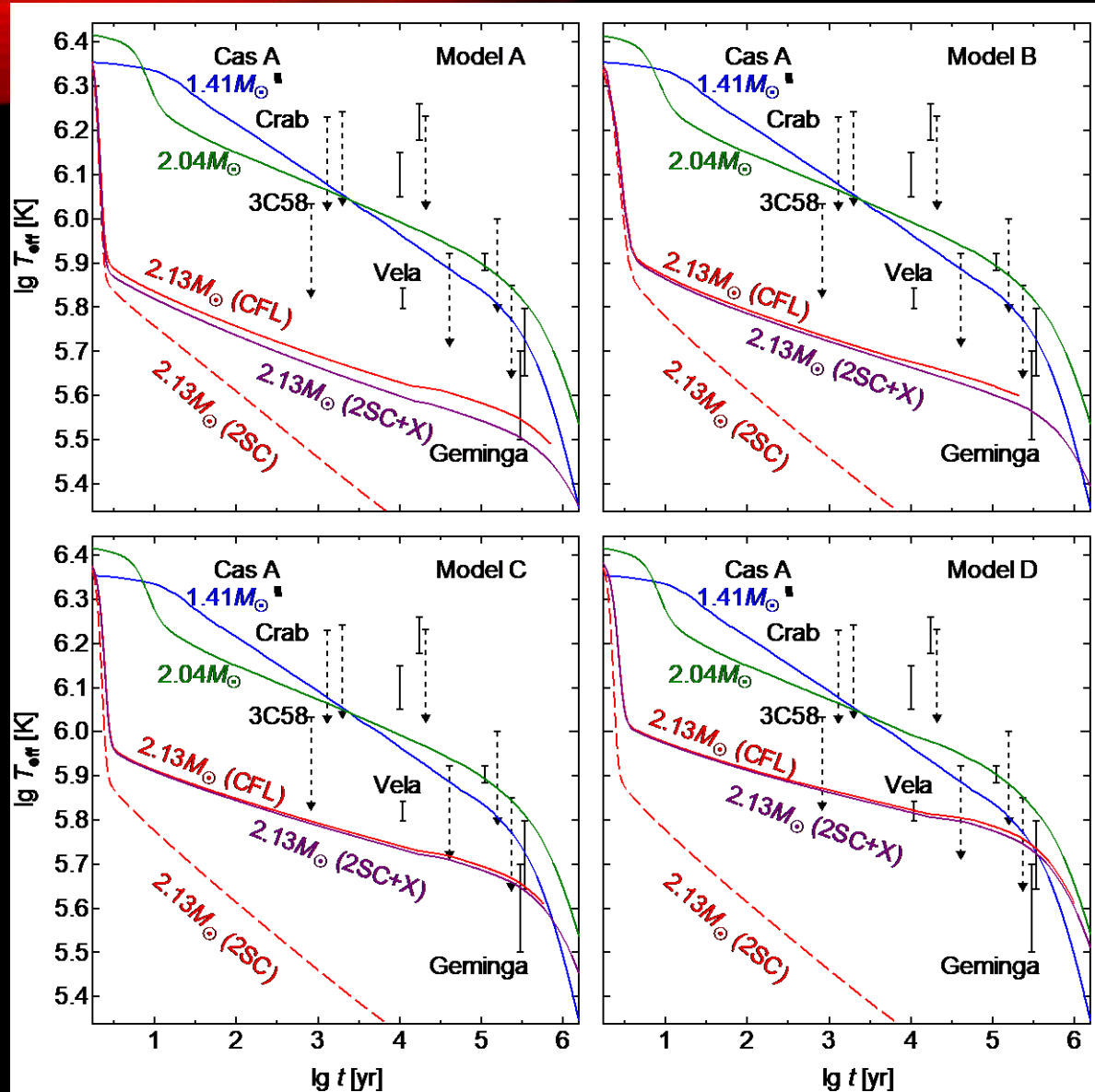
# Nucleon Superfluidity

- Neutrons and protons become superfluid  
Neutron:  ${}^1S_0, {}^3P_2$       Proton:  ${}^1S_0$
- Critical temperature ( $T_{cr}$ )
  - Functionated density dependence
- Effects on Cooling  
Superfluid transition: **Strong cooling (PBF)**  
(Page+ 2004)  
Superfluid state: **Suppresses other neutrino emission**
- n- ${}^3P_2$  critical temperature is continued by 2SC+X
- Calculating with changing the n, d- ${}^3P_2$  model

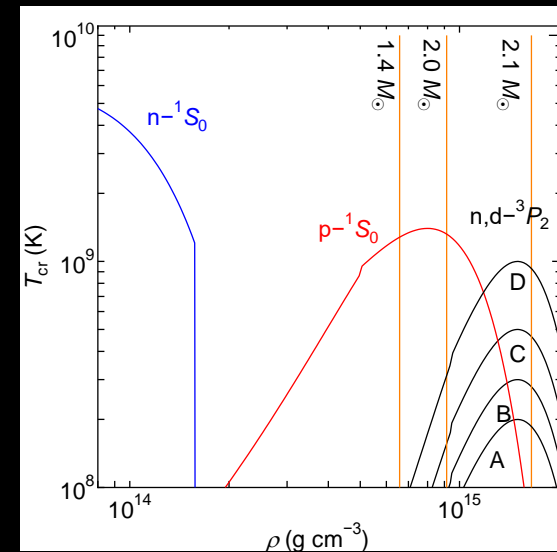


# Cooling Results

- Cooling behaviour depends on hadron superfluidity models
- Quark pairing affects cooling curves.



Checking superfluid cooling models with other superfluid phenomena.



# Summary of Neutron Star Cooling

- Rapid cooling in 2SC → Too cold **X**
  - Appearing density **✓**
- Marginal cooling with CFL (depending on  $n$ - ${}^3P_2$  critical temperature) → Observation **✓**
  - Appearing density does NOT match **X**
- **2SC+X is similar to CFL** (depending on  $n$ ,  $d$   ${}^3P_2$  critical temperature) → **Observation ✓**
  - Appearing density **✓**

- Show the effect on rotation following the method of Ho & Andersson (2012)

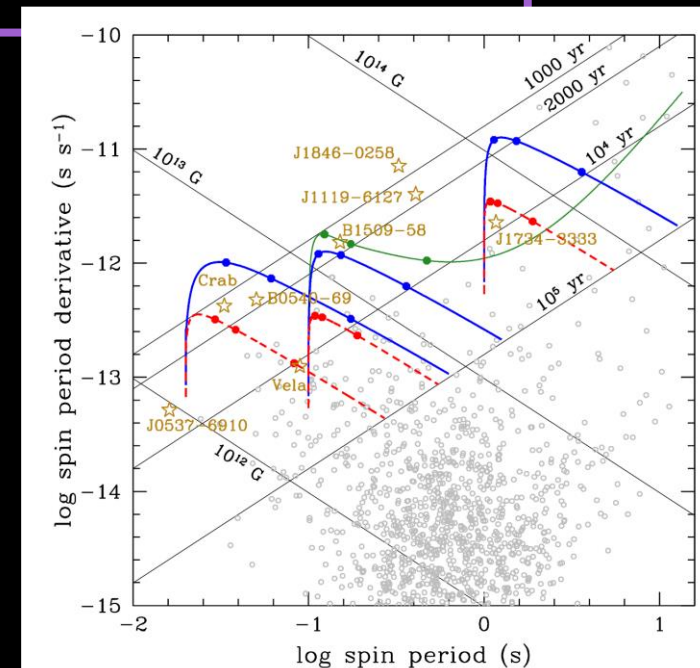
nature physics

LETTERS

PUBLISHED ONLINE: 30 SEPTEMBER 2012 | DOI: 10.1038/NPHYS2424

## Rotational evolution of young pulsars due to superfluid decoupling

Wynn C. G. Ho\* and Nils Andersson\*

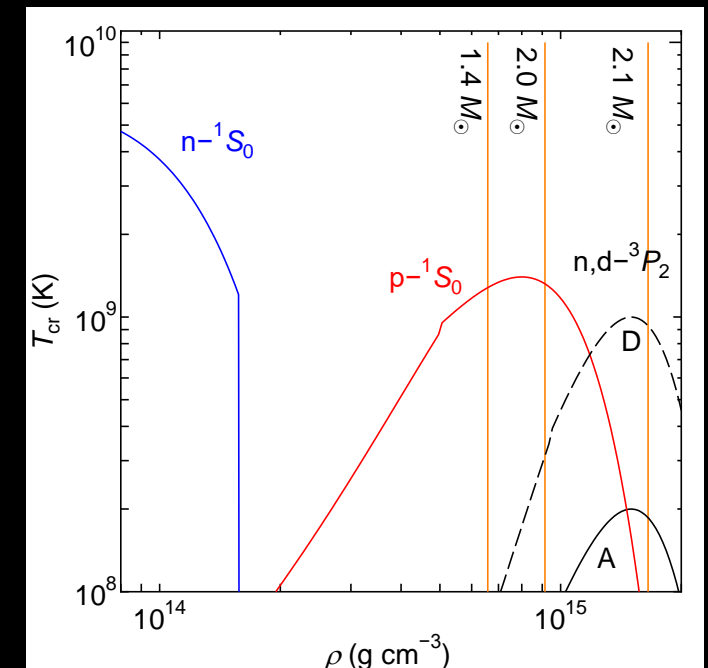


# Cooling Pulsar Evolution

1. As the neutron star cools, the region of superfluidity increases
2. Decreased moment of inertia in the normal state region
3. Spin-down due to magnetic dipole radiation changes with decreasing moment of inertia
4. Line in  $P - \dot{P}$  Diagram bends

- Settings

- Superfluid Model **A(Solid)** and **D(Dashed)**
- Superfluid state region does not affect rotation
- Friction is small and constant between super and normal
- Color superconducting regions are treated in the same way as nucleon superfluidity



# Variation of moment of inertia / rotation

## Moment of inertia

- General relativistic sphere moment of inertia
  - Exact relation: Ravenhall & Pethick (1994)

$$I = \frac{8\pi}{3} \int_0^R (\rho + P/c^2) \Lambda r^4 dr$$
$$\Lambda = \left(1 - \frac{2Gm}{c^2 r}\right)^{-1}$$

Calculate normal layer only

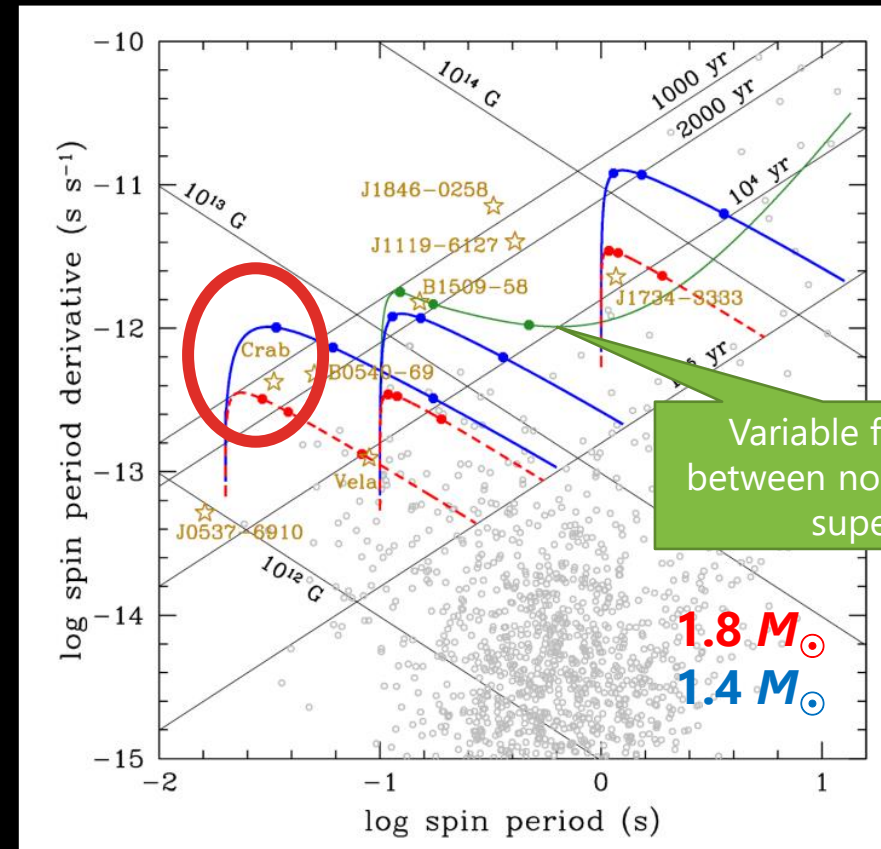
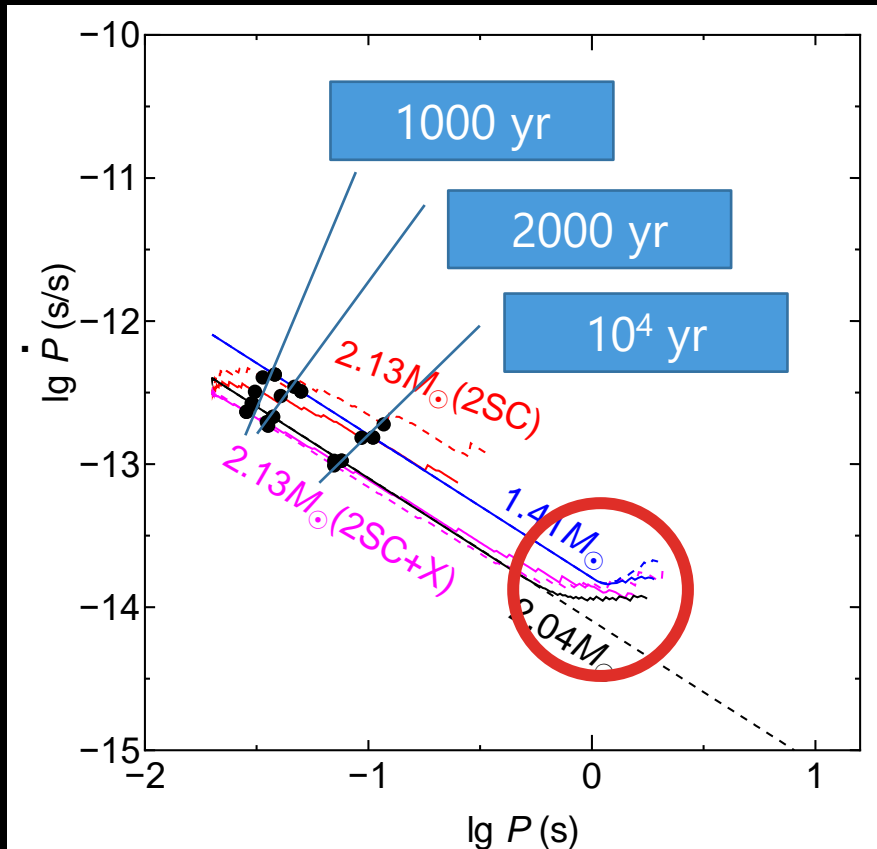
## Rotation

- Angular velocity variation considering magnetic dipole radiation

$$\frac{d\Omega}{dt} = (\Omega_{SF} - \Omega) \frac{1}{I} \frac{dI}{dt} - \frac{\beta \Omega^3}{I}$$
$$\beta \simeq B^2 R^6 / 6c^3$$
$$\Omega_{SF} - \Omega \leq 1.0 \times 10^{-6}$$

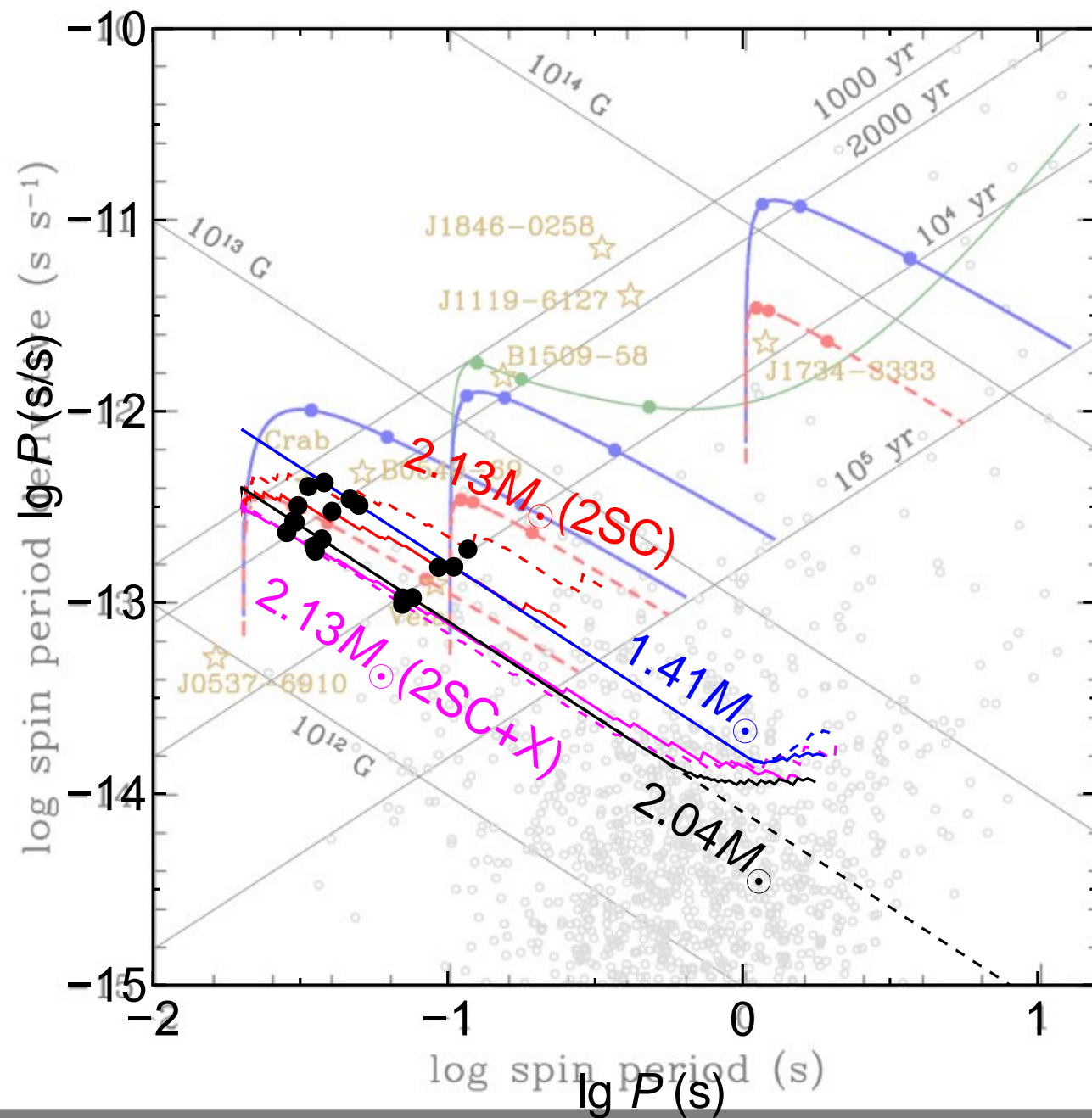
Parameters: Initial value of  $B$  &  $\Omega$   
 $(P, B) = (0.02 \text{ s}, 5 \times 10^{12} \text{ G})$

# $P - \dot{P}$ Diagram

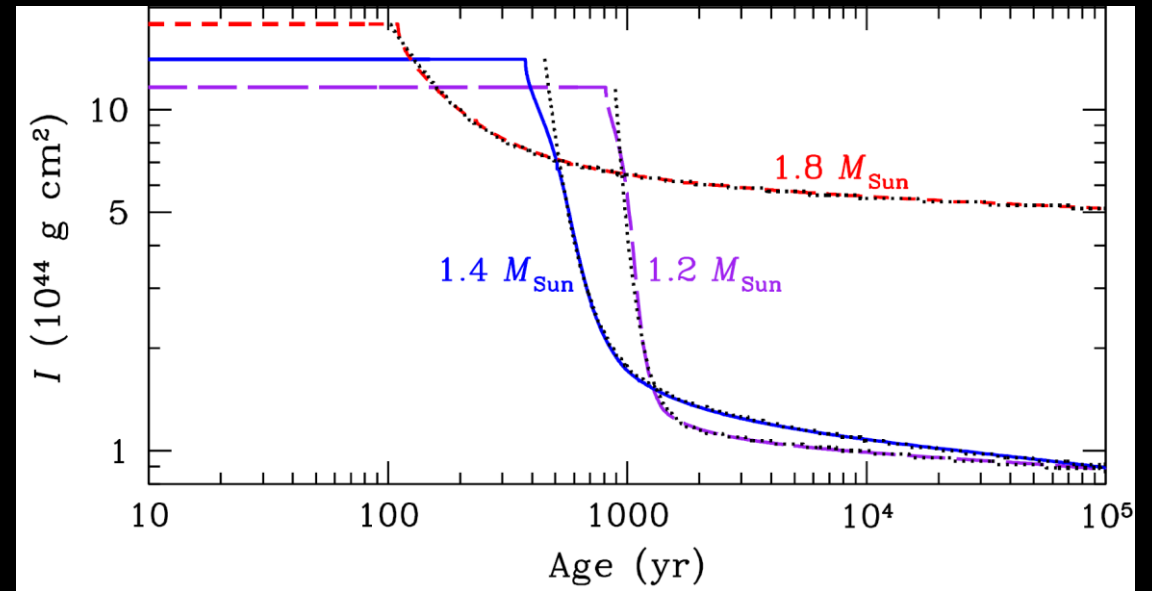
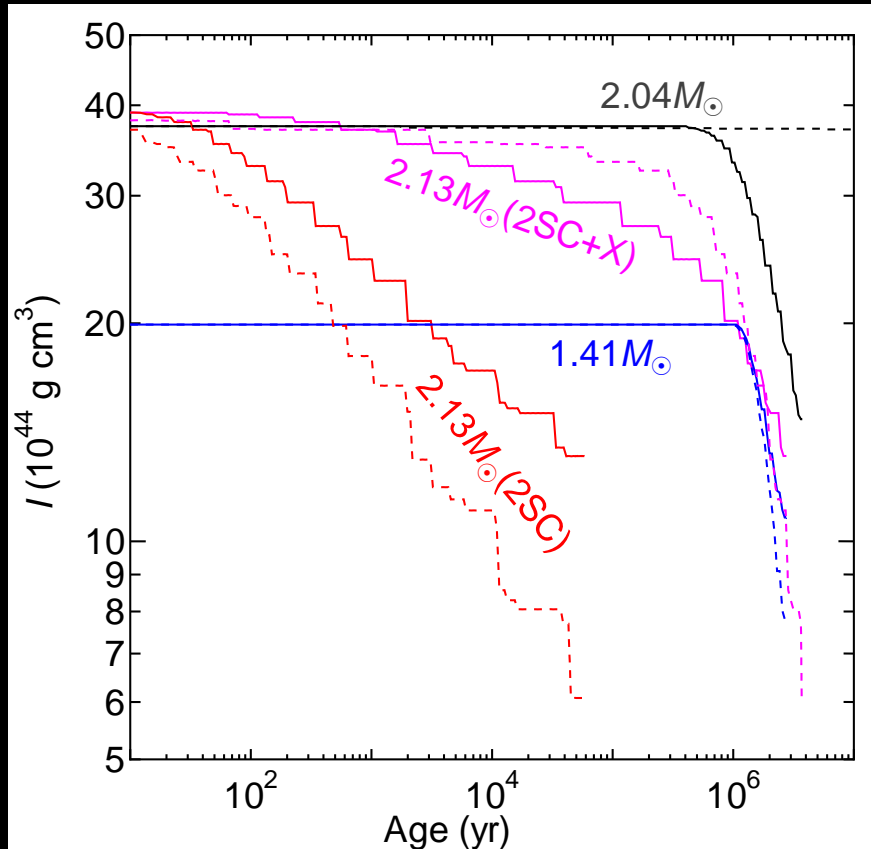


$P - \dot{P}$  Diagram differs to previous research (Ho&Andersson)

- Rise timing of  $\dot{P}$



# Time variation of moment of inertia



Evolutional tracks of the momentum inertia differ to the previous research (Ho&Andersson)



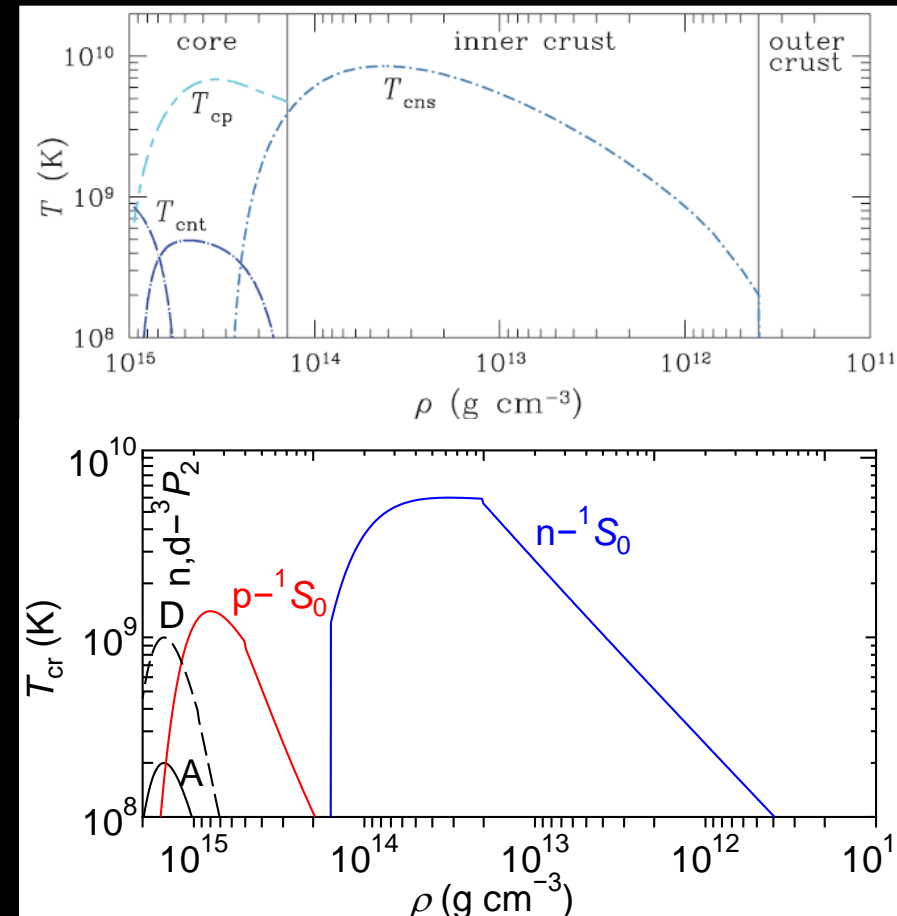
# The reason of the time variation of moment of inertia

Ho & Andersson (2012)

- Superfluid model with "shallow" component
- Outer layer of core: **Superfluid from the beginning**
  - **Moment of inertia decreases in early stage**
- Curves rise in  $P - \dot{P}$  Diagram in early stage

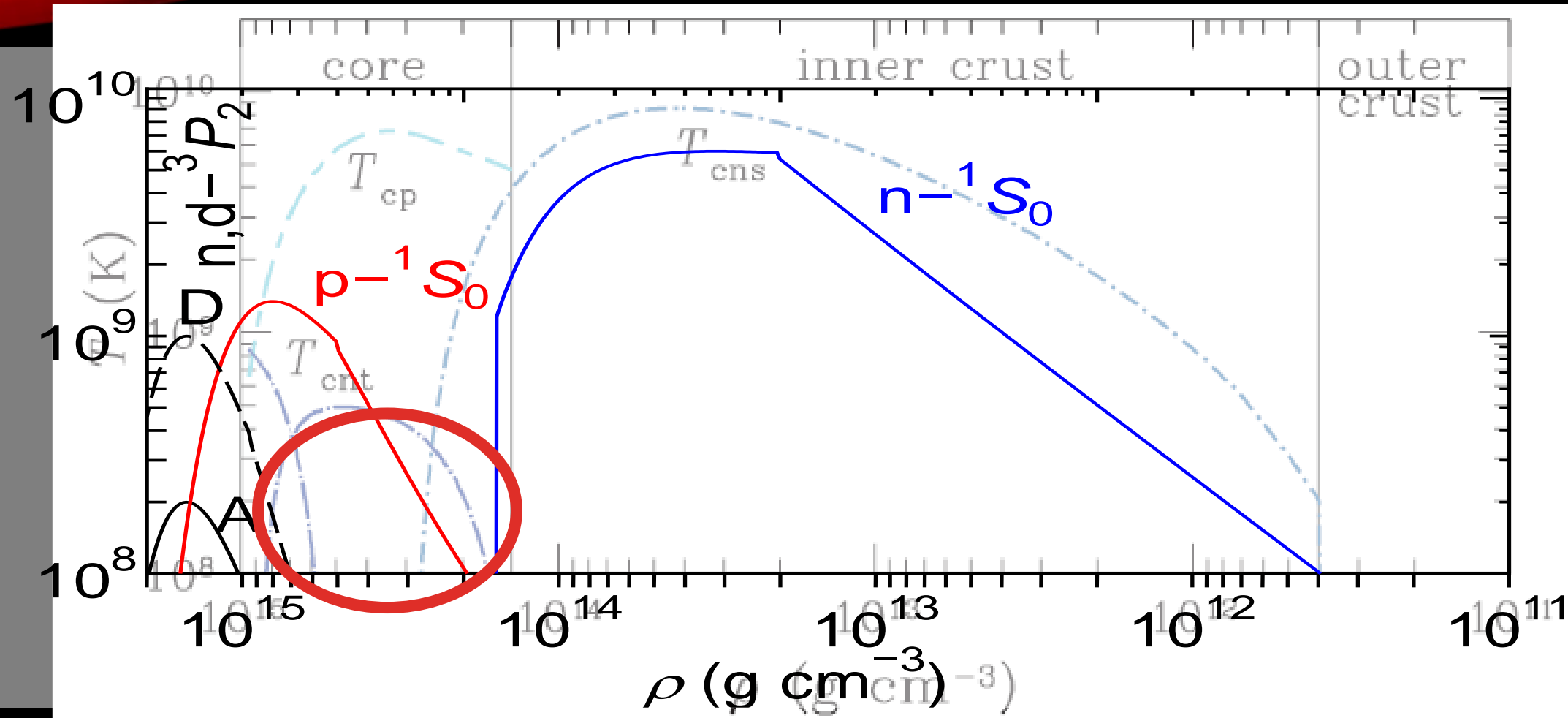
Ours

- Superfluid transition starts from the centre of the core
- Superfluid region increases by time
- When the superfluid region expands to a certain level, the moment of inertia is greatly reduced.
- Curves rise in  $P - \dot{P}$  Diagram in late stage



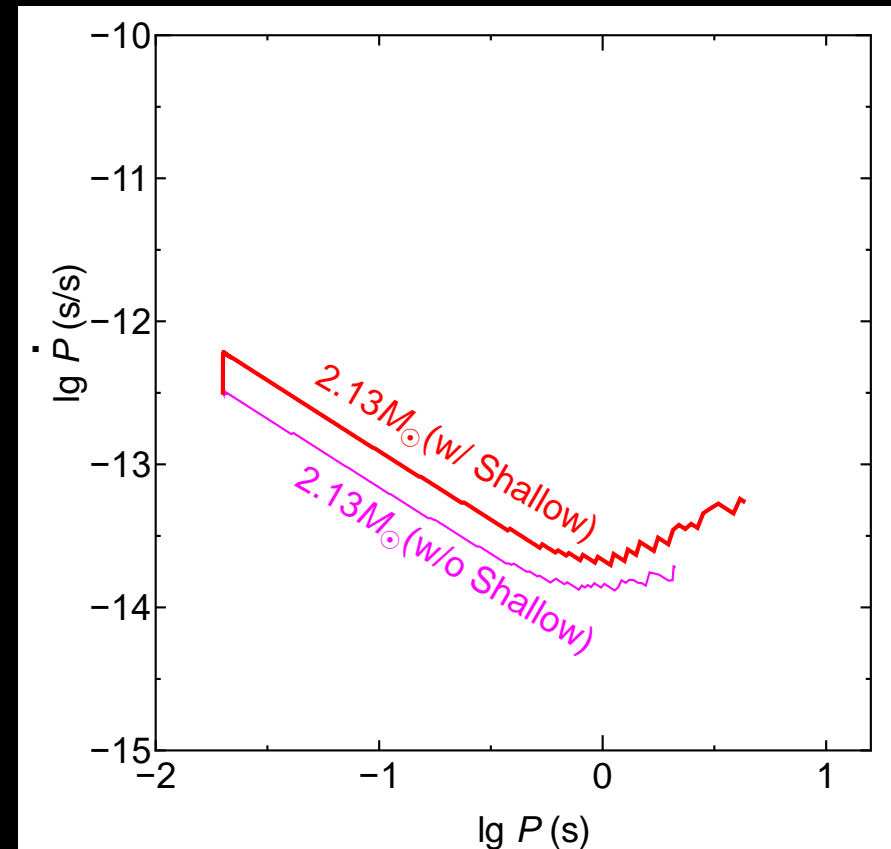
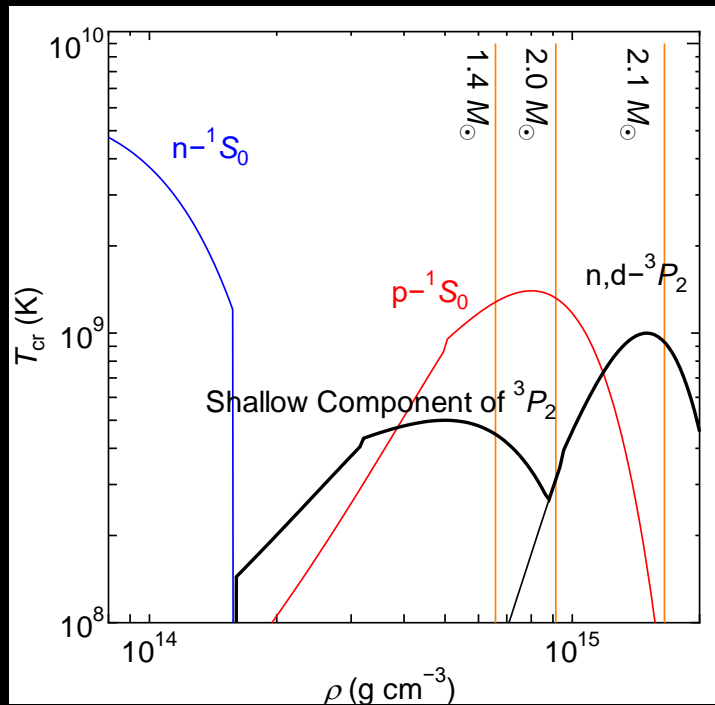
Critical temperatures slightly outside the centre are effective on moment of inertia

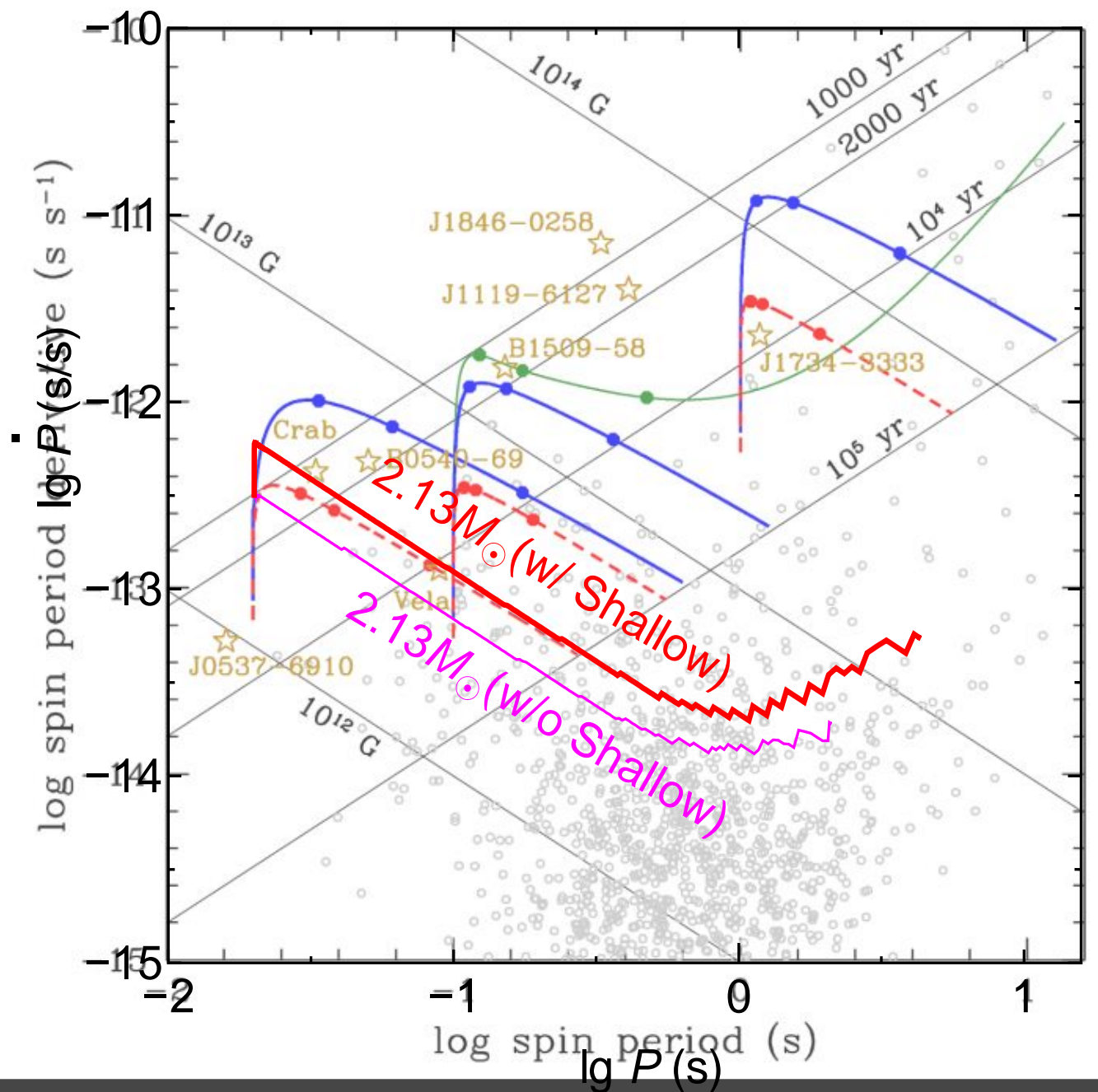
$T_{cr} (K)$



# Shallow Component of ${}^3P_2$

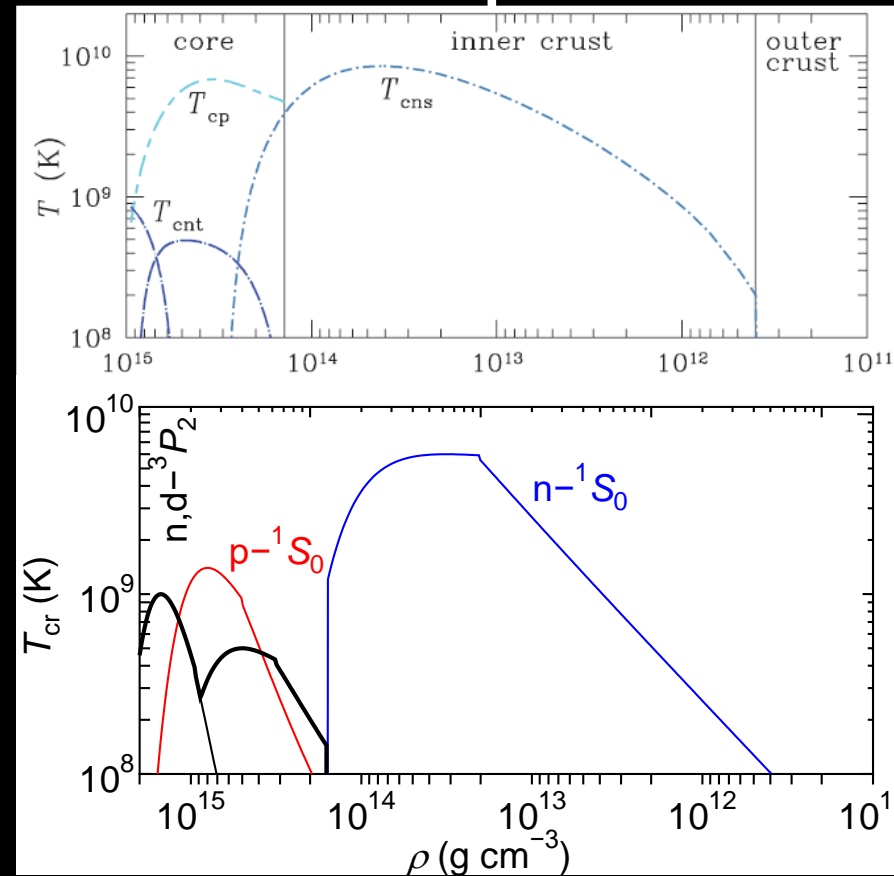
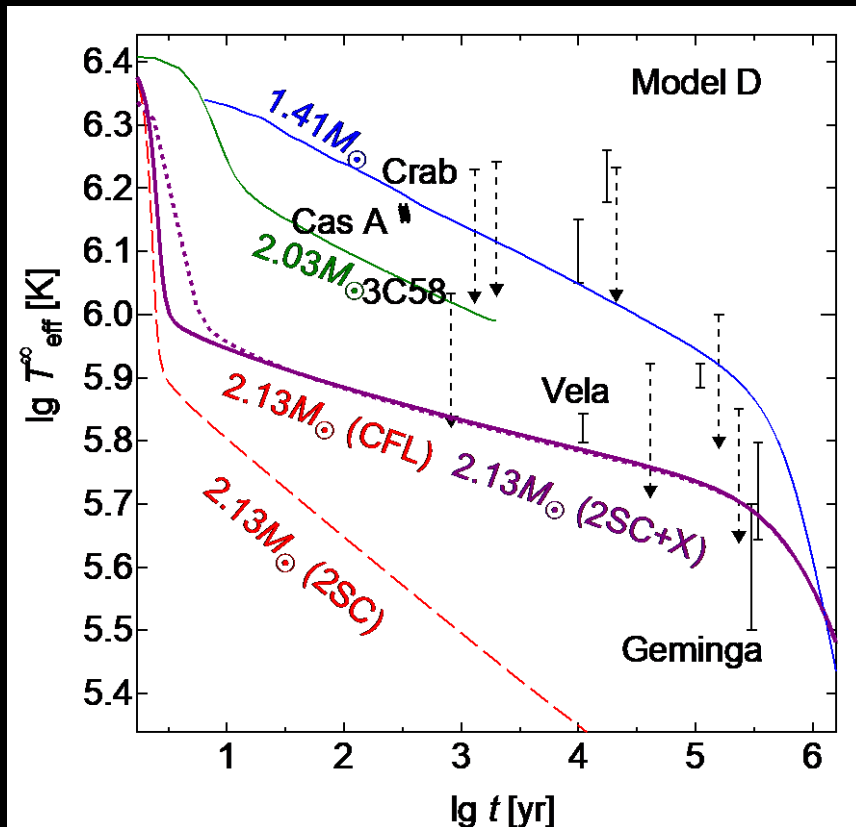
- Including "Shallow Component" to  ${}^3P_2$  superfluidity
- Rise of  $P - \dot{P}$  curve can be replicated





# Shallow Component of ${}^3P_2$

- With Shallow component,  $P - \dot{P}$  curve rises at early stage.
- Cooling curves also change



# Summary of Pulsar Evolution

- Calculate **the pulsar evolution** based on the cooling calculation of neutron stars
  - Magnetic dipole radiation
  - Time variation of Momentum of Inertia by superfluid transition  $\Rightarrow$   **$P - \dot{P}$  Diagram**
- Results of current model slightly differ from Ho & Andersson (2012)
  - $P - \dot{P}$  Diagram: Variation timing of  $\dot{P}$
  - Momentum of Inertia: Decrease timing of the momentum of inertia  $\rightarrow P - \dot{P}$  Diagram
  - **Difference of superfluid model**
- Variation timing of the momentum of inertia of entire star
  - Depends on neutron  ${}^3\text{P}_2$  superfluid critical temperature in outer layer, not centre
  - **Observation of rise timing of  $P - \dot{P}$  Diagram to upper right may constrain the "tail" of the critical temperature of neutron  ${}^3\text{P}_2$  ?**