

1D GRPIC Simulations of High-Energy Gamma Rays from BH Magnetospheres

一般相対論的プラズマ粒子シミュレーションによる
ブラックホール磁気圏強電場領域由来
ガンマ線放射の研究

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TOHOKU
UNIVERSITY

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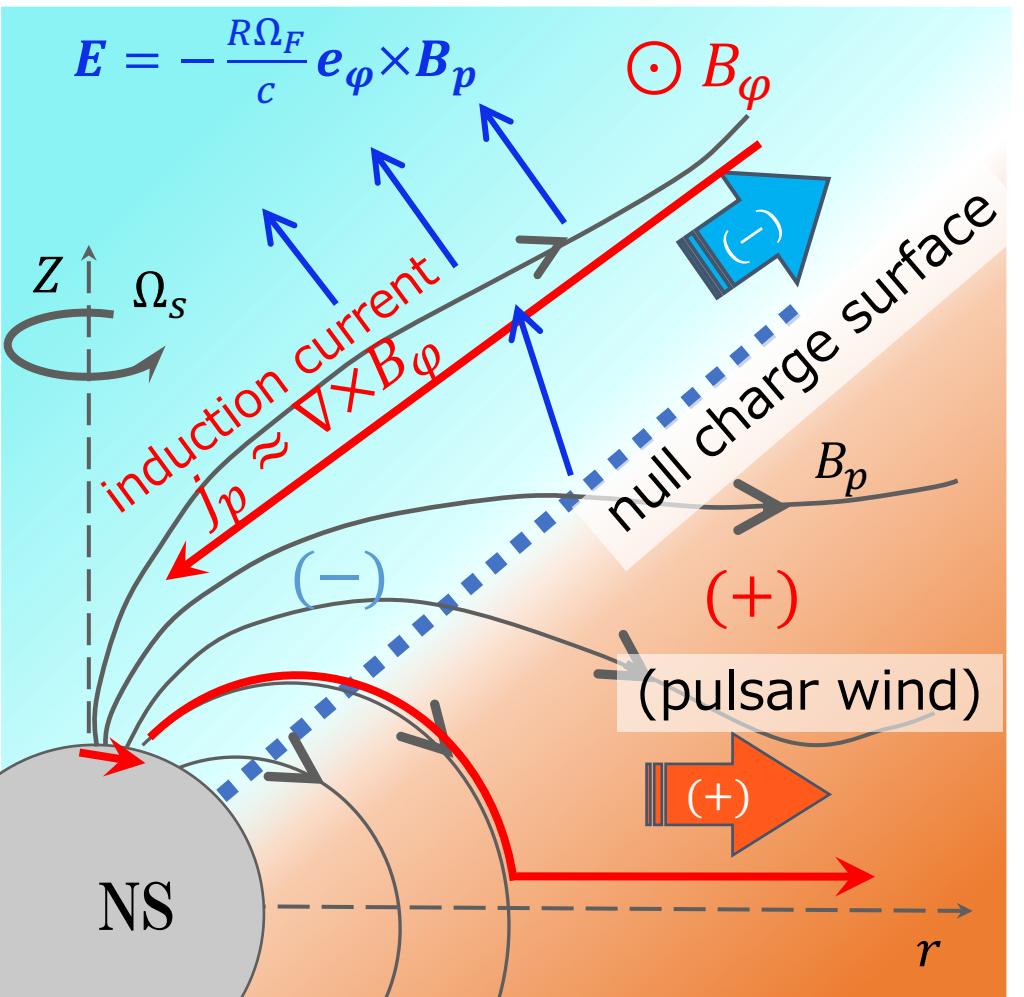
Theoretical Astrophysics
Tohoku University

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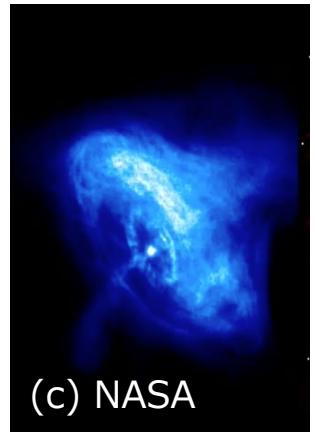
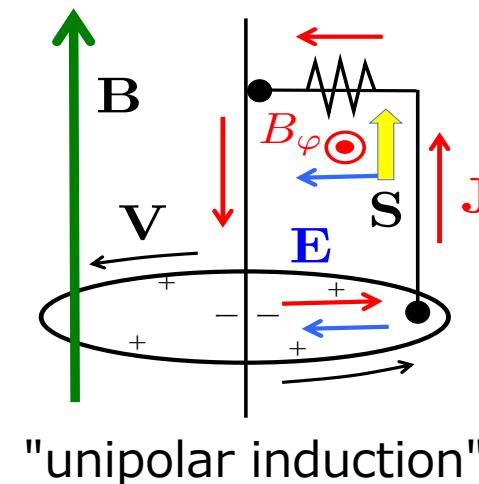
Current Structure • Charge Distribution of Pulsar Magnetospheres

Goldreich & Julian 1969



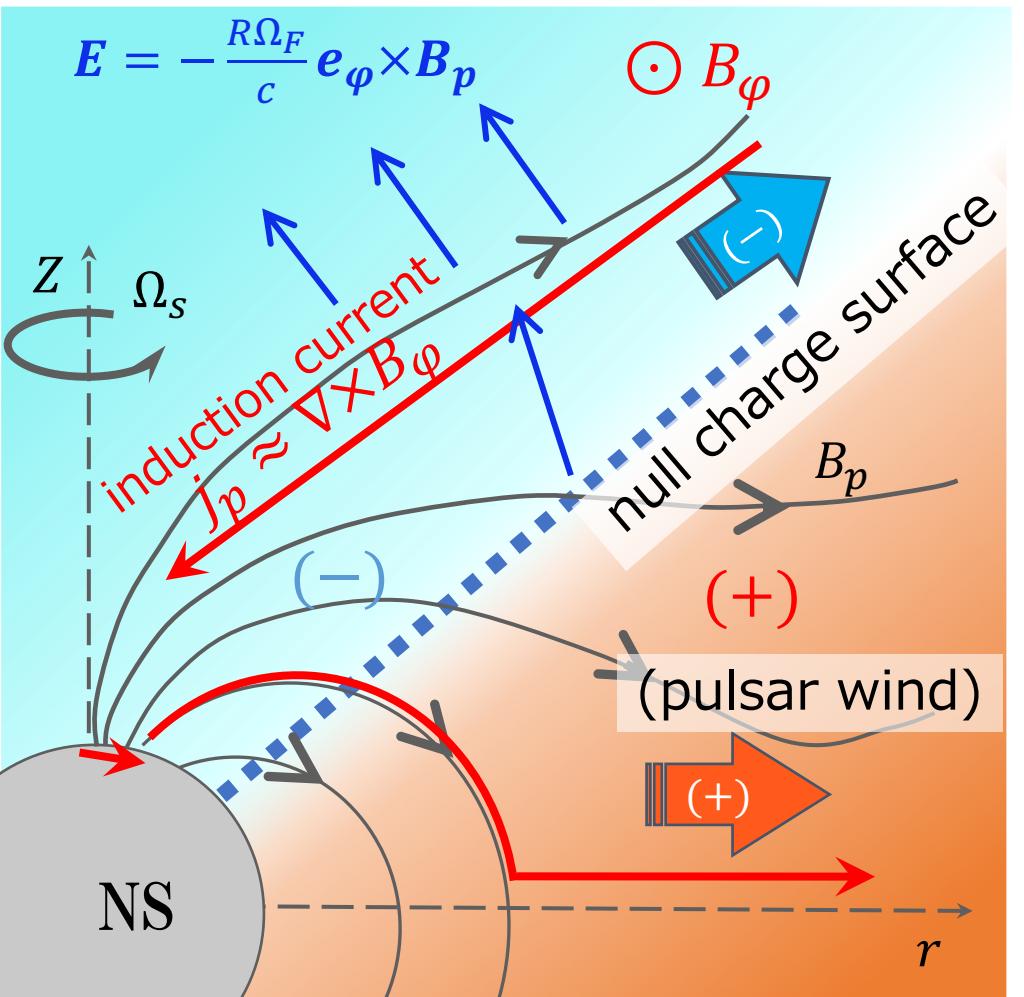
- **plasma injection** from NS / via $\gamma B \cdot \gamma \gamma$ interactions
- **co-rotation** of NS, axisymmetrical B-field, plasma ($\Omega_s = \Omega_F$)
NS rotation energy → **steady plasma flow**
maintaining j_p

$$S_p = \frac{1}{4\pi} E \times B_\phi = - \frac{1}{4\pi} R\Omega_F B_\phi B_p \quad \text{"pulsar wind"}$$



Current Structure • Charge Distribution of Pulsar Magnetospheres

Goldreich & Julian 1969



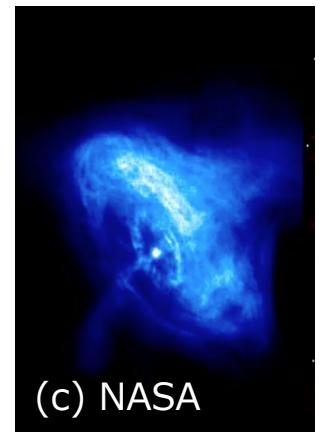
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- Goldreich-Julian charge distribution

$$\rho_e = \frac{1}{4\pi} \nabla \cdot \mathbf{E} \approx - \frac{1}{2\pi c} \Omega_s \cdot \mathbf{B}$$

→ cone-like null charge surface
(where $\Omega_s \cdot \mathbf{B}$ vanishes)

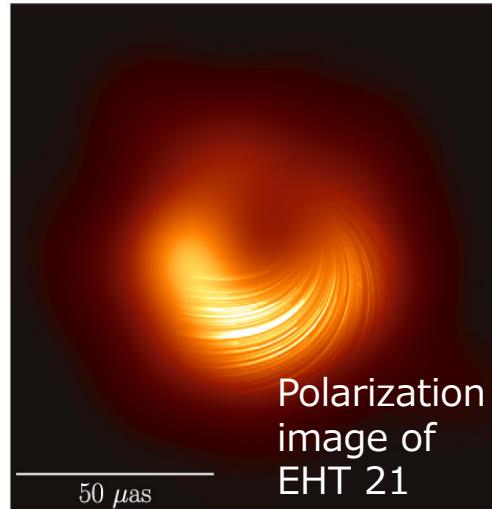


Formation of BH magnetospheres

<B-fields transportation>

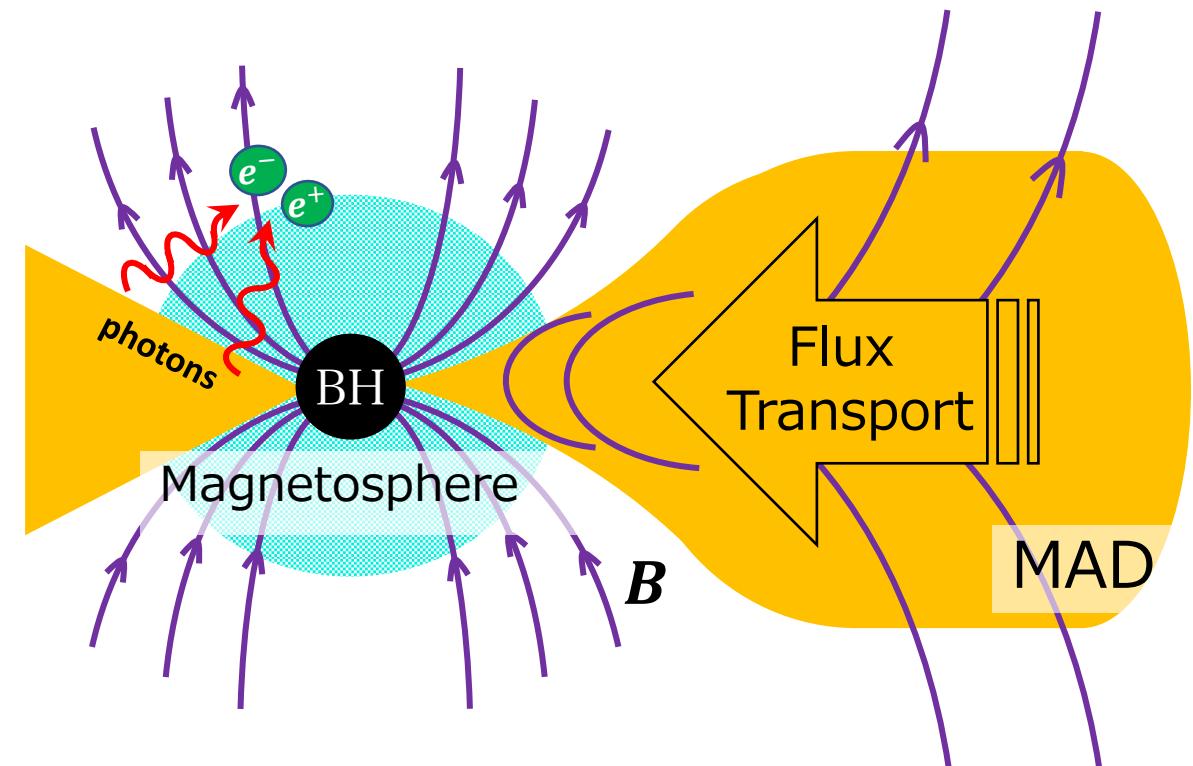
Theory:
 infalling gas bring magnetic flux
 → highly-magnetized gas disk around BHs
(Magnetically Arrested Disks, MADs)

EHT observations:
 highly-magnetized, poloidal B structure
 are expected at M87 vicinity



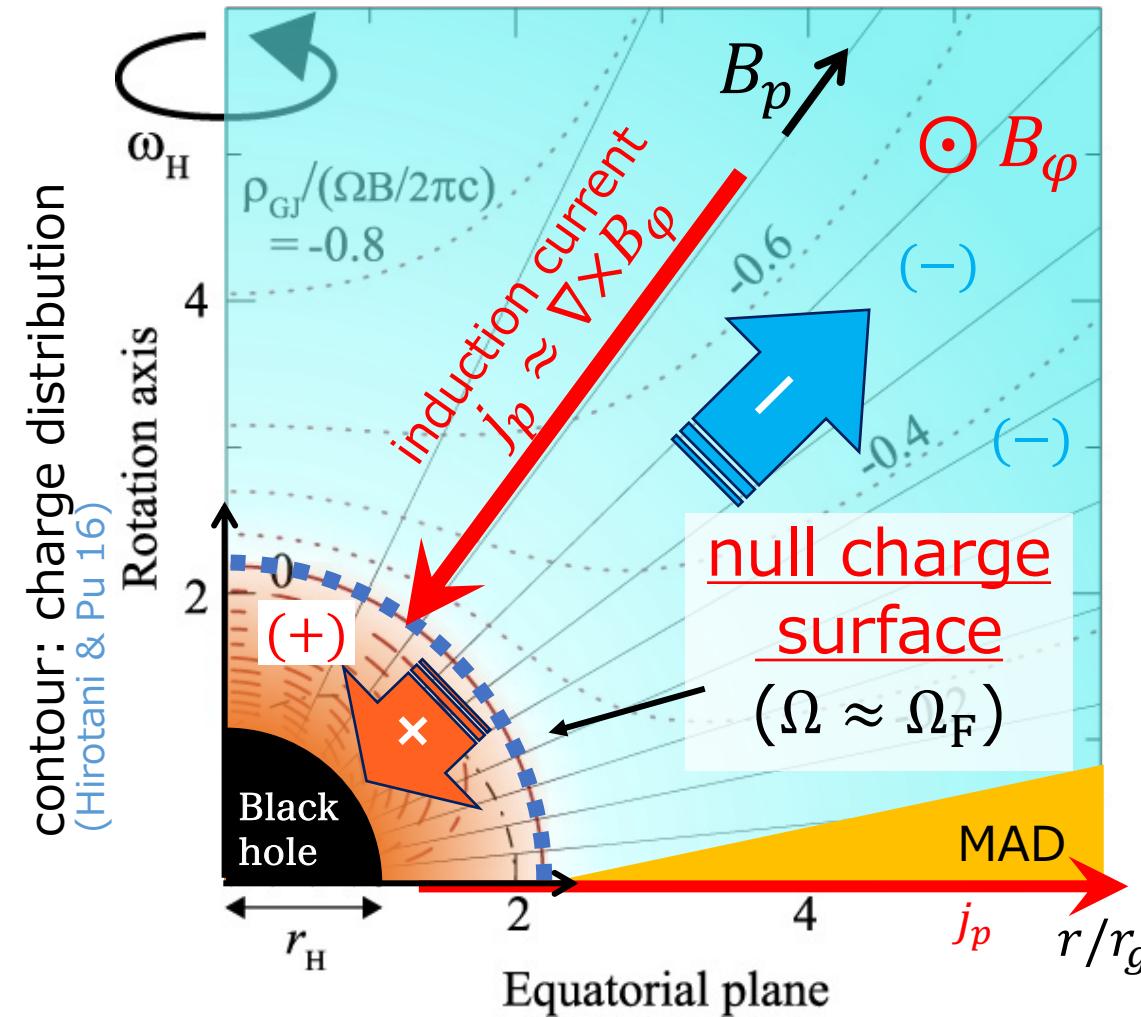
<Plasma injection>

Theory:
 main source = $\gamma\gamma$ interactions of disk photons



Current Structure • Charge Distribution of BH Magnetosphere

- ◎ sufficient plasma → steady EM structure (**BZ process**) (e.g. Blandford & Znajek 77)



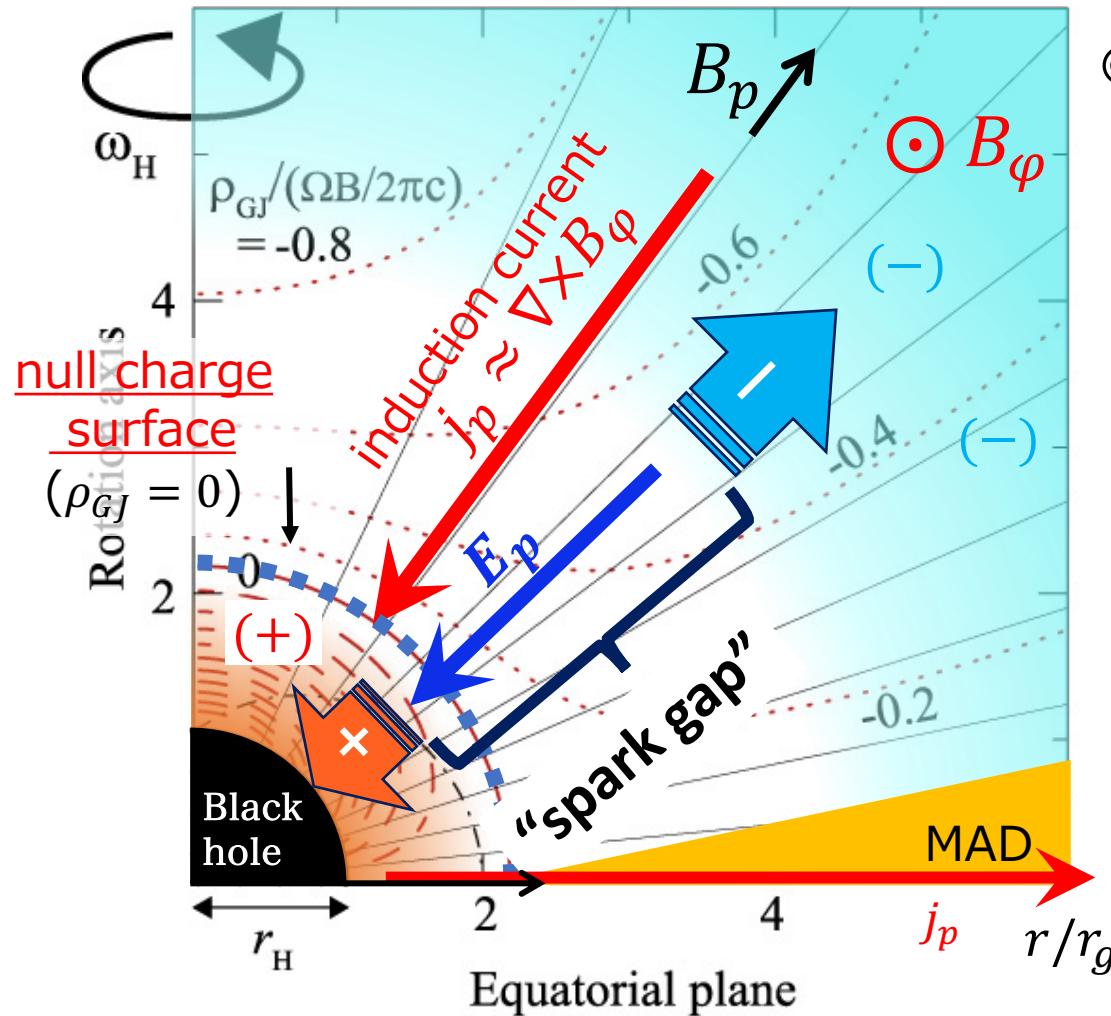
- Difference
- OPEN B-field structure
- $\Omega(r, \theta) > \Omega_F$ in nearby, $\Omega < \Omega_F$ at far
- ◎ **far zone:** negatively charged outflow maintains the stationary current (connected to the jet)
 - ◎ **near horizon:** positively charged inflow (consistent w/ infall due to the gravity, rapid rotation)



almost spherical null charge surface ($r_{null} \sim 2r_g$)

Formation of “Spark Gap” in Charge-Starved BH Magnetospheres

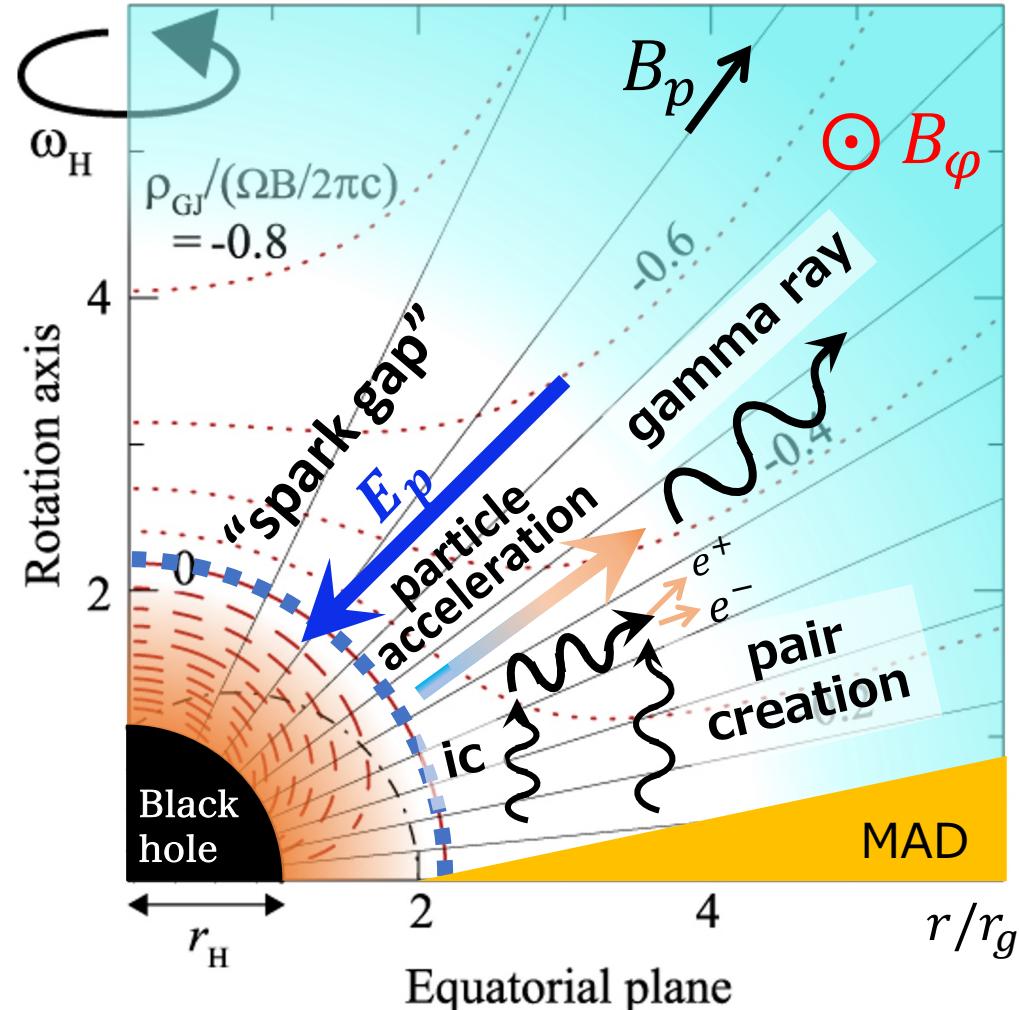
◎ **charge starved** → time-dependent E-field (analogical to pulsar polar cap/outer gap)



- ◎ charge starved due to the low plasma injection (Levinson & Rieger 11; Levinson & Segev 17; Hirotani & Pu 16 etc...)
- local charge deficiency around null surface $n < n_{GJ} (= |\rho_{GJ}|/e)$
- displacement current develops, maintaining electric current $\partial_t(E_p) \approx -4\pi(j^r - J_0/r^2)$
- local, intermittent E-field region (“**spark gap**”)

Formation of “Spark Gap” in Charge-Starved BH Magnetospheres

◎ charge starved → time-dependent E-field



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$$\partial_t(E_p) \approx -4\pi(j^r - J_0/r^2)$$

local, intermittent E-field region (“spark gap”)



efficient acceleration, gamma-ray emission
 secondary pair creation

Examining BH Spark Gap Natures...

- ◎ ... by **semi-analytic modeling** (Beskin et al. 92; Hirotani & Pu 16; Levinson & Segev 17 etc...)

steady gap approximations (likewise pulsar outer gap model Cheng et al. 1986a,b)

- ◎ ... by **Particle-In-Cell simulations (mainly for SMBHs)**

introducing time-dependence

- 1D local model

(Levinson & Cerutti 18; Chen et al. 20; Kisaka et al. 20, 22)

fixed global B-field structure

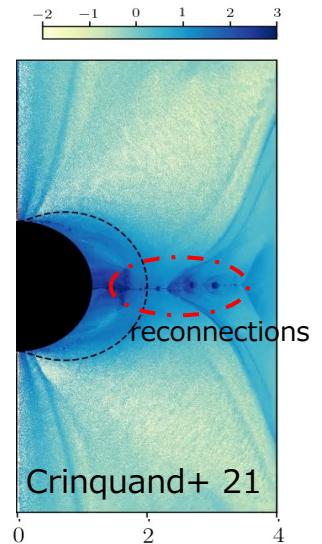
solving E-field & plasma evolutions



- 2D global model

(Parfrey et al. 19; Crinquand et al. 20,21;
Hirotani et al. 22,23; Niv et al. 23)

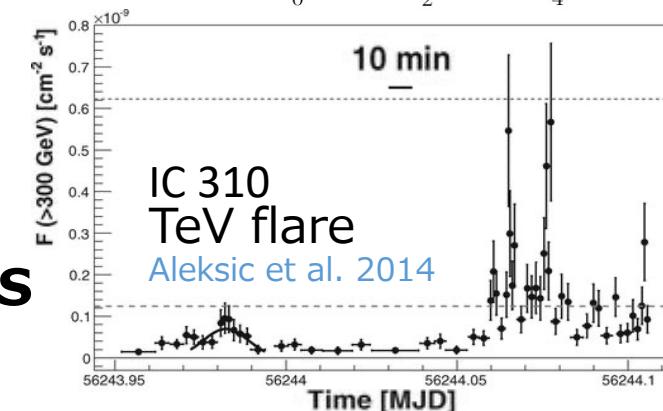
considering time-dependent B-field



escaping gamma rays for SMBH magnetospheres

= detectable!

possible connection w/ TeV flares from known AGNs



My Current Research:

High-energy Gamma Rays from **Stellar-Mass BH Magnetospheres** in 1D GRPIC Simulations

(Kin et al. 23 in prep.)

Motivation: Detecting Isolated stellar-mass Black Holes via Gamma-Rays?

- ◎ **$\sim 10^8$ undetected IBHs in the Galaxy**

$$SFR \times V_{gal} \times t_{galaxy}$$

$$\sim 10^{-14} \text{ pc}^{-3} \text{ yr}^{-1} \quad \sim 10^{11} \text{ pc}^3 \quad \sim 10 \text{ Gyr}$$

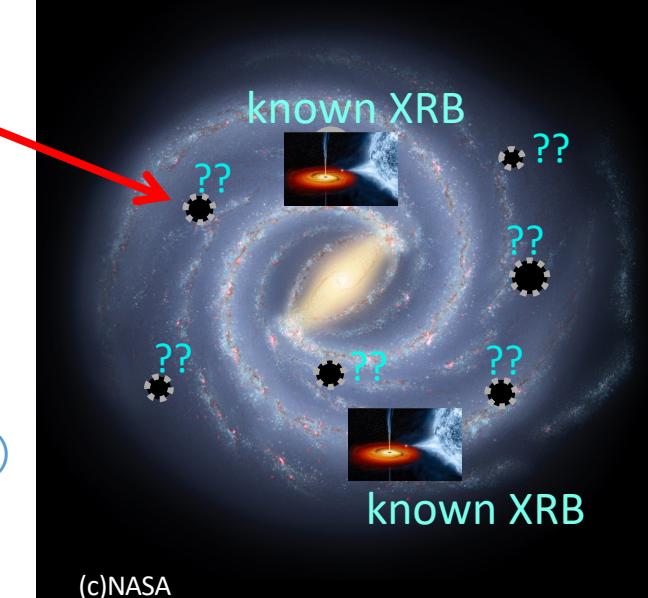
for $10M_\odot$

→ possible interactions w/ ISM clouds

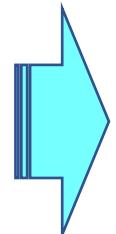
- ◎ **MAD formation around IBHs** (e.g. Ioka et al. 17; Kimura et al. 21)

efficient magnetic flux transportation

→ become MAD for low- β ISM accretion



(c)NASA



Gamma rays from IBH magnetospheric "spark gap" can be detected?

Simulation Setting

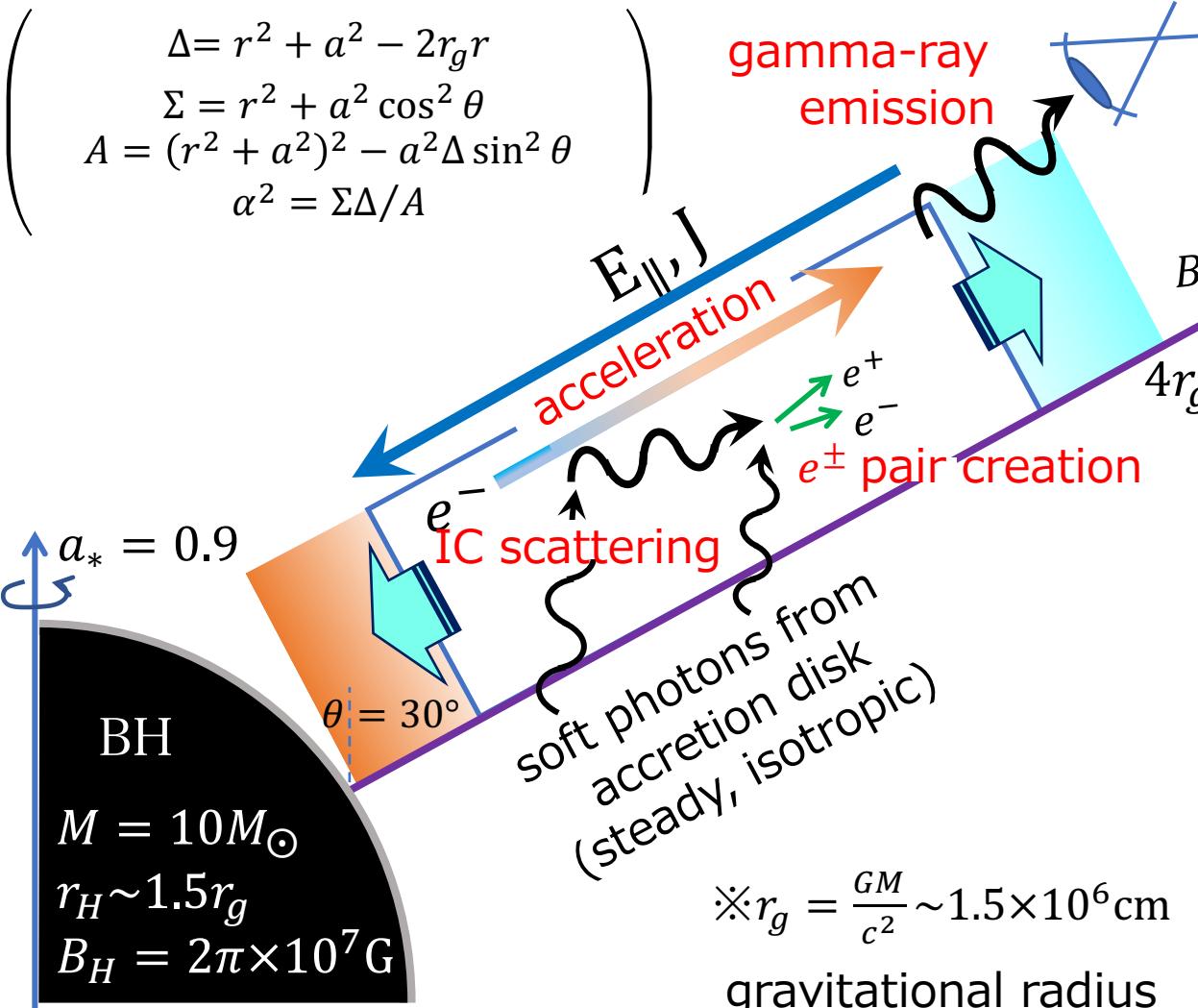
Metric term

$$\Delta = r^2 + a^2 - 2r_g r$$

$$\Sigma = r^2 + a^2 \cos^2 \theta$$

$$A = (r^2 + a^2)^2 - a^2 \Delta \sin^2 \theta$$

$$\alpha^2 = \Sigma \Delta / A$$



Particle-In-Cell

◎ 1D • GRPIC simulation code

(Levinson & Cerutti18; Kisaka et al.20;22)

- considering General Relativistic effect
- solve particle motion w/ real mass

$$\frac{du_\pm}{dt} = -\sqrt{g_{rr}} \gamma_\pm \partial_r(\alpha) + \alpha \left(\frac{q_\pm}{m_e} E_r - \frac{P}{m_e v_\pm} \right) : e^\pm \text{ EoM}$$

gravity (inertia term) acceleration back reaction of radiation

$$\frac{dp^r}{dt} = -\sqrt{g^{rr}} p^t \partial_r(\alpha) : \text{ic photon trajectory}$$

→ evaluate the charge/current distribution,
solve the Maxwell's eq. at each grid

$$\partial_t (\sqrt{A} E_r) = -4\pi (\Sigma j^r - J_0) : \text{Ampere's law}$$

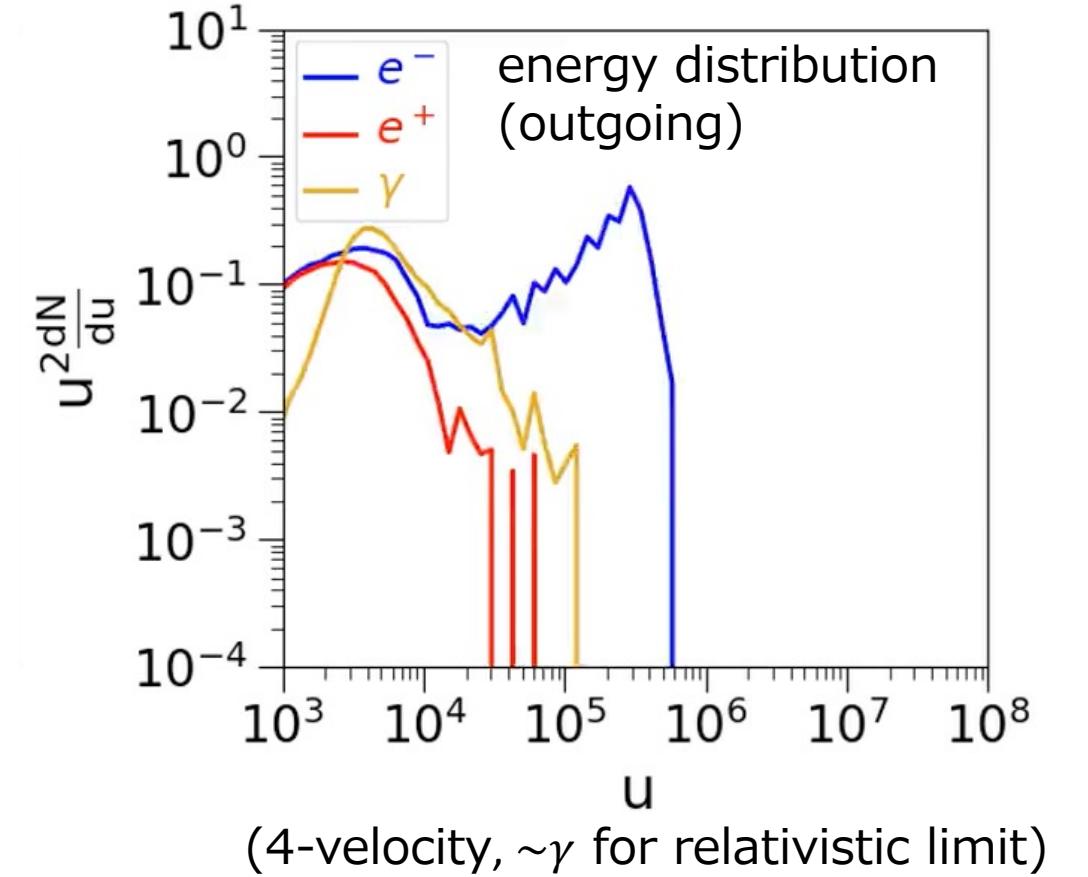
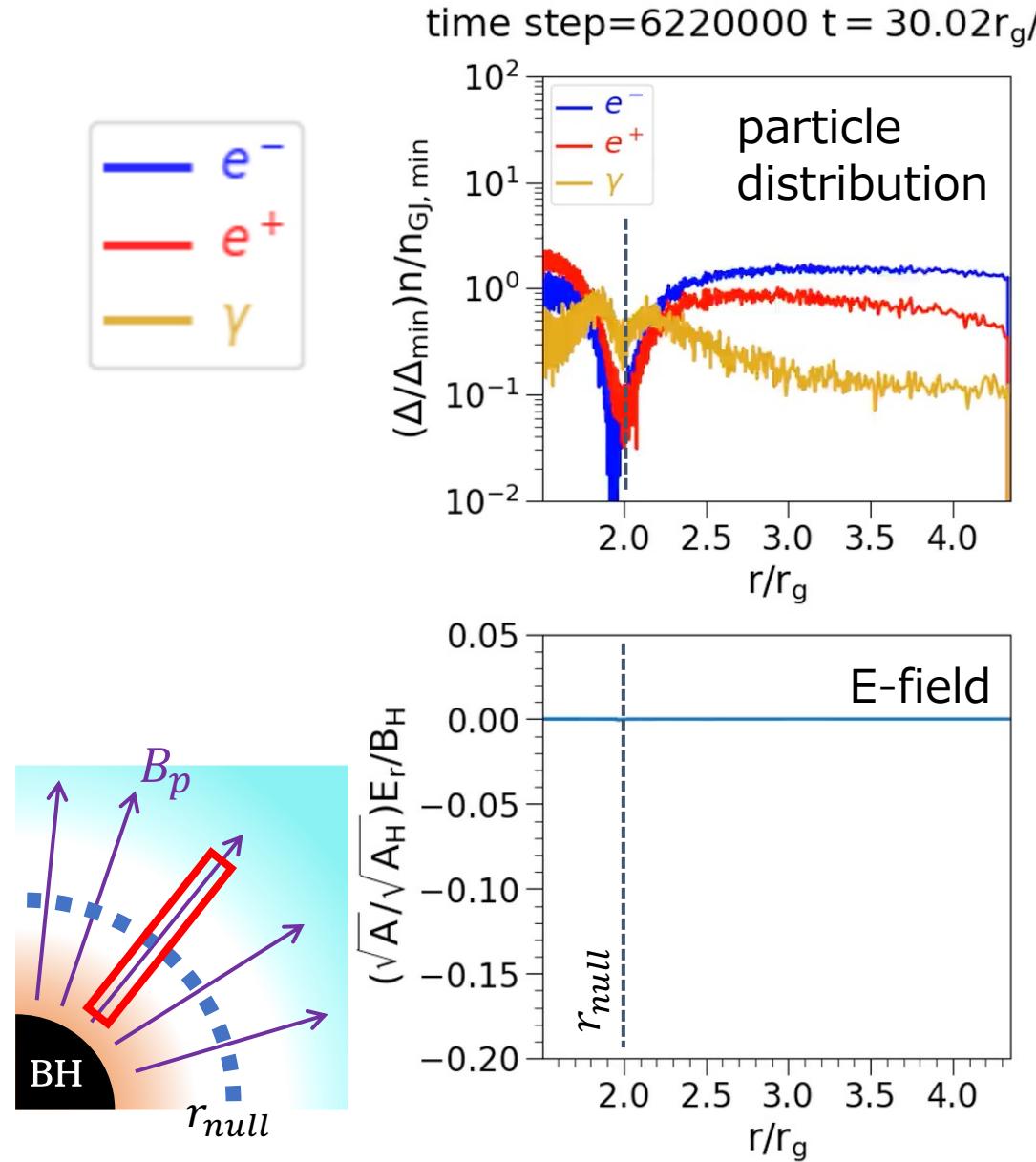
$$\partial_r (\sqrt{A} E_r) = 4\pi \Sigma (j^t - \rho_{GJ}) : \text{Gauss' law}$$

- include IC • pair creation interactions

◎ Kerr spacetime

◎ steady, axisymmetric B-field (split-monopole)

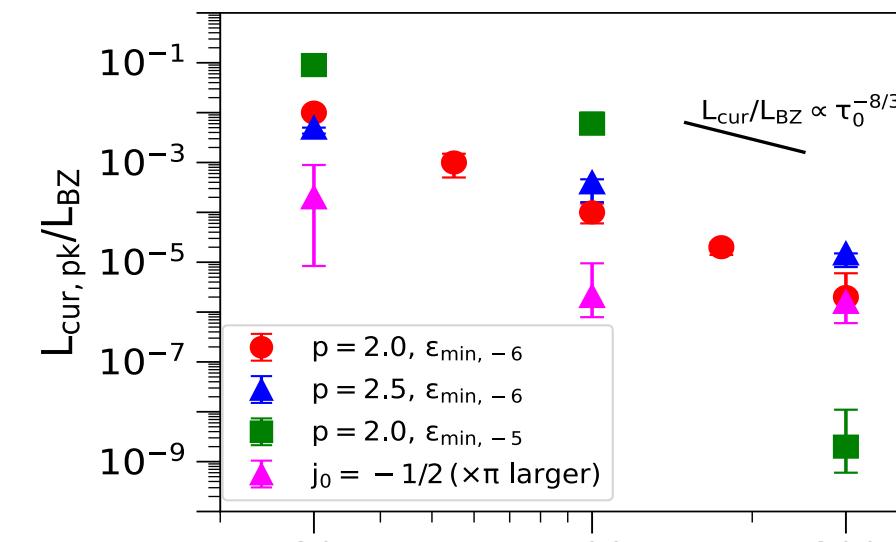
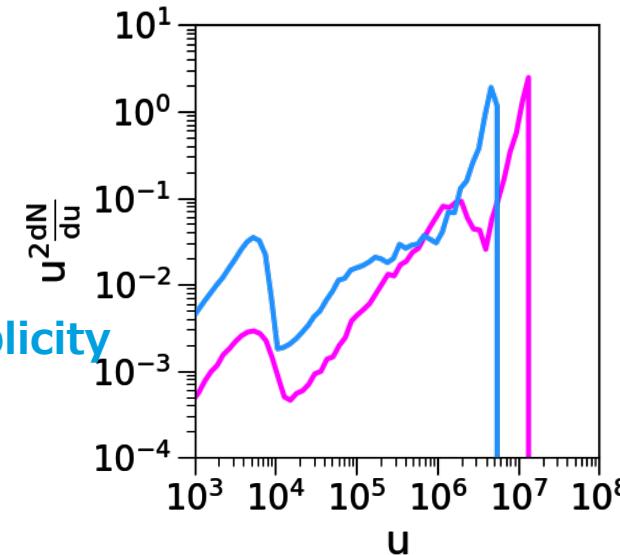
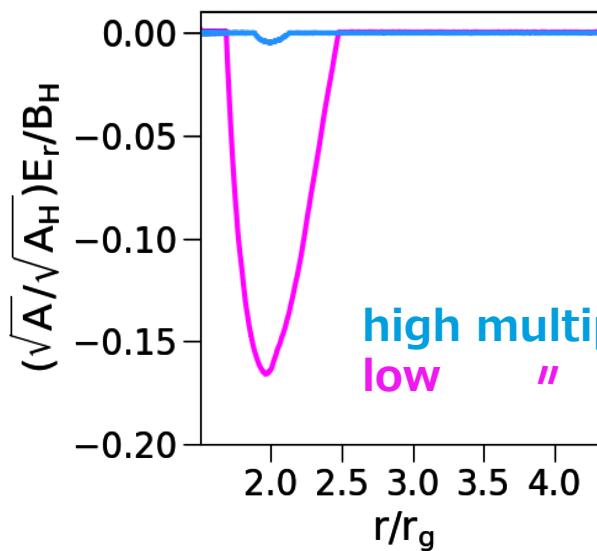
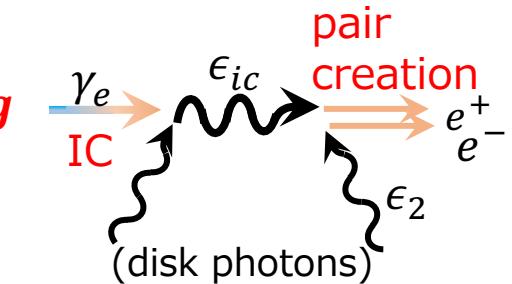
Simulation Result: Overall Evolution



- ◎ **quasi-periodic gap**
- ◎ $\gamma_{e^-, pk} \sim 10^7$
→ **GeV-TeV gamma-rays**

Simulation Result: Disk Photon Intensity Dependence

$\tau_0 \approx n_\gamma \sigma_T r_g \propto$ disk photon intensity : Thomson depth for r_g
 controlling e^\pm pair multiplicity ($\tau_{pair} \sim 0.1 \tau_0 \times (\epsilon_{ic} \epsilon_2)^{-1}$)



◎ multiplicity affects gap dynamics (period, gap size)

◎ $L_{cur,pk}/L_{BZ} \sim 10^{-2} (\tau_0/30)^{-\alpha}$

($\alpha \sim 8/3$ for $p = 2.0, \epsilon_{min} = 10^{-6}$)

higher intensity = higher multiplicity

L_{BZ} : BZ luminosity

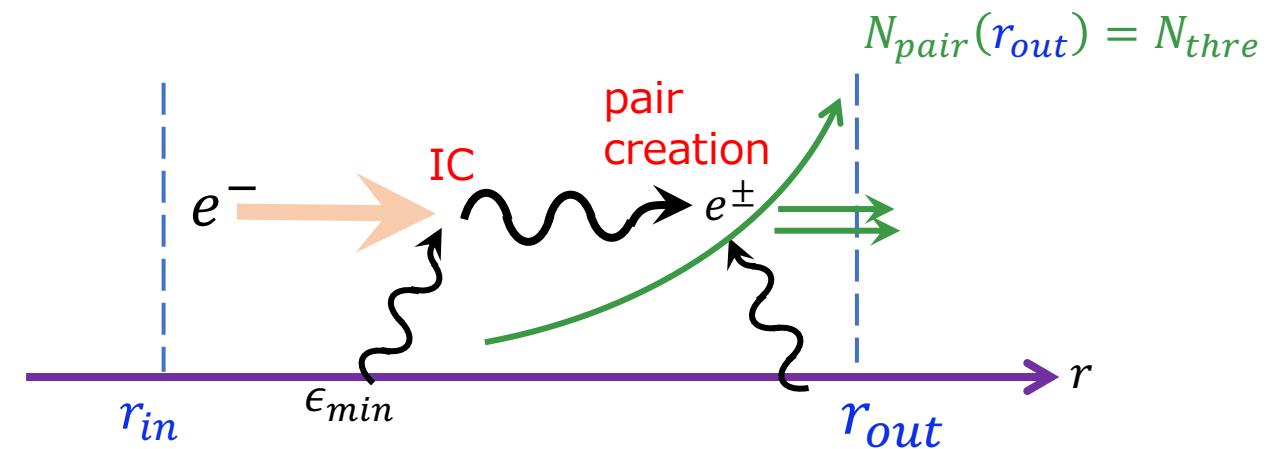
Semi-Analytic Model of Gamma-Ray Emission from Gap

◎predicting gamma-ray emissivity **for wide range of BH mass , gas density**



solve simplified EoM, photon transfer, pair creations
 → gap boundaries: enough pairs created

maximum Lorentz factor γ_{pk}
gamma-ray peak luminosity $L_{cur,pk}$



Semi-Analytic Model of Gamma-Ray Emission from Gap

- ◎ predicting gamma-ray emissivity **for wide range of BH mass , gas density**

ISM density, BH mass → one-zone IBH MAD model

(Kimura et al.21)

the gap disk photon B-field strength
opacity peak energy at the horizon

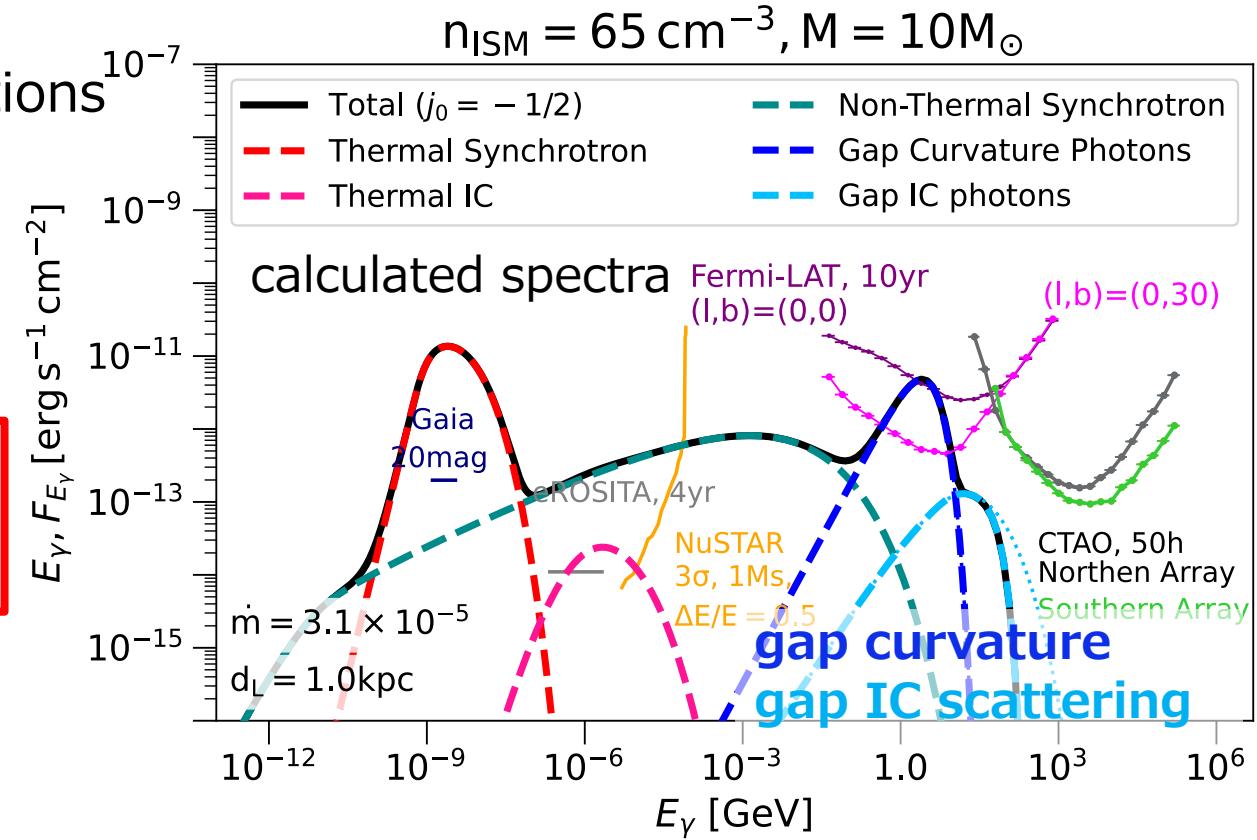
τ_0 、 ϵ_{min} 、 B_H

solve simplified EoM, photon transfer, pair creations
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maximum Lorentz factor γ_{pk}

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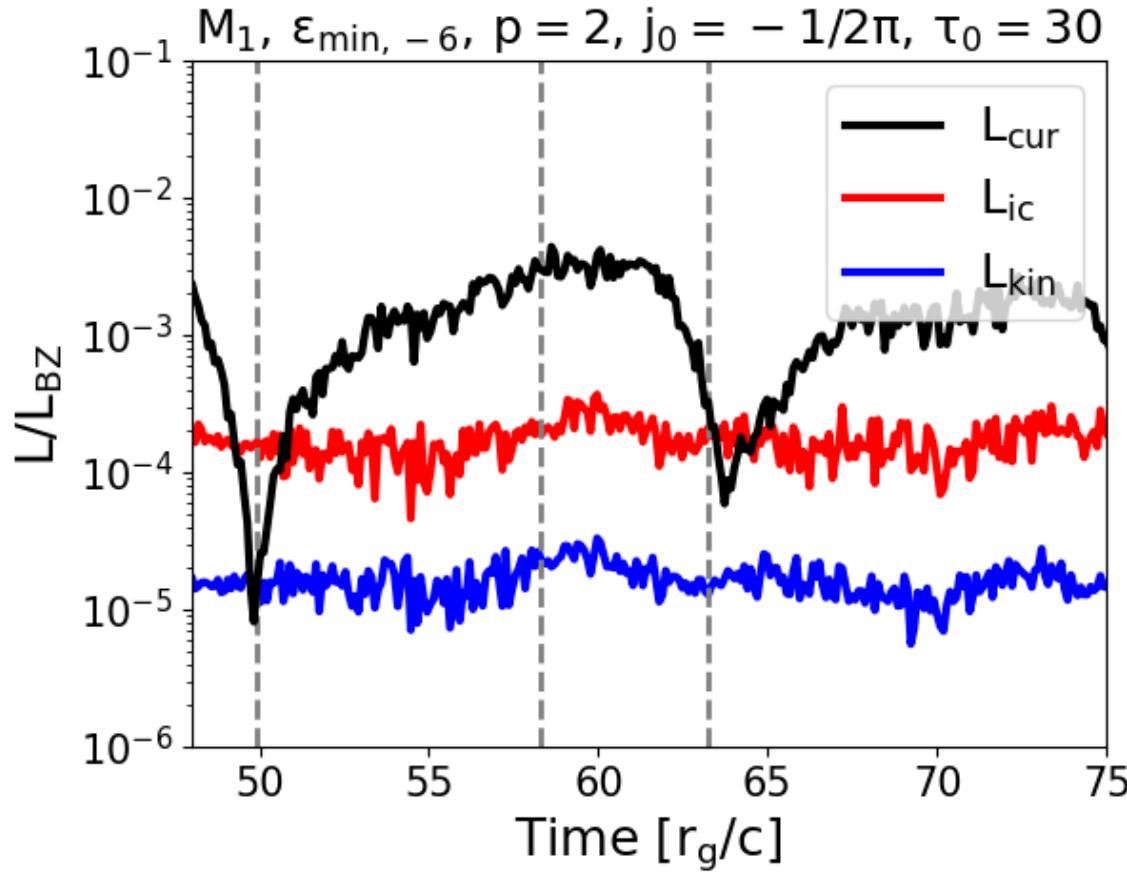
**GeV-TeV gamma-ray emission
detectable from ~kpc**



Summary

- ◎ Research Motivation: finding undetected isolated BHs through **gamma-ray** observation
gas infall → formation of **BH magnetosphere**, particle acceleration?
- ◎ Method: analyze plasma dynamics & gamma-ray characteristics using
1D • GRPIC simulation + semi-analytic modeling
⇒ **GeV-TeV gamma rays detectable from ~kpc, unID candidates**

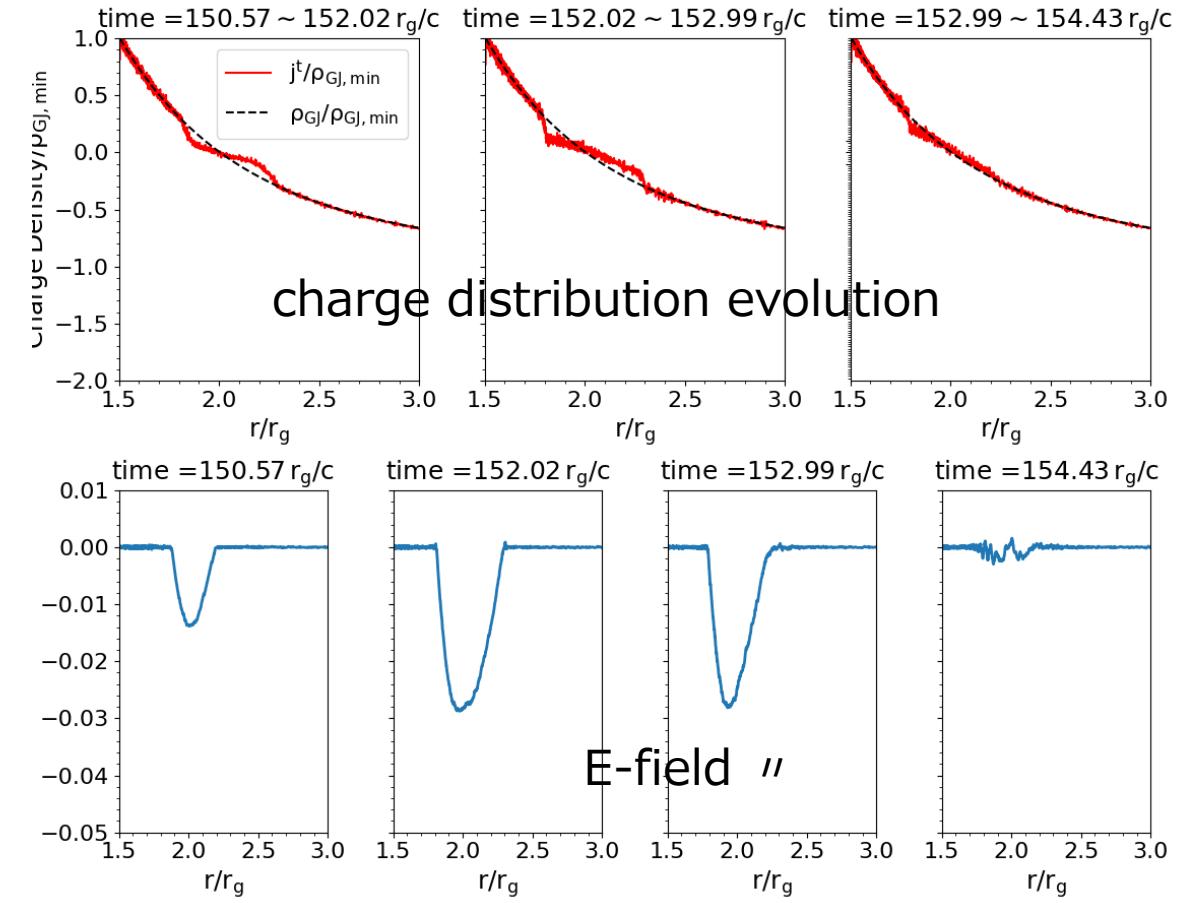
Back up



typical light curve at the outer boundary

$$j_0 = -1/2\pi$$

$\alpha = -1, p = 2.0, \epsilon_{\min} = 10^{-6}, \tau_0 = 55$



Discussion: strategy

\asymp pc (persec) $\sim 3 \times 10^{18}$ cm

◎ spectrum peak: 10-100GeV \rightarrow main target: **Fermi-LAT unIDs**

- hard spectral index for low energy side
 - has a peak in Fermi-LAT sensitivity range
 - association w/ gas clouds, point source
- ...will be the selection criteria

◎ determining the position: cross-match w/ optical~X-ray

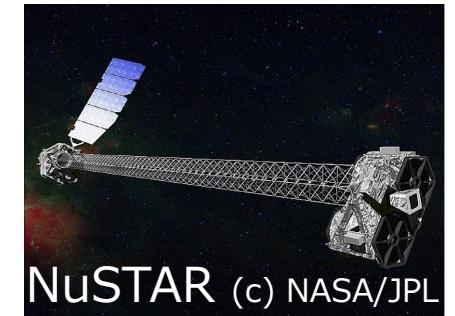
spatial resolution of gamma-ray detector: $\sim 1^\circ$

\rightarrow if \sim kpc distant, size of error circle: \sim pc!

cross-matching w/ data from other observation;

- Gaia(IR~optical)
- NuSTAR etc... (X-ray)

...will be crucial to narrow down & determine the location



Discussion: time variability

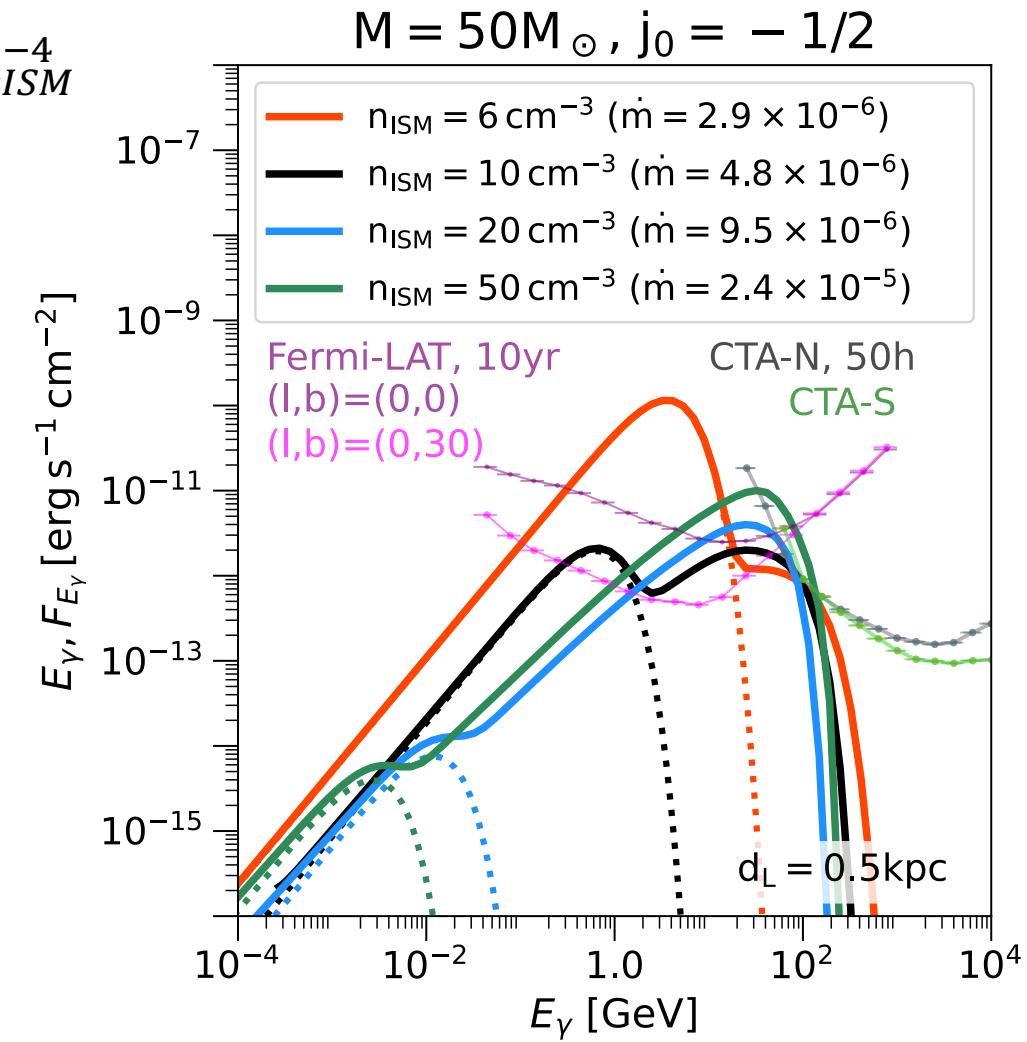
◎ L_{cur} sensitive to gas density: $L_{cur}/L_{BZ} \propto \tau_0^{-3.8} \propto n_{ISM}^{-4}$

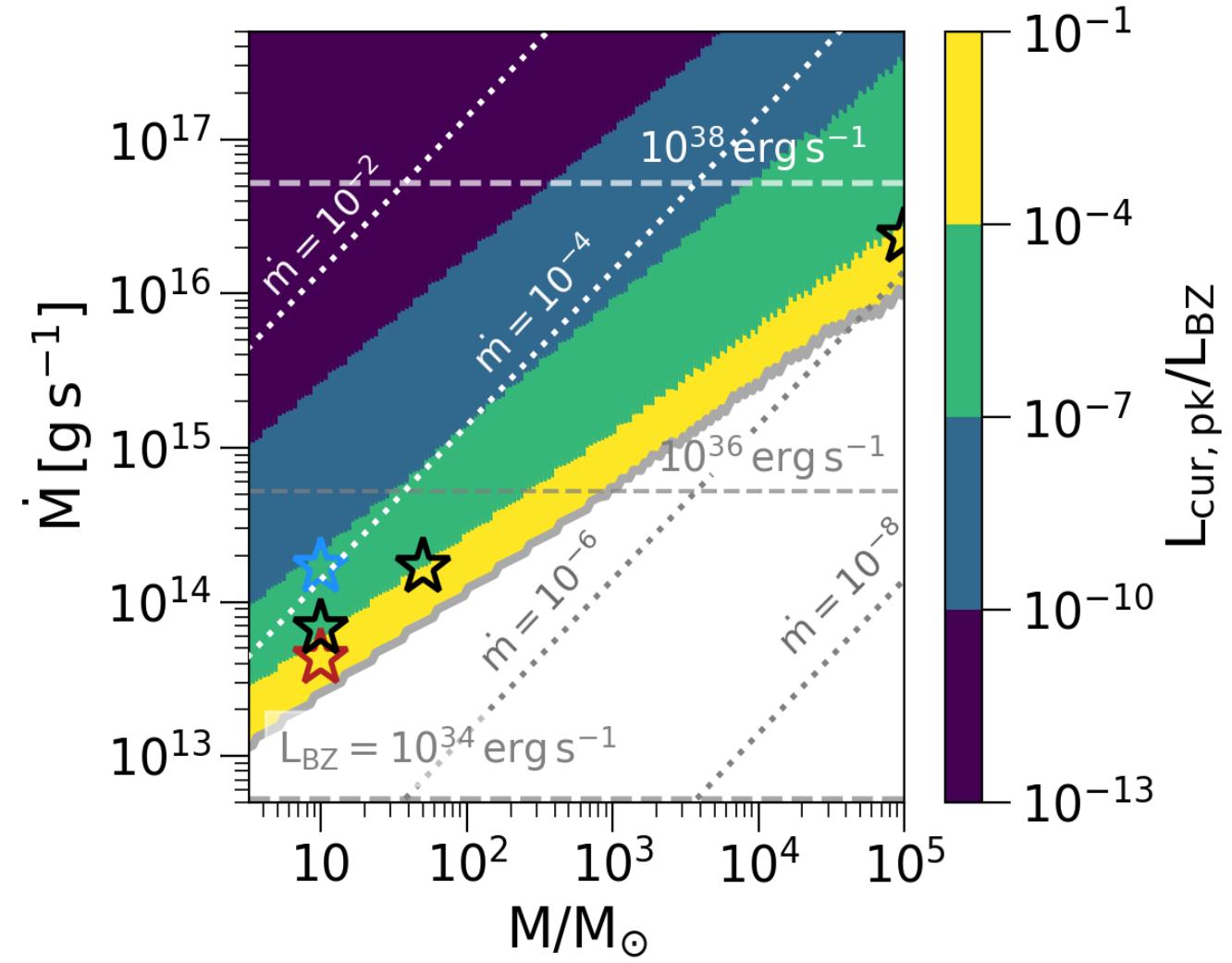
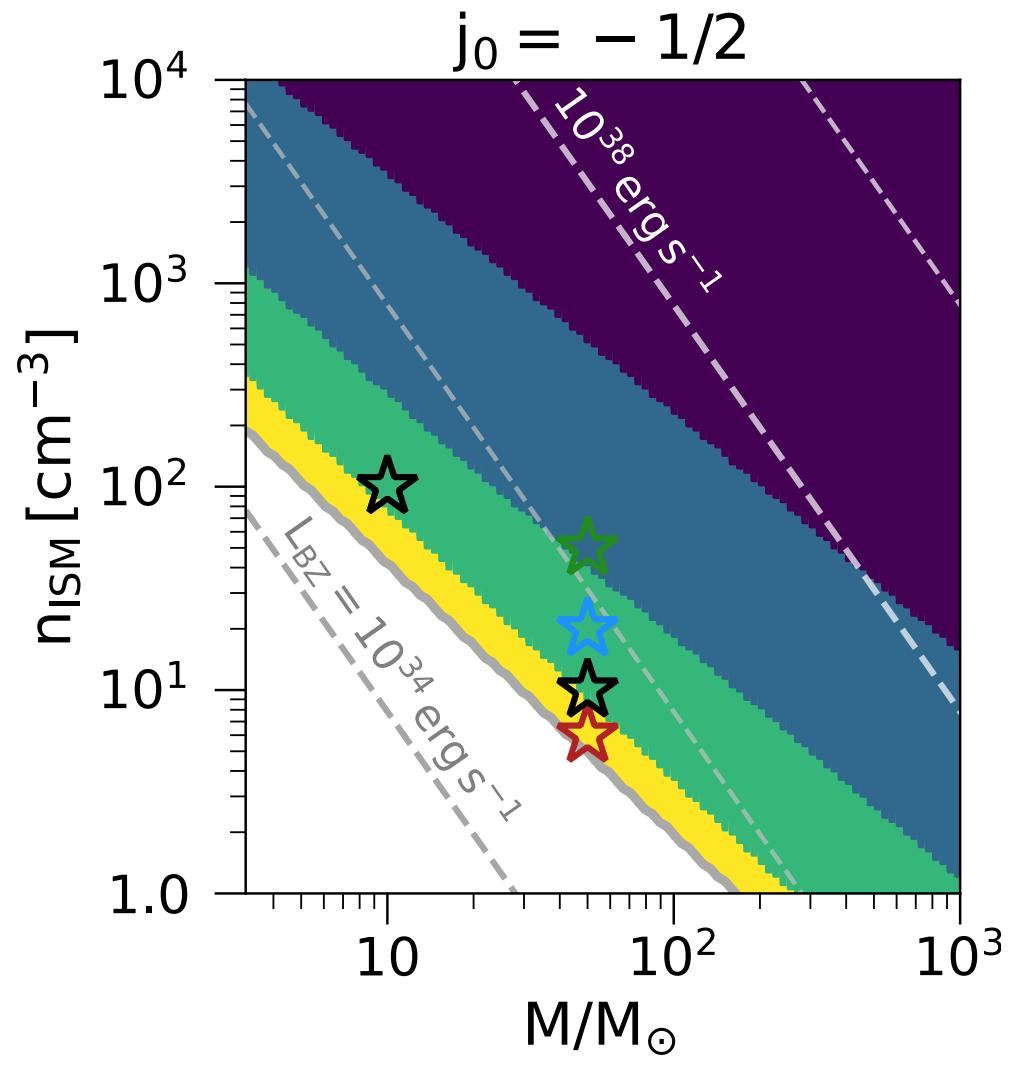
→ variable due to the gas turbulence

the timescale :

$$\tau \sim \frac{R_a (\text{length scale of gas turbulence})}{V_R (\text{infalling gas velocity})}$$

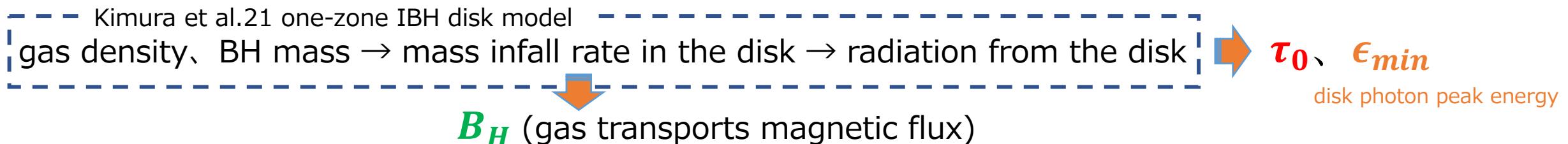
$$\approx \frac{8.3 \times 10^{13} \text{ cm}}{10^6 \text{ cm s}^{-1}} \approx 10^8 \text{ s}$$





Semi-Analytic Model of Gamma-Ray Emission from The “Gap”

- ◎ imitate the acceleration of e^- , gamma-ray emissivity during the oscillation



- 1) solve the simplified EoM of one e^- from inner boundary r_{in}

$$\frac{d(\gamma m_e c)}{dt} = e E_r(B_H, r_{in}) - \frac{P_{cur}(\gamma)}{c} - \frac{P_{ic}(\gamma)}{c}$$

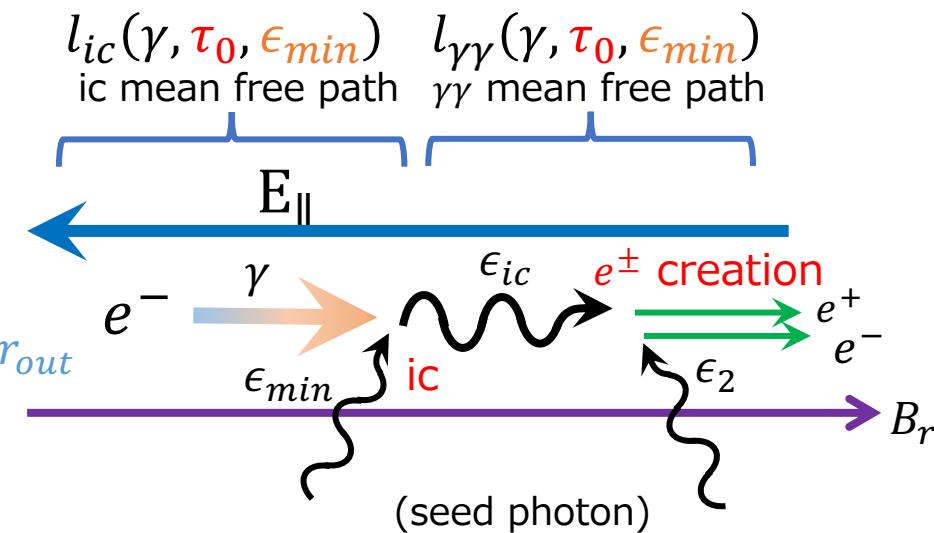
- 2) consider the pair creation

→ the position where enough pair created = outer boundary r_{out}

- 3) iteration, find r_{in} , r_{out} that satisfy $\frac{r_{in}+r_{out}}{2} \sim 2.0 r_g$

define the evolution of the e^- energy from $r_{in} \rightarrow r_{out}$

→ gamma-ray luminosity of the e^- also found unanimously

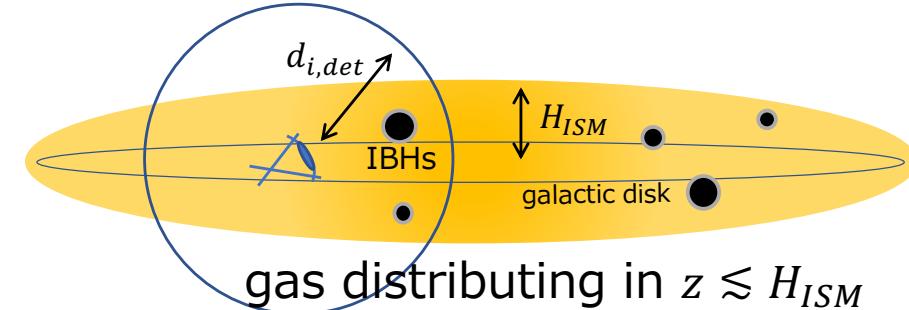


Discussion: expected number of detection in certain gas phase \mathcal{N}_{det}

- ◎ \mathcal{N}_{det} = number of IBHs in gas & sensitivity limit
sensitivity limit $d_{i,det}$: luminosity vs sensitivity

$$d_{i,det} = \sqrt{\frac{L_{obs}}{4\pi F_{sen}}} \sim 5 L_{obs,33}^{1/2} F_{sen,-12}^{-1/2} \text{ kpc}$$

$$\therefore \mathcal{N}_{det} \sim n_0 \xi_0 \frac{1-\gamma}{M_2^{1-\gamma} - M_1^{1-\gamma}} M^{1-\gamma} 2\pi H_{ISM} d_{i,det}^2 \simeq 3.7 \left(\frac{d_{i,det}}{5 \text{ kpc}} \right)^2 \left(\frac{M}{50 M_\odot} \right)^{1-\gamma} \text{ (for Cold HI)}$$



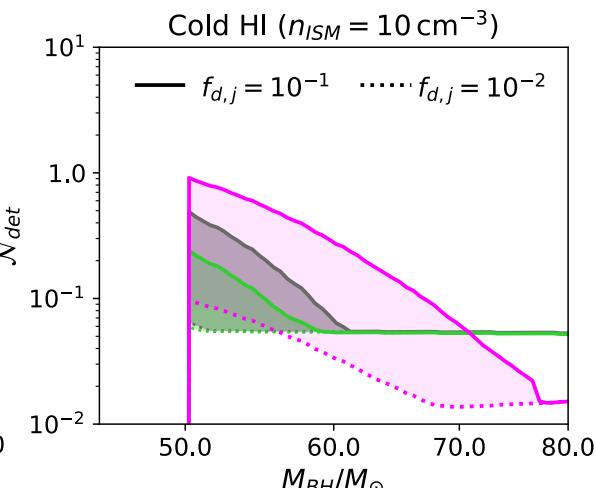
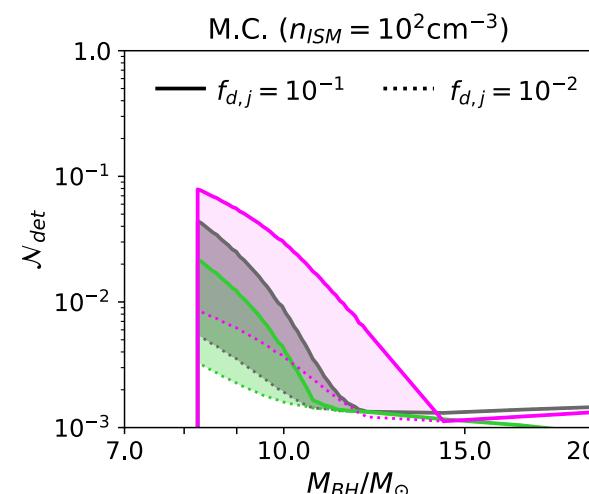
— CTAO-N,50h
— CTAO-S,50h
— Fermi-LAT,(0,0),10y

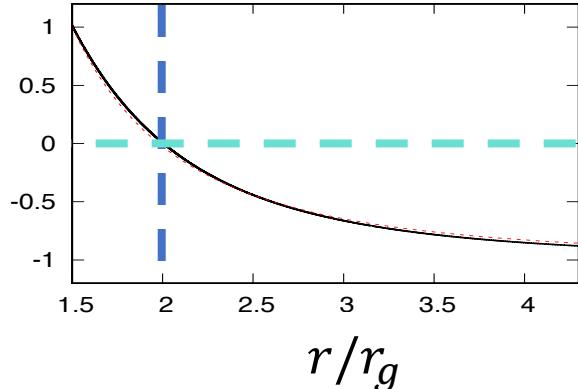
$$\frac{dN}{dM} \propto M^{-\gamma} \quad (\gamma \sim 2.6 \text{ Abbott et al.21})$$

ξ_0 : Volume filling factor

$n_0 \sim \mathcal{R}_{GW} n_{gal}^{-1} H_0^{-1} \sim 2 \times 10^2 \text{ kpc}^{-3}$: merged BH density

if n_0 is high, could be larger





ブラックホール磁気巻

◎パルサーとの大きな違い：磁場をアンカーする“表面”がない

→ブラックホール周りの時空の角速度 $\Omega (\equiv \frac{d\varphi}{dt})$ と“磁力線の角速度” Ω_F が一致しない

スプリットモノポール型の磁場を仮定すると
赤道面以外はモノポール

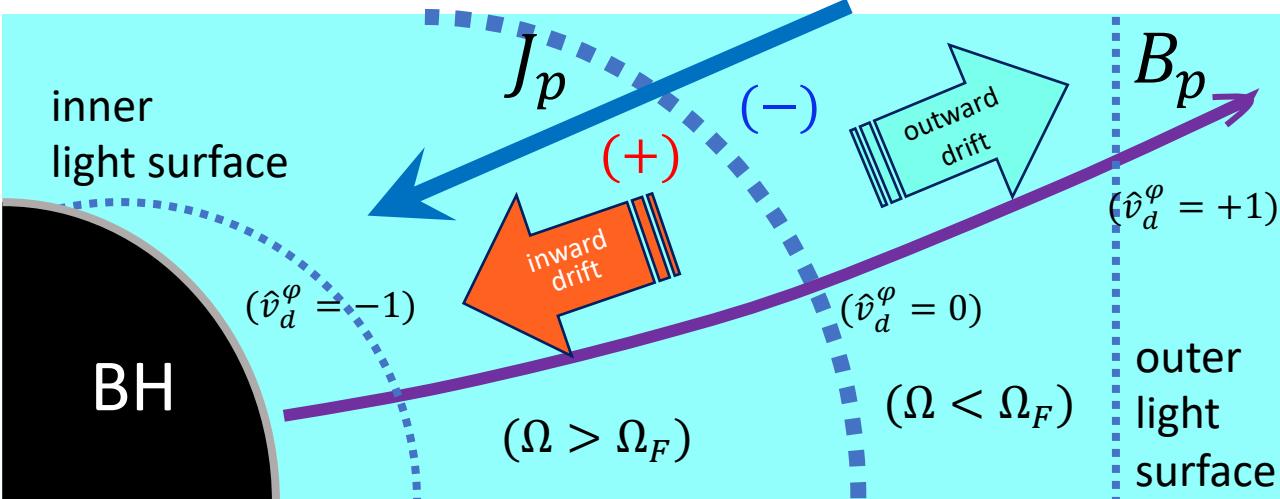
$$\Omega_F = \boxed{\frac{\Omega_H}{2}} + O(a^3)$$

$(\Omega_H:$ ホライズンでの $\Omega)$

$$\mathbf{D}_p = -\frac{1}{\alpha} (\Omega_F - \Omega) \mathbf{e}_\varphi \times \mathbf{B}_p \quad v \times B \text{電場}$$

$$\rho_{GJ} = \frac{1}{4\pi} \nabla \cdot \mathbf{D}_p = \dots = \frac{B_H \sqrt{A_H}}{4\pi \sqrt{-g}} \times \partial_\theta \left[\frac{\sin^2 \theta}{\alpha^2} \underline{(\Omega - \Omega_F)} \right] \dots \approx \boxed{-\frac{\Omega_F B_H}{2\pi c}}$$

$\rho_e \approx 0 \text{ at } \Omega \approx \Omega_F$



$(B_H, : \text{ホライズンでの磁場})$