Spectral modification of vector mesons

Philipp Gubler (JAEA)

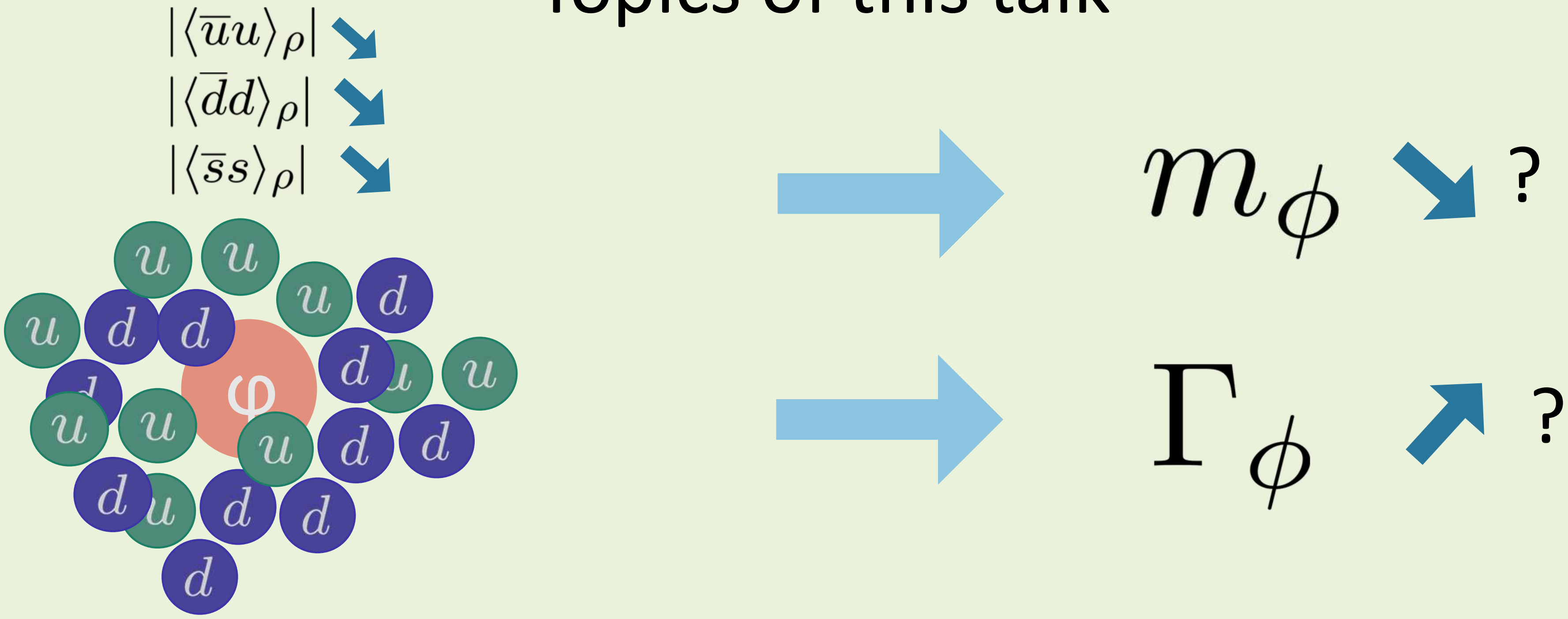
P. Gubler, M. Ichikawa, T. Song and E. Bratkovskaya, 2408.15364 [hep-ph].
L.M. Abreu, P. Gubler, K.P. Khemchandani, A. Martínez Torres and A. Hosaka, arXiv:2409.????? [hep-ph].

Talk at the E16 Collaboration Workshop
Academia Sinica, Taipei, Taiwan
September 9, 2024

Work done in collaboration with:

M. Ichikawa (JAEA)
T. Song (GSI)
E. Bratkovskaya (Goethe U. Frankfurt)
L. Abreu (U. Federal da Bahia)
K.P. Khemchandani (U. Federal de Sao Paulo)
A. Martínez Torres (U. de Sao Paulo)
A. Hosaka (Osaka U./RCNP)

Topics of this talk



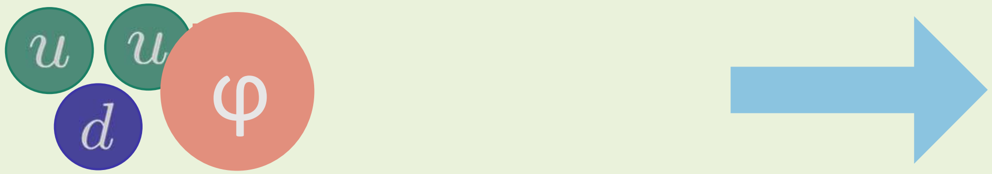
$|\langle \bar{u}u\rangle_\rho|$ ↘

$|\langle \bar{d}d\rangle_\rho|$ ↘

$|\langle \bar{s}s\rangle_\rho|$ ↘

m_ϕ ↘ ?

Γ_ϕ ↗ ?



How attractive is this interaction?

Can there be a φN bound state?

What is the relation to the φ meson mass shift in nuclear matter?

Contents

★ Recent news about the φN interaction

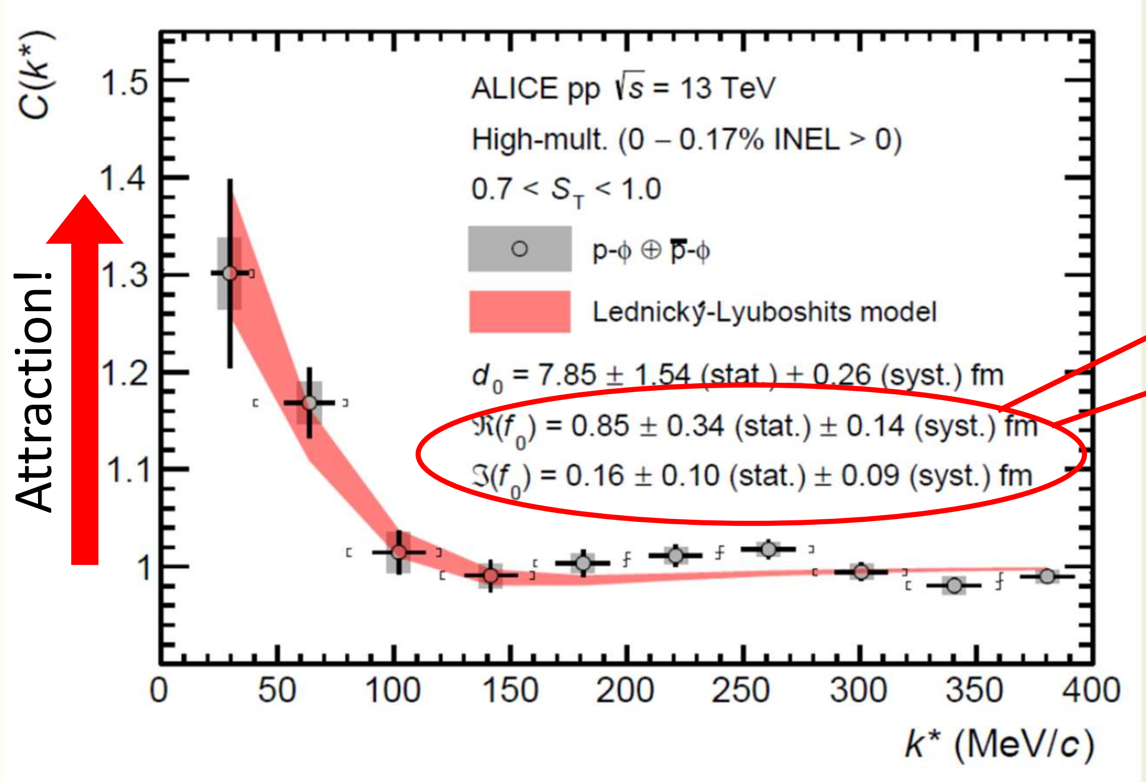
ALICE φN correlation measurement and its theoretical interpretation

★ Better understanding of φ meson production in pA reactions through transport simulations

Present status of ongoing research using PHSD

ALICE: pp

φN correlation function



S. Acharya et al. (ALICE Coll.), Phys. Rev. Lett. **127**, 172301 (2021).

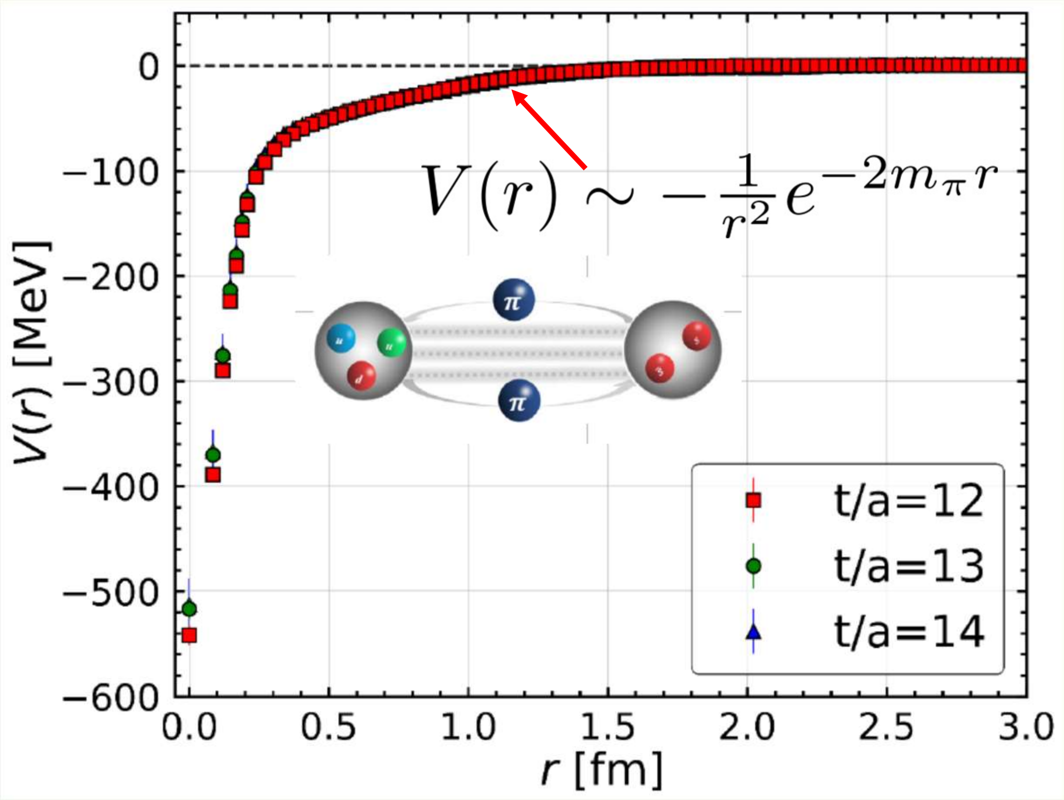
★ Strongly attractive

★ Small absorption

Caution:
Spin averaged value with
non-trivial weights
(thanks to Kamiya-san!)

φN potential from HAL QCD

Spin 3/2 channel



$$a_0^{3/2} = 1.43(23)_{\mathrm{stat.}}\left(^{+36}_{-06}\right)_{\mathrm{syst.}}\ \mathrm{fm}$$

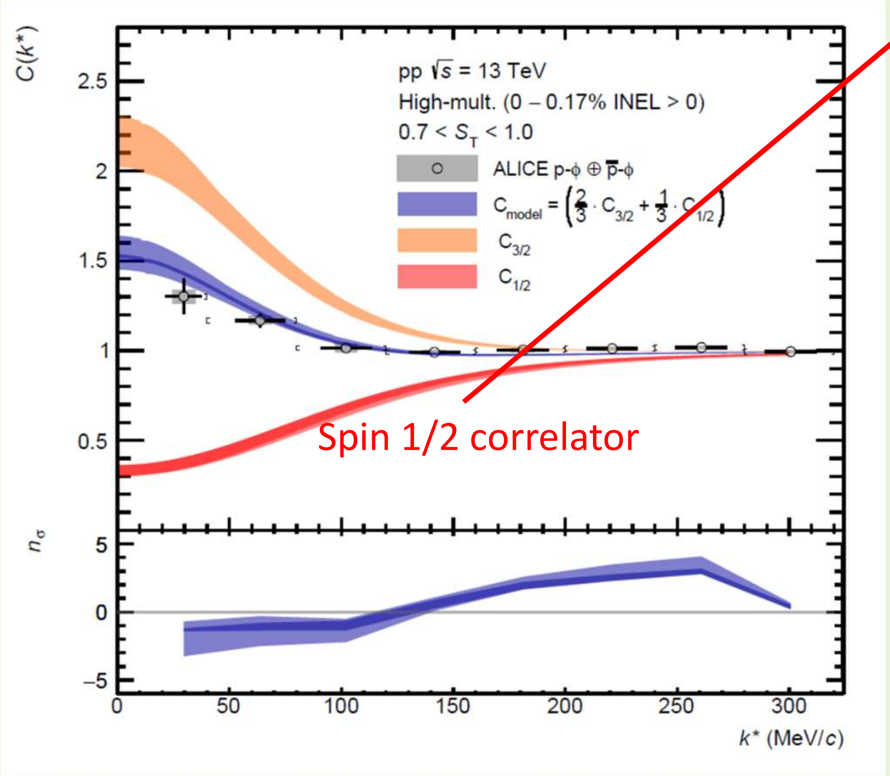
Indication for a quite strong and attractive interaction!

See talk by T. Doi for more details

Y. Lyu et al. (Lattice QCD, HAL QCD Collaboration), Phys. Rev. D **106**, 074507 (2022).

Even more recent results

Combination of ALICE pp-data and HAL QCD (spin 3/2) calculation



Spin 1/2 correlator

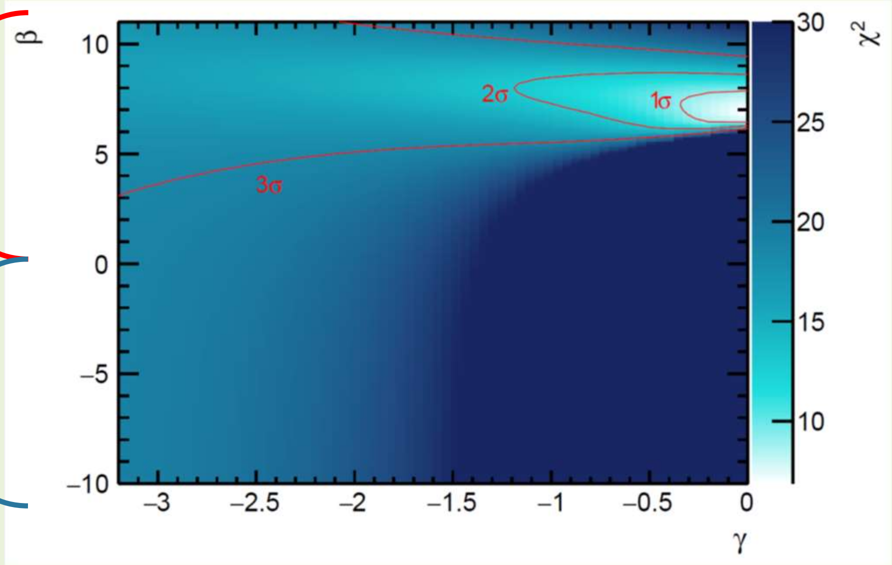
E. Chizzali et al., Phys. Lett. B **848**, 138358 (2024).

Bound state?

Repulsive?

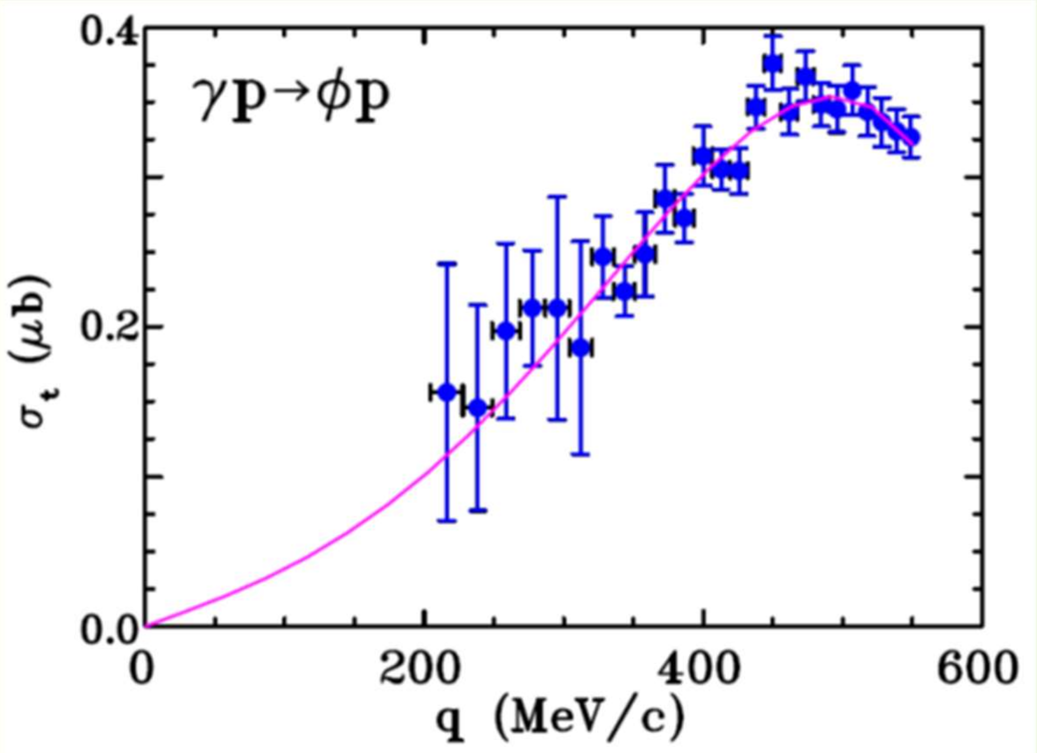
Attractive

Repulsive



Evidence for φ-N bound state!

Photoproduction measurement (CLAS)



Vector meson dominance is assumed

$$\frac{d\sigma^{\gamma p \to V p}}{d\Omega}\Big|_{\rm thr} = \frac{q}{k}\frac{1}{64\pi}|T^{\gamma p \to V p}|^2$$

$$= \frac{q}{k}\frac{\pi\alpha}{g_V^2}\frac{d\sigma^{V p \to V p}}{d\Omega}\Big|_{\rm thr} = \frac{q}{k}\frac{\pi\alpha}{g_V^2}|\alpha_{V p}|^2$$

$$|a_0| = 0.063 \pm 0.010\,\mathrm{fm}$$

Consistent with weak φN interaction

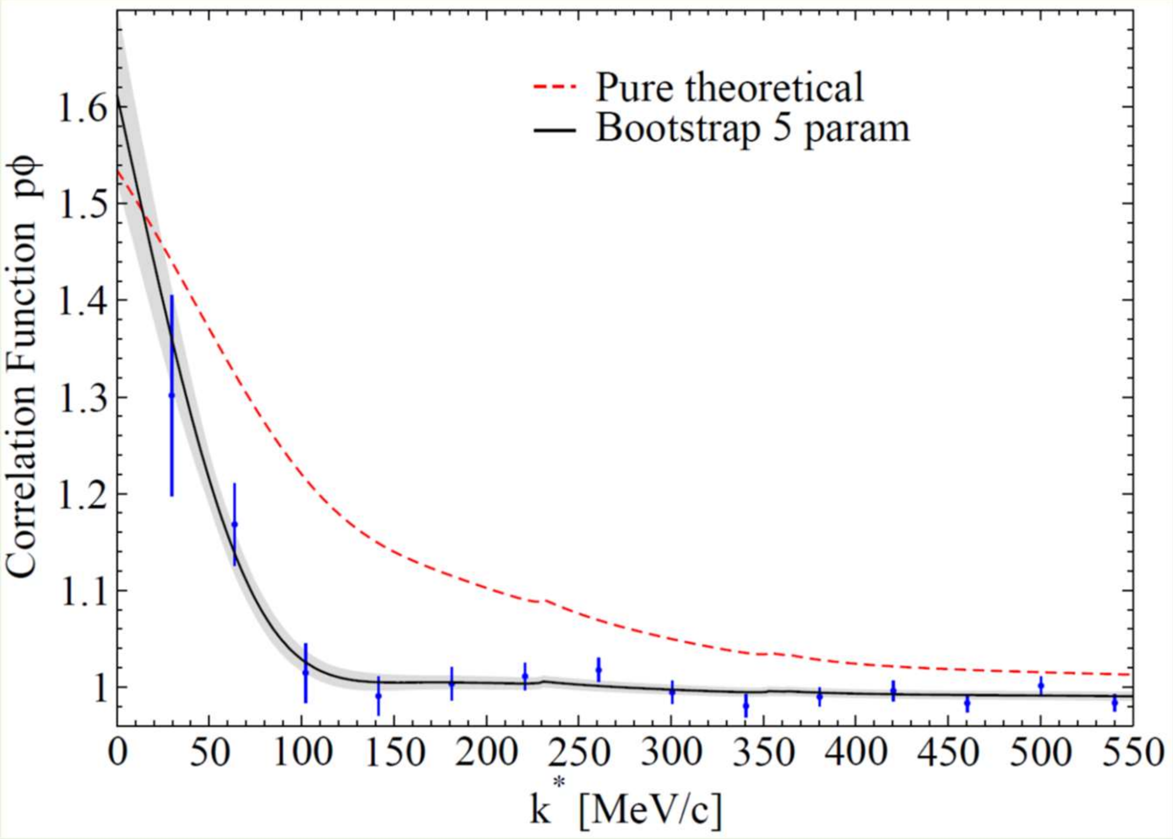
I.I. Strakovsky et al., Phys. Rev. C 101, 045201 (2020).

A lot of recent theoretical activity

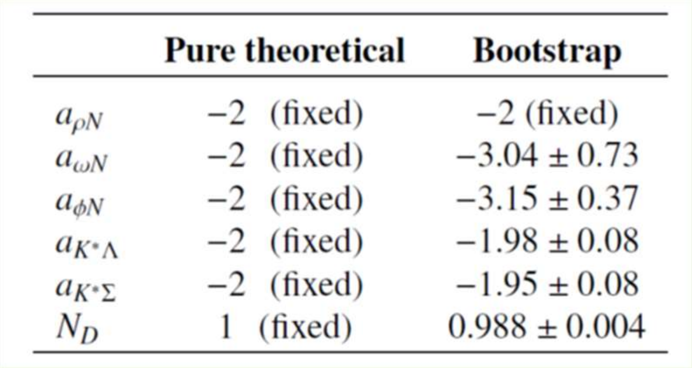
- Testing the φ-nuclear potential in pion-induced φ meson production on nuclei near threshold,
  E. Y. Paryev, Nucl. Phys. A **1032**, 122624 (2023).
- The φp bound state in the unitary coupled-channel approximation,
  B.-X. Sun, Y.-Y. Fan, Q.-Q. Cao, Commun. Theor. Phys. **75**, 055301 (2023).
- Possible ${}^{3}_{\phi}H$ hypernucleus with the HAL QCD interaction,
  I. Filikhin, R. Y. Kezerashvili, B. Vlahovic, Phys. Rev. D **110**, L031502 (2024).
- φ-p bound state and completeness of quantum states,
  A. Kuros, R. Maj, S. Mrowczynski, arXiv:2408.11941 [hep-ph].
- Bound states of ${}^{9}_{\phi}Be$ and ${}^{6}_{\phi\phi}He$ with φ+α+α and φ+φ+α cluster models,
  I. Filikhin, R. Y. Kezerashvili, B. Vlahovic, arXiv:2408.13415 [nucl-th].
- Relevance of the coupled channels in the φp and ρ^{0}p Correlation Functions ,
  A. Feijoo, M. Korwieser, L. Fabbietti, arXiv:2407.01128 [hep-ph].

New analysis of the ALICE data

A. Feijoo, M. Korwieser and L. Fabbietti, arXiv:2407.01128 [hep-ph].



Coupled channel approach, with subtraction constants as fittable parameters:



	Pure theoretical	**Bootstrap**
$a_{\rho N}$	-2 (fixed)	-2 (fixed)
$a_{\omega N}$	-2 (fixed)	-3.04 ± 0.73
$a_{\phi N}$	-2 (fixed)	-3.15 ± 0.37
$a_{K^*\Lambda}$	-2 (fixed)	-1.98 ± 0.08
$a_{K^*\Sigma}$	-2 (fixed)	-1.95 ± 0.08
N_D	1 (fixed)	0.988 ± 0.004

New analysis of the ALICE data

A. Feijoo, M. Korwieser and L. Fabbietti, arXiv:2407.01128 [hep-ph].

Obtained scattering length and effective range

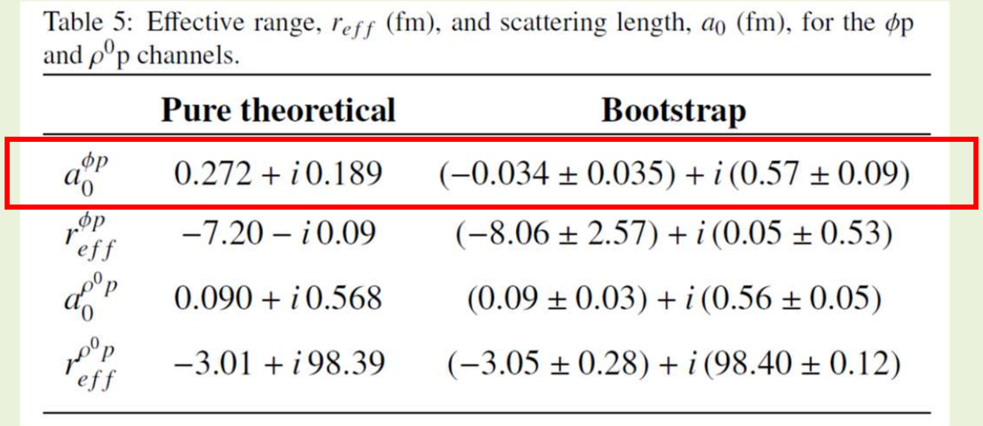
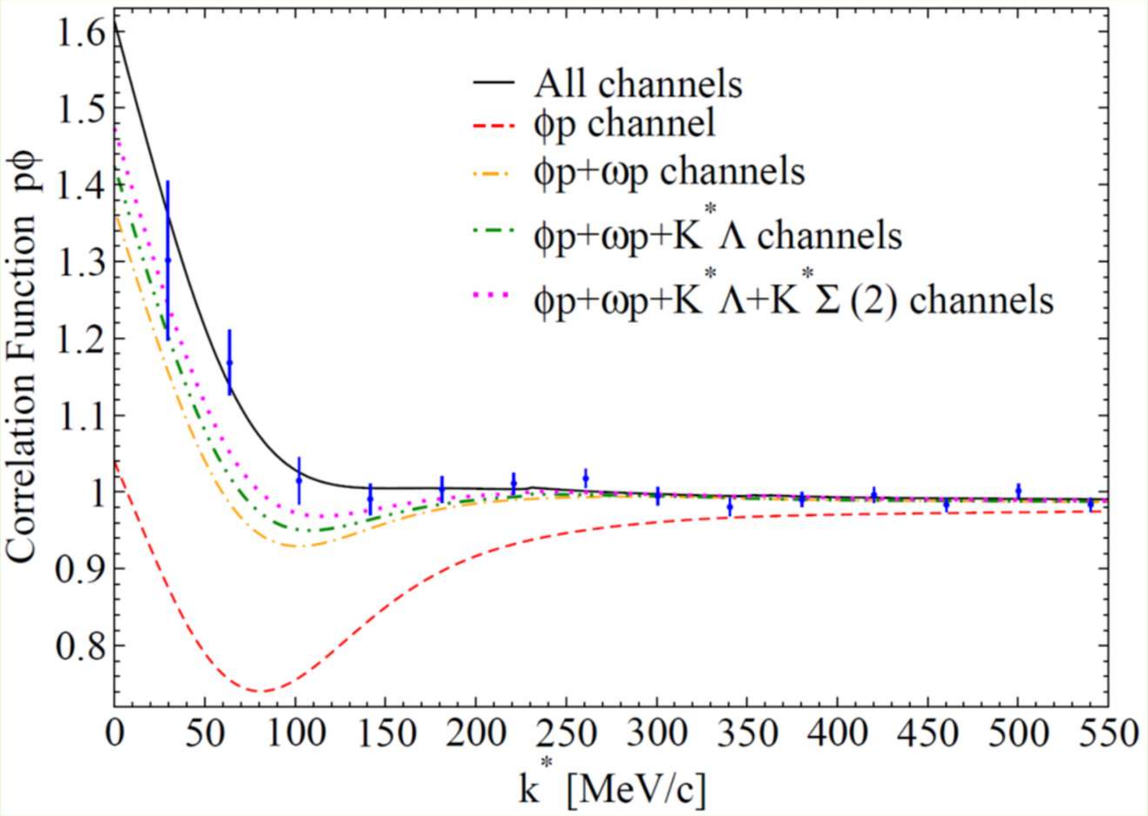


Table 5: Effective range, r_{eff} (fm), and scattering length, a_0 (fm), for the ϕp and ρ^0p channels.

	Pure theoretical	Bootstrap
$a_0^{\phi p}$	$0.272 + i\,0.189$	$(-0.034 \pm 0.035) + i\,(0.57 \pm 0.09)$
$r_{eff}^{\phi p}$	$-7.20 - i\,0.09$	$(-8.06 \pm 2.57) + i\,(0.05 \pm 0.53)$
$a_0^{\rho^0 p}$	$0.090 + i\,0.568$	$(0.09 \pm 0.03) + i\,(0.56 \pm 0.05)$
$r_{eff}^{\rho^0 p}$	$-3.01 + i\,98.39$	$(-3.05 \pm 0.28) + i\,(98.40 \pm 0.12)$

Decomposition into different hadronic channels



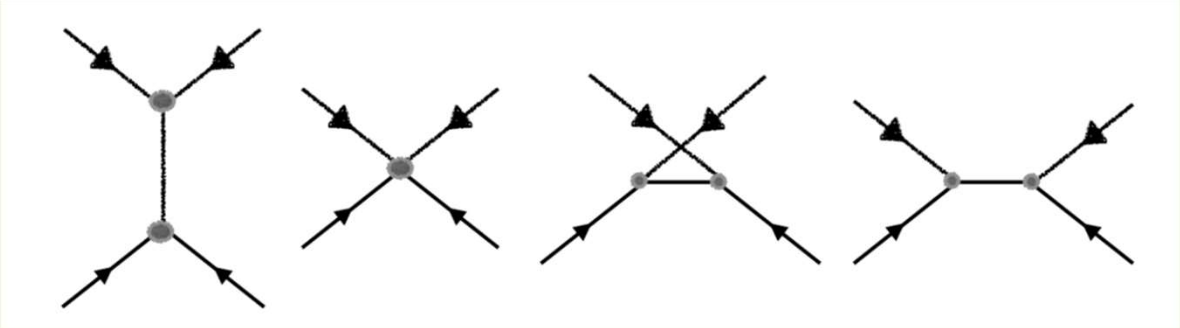
An even newer analysis of the ALICE data

L.M. Abreu, P. Gubler, K.P. Khemchandani, A. Martínez Torres and A. Hosaka, arXiv:2409.???? [hep-ph].
(will appear on the arXiv this week!)

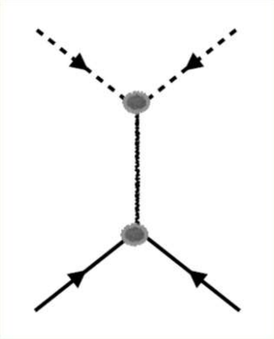
Starting point:
Hadronic Meson-Baryon interaction Lagrangian

1) Vector Meson-Baryon interaction: Based on Hidden Local Symmetry

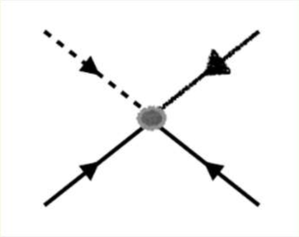
$$
\begin{aligned}
\mathcal{L}_{\mathrm{V}B} \quad &= -g\Big\{ \langle \bar{B}\gamma_\mu \left[V_8^\mu, B\right]\rangle + \langle \bar{B}\gamma_\mu B\rangle\langle V_8^\mu\rangle + \frac{1}{4M}\Big(F\langle \bar{B}\sigma_{\mu\nu}\left[V_8^{\mu\nu}, B\right]\rangle \\
&\quad + D\langle \bar{B}\sigma_{\mu\nu}\left\{V_8^{\mu\nu}, B\right\}\rangle\Big) + \langle \bar{B}\gamma_\mu B\rangle\langle V_0^\mu\rangle + \frac{C_0}{4M}\langle \bar{B}\sigma_{\mu\nu}V_0^{\mu\nu}B\rangle \Big\},
\end{aligned}
$$



2) Pseudoscalar Meson-Baryon interaction: Based on Chiral Symmetry

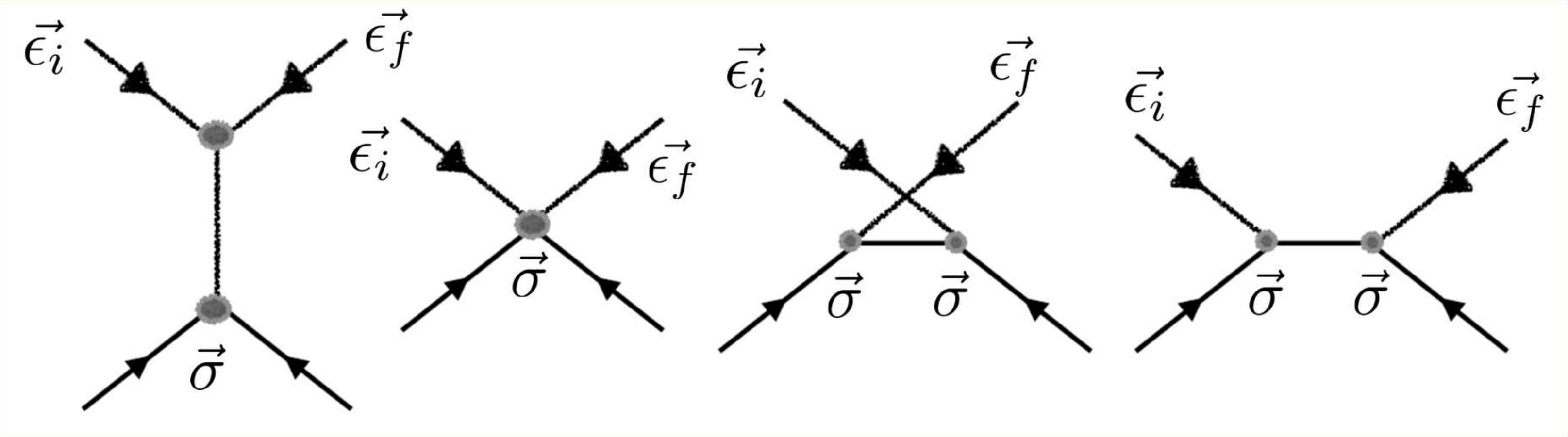


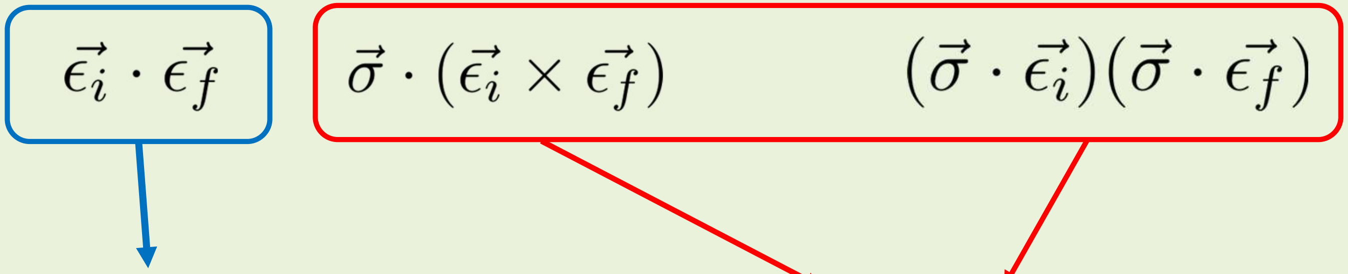
3) Transition between Pseudoscalar and Vector Mesons: Treating Vector meson as gauge boson in nonlinear sigma model (pion photoproduction + vector meson dominance)



L.M. Abreu, P. Gubler, K.P. Khemchandani, A. Martínez Torres and A. Hosaka, arXiv:2409.???? [hep-ph].

Crucial ingredient: spin dependent vector meson-baryon interactions





$\vec{\epsilon_i} \cdot \vec{\epsilon_f}$

$\vec{\sigma} \cdot (\vec{\epsilon_i} \times \vec{\epsilon_f})$ $\qquad$ $(\vec{\sigma} \cdot \vec{\epsilon_i})(\vec{\sigma} \cdot \vec{\epsilon_f})$

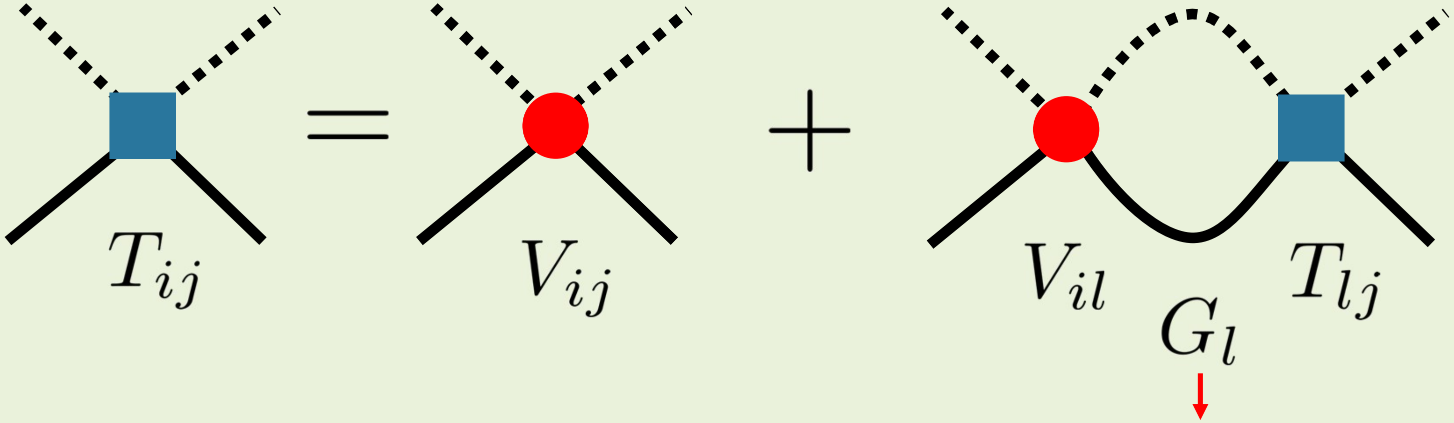
Considered so far
in Feijoo et al.

New!
Depend on spin configuration of the φN system (spin 1/2 or 3/2)

L.M. Abreu, P. Gubler, K.P. Khemchandani, A. Martínez Torres and A. Hosaka, arXiv:2409.???? [hep-ph].

Next step:
Solve the Bethe-Salpeter equation in the Vector Meson-Baryon channel of interest to obtain the full scattering amplitude T

$$T_{ij} = V_{ij} + V_{il} G_l T_{lj}$$



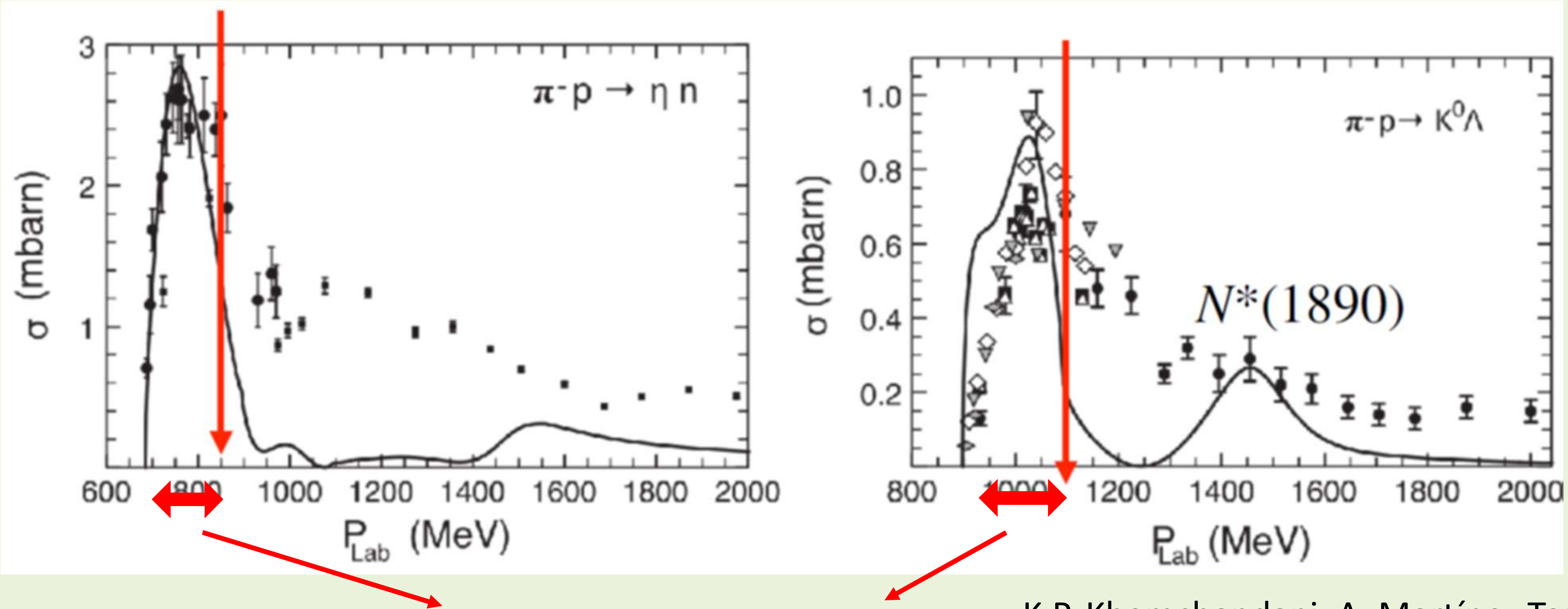
Emergence of subtraction constants, which are parameters of the model ←

Loop contains divergence that needs to be regulated. We use dimensional regularization

L.M. Abreu, P. Gubler, K.P. Khemchandani, A. Martínez Torres and A. Hosaka, arXiv:2409.???? [hep-ph].

Next step:
If possible, determine the parameters (subtraction constants) of the model from experimental data

For the spin 1/2 channel, scattering data are available for constraining the model parameters:

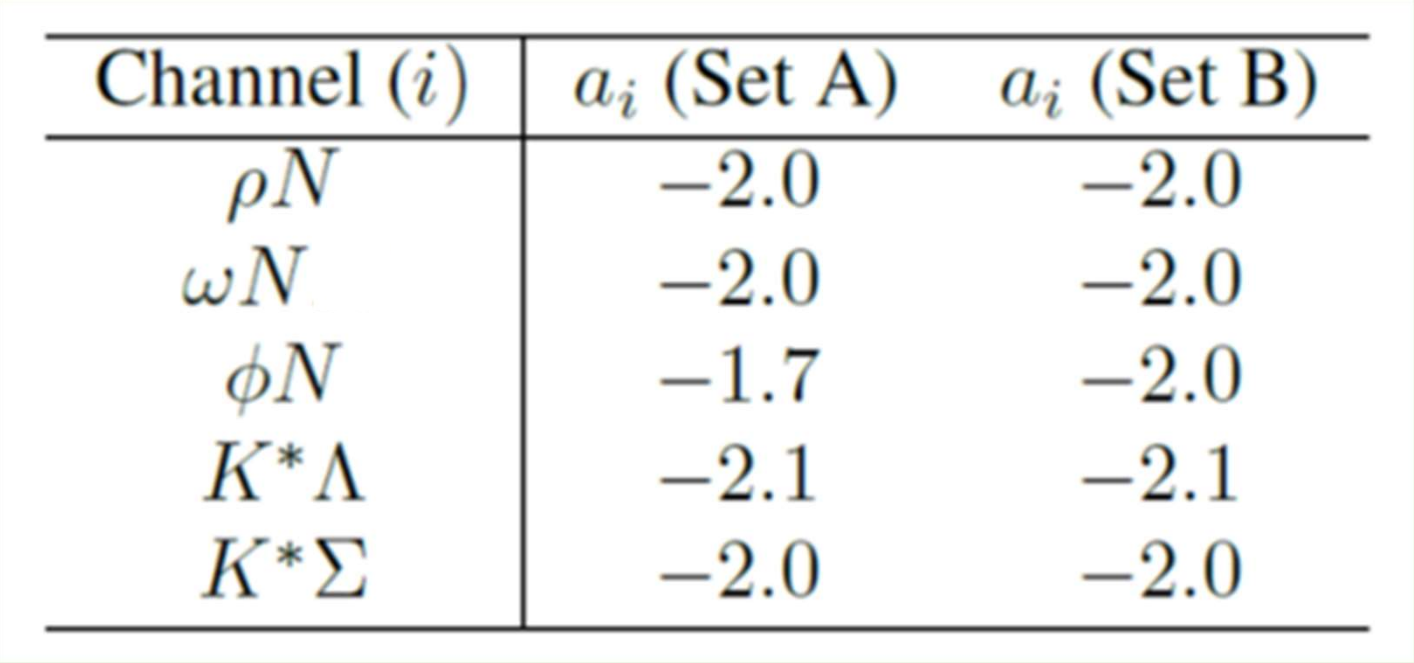


Regions used for the fit

K.P. Khemchandani, A. Martínez Torres, H. Nagahiro and A. Hosaka, Phys. Rev. D **88**, 114016 (2013).

L.M. Abreu, P. Gubler, K.P. Khemchandani, A. Martínez Torres and A. Hosaka, arXiv:2409.???? [hep-ph].

For the spin 3/2 channel, we use two data sets to evaluate the corresponding uncertainty:



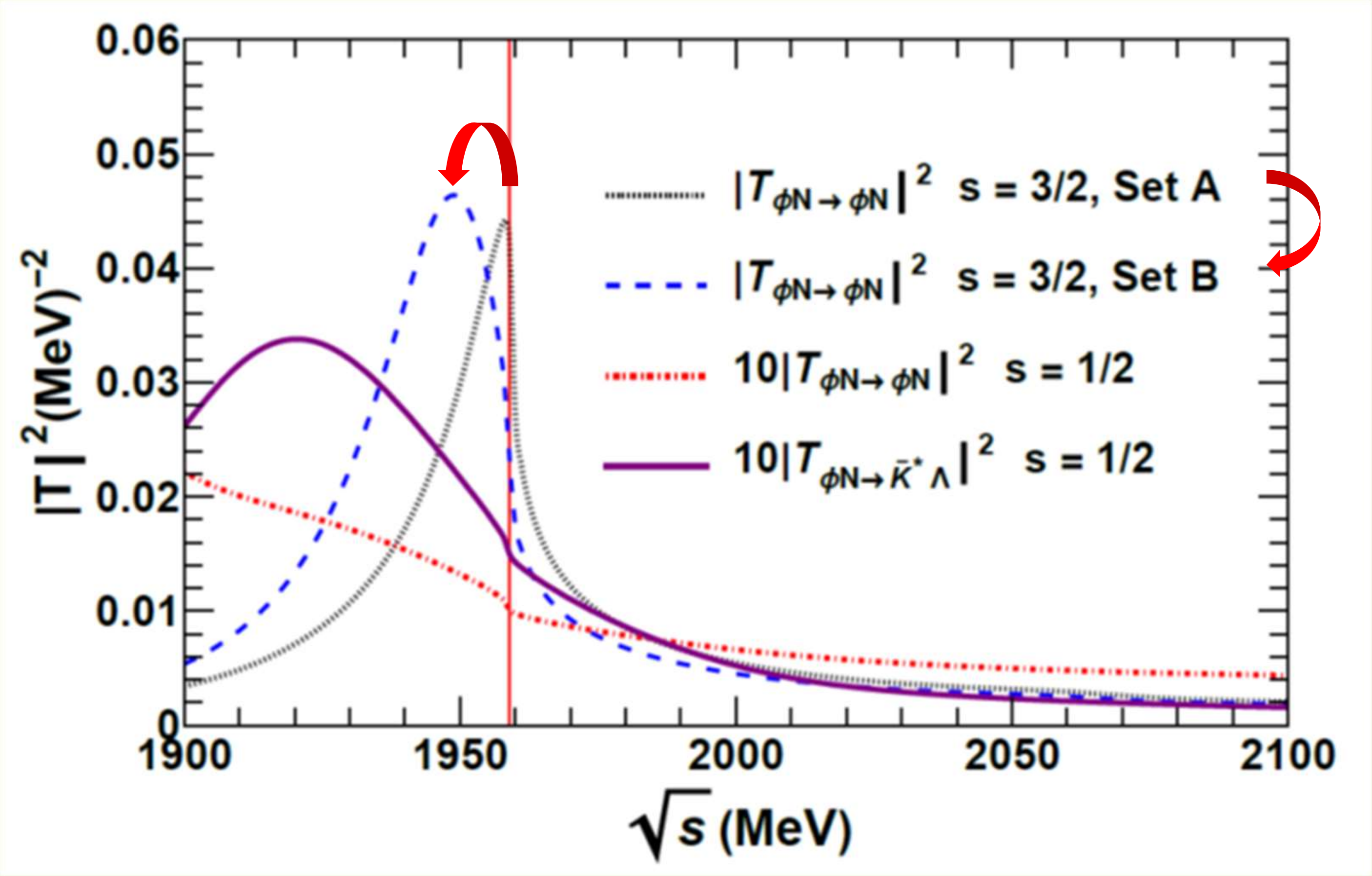
Channel (i)	a_i (Set A)	a_i (Set B)
ρN	-2.0	-2.0
ωN	-2.0	-2.0
ϕN	-1.7	-2.0
$K^*\Lambda$	-2.1	-2.1
$K^*\Sigma$	-2.0	-2.0

All are close to the “natural” value of $a_i = -2$

Only the value in the φN channel is modified, to study the change in the respective interaction strength

L.M. Abreu, P. Gubler, K.P. Khemchandani, A. Martínez Torres and A. Hosaka, arXiv:2409.???? [hep-ph].

The obtained scattering amplitudes



- ★ For spin 3/2, a state below the φN threshold can be found, depending on the values of the subtraction constants.
  However, above the threshold, the differences between the two cases are small.

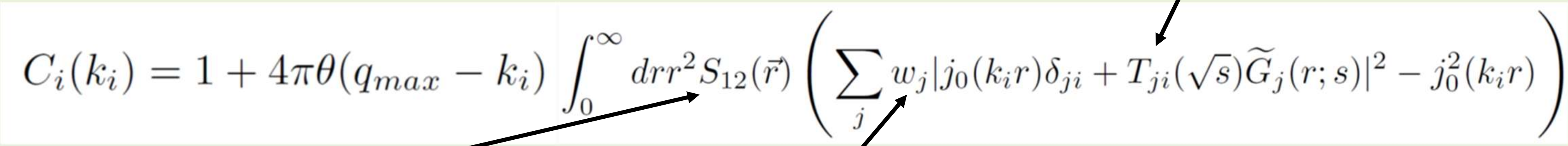
- ★ The spin 1/2 scattering amplitudes are much weaker than their spin 3/2 counterparts

L.M. Abreu, P. Gubler, K.P. Khemchandani, A. Martínez Torres and A. Hosaka, arXiv:2409.???? [hep-ph].

Next step:
Calculate the correlation function and compare with the ALICE data

We use a generalized Koonin-Pratt formula:

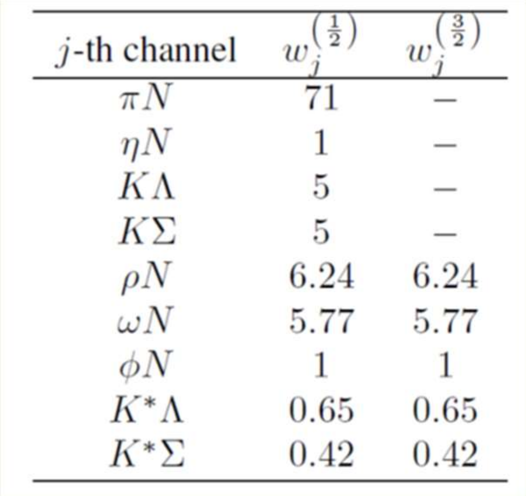
Scattering amplitudes:
See previous slide



$$C_i(k_i) = 1 + 4\pi\theta(q_{max} - k_i)\int_0^\infty drr^2 S_{12}(\vec{r})\left(\sum_j w_j |j_0(k_i r)\delta_{ji} + T_{ji}(\sqrt{s})\widetilde{G}_j(r;s)|^2 - j_0^2(k_i r)\right)$$

Source function:
Gaussian with radius of approx. 1 fm

Weights related to the multiplicity of pairs of primary particles created at the initial stage of the collision

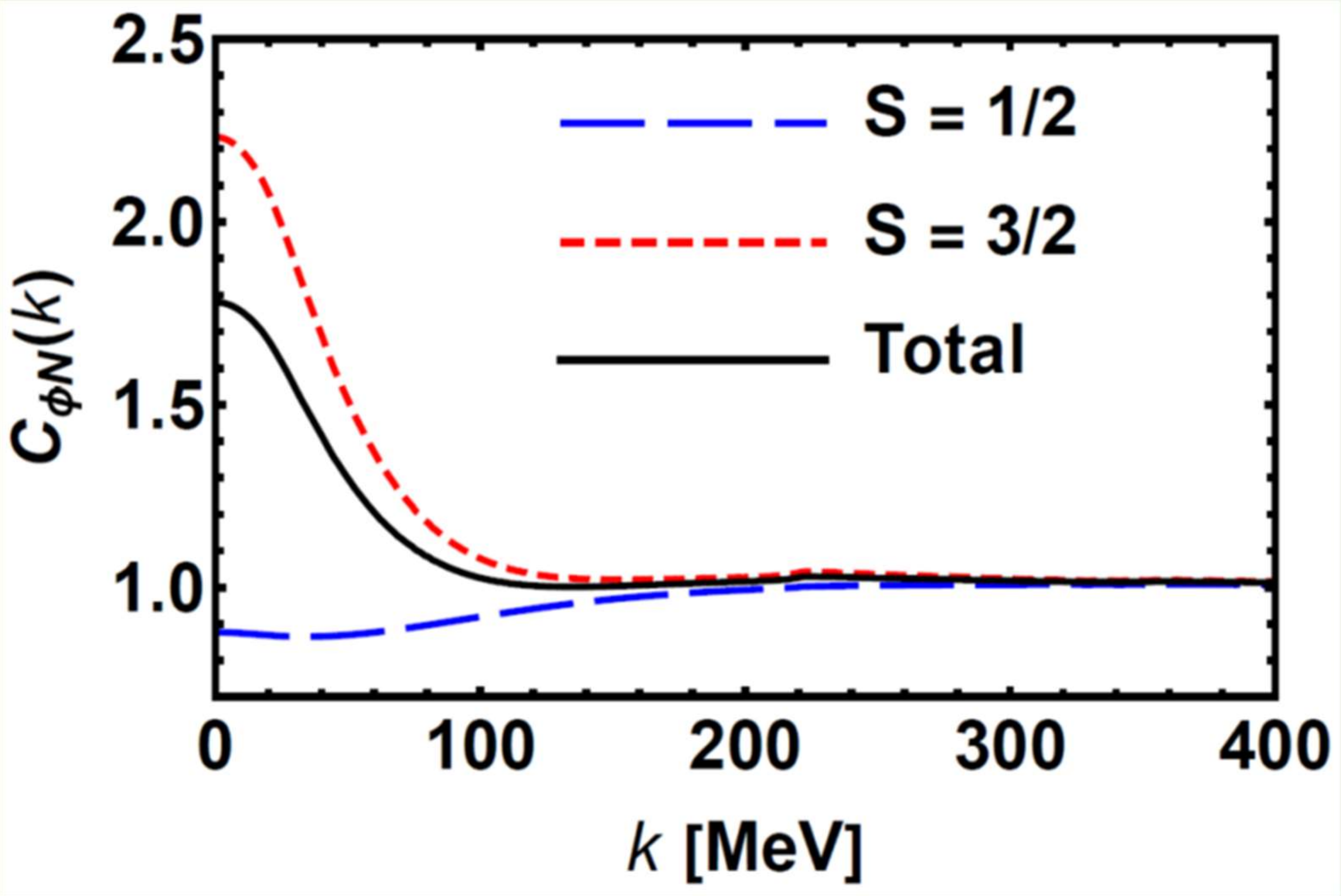


j-th channel	$w_j^{(\frac{1}{2})}$	$w_j^{(\frac{3}{2})}$
πN	71	–
ηN	1	–
$K\Lambda$	5	–
$K\Sigma$	5	–
ρN	6.24	6.24
ωN	5.77	5.77
ϕN	1	1
$K^*\Lambda$	0.65	0.65
$K^*\Sigma$	0.42	0.42

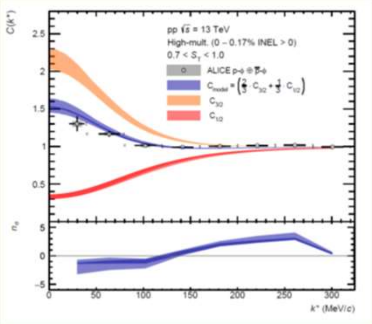
Obtained from
Thermal-Fist package

L.M. Abreu, P. Gubler, K.P. Khemchandani, A. Martínez Torres and A. Hosaka, arXiv:2409.???? [hep-ph].

The obtained correlation function (spin decomposed)

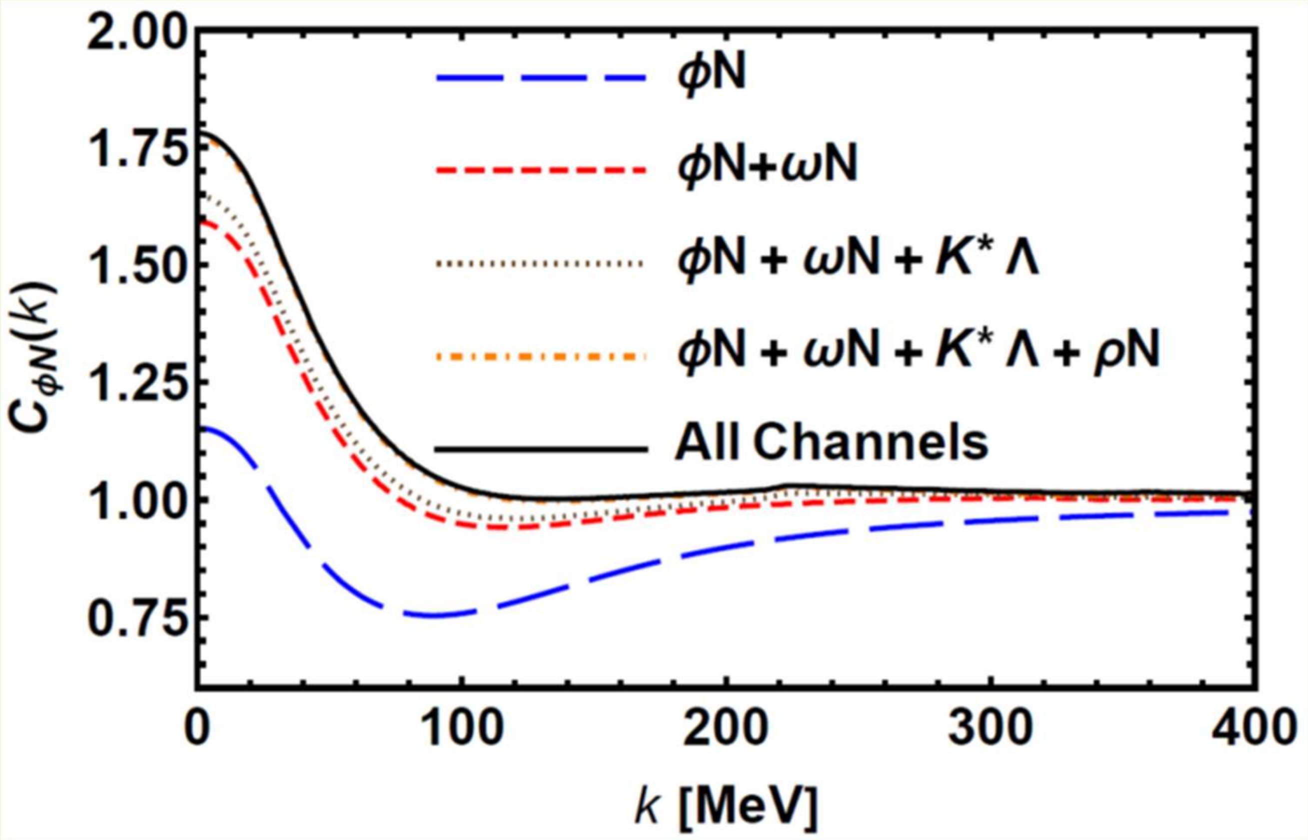


Qualitatively similar to
E. Chizzali et al.,
Phys. Lett. B **848**, 138358 (2024).



L.M. Abreu, P. Gubler, K.P. Khemchandani, A. Martínez Torres and A. Hosaka, arXiv:2409.???? [hep-ph].

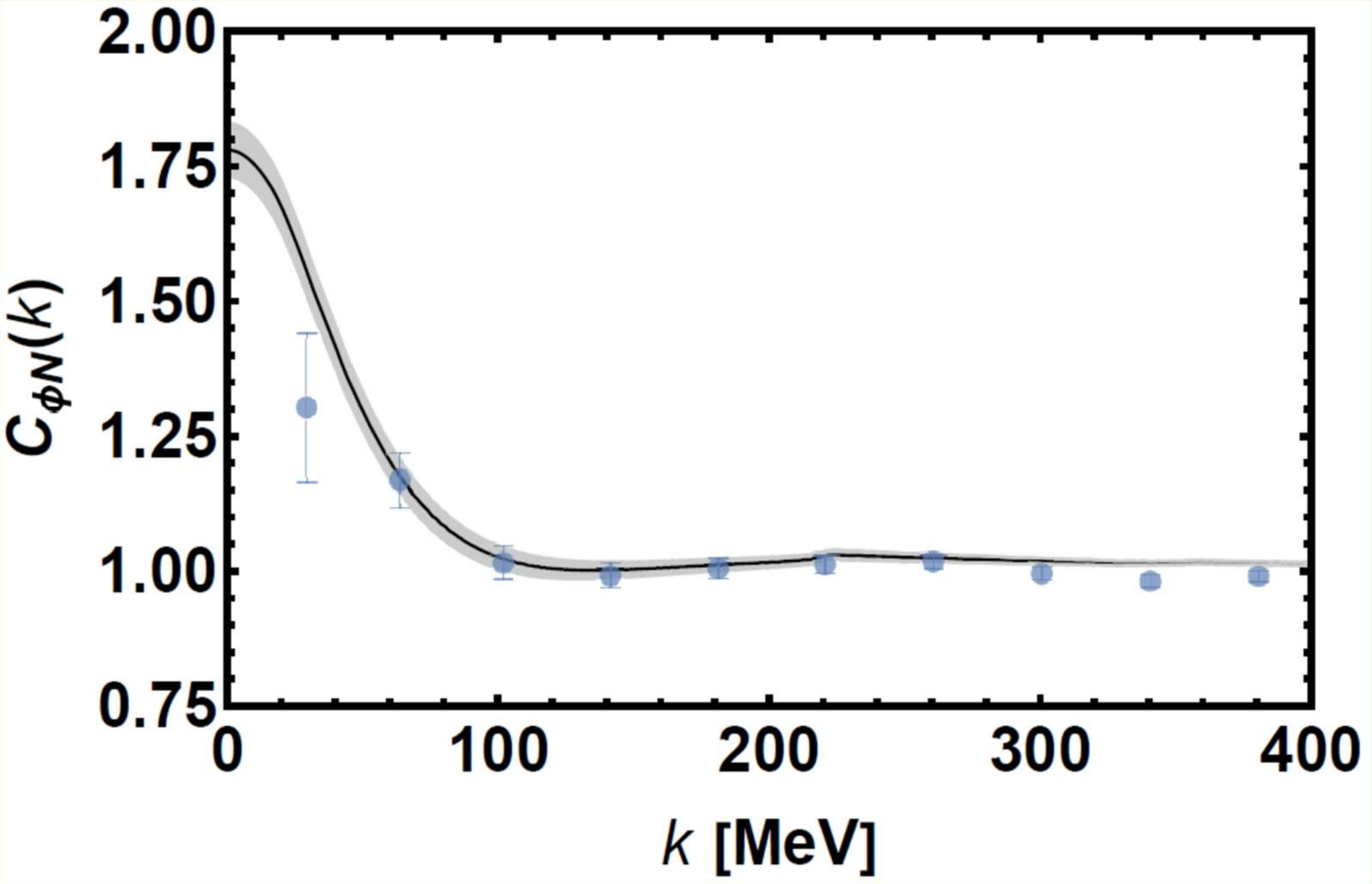
The obtained correlation function (channel decomposed)



Similar to
A. Feijoo et al.,
arXiv:2407.01128 [hep-ph].

L.M. Abreu, P. Gubler, K.P. Khemchandani, A. Martínez Torres and A. Hosaka, arXiv:2409.???? [hep-ph].

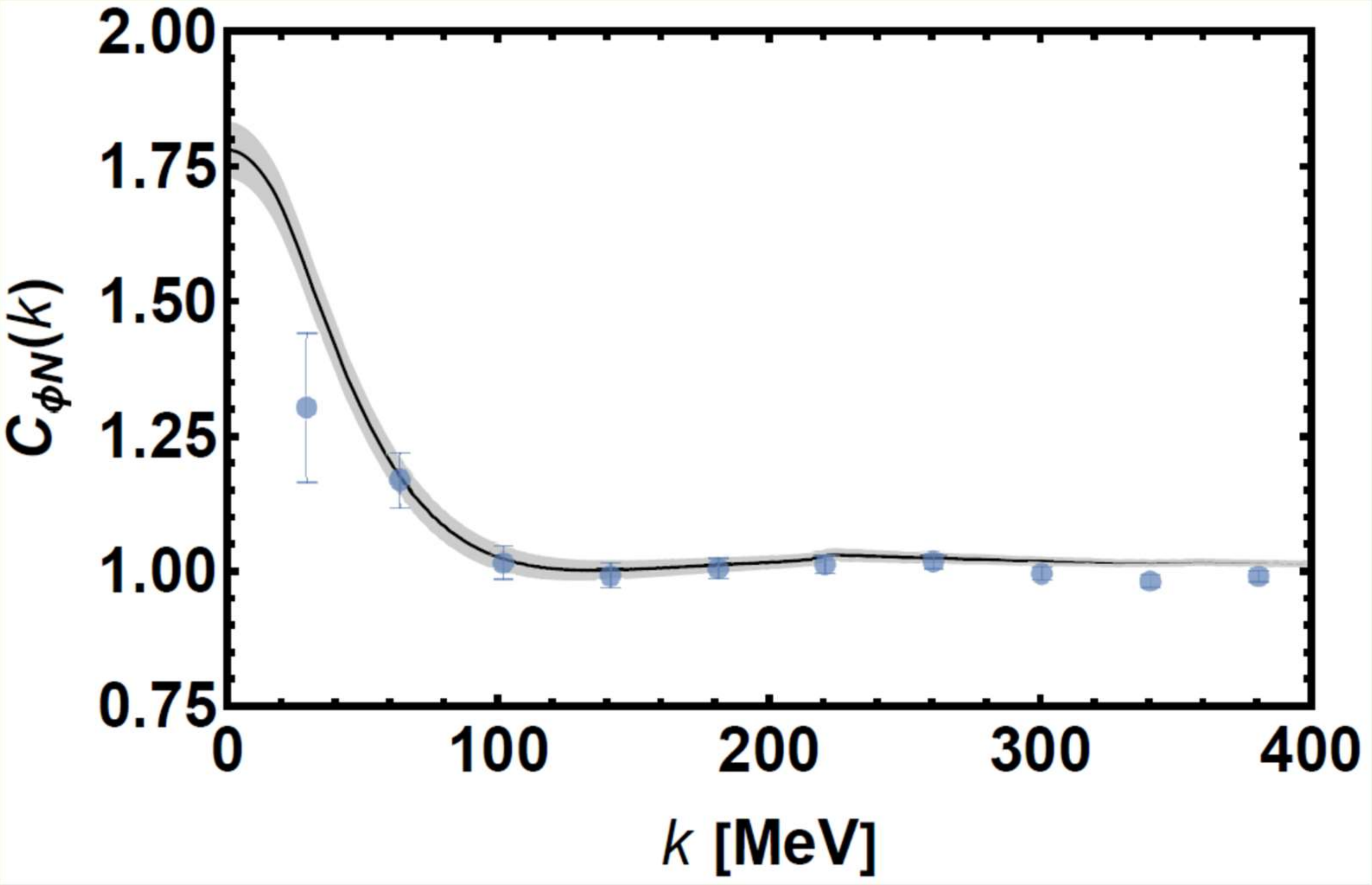
The obtained correlation function (compared with ALICE data)



Reasonably good agreement
without any parameter fitting!

L.M. Abreu, P. Gubler, K.P. Khemchandani, A. Martínez Torres and A. Hosaka, arXiv:2409.???? [hep-ph].

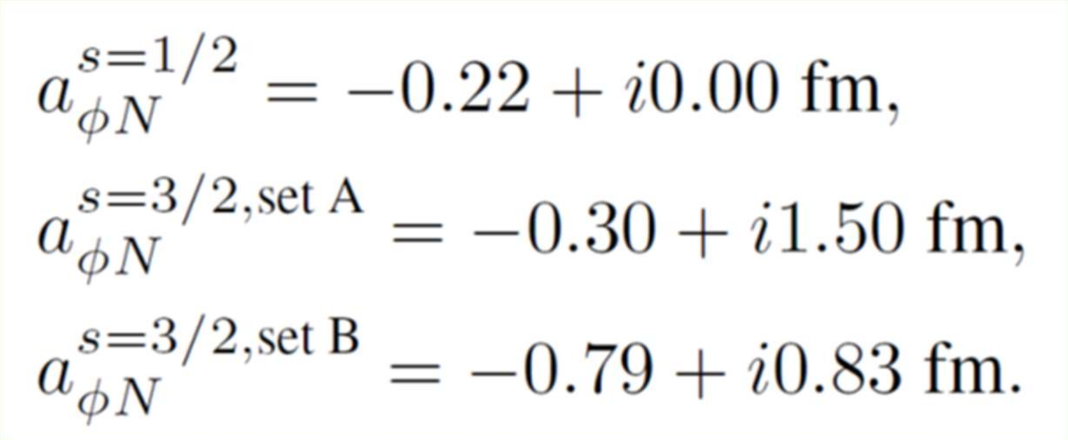
The obtained correlation function (compared with ALICE data)



Reasonably good agreement
without any parameter fitting!

L.M. Abreu, P. Gubler, K.P. Khemchandani, A. Martínez Torres and A. Hosaka, arXiv:2409.???? [hep-ph].

The obtained scattering lengths



$$a_{\phi N}^{s=1/2} = -0.22 + i0.00 \text{ fm},$$
$$a_{\phi N}^{s=3/2,\text{set A}} = -0.30 + i1.50 \text{ fm},$$
$$a_{\phi N}^{s=3/2,\text{set B}} = -0.79 + i0.83 \text{ fm}.$$

Large model parameter dependence!

Correlation function is not very sensitive to the scattering length

L.M. Abreu, P. Gubler, K.P. Khemchandani, A. Martínez Torres and A. Hosaka, arXiv:2409.???? [hep-ph].

Simple relation between φN scattering length and φ meson mass shift in nuclear matter

$$\left.\begin{aligned} V_\phi(\rho) &= -\frac{2\pi}{m_\phi}\rho\Big(1+\frac{m_\phi}{m_N}\Big)a_0 \\ &\simeq -85\frac{\rho}{\rho_0}\Big(\frac{a_0}{\mathrm{fm}}\Big)\,\mathrm{MeV} \end{aligned}\right\}$$

Valid within the linear density approximation

Larger than 100 MeV IF HAL QCD result is true for all spin configurations!

However, the above prescription seems problematic if a φN bound state (or resonance) is formed.

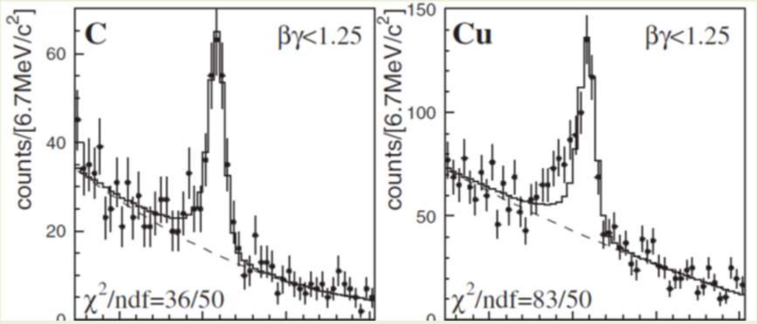
→ Need better theoretical understanding!

Comparison of theory and experiment

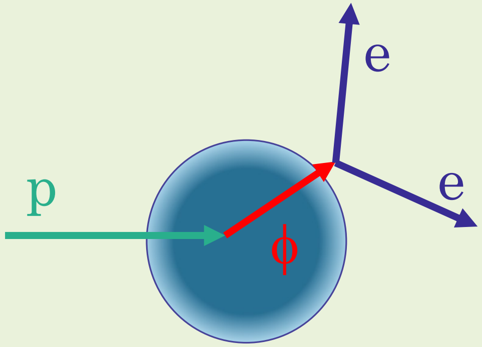
Information useful for theory

- Spectral function as a function of density
- Mass at normal nuclear matter density
- Decay width at normal nuclear matter density

Experimental data



Realistic simulation of pA reaction is needed!



Our tool: transport simulation
PHSD (Parton Hadron String Dynamics)

E.L. Bratkovskaya and W. Cassing, Nucl. Phys. A **807**, 214 (2008).
W. Cassing and E.L. Bratkovskaya, Phys. Rev. C **78**, 034919 (2008).

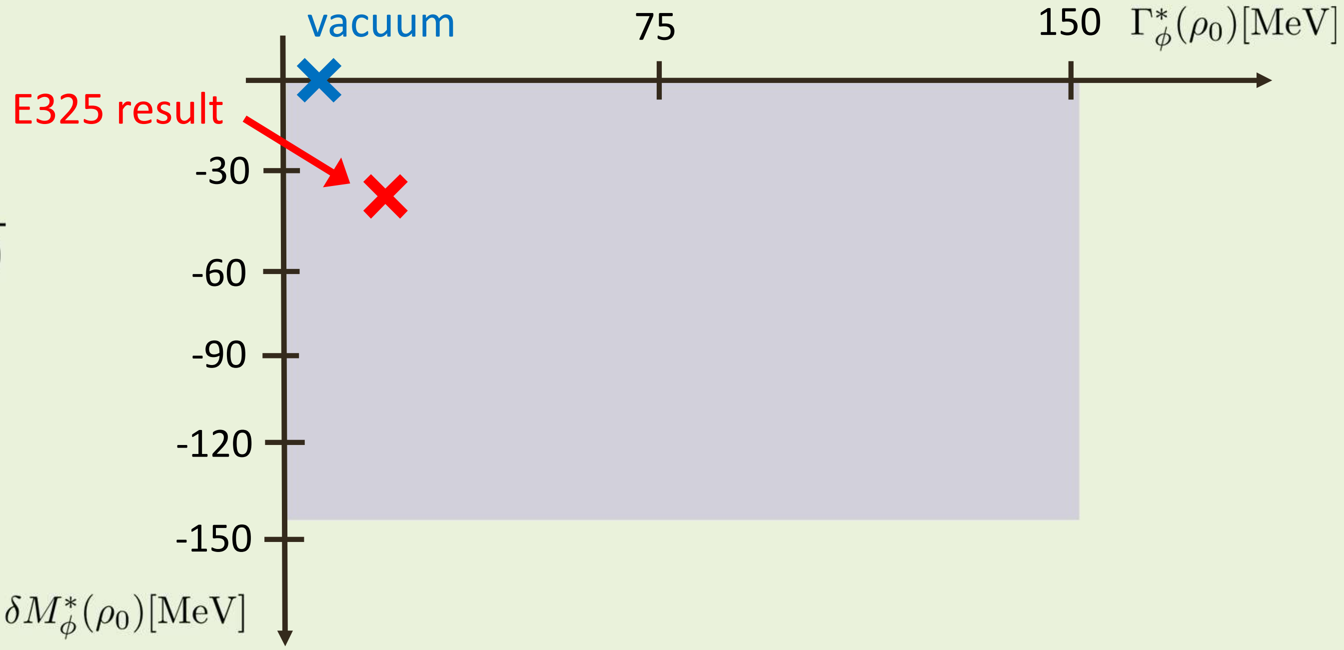
Off-shell dynamics of vector mesons and kaons
(dynamical modification of the mesonic spectral function during the simulated reaction)

Used spectral function:
Relativistic Breit-Wigner with density dependent mass and width

$$C\,\frac{2}{\pi}\,\frac{M^2\Gamma^*_\phi(M,\rho)}{[M^2-M^{*2}_\phi(\rho)]^2+M^2\Gamma^{*2}_\phi(M,\rho)}$$

$$\text{with}\ \left\{\begin{aligned} M^*_\phi(\rho) &= M^{\rm vac}_\phi\left(1-\alpha^\phi\,\frac{\rho}{\rho_0}\right),\\ \Gamma^*_\phi(M,\rho) &= \Gamma^{\rm vac}_\phi+\alpha^\phi_{\rm coll}\,\frac{\rho}{\rho_0}\end{aligned}\right.$$

Simulated scenarios:

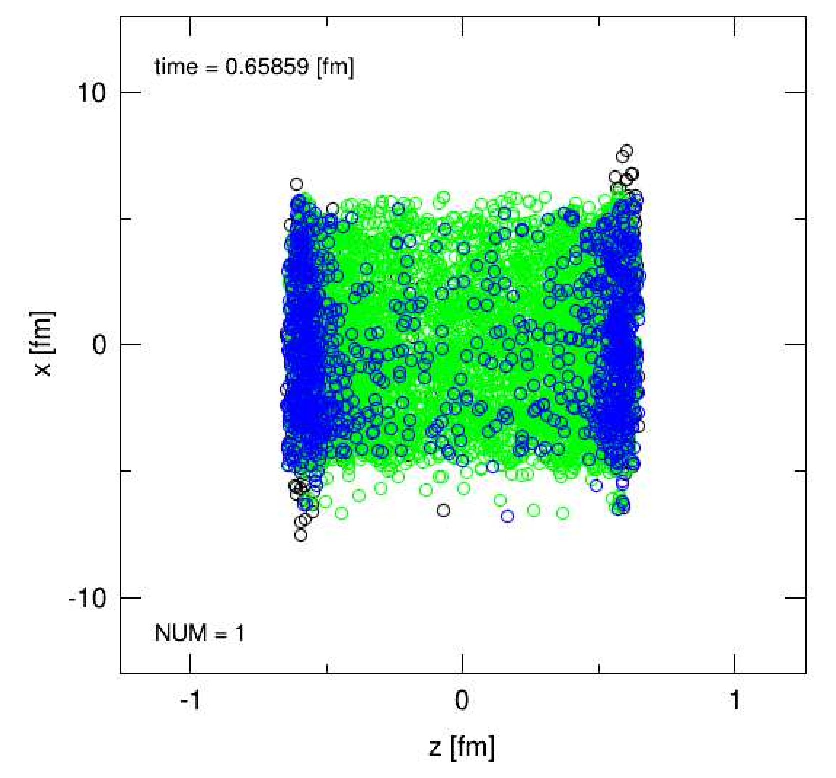


Example of a transport calculation

Au+Au collision at $s^{1/2} = 200$ GeV, b = 2 fm

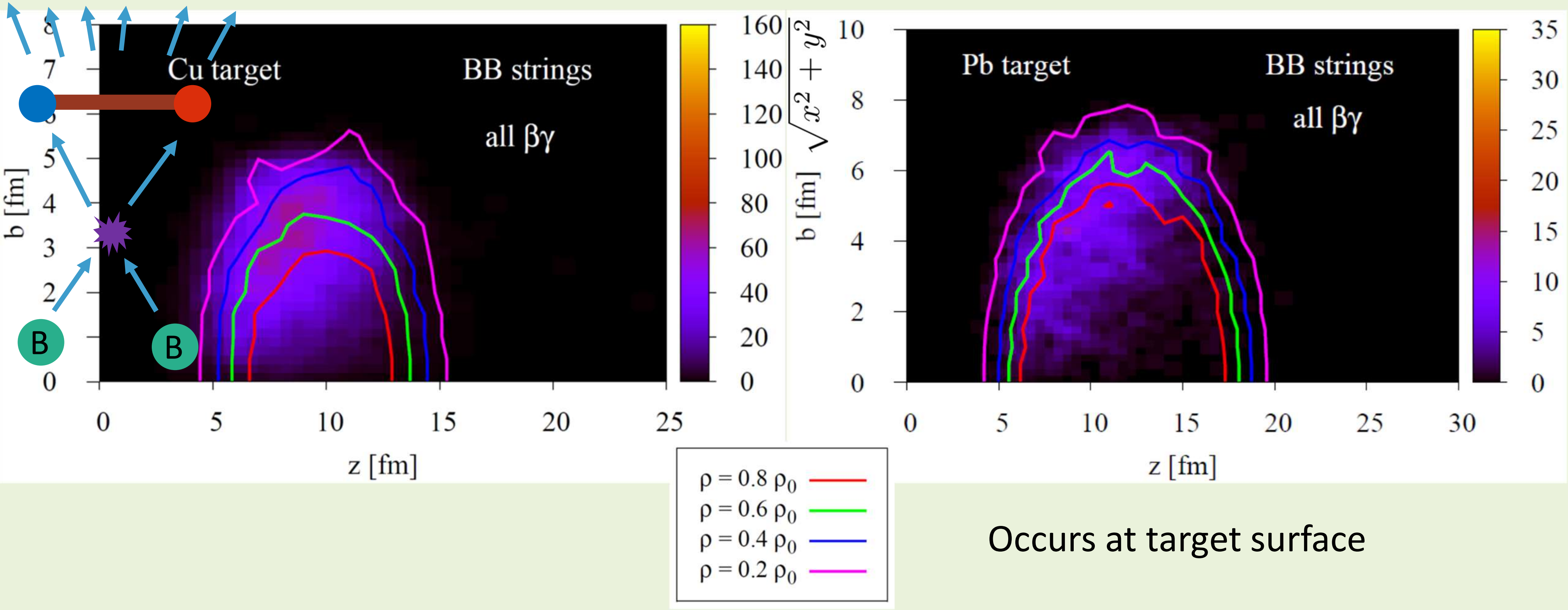
nucleons
quarks
gluons

will not be included in the simulations shown in this talk



How are φ mesons produced in 12 GeV pA collisions?

Production through initial high-energy collisions (via strings)

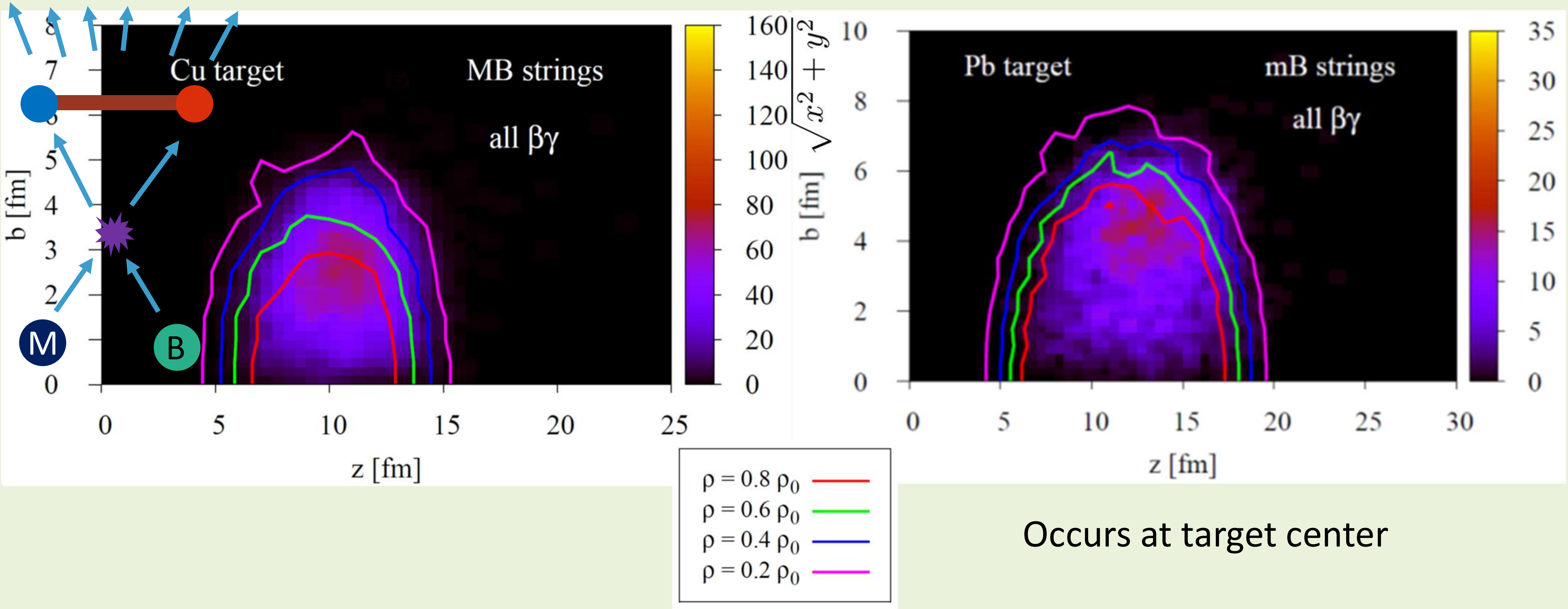


Occurs at target surface

P. Gubler, M. Ichikawa, T. Song and E. Bratkovskaya, 2408.15364 [hep-ph].

How are φ mesons produced in 12 GeV pA collisions?

Production through initial high-energy collisions (via strings)

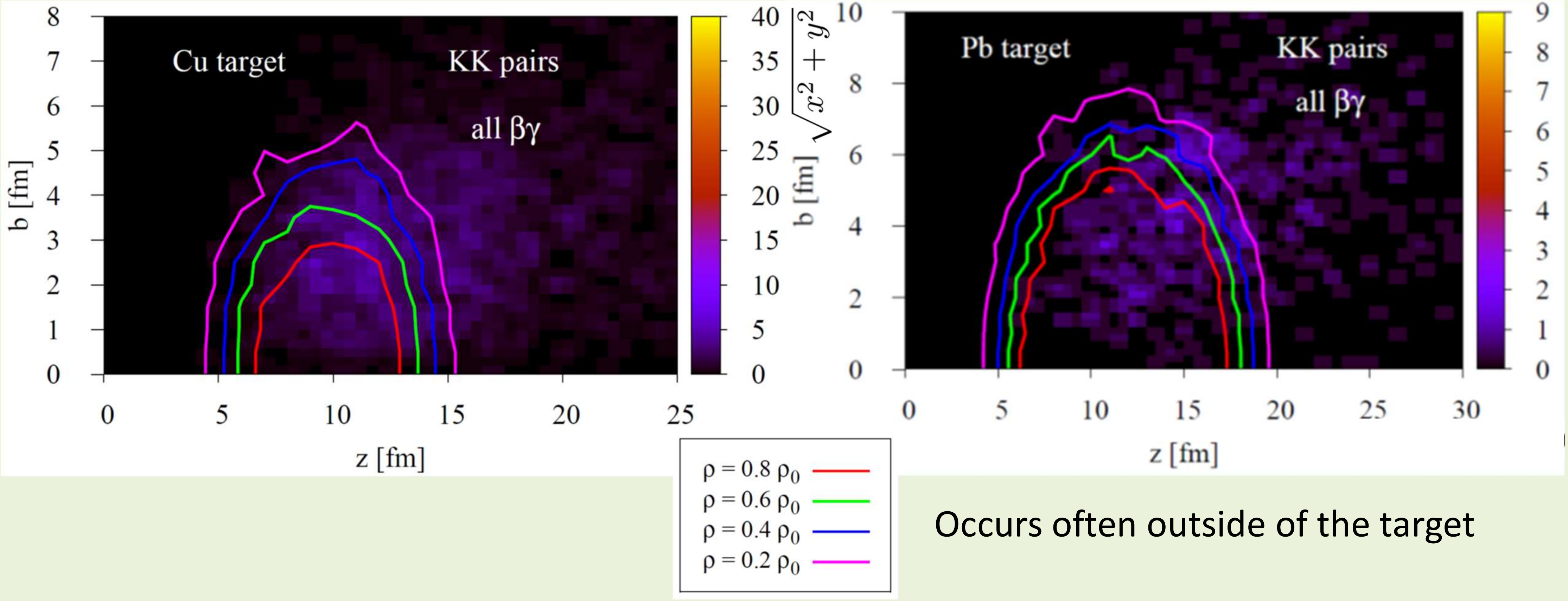


Occurs at target center

P. Gubler, M. Ichikawa, T. Song and E. Bratkovskaya, 2408.15364 [hep-ph].

How are φ mesons produced?

Production through secondary low-energy hadron collisions

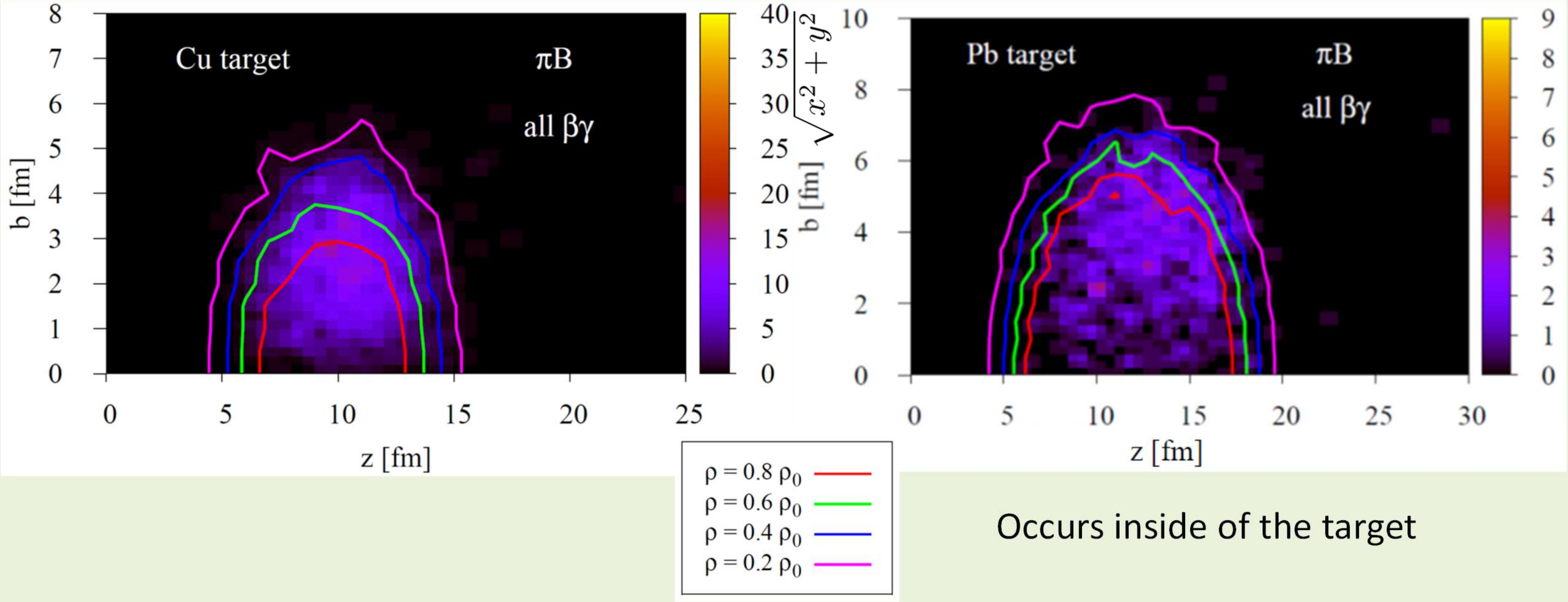


Occurs often outside of the target

P. Gubler, M. Ichikawa, T. Song and E. Bratkovskaya, 2408.15364 [hep-ph].

How are φ mesons produced in 12 GeV pA collisions?

Production through secondary low-energy hadron collisions

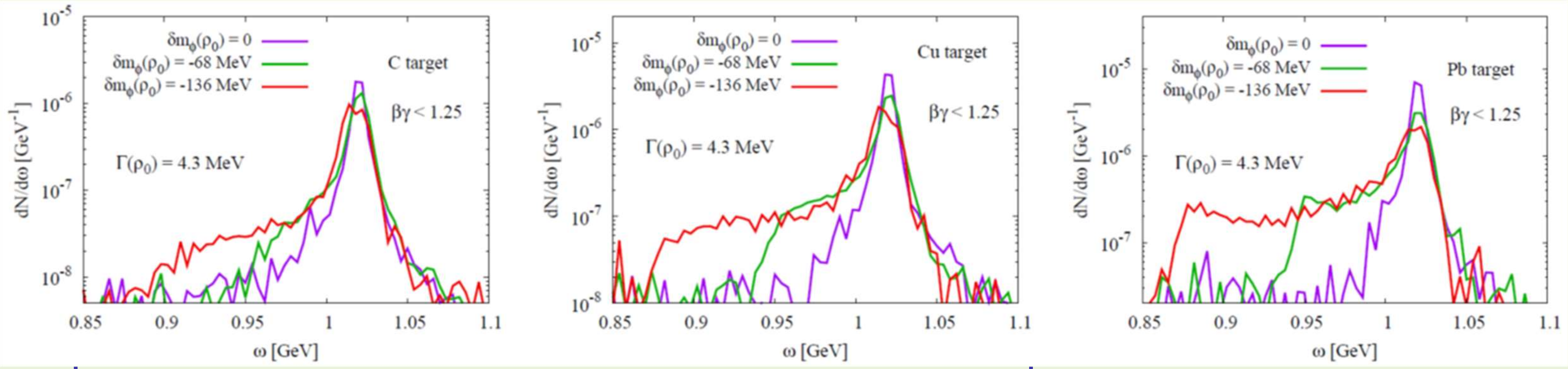


Occurs inside of the target

P. Gubler, M. Ichikawa, T. Song and E. Bratkovskaya, 2408.15364 [hep-ph].

The obtained dilepton spectrum (without experimental effects)

Pure mass shift scenarios (no broadening)



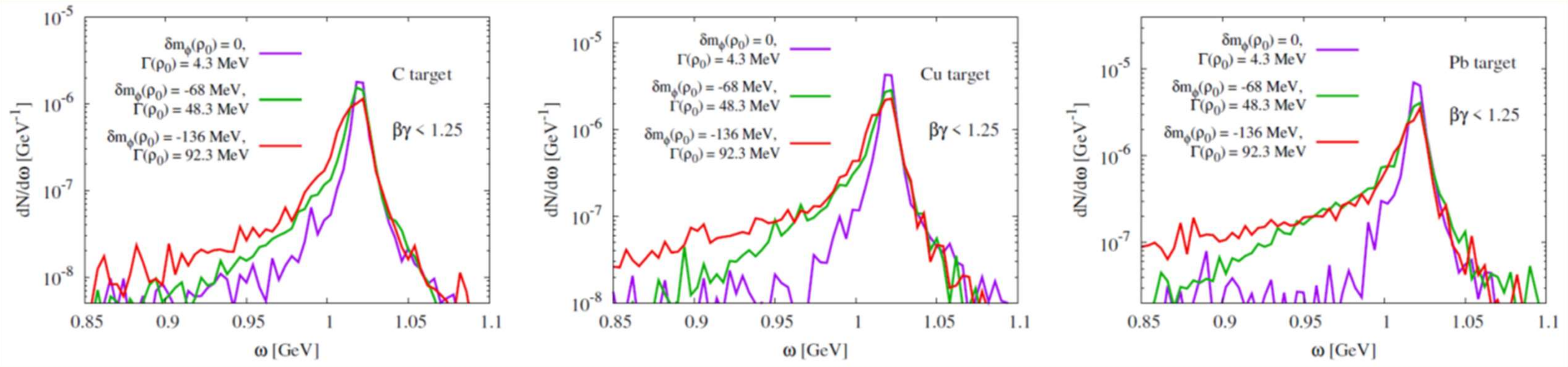
No second peak, but only shoulder structure for mass shift scenarios (even before considering experimental resolution effects)

Secondary peak can be generated for sufficient large mass shift scenario if the target is large enough (Pb here)

P. Gubler, M. Ichikawa, T. Song and E. Bratkovskaya, 2408.15364 [hep-ph].

The obtained dilepton spectrum (without experimental effects)

Mass shift + broadening scenarios



Second peak disappears even for the Pb target with only a shoulder structure left.

P. Gubler, M. Ichikawa, T. Song and E. Bratkovskaya, 2408.15364 [hep-ph].

Summary and conclusions

★ A lot of new theoretical and experimental information about the φN interaction is becoming available (LHC, HAL QCD)

→ **Strong hadronic medium effect?**

★ Several works have by now studied the ALICE Correlation Function data, but the results largely disagree

→ **Need better data?**

More reliable theory?

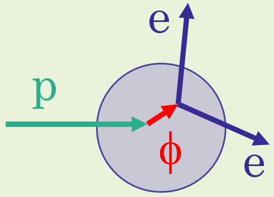
★ With the state-of-the-art PHSD transport approach, we can now study pA reactions more reliably

→ **Many opportunities for new studies and projects!**

Backup slides

Previous experimental results

KEK
E325



12 GeV
pA-reaction

Pole mass:
$$\frac{m_\phi(\rho)}{m_\phi(0)} = 1 - k_1 \frac{\rho}{\rho_0}$$

0.034 ± 0.007

Pole width:
$$\frac{\Gamma_\phi(\rho)}{\Gamma_\phi(0)} = 1 + k_2 \frac{\rho}{\rho_0}$$

2.6 ± 1.5

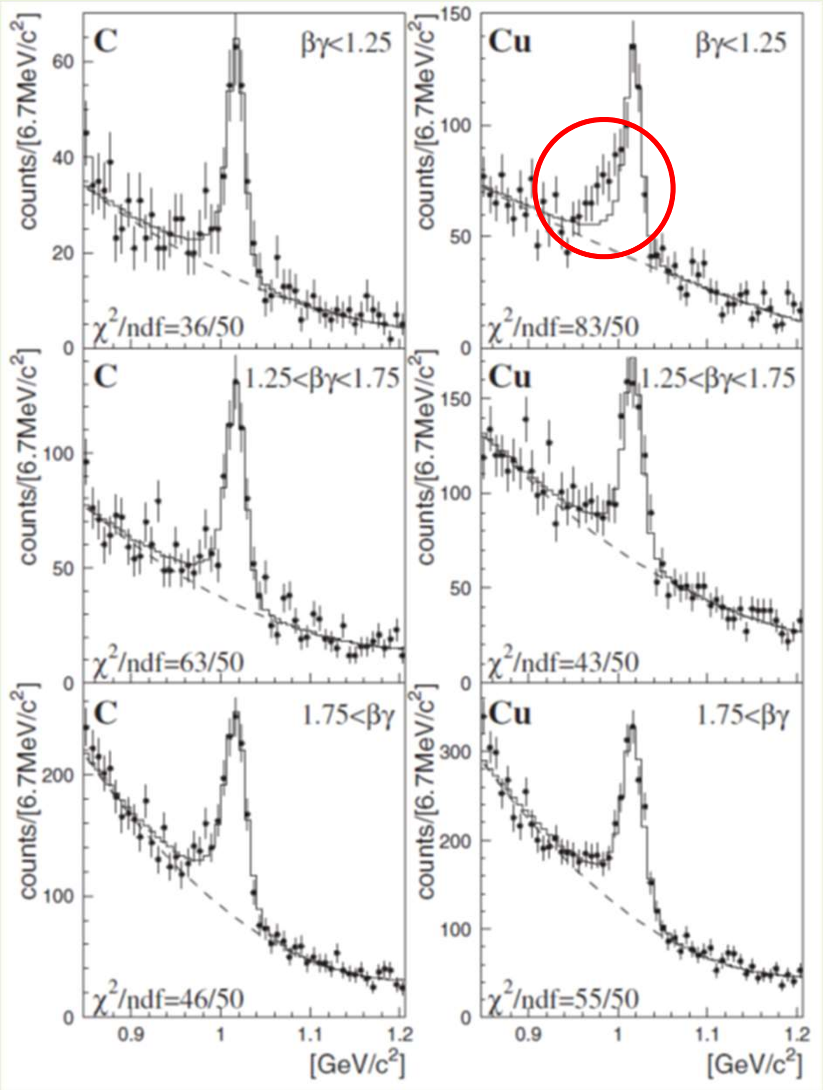
Measurement is being repeated with
~100x increased statistics at the
J-PARC E16 experiment!

See talks by H. Sako, K. Aoki and M. Ichikawa

slow φs

intermediate
φs

fast φs



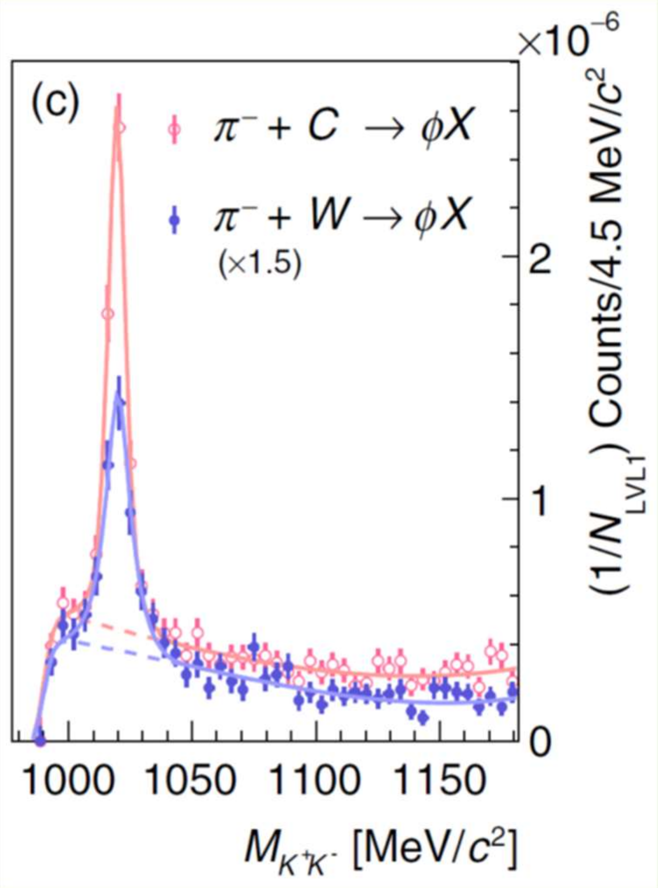
$$\beta\gamma = \frac{|\vec{p}|}{m_\phi}$$

R. Muto et al. (E325 Collaboration), Phys. Rev. Lett. **98**, 042501 (2007).

More recent results

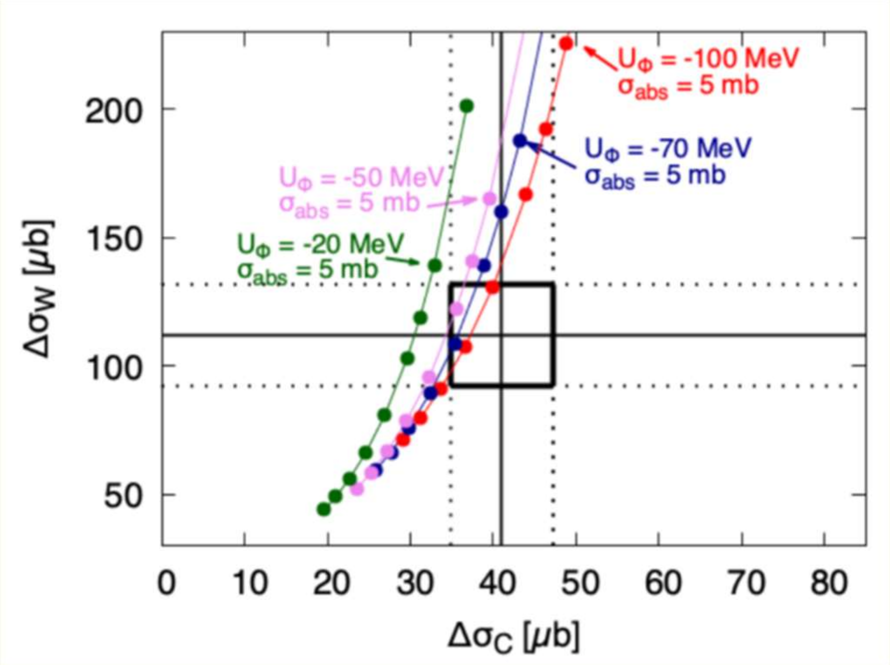
HADES: 1.7 GeV π^-A-reaction

K^+K^- - invariant mass spectrum



J. Adamczewski-Musch et al. (HADES Coll.), Phys. Rev. Lett. **123**, 022002 (2019).

Theoretical analysis of the of the total φ meson production cross section:

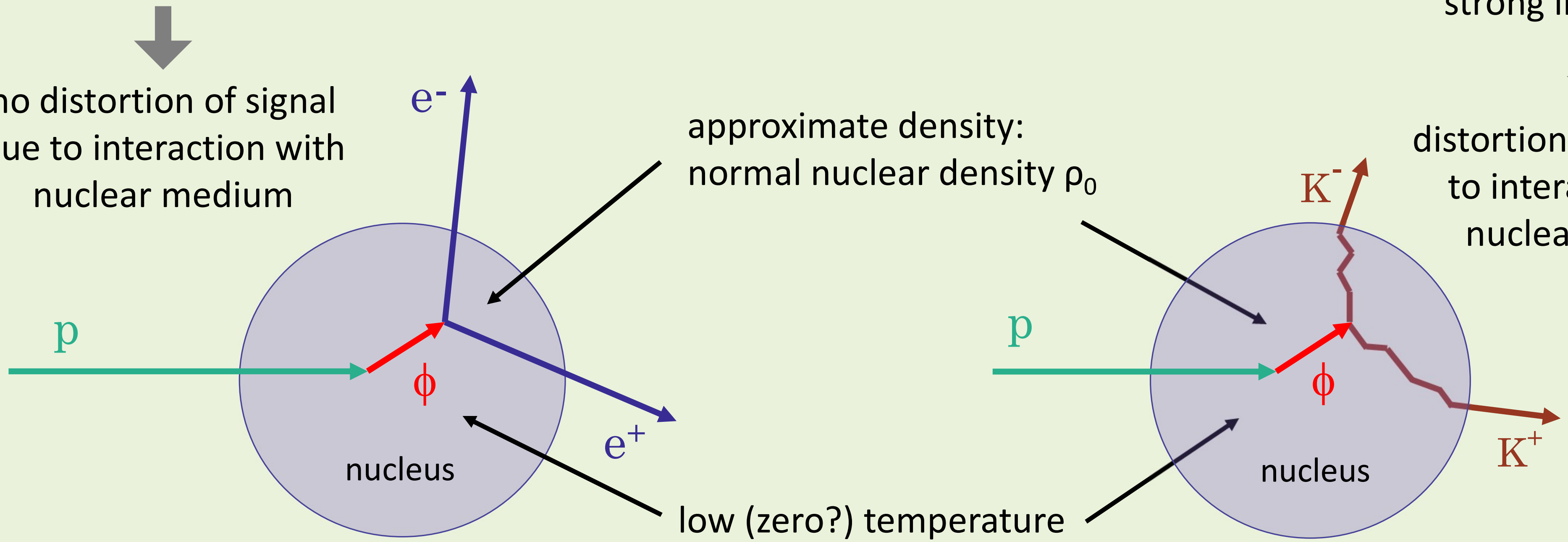


E. Ya. Paryev, Nucl. Phys. A **1032**, 122624 (2023).

- **Attractive φ-nucleus potential: -(50 - 100) MeV**
- **Relatively small imaginary part: 20 – 25 MeV**

The φ meson in pA collisions

Experiments to be discussed in this talk



no strong interaction

no distortion of signal
due to interaction with
nuclear medium

approximate density:
normal nuclear density ρ_0

strong interaction

distortion of signal due
to interaction with
nuclear medium

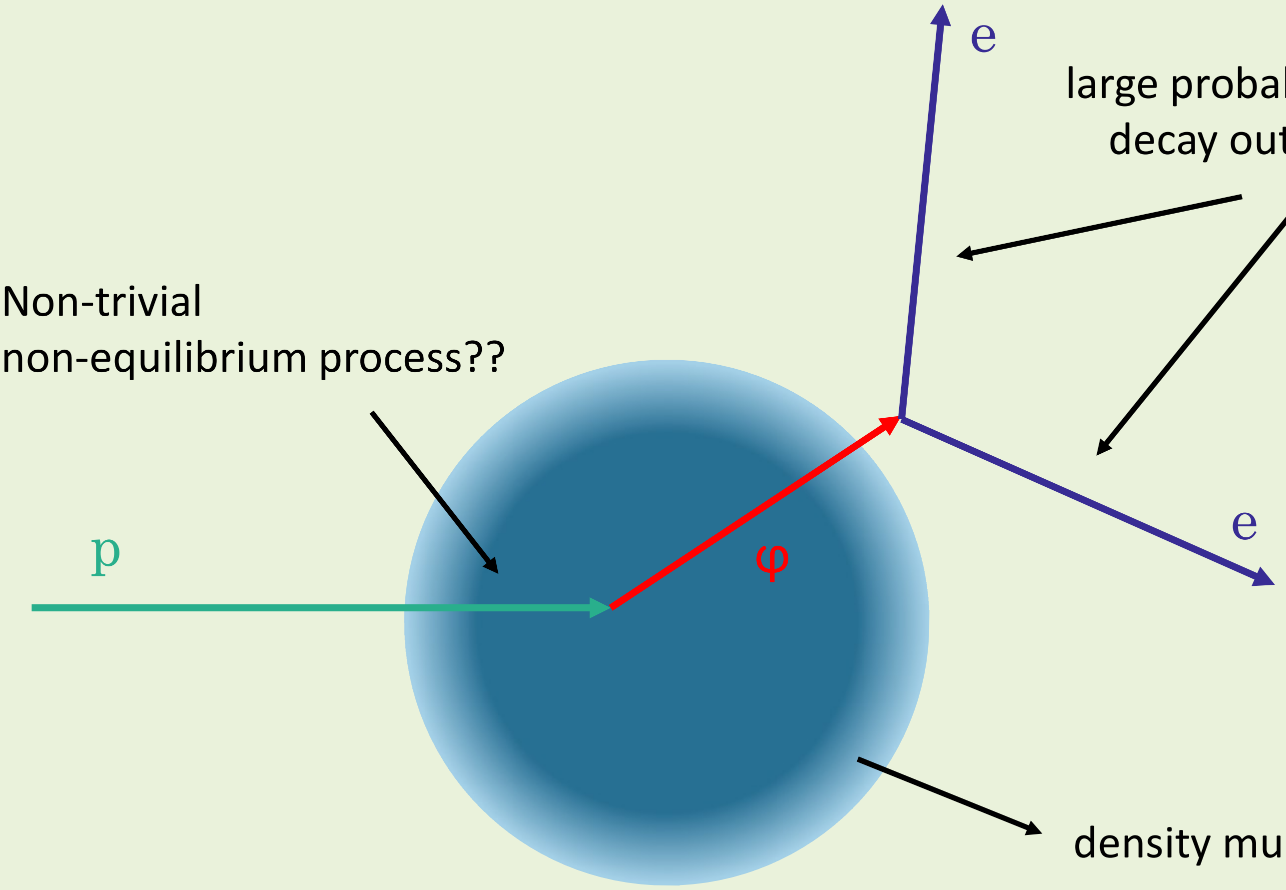
low (zero?) temperature

E325 (KEK)
E16 (J-PARC)

pion beam (HADES)
E88 (J-PARC)

In reality, things are more complicated…

Proton induced generation of vector mesons in nuclei



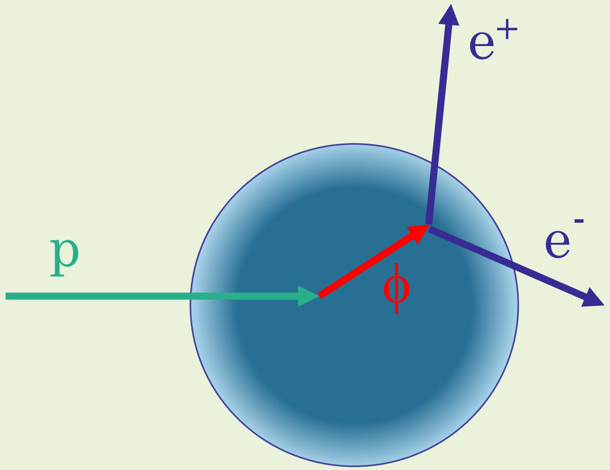
large probability of vector meson decay outside of the nucleus

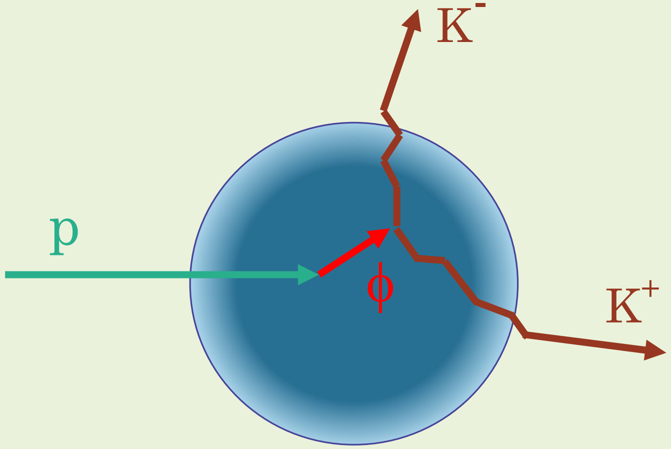
Non-trivial
non-equilibrium process??

density much below ρ_0

Further tasks for theory

Have a good understanding of the production mechanisms of the φ mesons in nuclei from pA reactions.





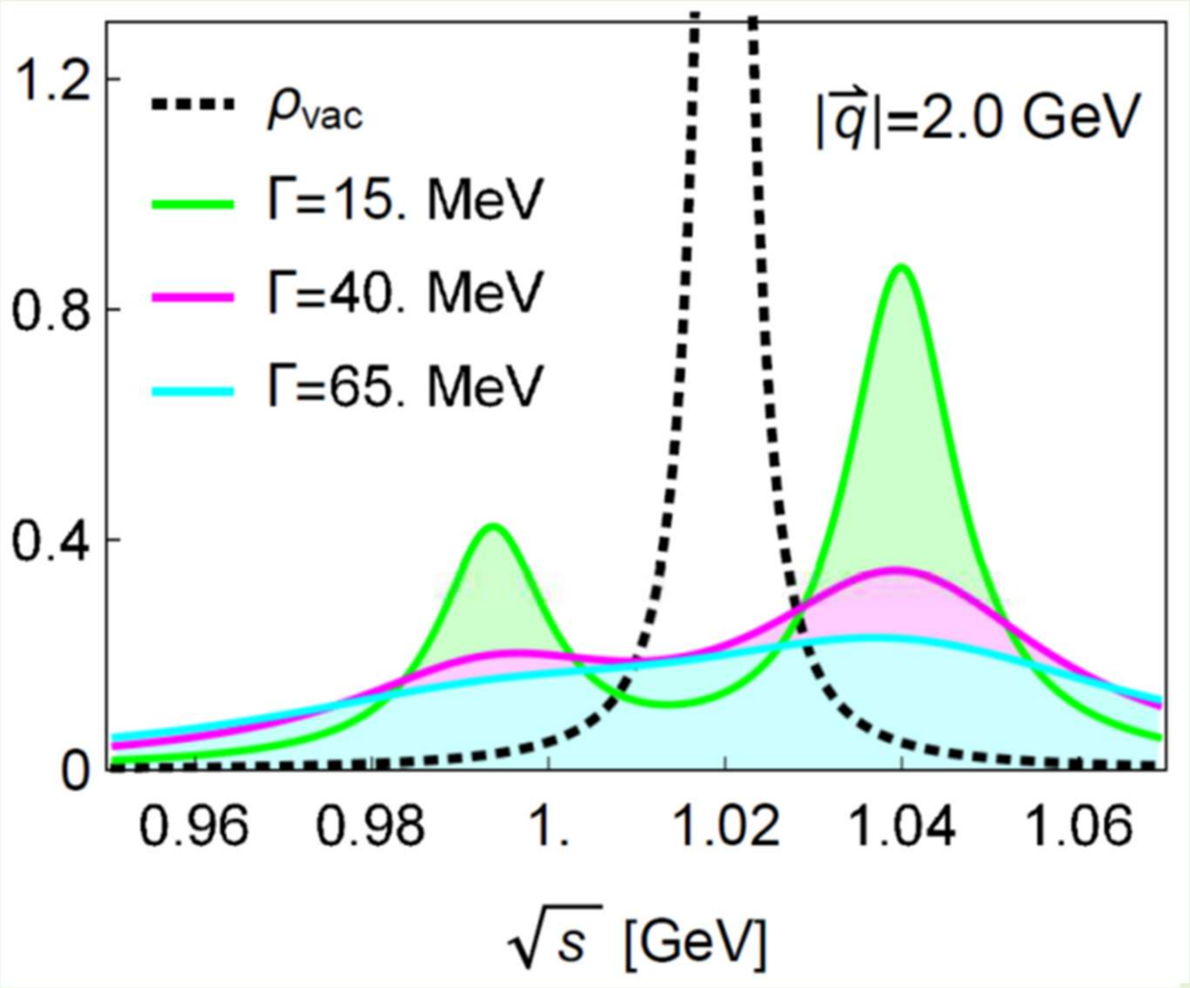
- ★ Where (and at what densities) is the φ meson produced and where does it decay?
- ★ How do the final state interactions of the decay particles influence the decay spectrum (especially for K^+K^-)?

➡ Realistic transport simulations using a transport approach (calculations using the PHSD code are ongoing)

➡ See talk by L. Oliva

The angle-averaged di-lepton spectrum

A double peak?



Computed at normal nuclear matter density

H.J. Kim and P. Gubler, Phys. Lett. B **805**, 135412 (2020).

First application of our formalism

Thermal model with single freeze-out to describe soft hadron production for Au + Au collisions at top RHIC energies:

$$f_V(q, X) = e^{-q^\mu \beta_\mu(x) - \xi(x)}$$

(Jüttner distribution)

$$\beta^\mu = \frac{u^\mu}{T}, \qquad \xi = \frac{\mu}{T}$$

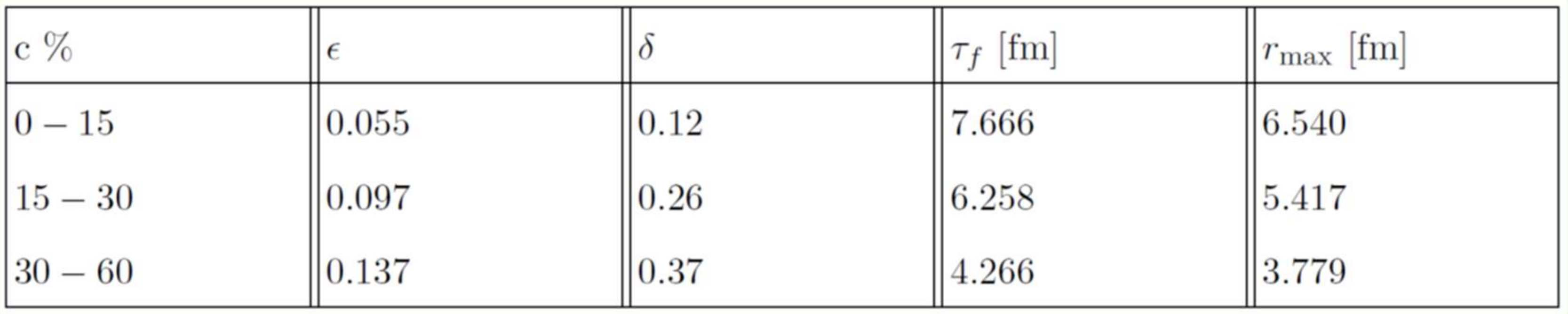
elliptical asymmetry:

$$\left\{ \begin{array}{l} x = r_{\max}\sqrt{1-\epsilon}\cos\phi, \\ y = r_{\max}\sqrt{1-\epsilon}\sin\phi \end{array} \right.$$

elliptical flow: $u^\mu = \frac{1}{N}(t, x\sqrt{1+\delta}, y\sqrt{1-\delta}, z)$

PHOENIX data at

$\sqrt{s_{NN}} = 130$ GeV



c %	ϵ	δ	τ_f [fm]	$r_{\max}$ [fm]
$0-15$	0.055	0.12	7.666	6.540
$15-30$	0.097	0.26	6.258	5.417
$30-60$	0.137	0.37	4.266	3.779

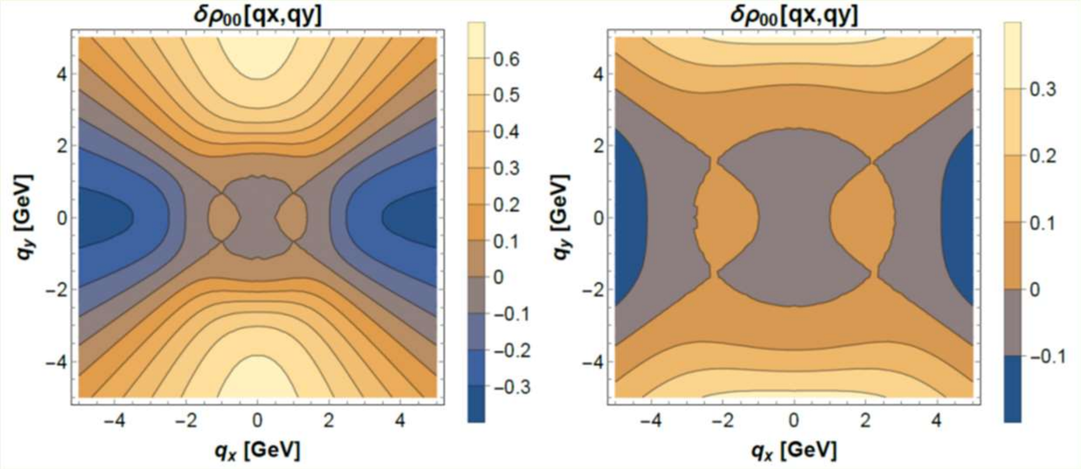
A. Kumar, D.-L. Yang and P. Gubler, 2312.16900 [nucl-th], to be published in PRD.

First application of our formalism

K^{*0} ϕ

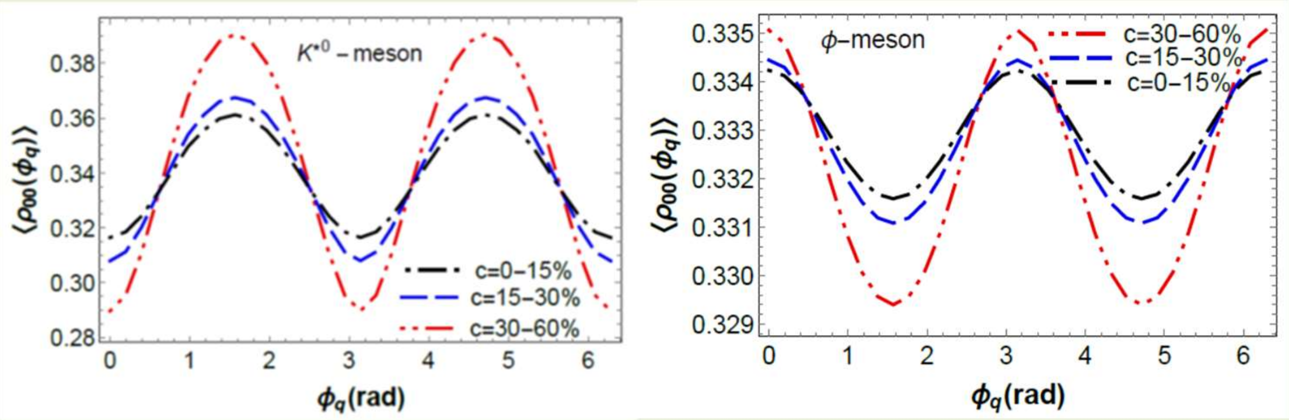
$\rho_{00} - \frac{1}{3}$

$c = 30 - 60\%$



ρ_{00}

Azimuthal angle dependence



A. Kumar, D.-L. Yang and P. Gubler, 2312.16900 [nucl-th], to be published in PRD.

φ meson at rest in nuclear matter

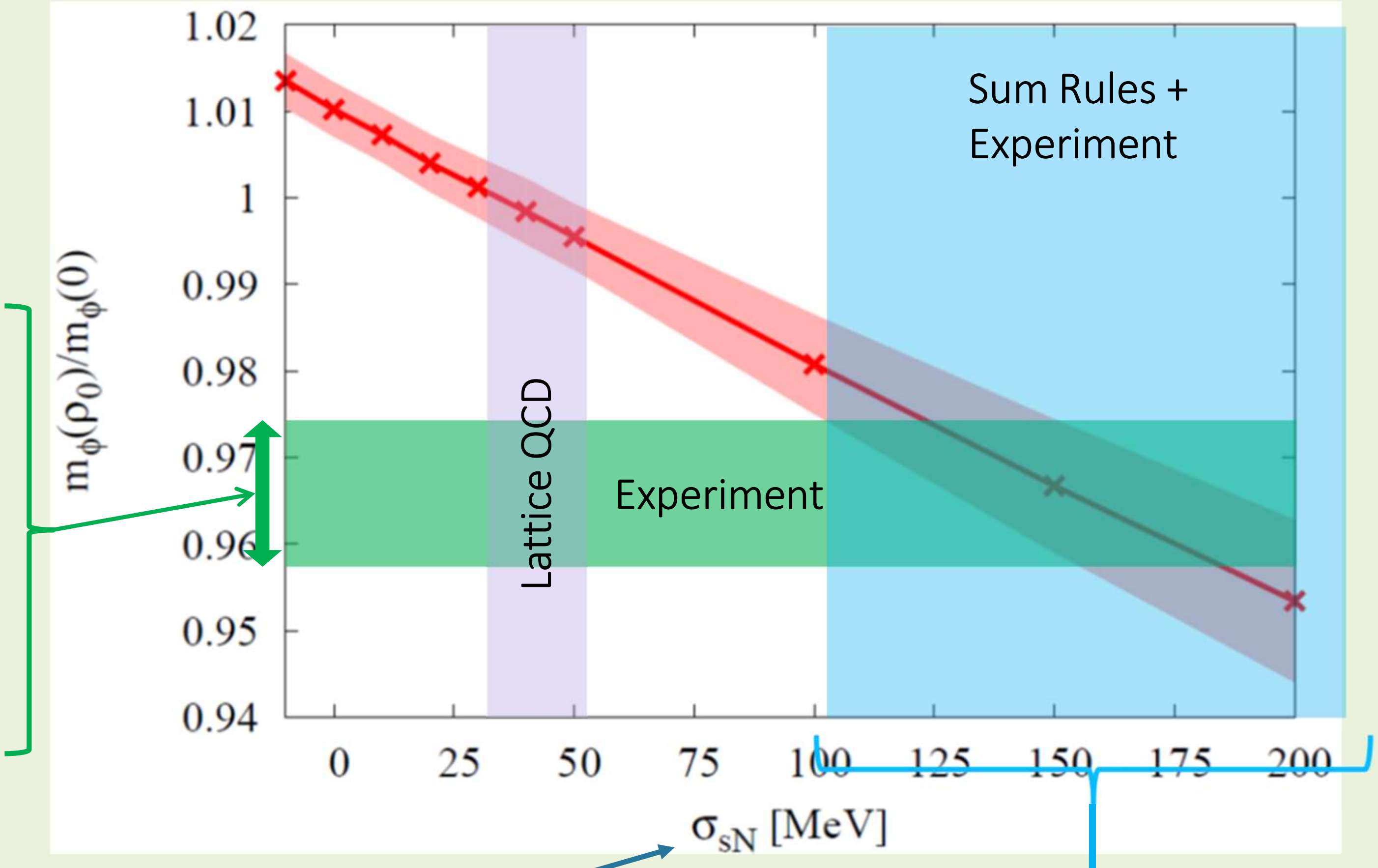
The φ meson mass in nuclear matter probes the strange quark condensate at finite density!

Not consistent?

R. Muto et al.
(KEK, E325 Collaboration),
Phys. Rev. Lett. **98**,
042501 (2007).

Measurement will be
repeated at the
J-PARC E16 experiment
(with 100 times
increased statistics!)

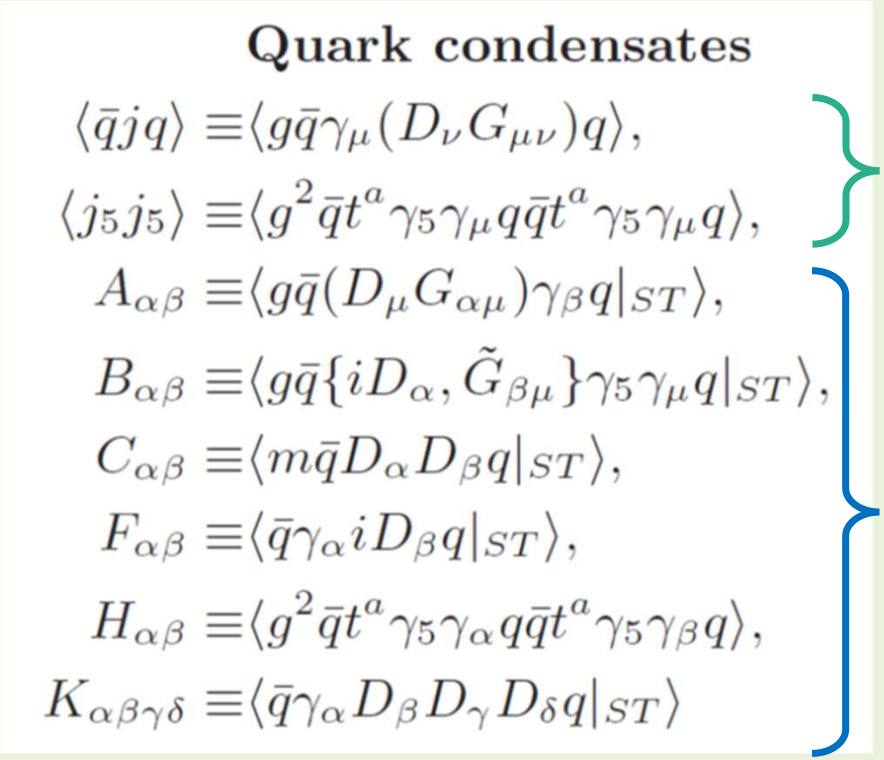
$$\frac{m_\phi(\rho)}{m_\phi(0)} = 0.966 \pm 0.007$$



$$\sigma_{sN} = m_s\langle N|\bar{s}s|N\rangle$$

$$\sigma_{sN} \sim 160 \pm 50 \text{ MeV}$$

Condensates that appear in the vector channel

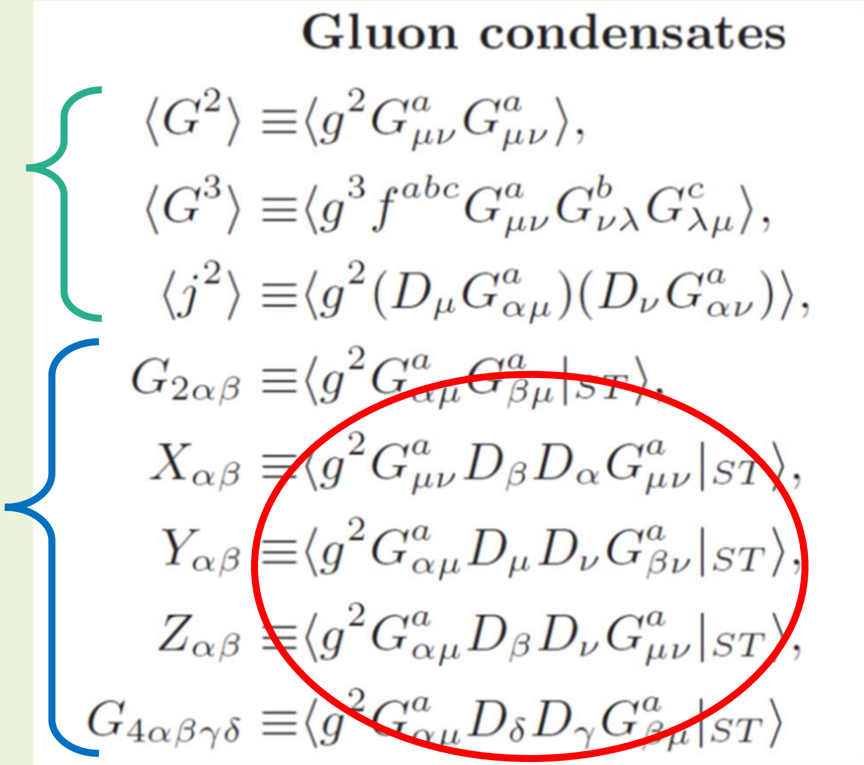


Quark condensates

$$\langle \bar{q}jq\rangle \equiv \langle g\bar{q}\gamma_\mu(D_\nu G_{\mu\nu})q\rangle,$$
$$\langle j_5 j_5\rangle \equiv \langle g^2\bar{q}t^a\gamma_5\gamma_\mu q\bar{q}t^a\gamma_5\gamma_\mu q\rangle,$$
$$A_{\alpha\beta} \equiv \langle g\bar{q}(D_\mu G_{\alpha\mu})\gamma_\beta q|ST\rangle,$$
$$B_{\alpha\beta} \equiv \langle g\bar{q}\{iD_\alpha, \tilde{G}_{\beta\mu}\}\gamma_5\gamma_\mu q|ST\rangle,$$
$$C_{\alpha\beta} \equiv \langle m\bar{q}D_\alpha D_\beta q|ST\rangle,$$
$$F_{\alpha\beta} \equiv \langle \bar{q}\gamma_\alpha iD_\beta q|ST\rangle,$$
$$H_{\alpha\beta} \equiv \langle g^2\bar{q}t^a\gamma_5\gamma_\alpha q\bar{q}t^a\gamma_5\gamma_\beta q\rangle,$$
$$K_{\alpha\beta\gamma\delta} \equiv \langle \bar{q}\gamma_\alpha D_\beta D_\gamma D_\delta q|ST\rangle$$

scalar

non-scalar

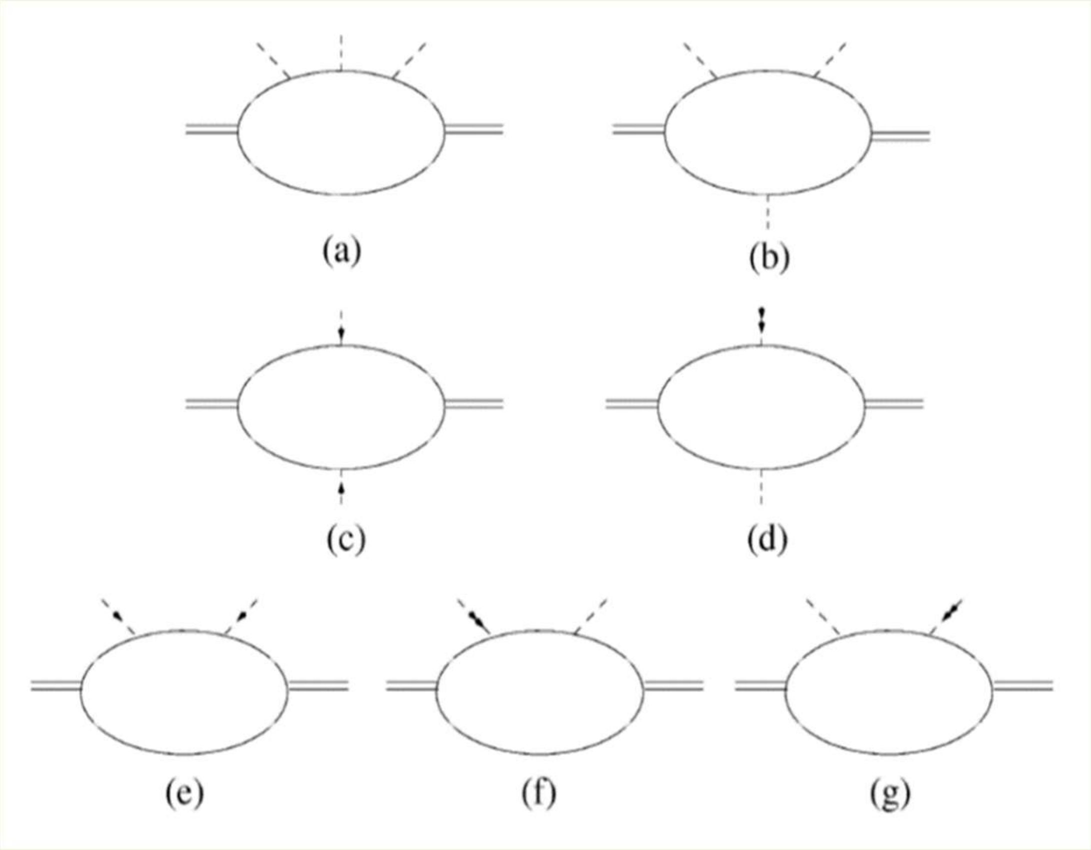


Gluon condensates

$$\langle G^2\rangle \equiv \langle g^2 G^a_{\mu\nu}G^a_{\mu\nu}\rangle,$$
$$\langle G^3\rangle \equiv \langle g^3 f^{abc}G^a_{\mu\nu}G^b_{\nu\lambda}G^c_{\lambda\mu}\rangle,$$
$$\langle j^2\rangle \equiv \langle g^2(D_\mu G^a_{\alpha\mu})(D_\nu G^a_{\alpha\nu})\rangle,$$
$$G_{2\alpha\beta} \equiv \langle g^2 G^a_{\alpha\mu}G^a_{\beta\mu}|ST\rangle,$$
$$X_{\alpha\beta} \equiv \langle g^2 G^a_{\mu\nu}D_\beta D_\alpha G^a_{\mu\nu}|ST\rangle,$$
$$Y_{\alpha\beta} \equiv \langle g^2 G^a_{\alpha\mu}D_\mu D_\nu G^a_{\beta\nu}|ST\rangle,$$
$$Z_{\alpha\beta} \equiv \langle g^2 G^a_{\alpha\mu}D_\beta D_\nu G^a_{\mu\nu}|ST\rangle,$$
$$G_{4\alpha\beta\gamma\delta} \equiv \langle g^2 G^a_{\alpha\mu}D_\delta D_\gamma G^a_{\beta\mu}|ST\rangle$$

Wilson coefficients were not yet available until recently

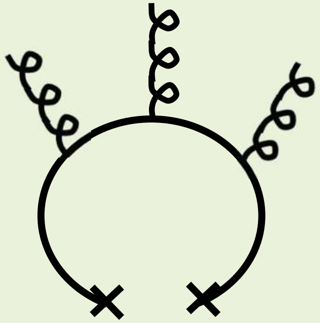
OPE calculation



Mass singularities in chiral limit!

$$\frac{1}{m^2},\ \log\left(\frac{\mu^2}{m^2}\right),\ \ldots$$

Subtract corresponding quark condensate contribution

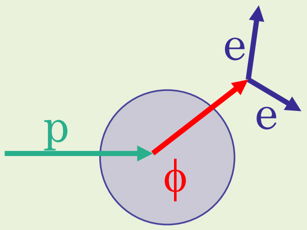


$$\langle \bar{q}\,\Gamma(D,G)q \rangle$$

S. Kim and S.H. Lee, Nucl. Phys. **679**, 517 (2001).

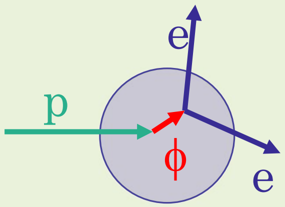
H.J. Kim, P. Gubler and S.H. Lee, Phys. Lett. B **772**, 194 (2017).

Experimental di-lepton spectrum



outside decay

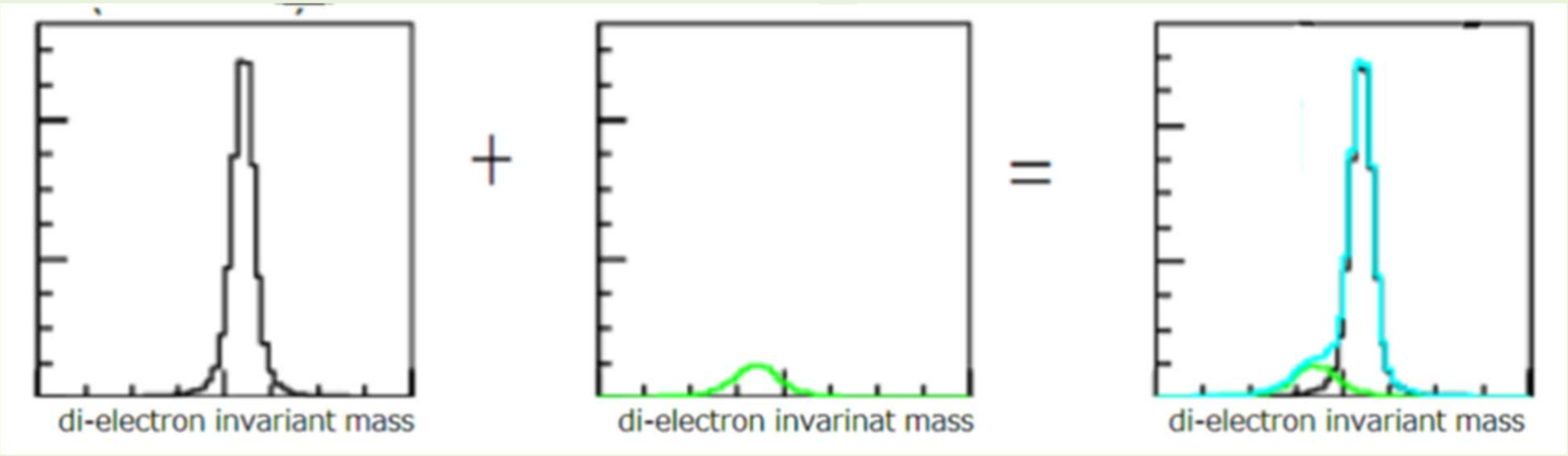
+



inside decay

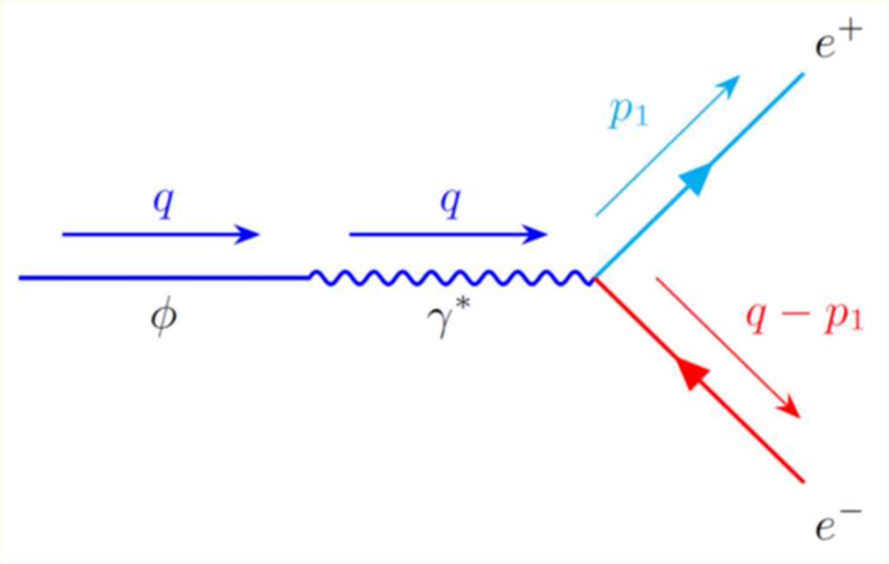
=

Experimentally
observed spectrum

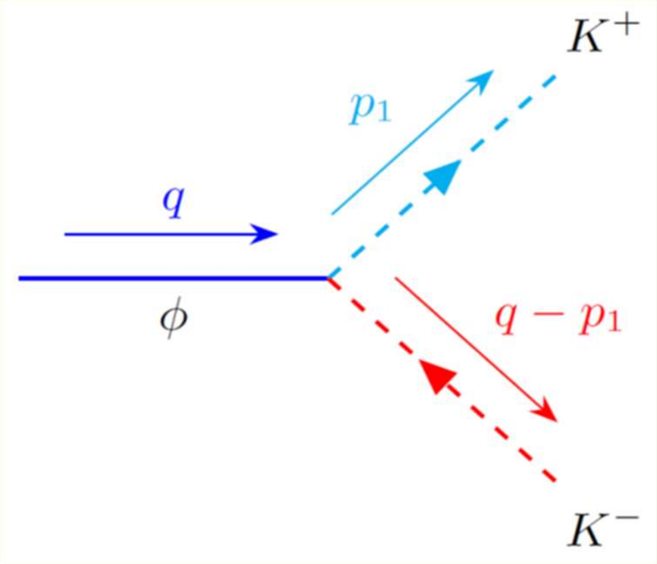


Can the two polarizations be disentangled?

→ Look at the angular distributions of various decay channels

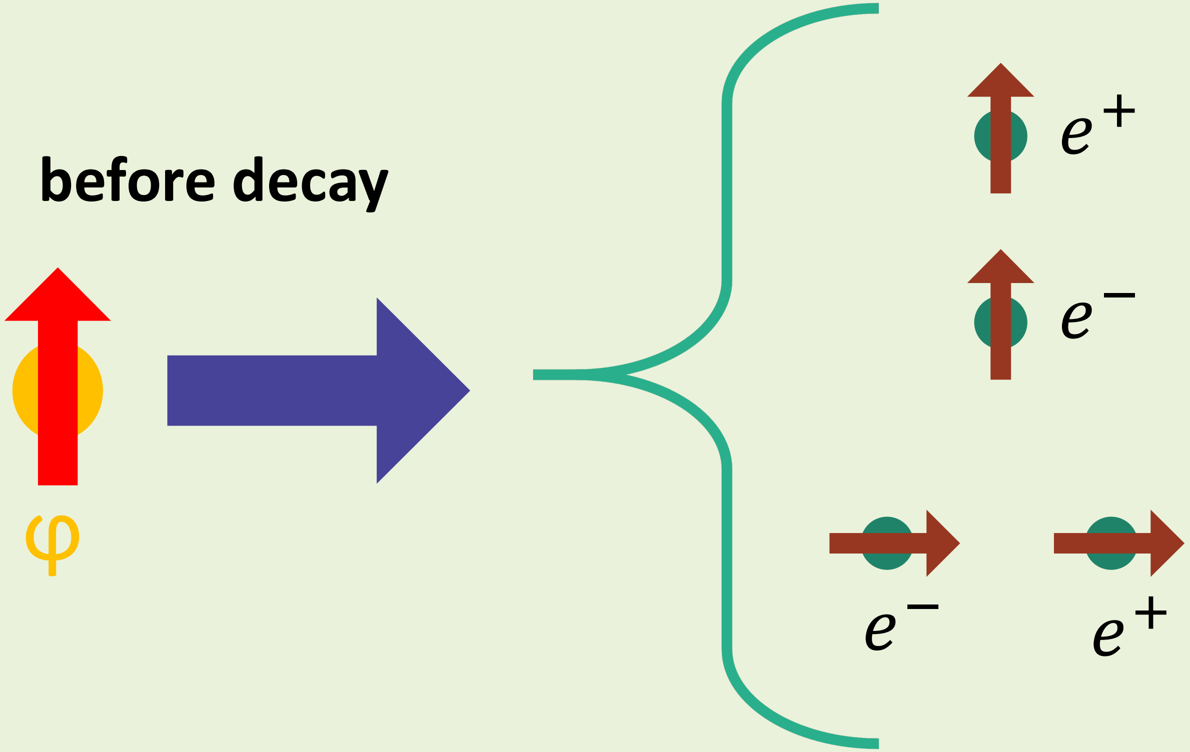


To be measured soon at
the J-PARC E16 experiment



New E88 experiment
at J-PARC (in a few years)

A simple example of dilepton decay of a longitudinally polarized φ



before decay

after decay

decay <u>perpendicular</u> to
the φ meson momentum:
allowed

decay <u>parallel</u> to the
φ meson momentum:
not allowed

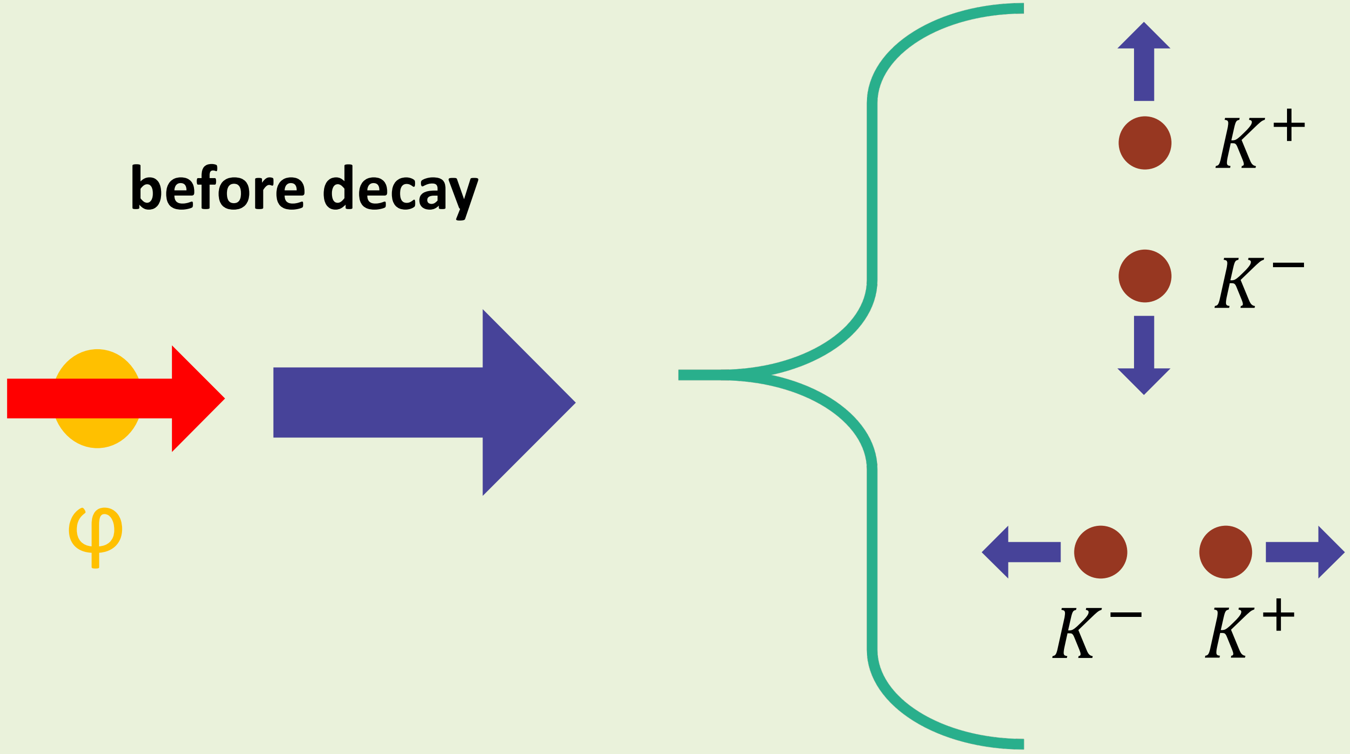
measurement in this direction
can be used to observe purely
transverse mode

A simple example of K^+K^- decay of a transversely polarized φ

$$|Y_1^1(\theta,\phi)|^2 \sim \sin^2\theta$$

before decay

after decay



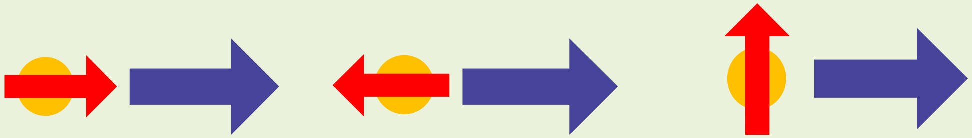
decay perpendicular to
the φ meson momentum:
allowed

decay parallel to the
φ meson momentum:
not allowed

↓

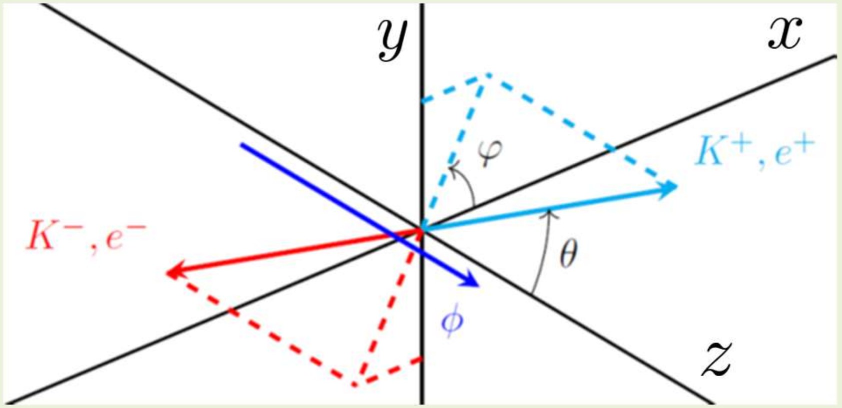
measurement in this direction
can be used to observe purely
longitudinal mode

Full angular distribution of dilepton decay



Initial polarization: $$|V\rangle = \underbrace{a_{+1}|+1\rangle + a_{-1}|-1\rangle}_{\text{Transverse polarization}} + a_0|0\rangle$$

Longitudinal polarization



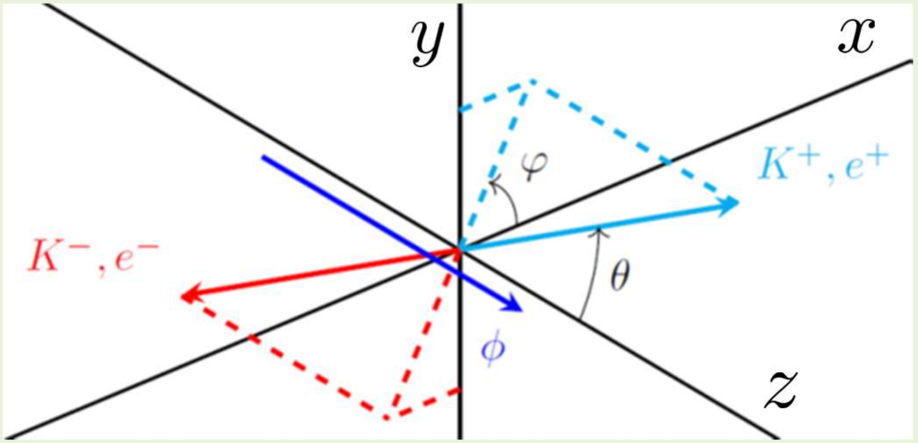
θ: polar angle

ϕ: azimuthal angle

$$\frac{1}{\Gamma}\frac{d\Gamma}{d\Omega} = \frac{3}{16\pi}\Big[(|a_{+1}|^2 + |a_{-1}|^2)(1+\cos^2\theta) + 2|a_0|^2(1-\cos^2\theta) \\ + 2\mathrm{Re}(a_{+1}a_{-1}^*)\sin^2\theta\cos 2\phi + \ldots\Big]$$

other ϕ-dependent terms

Full angular distribution of dilepton decay



θ: polar angle
ϕ: azimuthal angle

With

$$|a_{+1}|^2 + |a_{-1}|^2 + |a_0|^2 = 1, \;\; |a_0|^2 = \rho_{00}$$

00-component of spin-density matrix

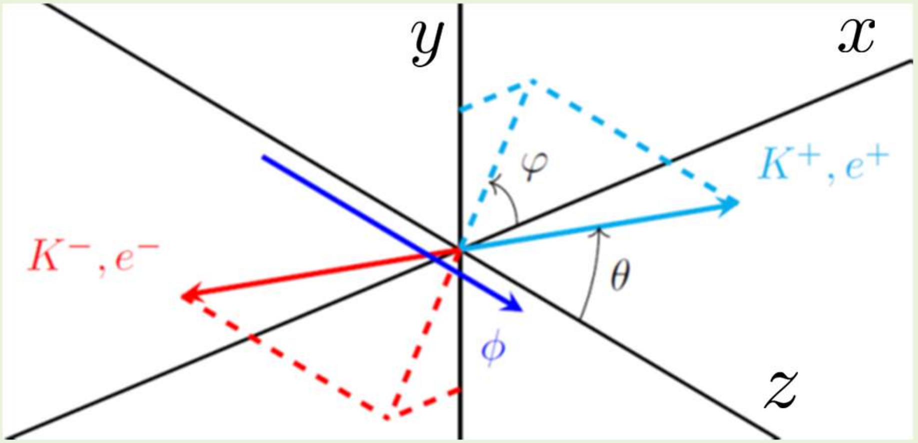
$$\Rightarrow \frac{1}{\Gamma}\frac{d\Gamma}{d\Omega} = \frac{3}{16\pi}\left[1 + \cos^2\theta + \rho_{00}(1 - 3\cos^2\theta) + \ldots\right]$$

ϕ-dependent terms

$$\Rightarrow \rho_{00} = \frac{1}{3}$$

Unpolarized case: vanishing θ-dependence

Full angular distribution of K^+K^- decay



θ: polar angle
ϕ: azimuthal angle

$$\frac{1}{\Gamma}\frac{d\Gamma}{d\Omega} = \frac{3}{16\pi}\Big[\overbrace{(|a_{+1}|^2 + |a_{-1}|^2)}^{\text{Transverse modes}}\sin^2\theta + \overset{\text{Longitudinal mode}}{2|a_0|^2\cos^2\theta}$$

$$-2\mathrm{Re}(a_{+1}a^*_{-1})\sin^2\theta\cos 2\phi + \ldots\Big]$$

$$= \frac{3}{16\pi}\left[1 - \cos^2\theta - \rho_{00}(1 - 3\cos^2\theta) + \ldots\right]$$

ϕ-dependent terms