

# Topological symmetries and their gaugings in 2d CFT and 3d TFT

Ingo Runkel (Hamburg Univ.)

## Outline

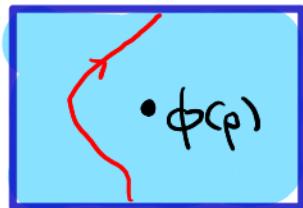
Part I Topological line defects in 2d QFT & 3d TFT

Part II Applications

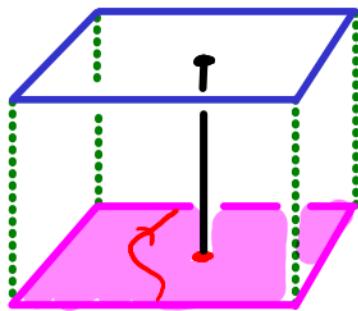
- Gauging topol. symm. in 2d QFT
- Gauging as an equiv. rel.
- Non-topol. defects & integrable flows

## Part I

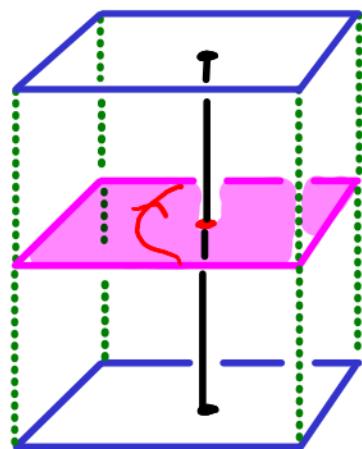
2d QFT with  
topological line  
defects



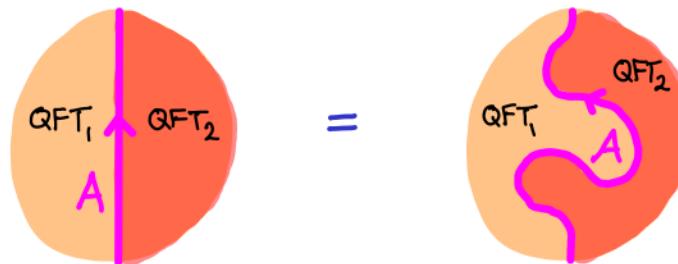
Symmetry topol.  
field theory  
("SymTFT")



Chiral TFT



# Topological defects

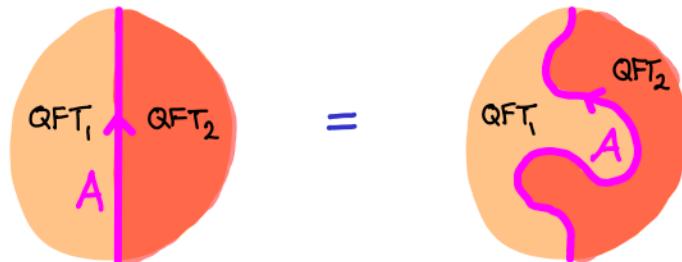


For now :  $QFT_1 = QFT_2$

The symmetries of a quantum mech. system form a group.

The topol. line defects of a 2d QFT form a  .

# Topological defects



For now :  $QFT_1 = QFT_2$

The symmetries of a quantum mech. system form a group.

The topol. line defects of a 2d QFT form a ↗.

pirital monoidal category

Fuchs, Schweigert, IR '02

Fröhlich, Fuchs, Schweigert, IR '09

Davydov, Kong, IR '11

Thorngren, Wang '19

## Fusion categories

Assume :

- finite list of elementary top. def.  $\{1, a_2, \dots, a_n\}$

- closed under fusion

$$a \otimes b \simeq \bigoplus c$$

$$\text{Diagram: } \begin{array}{c} \text{Two vertical pink arrows labeled } a \text{ and } b \text{ pointing up through a light blue rounded rectangle.} \\ = \sum_c \end{array}$$

- quantum dimensions

$$\text{Diagram: } \begin{array}{c} \text{Two light blue rounded rectangles, each containing a circle with a pink arrow labeled } a. \\ = \end{array}$$

spherical  
fusion category

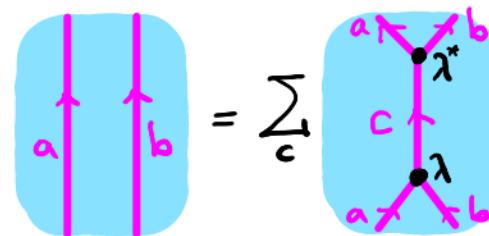
# Fusion categories

Assume :

- finite list of elementary top. def.  $\{1, a_2, \dots, a_n\}$

- closed under fusion

$$a \otimes b \simeq \bigoplus c$$



- quantum dimensions

assume 0 or 1 dim.  
coupling spaces

spherical  
fusion category

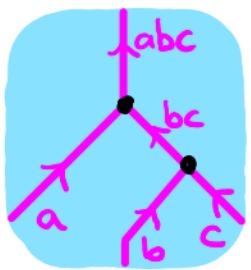
## ... fusion categories

F-matrices : Change of basis in  $a \otimes b \otimes c \rightarrow d$

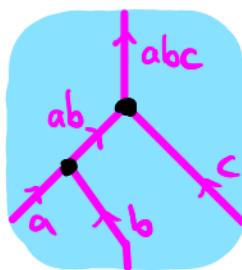
$$\begin{array}{c} \text{Diagram 1: } \begin{array}{c} \text{Top node: } d \\ \text{Bottom nodes: } a, b, c \\ \text{Arrows: } a \rightarrow p, b \rightarrow q, c \rightarrow r \end{array} \\ \xrightarrow{\quad \quad} \\ \text{Diagram 2: } \begin{array}{c} \text{Top node: } d \\ \text{Bottom nodes: } p, q, r \\ \text{Arrows: } p \rightarrow a, q \rightarrow b, r \rightarrow c \end{array} \\ = \sum_q F_{pq}^{(abc)d} \\ \xleftarrow{\quad \quad} \\ \text{Diagram 3: } \begin{array}{c} \text{Top node: } d \\ \text{Bottom nodes: } a, b, c \\ \text{Arrows: } a \rightarrow p, b \rightarrow q, c \rightarrow r \end{array} \end{array}$$

## ... fusion categories

Group symmetry :  $a, b, \dots \in G$  finite group  
 $a \otimes b \simeq c$  (no sum)



$$= \omega(a, b, c) \cdot$$



$$[\omega] \in H^3(G, \mathbb{C}^\times)$$

obstruction to setting  $\omega = 1$

↪ obstruction to gauging  $G$

## ... fusion categories

E.g.  $G = \mathbb{Z}_2$      $H^3(\mathbb{Z}_2, \mathbb{C}^\times) \simeq \mathbb{Z}_2$      $\{d: \mathbb{Z}_2 \rightarrow \mathbb{C}^\times\} \simeq \mathbb{Z}_2$     } 4 choices

Write multiplicatively :  $\mathbb{Z}_2 \simeq \{\pm 1\}$

	$\omega(\perp, \perp, \perp) = 1$	$\omega(\perp, \perp, \perp) = -1$
$d(\perp) = 1$	$\text{Vect}_{\mathbb{Z}_2}$	$\text{Vect}_{\mathbb{Z}_2}^\omega$
$d(\perp) = -1$	$\text{SVect}$	$\text{SVect}^\omega$

## ... fusion categories

$$\text{E.g. } G = \mathbb{Z}_2 \quad H^3(\mathbb{Z}_2, \mathbb{C}^\times) \simeq \mathbb{Z}_2 \quad \left\{ d : \mathbb{Z}_2 \rightarrow \mathbb{C}^\times \right\} \simeq \mathbb{Z}_2 \quad \left. \right\} \text{4 choices}$$

Write multiplicatively :  $\mathbb{Z}_2 \simeq \{\pm 1\}$

	$\mathbb{Z}_2$ can be gauged $\omega(\perp, \perp, \perp) = 1$	$\mathbb{Z}_2$ cannot be gauged $\omega(\perp, \perp, \perp) = -1$
$d(\perp) = 1$	$\text{Vect}_{\mathbb{Z}_2}$ gauging well-defined on oriented surfaces	$\text{Vect}_{\mathbb{Z}_2}^\omega$
$d(\perp) = -1$	$\text{SVect}$ gauging needs a spin structure	$\text{SVect}^\omega$

## ... fusion categories

E.g.: simplest non-group case : Fib

- simple obj  $\mathbb{1}, \varphi$
- tensor  $\varphi \otimes \varphi \simeq \mathbb{1} \oplus \varphi$
- F-matrix : 2 choices
- dim's : 1 choice

Spherical fusion cat. classified for :

$$\#(\text{simple obj}) =$$

## ... fusion categories

E.g.: simplest non-group case : Fib

- simple obj     $\mathbb{1}, \varphi$
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- F-matrix : 2 choices
- dim's            : 1 choice

Spherical fusion cat. classified for :

Ostrik '02, '13

$$\#(\text{simple obj}) = 1, 2, 3$$

# Symmetry topological field theory

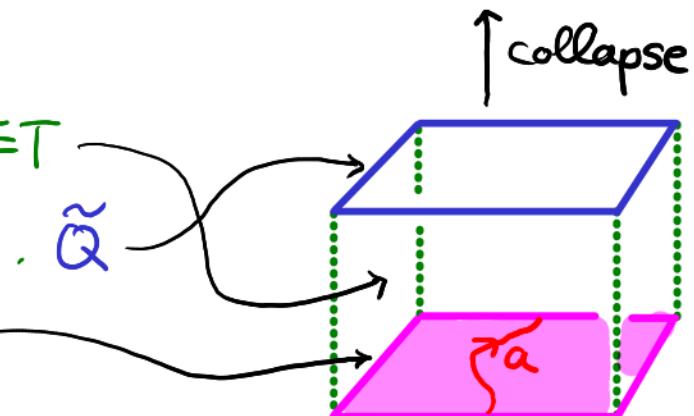
Gaiotto, Kulp '20  
Freed, Moore,  
Teleman '22  
...

Q 2dQFT with spherical fusion  
cat.  $\mathcal{F}$  of topol. line defects



Get

- $Z_{\mathcal{F}}$  3d state sum TFT
- non-top. bnd. cond.  $\tilde{\mathcal{Q}}$
- top. bnd. cond.  $\tilde{\mathcal{F}}$

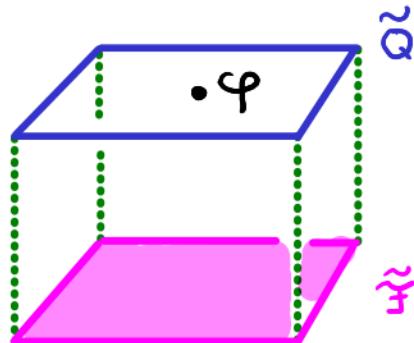


## ... Sym TFT

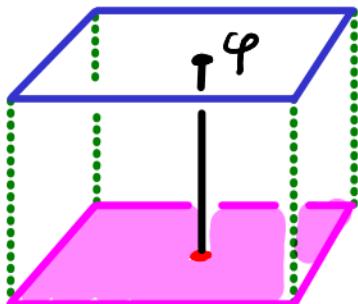
Bulk fields separate

- commute with all  $a \in \mathcal{F}$

$$\bullet \varphi^a = \bullet \varphi_a$$



- do not commute with all  $a \in \mathcal{F}$

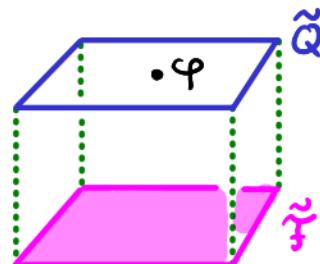


## Chiral TFT

Special to 2d CFT : E.g. Ising CFT

choice of $\mathcal{F}$	bulk field commutes ?		
	$1L$ $(0,0)$	$\varepsilon$ $(\frac{1}{2}, \frac{1}{2})$	$6$ $(\frac{1}{6}, \frac{1}{6})$
id	✓	✓	✓
$\mathbb{Z}_2$	✓	✓	✗
$\mathbb{Z}_2 \cup \{\text{KW}\}$	✓	✗	✗

only  $(\text{hol}) \otimes (\text{antihol})$  lives on  $\tilde{\mathcal{Q}}$ .



## ... chiral TFT

In rational CFT :

- 1) If only  $(\text{hol}) \otimes (\text{antihol})$  lives on  $\tilde{\mathbb{Q}}$ ,  
can "unfold" Sym TFT  $\rightarrow$  chiral TFT
- 2) There is a choice of  $\mathcal{F}$  s.t. 1) applies

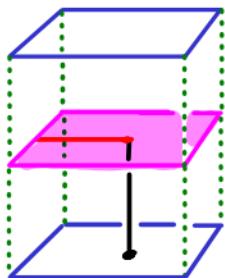
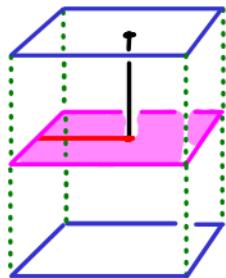
$Z_{\mathcal{F}}$  =  $C \boxtimes \bar{C}$

Fuchs, Schweigert, IR '02, ...

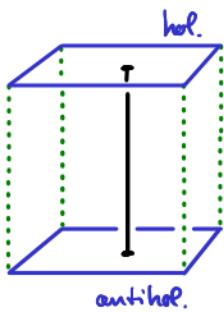
## ... chiral TFT

Nice :

- geometric separation  
into hol./antihol.

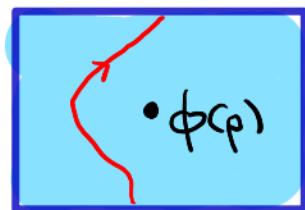


- very easy for identifying surf. def.

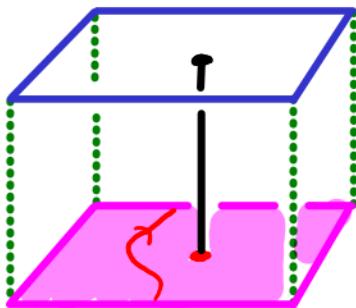


# Part I summary :

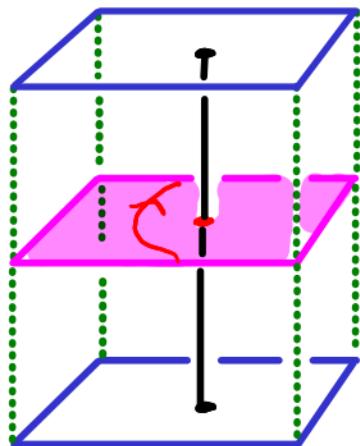
2d RCFT with  
topological line  
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Symmetry topol.  
field theory  
("SymTFT")



Chiral TFT



# Now : Part II – Applications

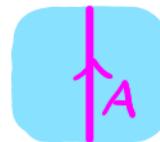
# Gauging topological symmetries in 2d

## aka generalised orbifolds

Fröhlich, Fuchs, Schweigert, IR '09  
Carqueville, IR '12

$Q$  : 2d QFT

$A$  : topol. def. in  $Q$   
(not nec. elementary)



Want to define gauged theory  $Q/A$ .

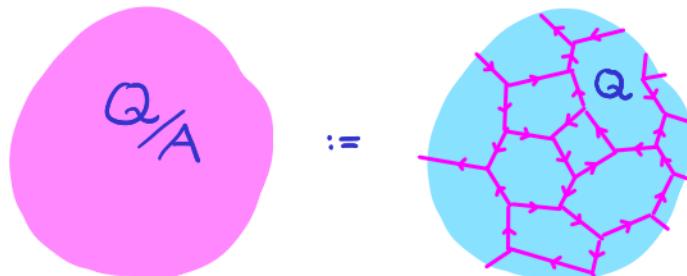
Also need:

$\mu, \Delta$  : topol. junctions

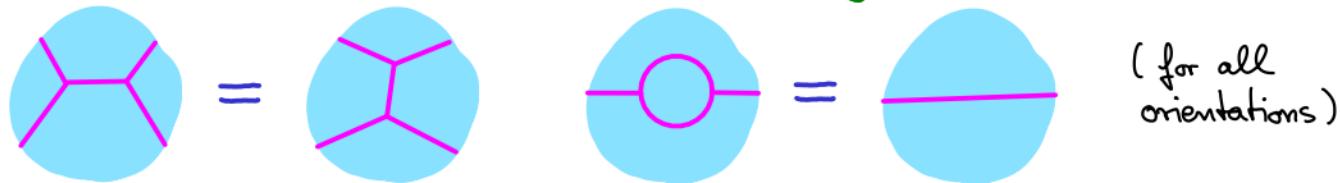


... gauging topol. sym.

Idea :  $Q/A :=$  "  $Q$  with  $A$ -network "



Need : Invariance under change of  $A$ -network

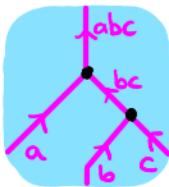


$\rightsquigarrow A$  is symmetric  $\Delta$ -separable Frobenius alg. (assume unit)

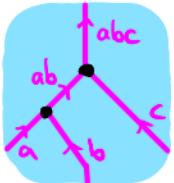
... gauging topol. Sym.

Example: Gauge a symmetry group  $G$

$$A = \bigoplus_{a \in G} a$$



$$= \omega(a, b, c) \cdot$$



$$[\omega] \in H^3(G, \mathbb{C}^\times)$$

$\Delta$ -sep. Trob. alg exists on  $A$  if and only if  
 $[\omega] = 1$  (obstruction vanishes)

$$\left\{ \begin{array}{l} \text{$\Delta$-sep. Trob. alg on } A \\ \text{up to isom.} \end{array} \right\} \simeq H^2(G, \mathbb{C}^\times)$$

## ... gauging topol. sym.

Fröhlich, Fuchs, Schweigert, IR '09

Thm.:

Any two rational 2d conformal field theories containing the chiral symmetry  $V \otimes V$   
(with unique vacuum and non-deg. 2pt correlators)

are obtained from each other by gauging  
top. sym.

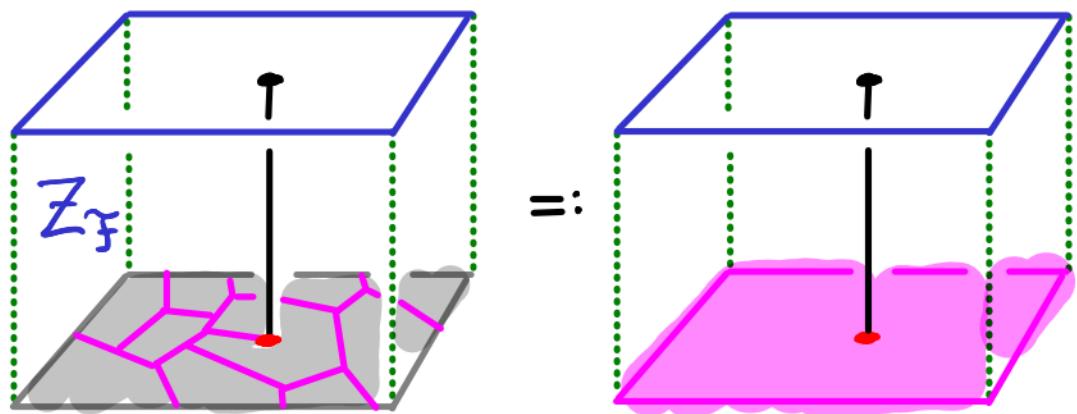
implies same central charge



Wrong for just group symmetries!

... gauging topol. Sym.

Sym TFT picture :



Any two top. bnd. of  $Z_F$  are related to each other by gauging.

# Invertible bubbles and gauging-

Carqueville, IR '12

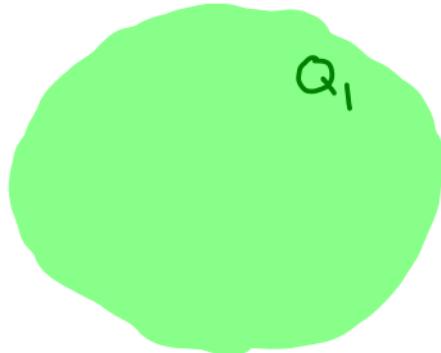
Given : topol. defect separating d-dim QFTs  $Q_1$  and  $Q_2$



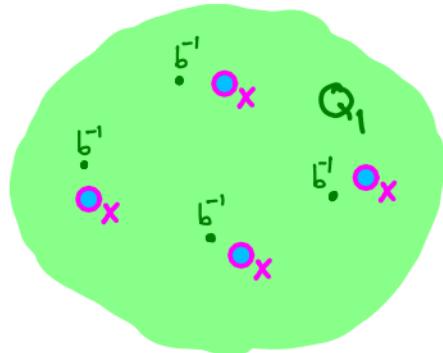
is invertible as a point defect in  $Q_1$  and vice versa.

Then :  $Q_1$  is obtained from  $Q_2$  by gauging a top.sym. and vice versa.

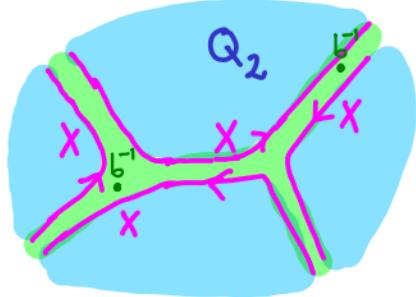
... invertible bubbles



=



=



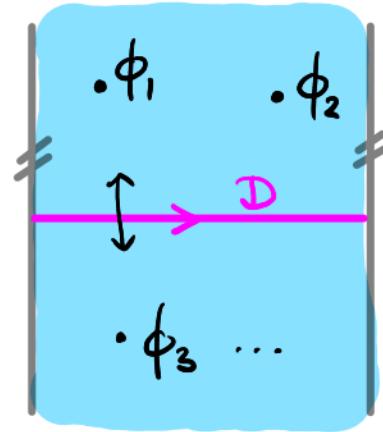
| Get equivalence relation on d-dim. QFTs .

# Perturbed defects and integrability

Ambrosino, Watts,  
IR: in prep.

$\mathcal{T}(Q)$ : Rigidly translation  
invariant defects in  
2dQFT  $Q$

$$[H_a, D] = 0$$



$\Rightarrow$  defect fusion is well-defined

$\mathcal{T}(Q)$  is a monoidal category

(morphisms : topological junctions )

## ... pert. def. & integrability

$\mathcal{T}(Q)$  extends monoidal cat.  $\mathcal{T}_{\text{top}}(Q)$  of topol. defects

$$\mathcal{T}_{\text{top}}(Q) \subset \mathcal{T}(Q) \quad (\text{full monoidal subcategory})$$

Idea : Consider  $\mathcal{T}(Q)$ , not just  $\mathcal{T}_{\text{top}}(Q)$  to get information about  $Q$  and flows  $Q \rightarrow Q'$

# ... pert. def. & integrability

Bazhanov, Lukyanov,  
Zamolodchikov '96  
IR '07

Apply to 2d CFT  $G$

$D$  : topol. defect     $\psi$  : holom. field on  $D$   
 $\bar{\psi}$  : antihol. —||—

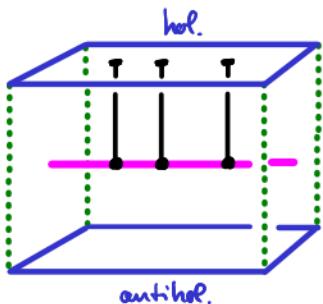
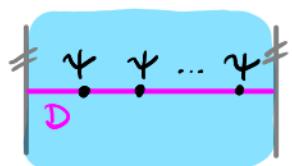
Perturbed defects ( $\lambda \in \mathbb{C}$ )

$$D(\lambda \psi) = \exp(\lambda \int_0^L \psi) \text{ on } D$$

$$D(\lambda \bar{\psi}) = \exp(\lambda \int_0^L \bar{\psi}) \text{ on } D$$

are rigidly transl. inv. since

$$\frac{\partial}{\partial t} \psi(s+it) = -i \frac{\partial}{\partial s} \psi(s+it)$$



# ... pert. def. & integrability

Manolopoulos, IR '09  
Bücher, IR '12

Get monoidal subcat.  $\mathcal{T}_{\text{pert}}(\mathcal{C}) \subset \mathcal{T}(\mathcal{C})$



- not semisimple
- abelian (kernels, quotients, ...)
- continuously many simple obj
- rigid (has duals), not spherical
- $\otimes$  of simples generically simple

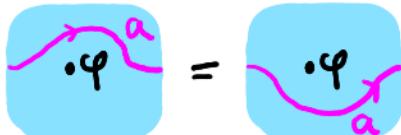
# ... pert. def. & integrability

Bazhanov, Lukyanov,  
Zamolodchikov '99  
IR '10, Bücher, IR '12  
Ambrosino, Watts, IR: in prep.

Perturb CFT by bulk field  $\varphi \sim C(\mu \varphi)$

What can we say about  $J(\mu)$ ?

$J_{\text{top}}(\mu)$  : topol. def a s.th.



$J_{\text{pert}}(\mu)$  : if **commutation condition** ( $\rightarrow$  next slide) holds

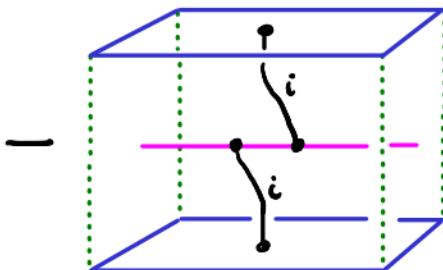
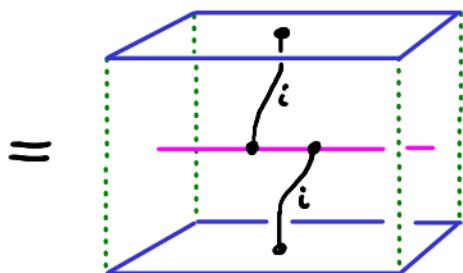
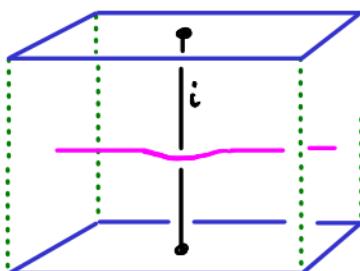
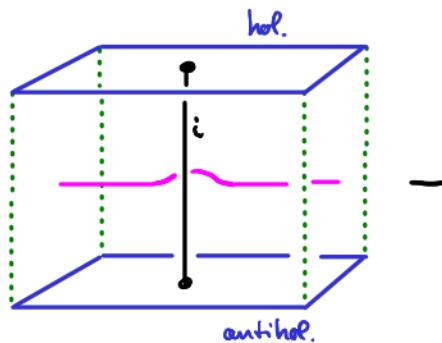
$$\underbrace{\left[ H_{\text{CFT}} + 2i\mu \int_0^L \varphi(s) ds \right]}_{\text{perturbed Hamiltonian}}$$

$$\underbrace{D(\lambda \varphi + \frac{\mu}{\lambda} \bar{\varphi})}_{\in \mathbb{C}^\times} = 0$$

D perturbed by  $\varphi$  and  $\bar{\varphi}$

# ... pert. def. & integrability

The commutation condition (in chiral TFT with trivial surf def.)



## ... pert. def. & integrability

If commutation condition holds :

$D(\lambda \varphi + \frac{\mu}{\lambda} \bar{\varphi})$  is one-param. family in  $T(u)$

→ family of non-local conserved charges  
in perturbed CFT  $C(\mu^q)$ .

## ... pert. def. & integrability

$$D(\lambda \psi + \frac{\mu}{\lambda} \bar{\psi}) \in J_{\text{pert}}(\mu) \subset J(\mu)$$

- monoidal cat.  $J_{\text{pert}}(\mu)$  can be computed in UV CFT  $\mathcal{C}$
- In examples :  $D$ 's commute for different  $\lambda \in \mathbb{C}^\times$

sign of integrability

# ... pert. def. & integrability

Example : A-type Virasoro minimal models

$$M(p, q)$$

Perturbing bulk field  $\varphi_{(r,s)}$

- $(1, 2)$       }      • solutions to comm. cond. exist
  - $(1, 3)$       }      •  $D(\lambda \Psi + \frac{u}{\lambda} \bar{\Psi})$  mutually commute
  - $(1, 5)$       }
  - $(1, 4)$       }      • solutions sporadically exist
  - $(1, 6)$       }
  - $(1, 7)$       }
  - ...      }      • pert. defect often would need reg.
- e.g. in  $M(p, q)$ ,  $q = 9, 10, 18$

## Outlook

Further applications in 2d CFT

- functional relations & integral eqn for D's
  - spin from oriented QFTs by gauging
  - lattice models with fusion cat. topol. sym.  
from 2d CFT
  - gauging non-inv. surface defects in 3d TFT
  - non-semisimple ("logarithmic") 2d CFT  
and chiral / SymTFT
- :