

Boundary Scattering and Non-invertible Symmetries in (1+1) Dimensions

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Based on the collaboration with Satoshi Yamaguchi (Osaka U.)

[To appear]

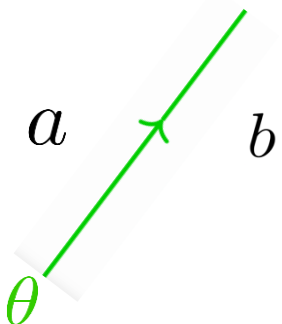
Kyushu IAS-iTHEMS conference: Non-perturbative methods in QFT @ Kyushu

Outline

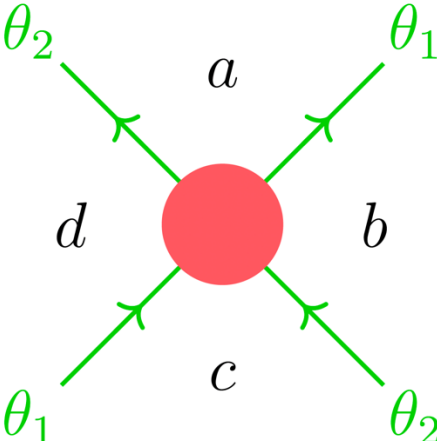
1. Introduction
2. Non-invertible symmetries and kinks
3. Boundary scattering
4. Modified boundary crossing relation
5. Summary and prospects

1. Introduction

- **Scattering amplitude** is a key physical quantity in QFTs .
- In general, it is quite hard to derive the exact S-matrix .
non-perturbatively
- However, if a QFT in **(1+1)d** has the **integrability**, we can determine the S-matrix by the axiomatic way .
- Our interest in this talk is the **kink** scattering .
- Kink : massive excitation interpolating two adjacent vacua $|a\rangle$ and $|b\rangle$

$$|K_{ab}(\theta)\rangle = \begin{array}{c} a \quad \nearrow \quad b \\ \theta \end{array} \quad (\theta : \text{rapidity})$$
A diagram showing a green arrow pointing upwards and to the right. The arrow starts at a point labeled 'theta' at the bottom left and ends at a point labeled 'b' at the top right. The letter 'a' is positioned to the left of the arrow's midpoint. The arrow is surrounded by a light gray, semi-transparent rectangular glow.

- We can introduce the S-matrix for two-body kink scattering :

$$S_{dc}^{ab}(\theta) = \text{diagram}$$


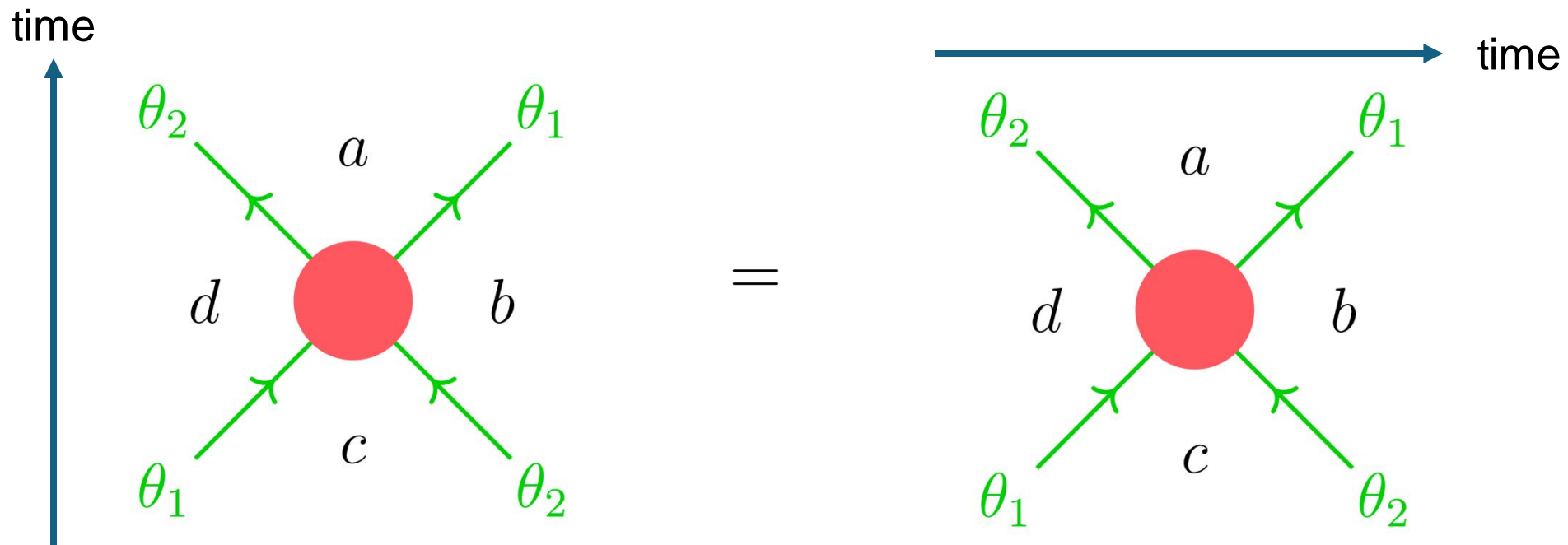
- Integrability allows us to factorize all scattering amplitudes into the product of two-body scattering defined above.
- There are three consistency conditions for the S-matrix.
(unitarity, Yang-Baxter equation, crossing relation)

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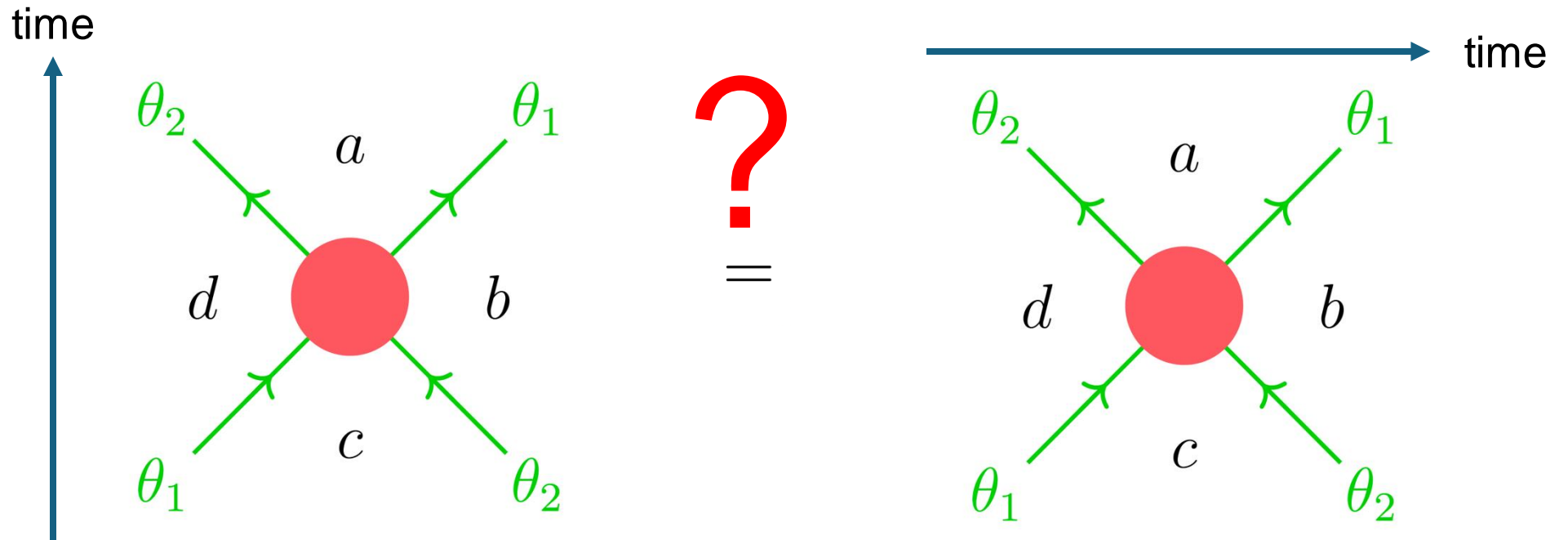
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(unitarity, Yang-Baxter equation, **crossing relation**)

- Crossing relation : $S_{dc}^{ab}(\theta) = S_{ad}^{bc}(i\pi - \theta)$



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- When **non-invertible symmetries** are spontaneously broken in the IR gapped phase, the crossing relation must be modified as

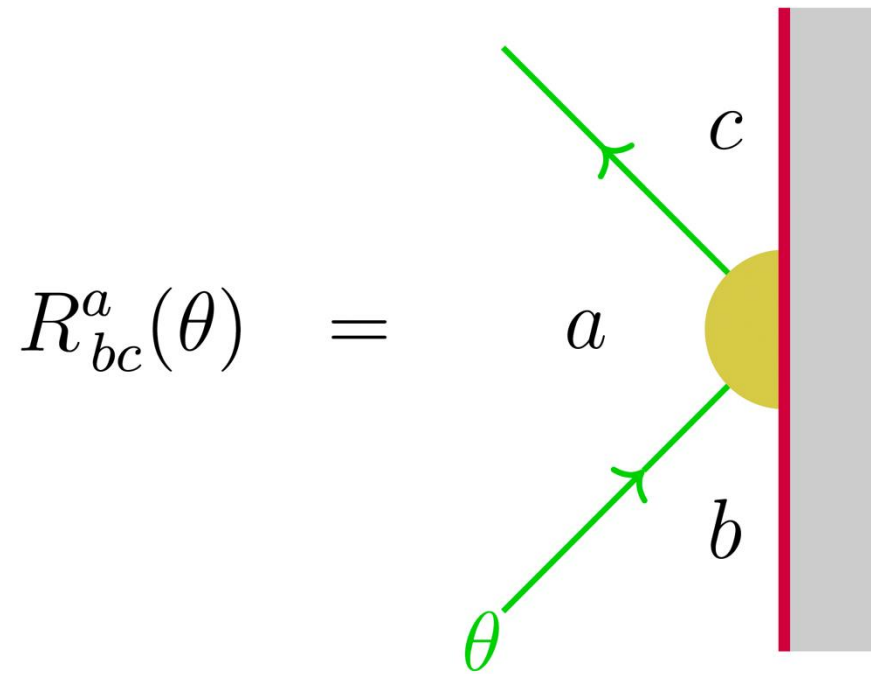
$$S_{cd}^{ab}(\theta) = \sqrt{\frac{d_a d_c}{d_b d_d}} S_{ad}^{bc}(i\pi - \theta) \quad [\text{Copetti, Cordova, Komatsu, '24}]$$

d_a : quantum dimension of topological defect

If we have a **boundary** in (1+1)d spacetime,

we can also define the **boundary S-matrix** .

[Ghoshal, Zamolodchikov, '94]



- Ghoshal and Zamolodchikov proposed the boundary crossing relation:

$$R_{ca}^b \left(\frac{i\pi}{2} - \theta \right) = \sum_d S_{cd}^{ba}(2\theta) R_{ca}^d \left(\frac{i\pi}{2} + \theta \right)$$

Main question of this talk

How should we modify this boundary crossing relation ?

Our work [S.S , S.Yamaguchi, to appear]

Assumption

1. Bulk integrability
2. Boundary is “**weakly-symmetric**” under the non-invertible symmetry.

→ Boundary crossing relation must be also modified as

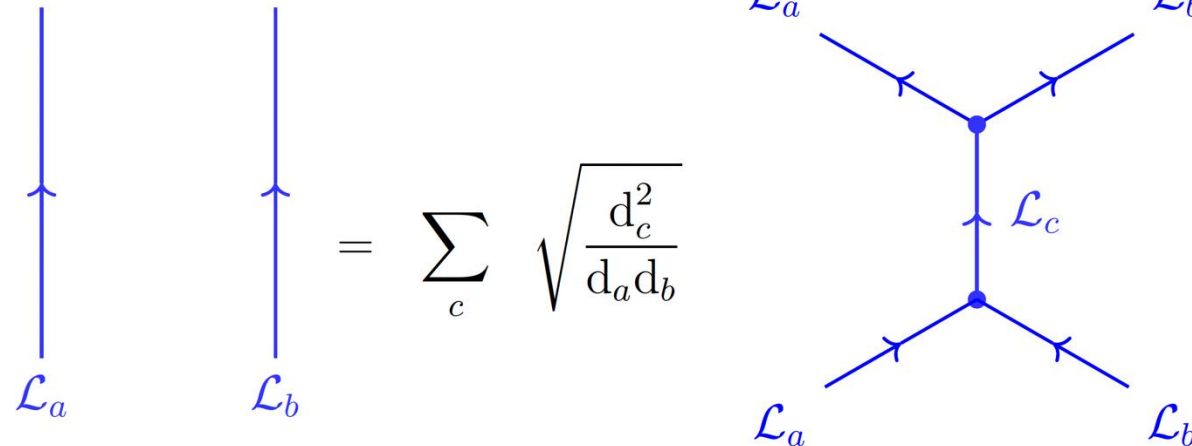
$$R_{ca}^b \left(\frac{i\pi}{2} - \theta \right) = \sum_d \sqrt{\frac{d_d}{d_b}} S_{cd}^{ba}(2\theta) R_{ca}^d \left(\frac{i\pi}{2} + \theta \right)$$

2. Non-invertible symmetries and kinks

- Non-invertible symmetry : **non-group** fusion rule

$$\mathcal{L}_a \otimes \mathcal{L}_b = \sum_c N_{bc}^a \mathcal{L}_c, \quad N_{bc}^a \in \{0, 1\} \quad (\text{multiplicity-free case})$$

- Partial fusion :



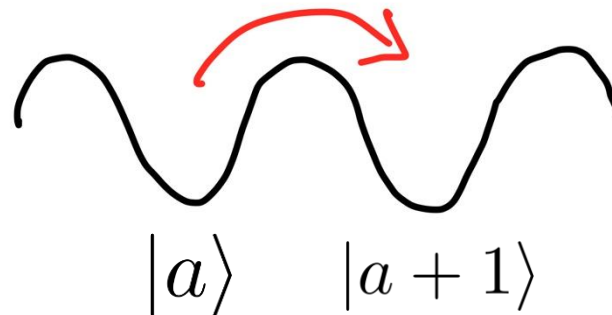
We assume that this fusion category is preserved along with RG flow to **IR gapped phase**.

Non-invertible symmetry is spontaneously broken in the IR.

→ We have degenerate vacua $|a\rangle$

$$|a\rangle = \mathcal{L}_a |0\rangle \quad (\text{regular module category})$$

→ We can define the kink between $|a\rangle$ and $|a \pm 1\rangle$.



Ex.

Tricritical Ising CFT



$-\varepsilon'$

Gapped Phase
with 3 vacua

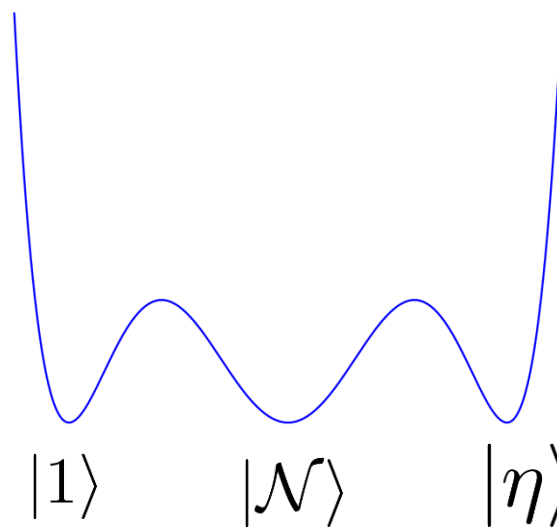
[Huse, '84;,
Chang, Lin, Shao, Wang, Yin, '19]

TY category is preserved

$$\mathcal{N} \otimes \mathcal{N} = 1 \oplus \eta$$

$$\mathcal{N} \otimes \eta = \eta \otimes \mathcal{N} = \mathcal{N}$$

$$\eta^2 = 1$$



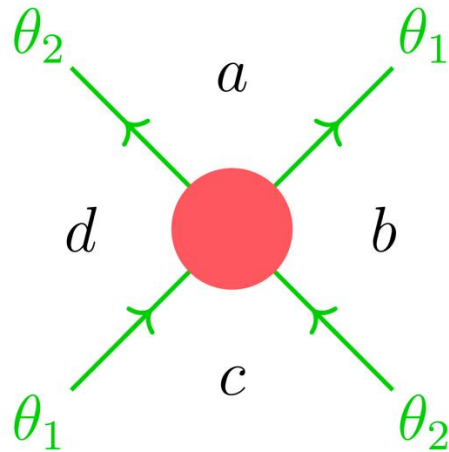
3. Boundary Scattering

[Ghoshal, Zamolodchikov, '94]

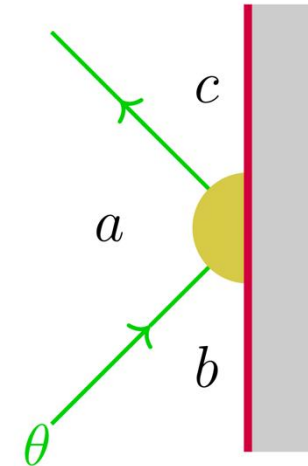
[Scattering data]

• Bulk S-matrix

not be affected
by the boundary



• Boundary S-matrix



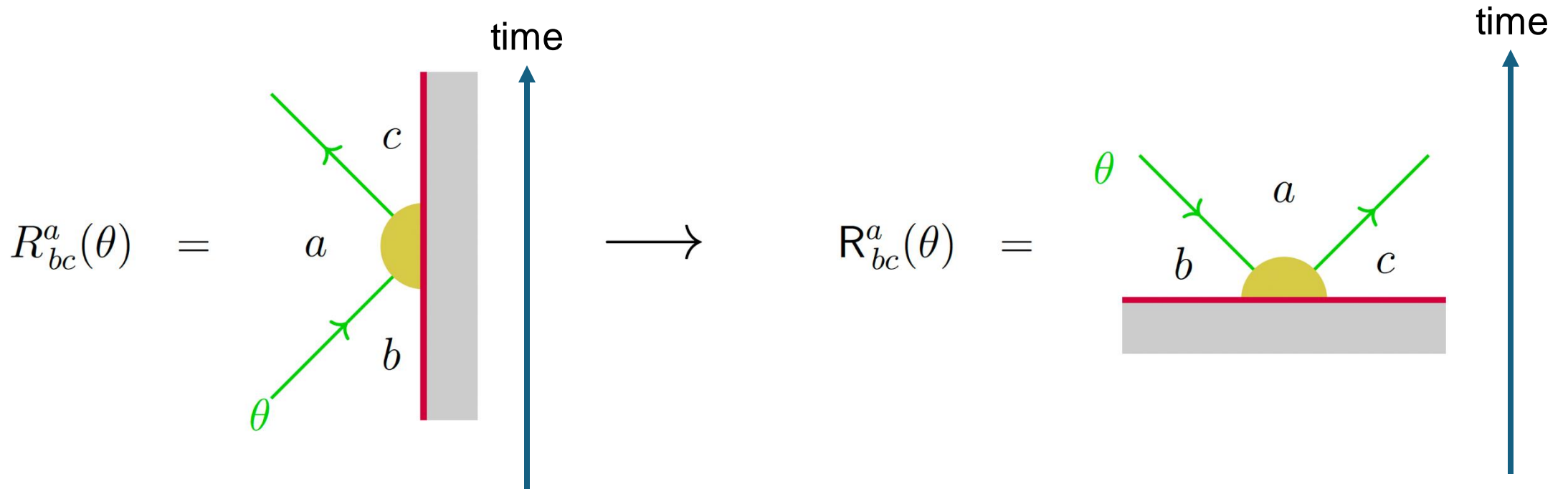
Recall that the bulk S-matrix must be subject to several conditions .

Ghoshal and Zamolodchikov proposed the three conditions for boundary S-matrix .

■ Boundary crossing relation:

Boundary analog of crossing relation is more non-trivial.

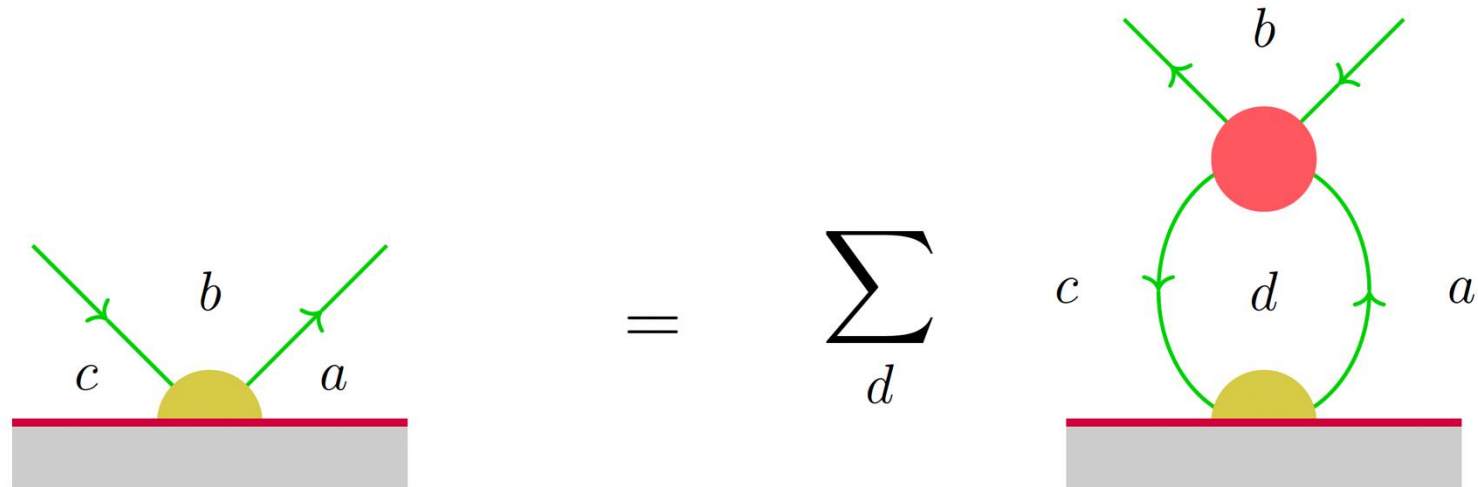
We first move to the rotated picture :



In this picture, the asymptotic kink state belongs to **bulk** Hilbert space.

In and out asymptotic states in the bulk Hilbert space must be related via **bulk S-matrix** .

$$R_{ca}^b \left(\frac{i\pi}{2} - \theta \right) = \sum_d S_{cd}^{ba}(2\theta) R_{ca}^d \left(\frac{i\pi}{2} + \theta \right)$$



Remark: we used the bulk integrability.

To sum up, boundary crossing relation can be put as follows:

$$R_{bc}^a(\theta) = R_{bc}^a(\theta)$$

$$R_{ca}^b\left(\frac{i\pi}{2} - \theta\right) = \sum_d S_{cd}^{ba}(2\theta) R_{ca}^d\left(\frac{i\pi}{2} + \theta\right)$$

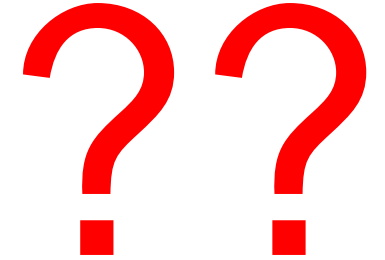
→ $R_{ca}^b\left(\frac{i\pi}{2} - \theta\right) = \sum_d S_{cd}^{ba}(2\theta) R_{ca}^d\left(\frac{i\pi}{2} + \theta\right)$

[Ghoshal, Zamolodchikov, '94]

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[Ghoshal, Zamolodchikov, '94]

Necessary to be modified when non-invertible symmetry is concerned .

4. Modified boundary crossing relation [\[S.S, S.Yamaguchi, to appear\]](#)

How should we modify ?

Our assumption :

1. Bulk integrability

4. Modified boundary crossing relation [\[S.S, S.Yamaguchi, to appear\]](#)

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1. Bulk integrability \longrightarrow necessary to suppress multi-kink contributions

4. Modified boundary crossing relation [S.S, S.Yamaguchi, to appear]

How should we modify ?

Our assumption :

1. Bulk integrability \longrightarrow necessary to suppress multi-kink contributions
2. Boundary is weakly-symmetric under the non-invertible symmetry.

topological junction $\text{Hom}(\mathcal{L}_a \otimes \mathcal{B}, \mathcal{B})$

[Choi, Rayhaun, Sanghavi, Shao, '23]

Firstly, we notice that in the rotated picture, the correctly normalized matrix R_{bc}^a must satisfy

$$R_{ca}^b \left(\frac{i\pi}{2} - \theta \right) = \sum_d S_{cd}^{ba}(2\theta) R_{ca}^d \left(\frac{i\pi}{2} + \theta \right)$$

But, we emphasize

$$R_{bc}^a(\theta) \neq R_{bc}^a(\theta)$$

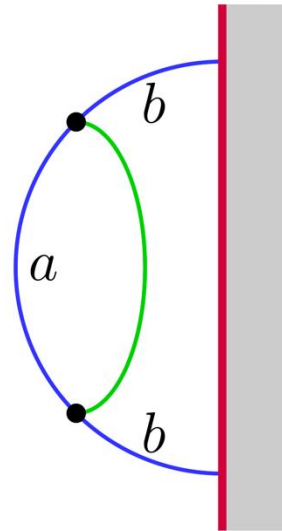
To find the proper relation, we must carefully look at the normalization!

- Normalization of $R_{bc}^a(\theta)$

Correctly normalized boundary S-matrix is given by

$$R_{bc}^a(\theta) = \frac{1}{\sqrt{C_{ac}^{\mathcal{B}} \cdot C_{ab}^{\mathcal{B}}}} \text{out}, \mathcal{B} \langle K_{ac}(-\theta) | K_{ab}(\theta) \rangle_{\mathcal{B}, \text{in}}$$

$$C_{ab}^{\mathcal{B}} := \mathcal{B} \langle K_{ab}(\theta) | K_{ab}(\theta) \rangle_{\mathcal{B}} =$$



By using the fusion property, we can evaluate as

$$C_{ab}^{\mathcal{B}} := \text{Diagram 1} = \sum_{d \in b \otimes b} \frac{\sqrt{d_d}}{d_b} \text{Diagram 2} \cong \frac{1}{d_b} \text{Diagram 3}$$

Assuming

$$\text{Diagram 4} = \frac{1}{F_b} \text{Diagram 5} \longrightarrow C_{ab}^{\mathcal{B}} = \frac{1}{G_b} \langle K_{ab}(\theta) | K_{ab}(\theta) \rangle \cdot {}_{\mathcal{B}}\langle 0 | 0 \rangle_{\mathcal{B}}$$

($G_b := d_b F_b$)

- Normalization of R_{bc}^a

Correctly normalized S-matrix in the rotated picture is given by

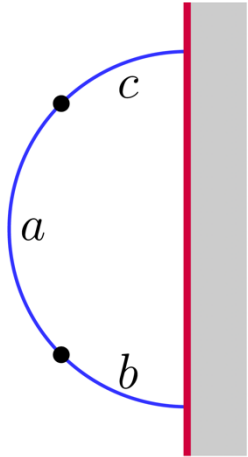
$$R_{bc}^a(\theta) = \frac{1}{\sqrt{C_{bac} \cdot \langle \mathcal{B} | \mathcal{B} \rangle}} \text{out} \langle K_{ba}(-\theta) K_{ac}(\theta) | \mathcal{B} \rangle$$

$$C_{bac} = \text{diagram} = \sum_{d \in a \otimes a} \frac{\sqrt{d_d}}{d_a} \text{diagram} \cong \frac{1}{d_a} \text{diagram} \text{ diagram}$$

The diagram shows the decomposition of a circle with two arcs labeled b and c into a sum over intermediate states d of two circles connected by a line labeled d . The first circle has arcs b and a , and the second has arcs a and c . This is then simplified to a product of two separate circles, one with arcs b and a , and the other with arcs a and c .

$$\longrightarrow C_{bac} = \frac{1}{d_a} \langle K_{ba}(\theta) | K_{ba}(\theta) \rangle \langle K_{ac}(\theta) | K_{ac}(\theta) \rangle$$

By using the boundary LSZ formula and analytic continuation, we can relate the boundary S-matrices R_{bc}^a and R_{bc}^a in terms of two-point Euclidean correlators .

$$R_{bc}^a(\theta) = \frac{1}{\sqrt{C_{ac}^{\mathcal{B}} \cdot C_{ab}^{\mathcal{B}}}} \text{ (diagram) } = \sqrt{\frac{C_{bac} \cdot \langle \mathcal{B} | \mathcal{B} \rangle}{C_{ac}^{\mathcal{B}} \cdot C_{ab}^{\mathcal{B}}}} R_{bc}^a = \sqrt{\frac{G_b G_c}{d_a}} R_{bc}^a$$


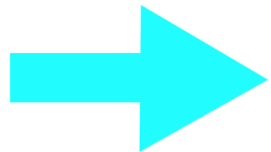
$$\longrightarrow R_{bc}^a(\theta) = \sqrt{\frac{d_a}{G_b G_c}} R_{bc}^a(\theta)$$

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Plug in

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Modified boundary crossing relation

$$R_{ca}^b\left(\frac{i\pi}{2} - \theta\right) = \sum_d \sqrt{\frac{d_d}{d_b}} S_{cd}^{ba}(2\theta) R_{ca}^d\left(\frac{i\pi}{2} + \theta\right)$$

5. Summary and prospects

- We revisit the boundary scattering theory.
- We found that the boundary crossing relation should be modified when non-invertible symmetry is concerned.
(see our paper for concrete models)
- How about monopole scattering?
[van Beest, Smith, Delmastro, Komargodski, Tong, '23]
- Non-regular vacuum module category?