

On Energy and Entropy of Surface Defect in Six Dimensions

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“Non-perturbative methods in QFT”, Kyushu University



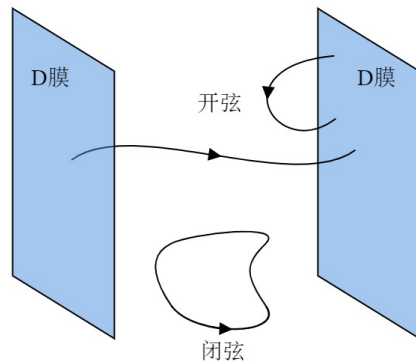
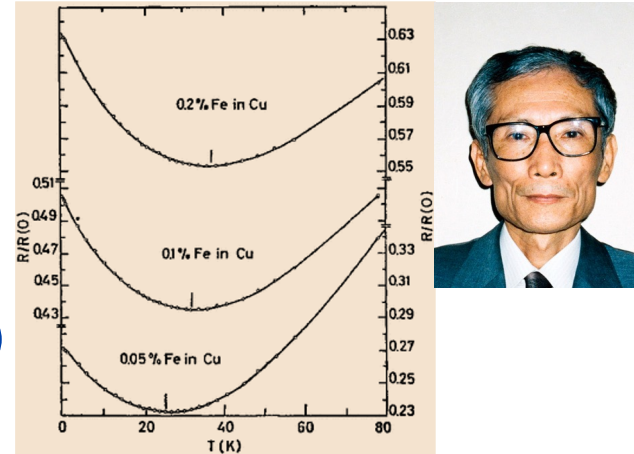
[with Ma-Ke Yuan, arXiv:2310.02096][with Zixiao Huang and Ma-Ke Yuan, arXiv:2501.09498]

Defects

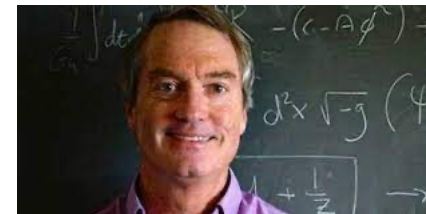
infinitely heavy
charge probe



impurity
(Kondo effect)

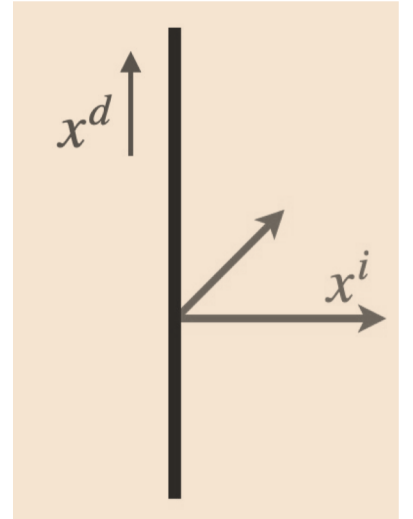


D-brane



Defect in QFT

- Defects are usually defined by **non-local** operators/B.C. in QFT .
- Defects are classified by their dimensions: **line defects**, **surface defects**..(Wilson line, Wilson surface).
- Physically, defect may arise from **p**-dimensional degrees of freedom coupled to **D**-dimensional bulk QFT



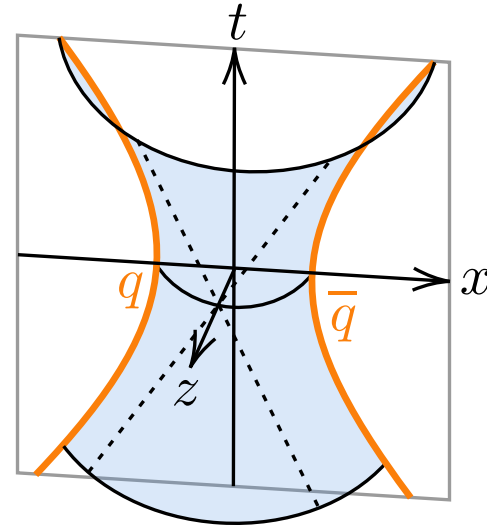
$$I_{\text{DCFT}} = \int d^D x \sqrt{g} \mathcal{L}_{\text{CFT}}[\phi] + \int d^p \hat{x} \sqrt{\gamma} \mathcal{L}_{\text{defect}}[\phi, \psi]$$

Defect landscape

- **Topological** defects, generalized symmetry.
- **Conformal** defects, impurities & critical system.
- Boundary, Cross-Cap, monodromy, entangling surface...

Fun with defects

- **Bootstrap** with defects
- **Localization** with defects
- Defect **RG**, Defect **ER=EPR**
- Defect **fusion**, **intersection**



[Yifan's talk & Tom's talk]

Conformal defect Kapustin 2006, Billò–Gonçalves–Lauria–Meineri 2016

- Defect breaks part of conformal symmetry.

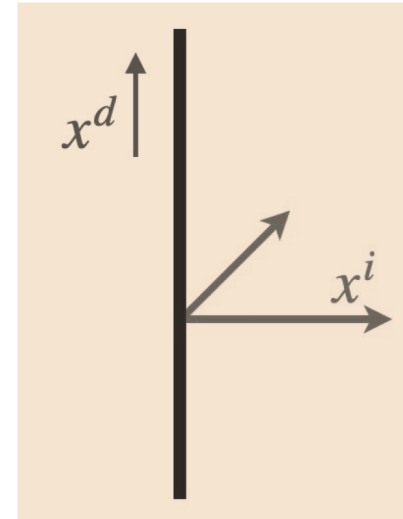
$$SO(1, D+1) \text{ to } SO(1, p+1) \times SO(D-p)$$

- Bulk stress tensor 1-pt function can be fixed up to a constant

$$\langle T^{ab} \rangle = -\frac{D-p-1}{D} \frac{h}{|x^i|^D} \delta^{ab}, \quad \langle T^{ai} \rangle = 0,$$

$$\langle T^{ij} \rangle = \frac{h}{|x^i|^D} \left(\frac{p+1}{D} \delta^{ij} - \frac{x^i x^j}{|x^i|^2} \right),$$

- Displacement operator



Weyl anomaly on surface defect

- For a surface defect, there is an anomaly, analogy to conformal anomaly in CFT2, but now we have b, d_1, d_2

$$\langle\langle \hat{T}_a^a \rangle\rangle = -\frac{1}{24\pi} \left[bR^\Sigma + d_1 \tilde{\Pi}_{ab}^\mu \tilde{\Pi}_\mu^{ab} - d_2 W_{ab}^{ab} \right]$$

- This can also be seen from log divergence of expectation value

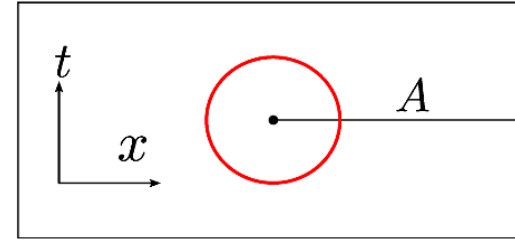
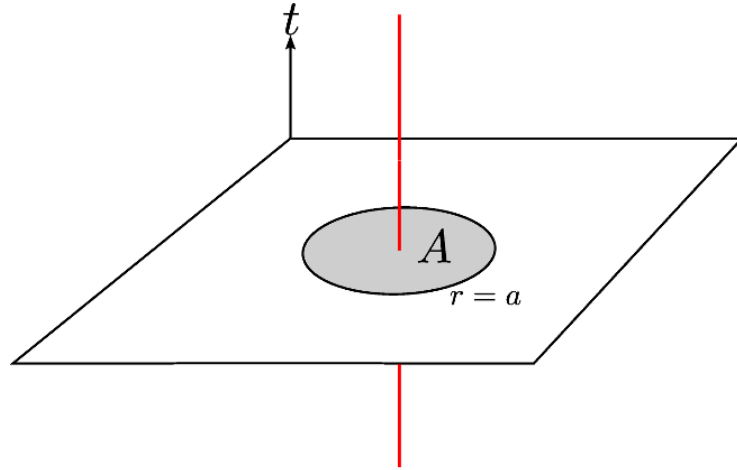
$$\log \langle D_\Sigma \rangle \supset \int_\Sigma \mathcal{A}_\Sigma \text{vol}_\Sigma \log \ell / \epsilon$$

$$\mathcal{A}_\Sigma = \frac{1}{24\pi} \left[bR^\Sigma + d_1 \tilde{\Pi}_{ab}^\mu \tilde{\Pi}_\mu^{ab} - d_2 W_{ab}^{ab} \right]$$

Extra entanglement due to defect

Jensen-O'bannon 2013,

Lewkowycz-Maldacena 2013



$$S_W = (1 - n\partial_n) \log \langle W \rangle|_{n=1} = \log \langle W \rangle + \int \langle T_{\tau\tau} \rangle_W$$

Interesting observable related to both $\log \langle W \rangle$ and h

Free fields

- free scalar and 2-form in **6d**

$$L_\phi = \frac{1}{2}(\partial\phi)^2 + \frac{1}{10}R\phi^2$$

$$L_B = \frac{1}{12}F_{\mu_1\mu_2\mu_3}F^{\mu_1\mu_2\mu_3}$$

- Surface defects in free fields

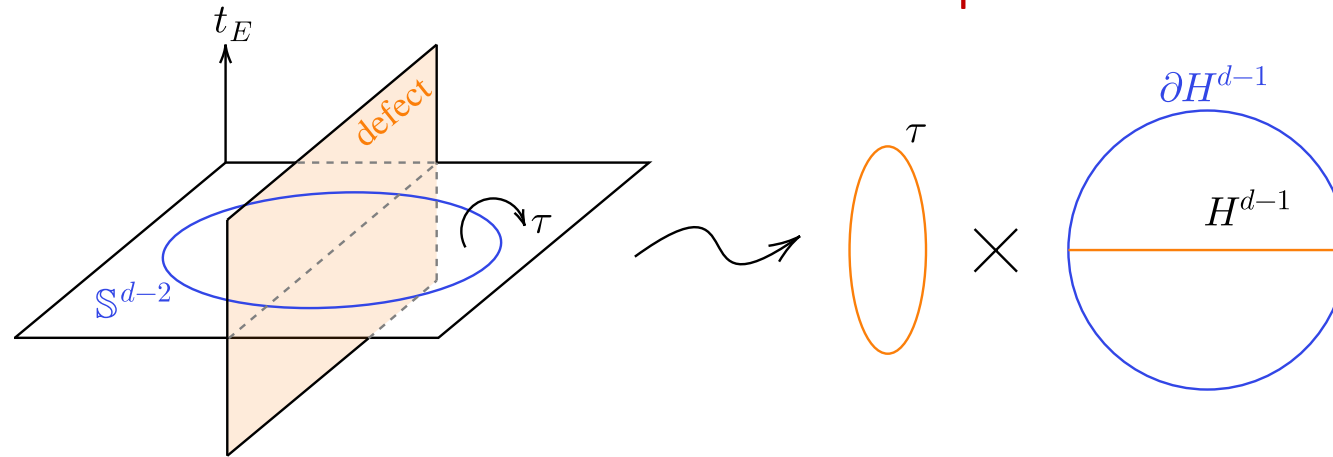
$$D_\phi = \exp\left(\int_\Sigma d^2\sigma\phi(\sigma)\right)$$

$$D_B = \exp\left(i\int_\Sigma B\right).$$

$$h_\phi = \frac{1}{20\pi^4}, \quad h_B = \frac{1}{4\pi^4}$$

Renyi entropy

- Map to $S_\beta * H^5$, with defect wrapped on $S_\beta * H^1$



$$ds_{\mathbb{R}^d}^2 = dt_E^2 + dr^2 + r^2 d\Omega_{d-2}^2$$

$$ds_{\mathbb{R}^d}^2 = \Omega^2 ds_{S^1 \times H^5}^2 = \Omega^2 \ell^2 (d\tau^2 + d\rho^2 + \sinh^2 \rho d\Omega_{d-2}^2)$$

$$t_E = \ell \frac{\sin \tau}{\cosh \rho + \cos \tau}, \quad r = \ell \frac{\sinh \rho}{\cosh \rho + \cos \tau}$$

Free energy

- Free scalar Kernel on $S_\beta * H^5$

$$K_{H^5}(t, \rho) = \frac{e^{-4t - \frac{\rho^2}{4t}} (\rho^2 + 2\rho t \coth \rho - 2t) \operatorname{csch}^2 \rho}{32\pi^{5/2} t^{5/2}}$$

$$K_{S_\beta^1}(t, \tau) = \sum_{n=-\infty}^{+\infty} \frac{\exp\left(-\frac{(\beta n + \tau)^2}{4t}\right)}{\sqrt{4\pi t}}$$

- Integrate propagator over $S_\beta * H^1$

$$G_\beta(\tau, \rho) = \int_0^{+\infty} dt K_{S_\beta^1}(t, \tau) * K_{H^5}(t, \rho) e^{4t}$$

Boundary term

- Additional term on the boundary of H^5 , from conformal coupling of the scalar to curvature

$$S = -\frac{4\pi}{10} \int_{\hat{\Sigma}} dA \langle\langle \phi^2 \rangle\rangle$$

- Evaluate it in the presence of defect,

$$S_\phi = -\frac{1}{10\pi} \log \ell/\epsilon$$

Defect EE

[Kobayashi-Nishioka-Watanabe,2018]

[Jensen-O'Bannon-Robinson-Rodgers 2018]

- Defect EE in terms of central charges was derived before

$$S = \frac{1}{3} \left(b - \frac{d-3}{d-1} d_2 \right) \log \ell / \epsilon$$

- Consistency check:

$$S_\phi = -\frac{1}{10\pi} \log \ell / \epsilon \quad d_2 = 6\pi \Omega_{d-3} \frac{d-1}{d} h_\phi \quad b = 0$$

Two form

- We use defect EE formula

$$b = \frac{3}{2\pi} \quad h_B = 1/4\pi^4 \quad S_B = \frac{1}{3} (b - 6\pi^3 h_B) \log \ell/\epsilon = 0$$

- From information theory, we conjecture that Renyi entropy also vanishes.

(2,0) free theory

- free tensor multiplet: 5 real scalars + 2 Weyl fermions + 1 two-form(with self-dual strength)

$$W = \exp \int_{\Sigma} (iB^+ - n^i \phi_i \text{vol}_{\Sigma})$$

- Defect contribution to Renyi entropy is a sum,

$$S = S_{\phi} + \frac{1}{2} S_B = -\frac{1}{10\pi} \log \ell/\epsilon$$

(2,0) in the large N

- CFT6/AdS7 [Maldacena 1997]

A_{N-1} (2, 0) SCFT

M-theory on $\text{AdS}_7 \times S^4$

$$ds_{11}^2 = L^2(ds_{\text{AdS}_7}^2 + \frac{1}{4}ds_{S^4}^2) , \quad F_4 = dC_3 = \pi^2 L^3 \text{vol}_{S^4} , \quad L^3 = 8\pi N \ell_p^3$$

- Holographic dual of surface operator in **fundamental rep.** is a single M2 brane,

$$S_{\text{M2}} = -\log\langle W \rangle$$

Super-Renyi entropy of 6d SCFT

[Nishioka-Yaakov 2013]
[YZ 2016, Yankielowicz-YZ 2017]

- Twist the theory on $S_\beta * H^5$ by R-symmetry leads to **supersymmetric Renyi entropy (SRE)**: a polynomial of $\gamma := 1/n$

$$S_\nu^{(1,0)} = \sum_{n=0}^3 s_n (\gamma - 1)^n,$$

$$s_0 = \frac{1}{6}(8\alpha - 8\beta + 8\gamma + 3\delta),$$

$$s_1 = \frac{1}{4}(2\alpha - 3\beta + 4\gamma + \delta),$$

$$s_2 = \frac{1}{24}(2\alpha - 5\beta + 8\gamma),$$

$$s_3 = \frac{1}{192}(\alpha - 4\beta + 16\gamma).$$

$$\mathcal{I}_8 = \frac{1}{4!} (\alpha c_2^2(R) + \beta c_2(R)p_1(T) + \gamma p_1^2(T) + \delta p_2(T))$$

[Ohmori-Shimizu-Tachikawa-Yonekura 2014]

Large N limit

- Large N limit is **charged** hyperbolic black hole with **$k=-1, m=0$**

$$ds_7^2 = - (H_1 H_2)^{-4/5} f dt^2 + (H_1 H_2)^{1/5} (f^{-1} dr^2 + r^2 d\Omega_{5,k}^2)$$

$$f(r) = k - \frac{m}{r^4} + \frac{r^2}{L^2} H_1 H_2, \quad H_i = 1 + \frac{q_i}{r^4},$$

[Cvetic-Duff-Hoxha-Liu-Lu-Lu-Martinez-Acosta-Pope-Sati-Tran 1999]

- One can check holographic SRE

$$\frac{S_\gamma[A_{N \rightarrow \infty}]}{N^3} = \frac{1}{192} (\gamma - 1)^3 + \frac{1}{12} (\gamma - 1)^2 + \frac{1}{2} (\gamma - 1) + \frac{4}{3}$$

Defect Super-Renyi

- M2 brane action

$$S_{\text{M2}} = T_2 \int d^3\sigma \sqrt{-\det[g]}$$

- Embedding

$$\sigma_0 = \tau, \quad \sigma_1 = \rho, \quad \sigma_2 = r$$

- On-shell action

$$S_{\text{M2}} = T_2 \int_0^\beta d\tau \int_{-\infty}^{\infty} d\rho \int_{r_H}^{\Lambda} r_H (H_1 H_2)^{-1/5} \sqrt{\Delta}$$

- Large N defect SRE

$$S_\gamma[A_{N \rightarrow \infty}] = N(r_1 r_2 (\gamma - 1) + 2) \log \ell / \epsilon \qquad S_n = \pi T_2 V_{H^1} \left(\frac{7n + 1}{16n} \right)$$

A closed formula

[arXiv:2501.09498 with Huang and Yuan]

$$S_\gamma[\mathfrak{g}] = \frac{2b_{\mathfrak{g}} - d_{2\mathfrak{g}}}{6} \left[\frac{r_1 r_2}{2} (\gamma - 1) + 1 \right] \log \ell / \epsilon$$

$$r_1 + r_2 = 1$$

- Large gamma defect SRE shares the same behavior with defect SCE (**supersymmetric Casimir energy on $S^5 * H^1$**)
- Defect SCE is determined from localization, further confirmed by anomaly polynomial
- The constant term is fixed by assuming a linear shift of $b/3$ by d_2

Supersymmetric Casimir energy

[Bobev-Bullimore-Kim 2015]

- 5d partition function from localization (large β)

$$\int \prod_{i=1}^N d\nu_i \exp \left[\frac{2\pi}{\omega_1 \omega_2 \omega_3} \left(-\frac{\pi}{\beta} \sum_{i=1}^N \nu_i^2 + \frac{\sigma_1 \sigma_2}{2} \sum_{j < i} (\nu_i - \nu_j) \right) \right]$$

- Saddle point solution

$$-\frac{2\pi}{\beta} \nu_i + \frac{\sigma_1 \sigma_2}{2} ((i-1) - (N-i)) = 0 \quad \Rightarrow \quad \nu_i = \frac{\beta \sigma_1 \sigma_2}{4\pi} (2i - N - 1)$$

- 6d supersymmetric Casimir energy of A-type (2,0)

$$\beta E_c = -\log \mathcal{Z} = -\beta \frac{N(N^2 - 1) \sigma_1^2 \sigma_2^2}{24 \omega_1 \omega_2 \omega_3}$$

Defect SCE

[arXiv:2501.09498 with Huang and Yuan]
[Mori-Yamaguchi 2014]

- Wilson loop insertion in localization $\text{Tr}_{\mathcal{R}} \exp [2\pi\nu/\omega_j]$

$$\langle W_{(k)} \rangle = \frac{1}{\mathcal{Z}} \int \prod_{i=1}^N d\nu_i \exp \left[\frac{2\pi}{\omega_1\omega_2\omega_3} \left(-\frac{\pi}{\beta} \sum_{i=1}^N \nu_i^2 + \frac{\sigma_1\sigma_2}{2} \sum_{j<i} (\nu_i - \nu_j) + \omega_2\omega_3 k\nu_N \right) \right]$$

- For symmetric representation

$$\begin{aligned} \beta E_{(k)} &= -\log \langle W_{(k)} \rangle \\ &= -\frac{\beta}{2\omega_1} (k(N-1)\sigma_1\sigma_2 + k^2\omega_2\omega_3) \end{aligned}$$

- For anti-symmetric rep.

$$\beta E_{[k]} = -\frac{\beta}{2\omega_1} (k(N-k)\sigma_1\sigma_2 + k\omega_2\omega_3)$$

Defect SCE from anomaly polynomial

- Anomaly polynomial from inflow [Shimizu-Tachikawa 2016, Wang 2021]

$$I_4 = \frac{1}{4}(\Lambda, \Lambda) (c_2(F_L) - c_2(F_R)) \\ + \frac{1}{2}(\Lambda, \rho) (c_2(F_I) - c_2(F_F))$$

- Equivariant integration

$$E_g = -\frac{1}{2\pi} \int I_4 \\ = -\frac{1}{4\omega_1} \left[\frac{1}{2} (\Lambda, \Lambda) (a_L^2 - a_R^2) + (\Lambda, \rho) (a_I^2 - a_F^2) \right]$$

- Identify 2d and 6d parameters and use the recent defect anomaly formula

$$b = 24 (\Lambda, \rho) + 3 (\Lambda, \Lambda) , \quad d_2 = 24 (\Lambda, \rho) + 6 (\Lambda, \Lambda)$$

$$E_g = -\frac{1}{\omega_1} \left[\frac{d_{2g} - b_g}{6} \omega_2 \omega_3 + \frac{2b_g - d_{2g}}{24} \sigma_1 \sigma_2 \right]$$

$$\omega_1 = \omega_2 = 1, \quad \sigma_1 \sigma_2 = 2\omega_3$$

$$E_g|_{\text{chiral limit}} = -\frac{1}{12} \omega_3 d_2$$

[Chalabi-O'bannon-Robinson-Sisti 2020]

Conclusion&Discussion

- We propose a closed formula for defect contribution to Super-Renyi entropy as well as supersymmetric Casimir energy in terms of central charges in M5 brane theories.
- A defect Cardy formula?
- Bounds on defect anomalies?
- (1,0)? Other dimensional defects? Gukov-Witten in 4d?