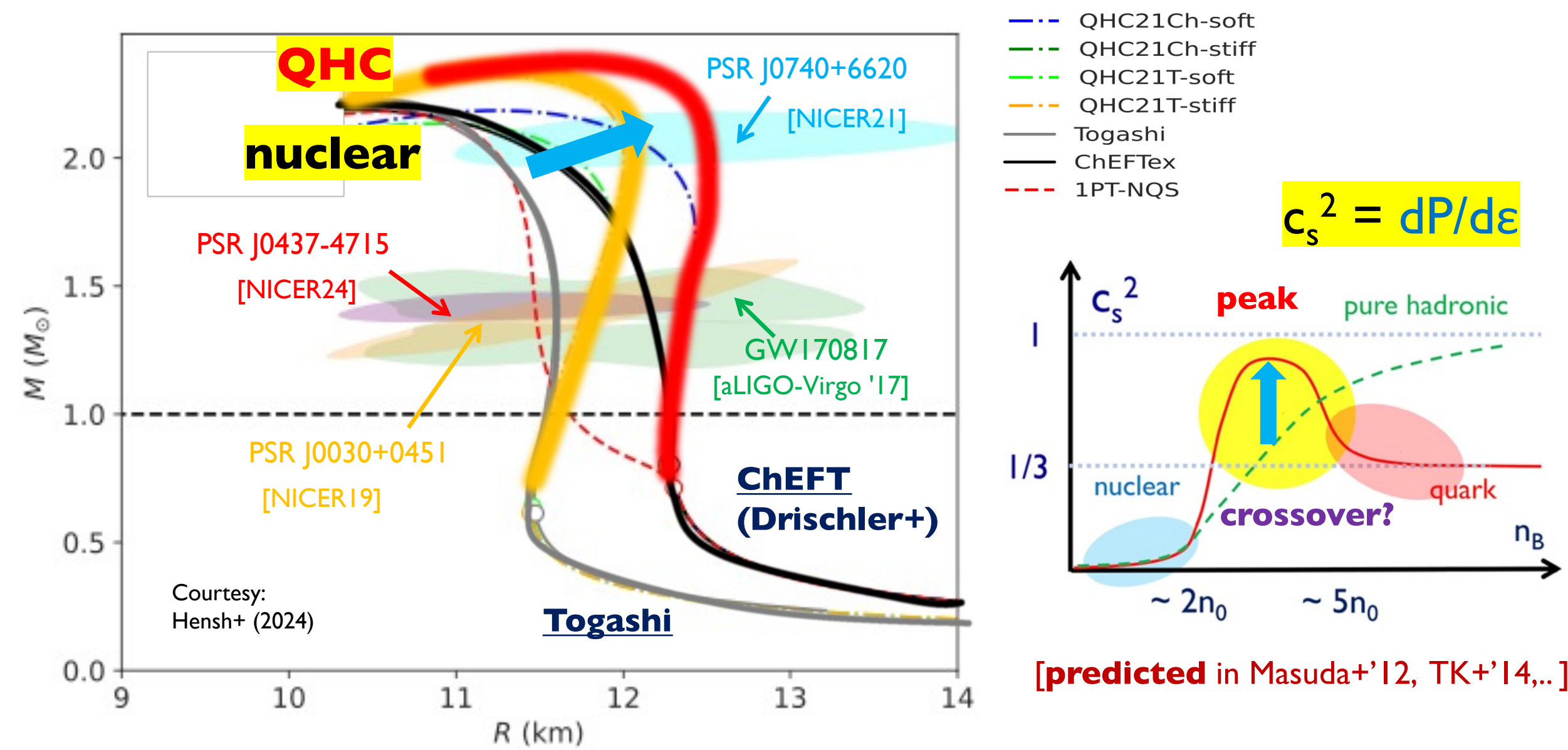


A quarkyonic matter model

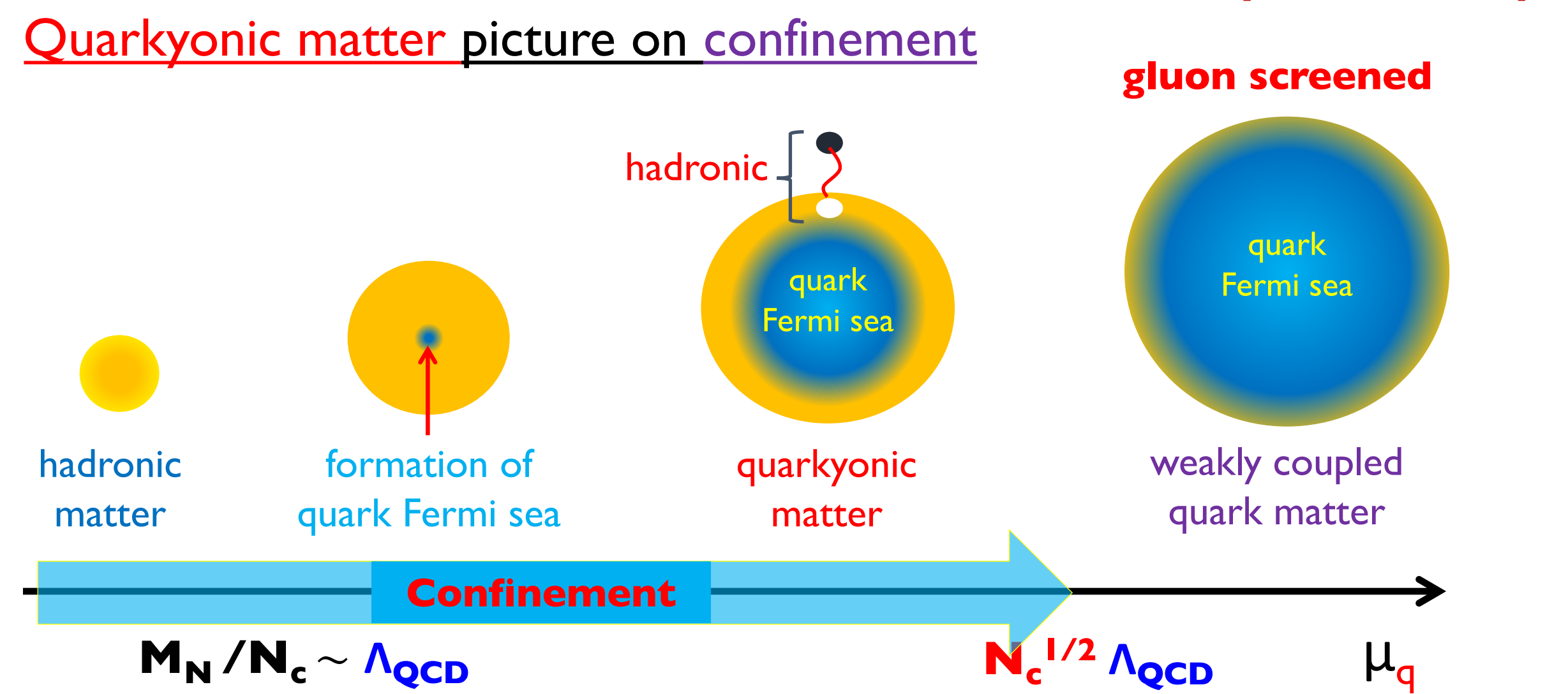
Toru Kojo (KEK), Yuki Fujimoto (iTHEMS), Larry McLerran (INT)

Summary: *Quark saturation* drives rapid stiffening of EOS at $\sim 2-3n_{\text{sat}}$. It also suppresses the hyperon production in dense matter with a shift $\mu_B^{\text{onset}} = M_Y \Rightarrow 2M_Y - M_N$. The new onset density is large, $n_B^{Y, \Xi\text{-onset}} \sim 6n_0$, even before NY repulsions are included.

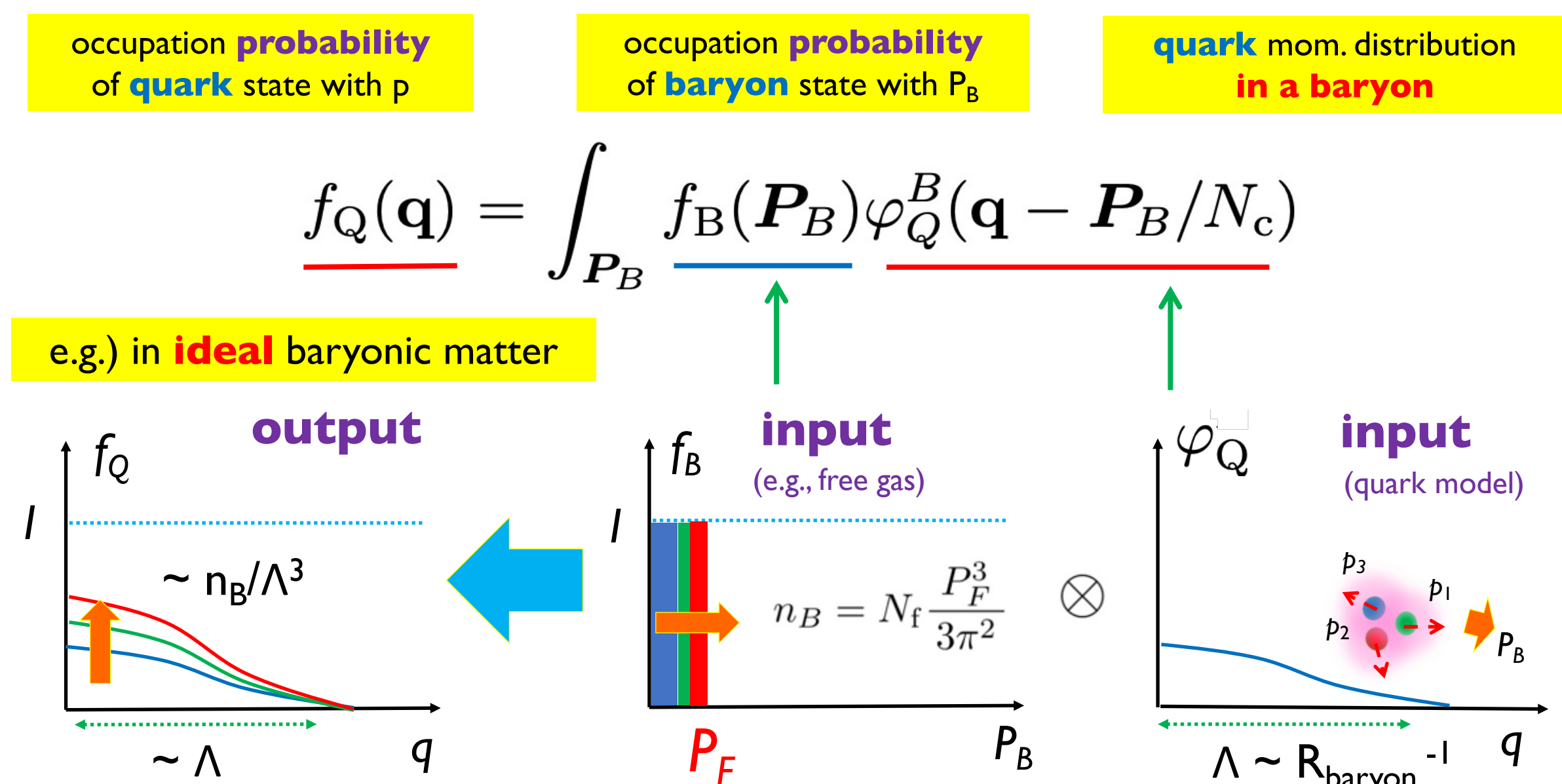
Soft-to-stiff EOS: rapid evolution of pressure



Evolution from nuclear to quark matter



Sum rules for occupation probabilities



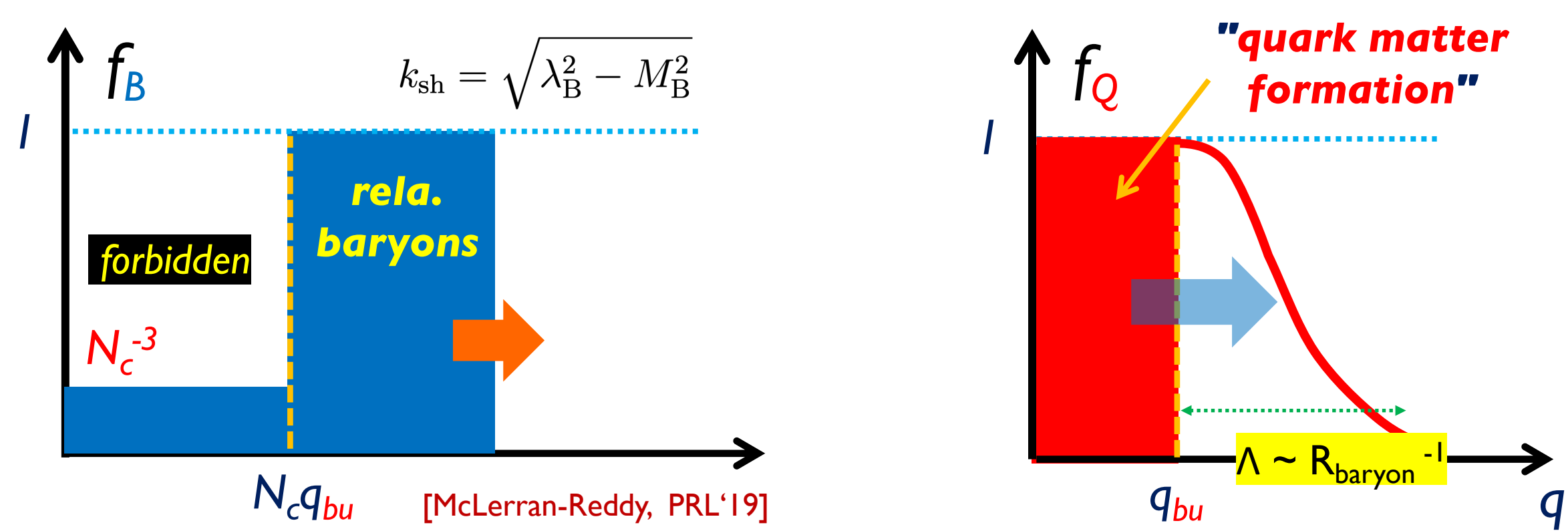
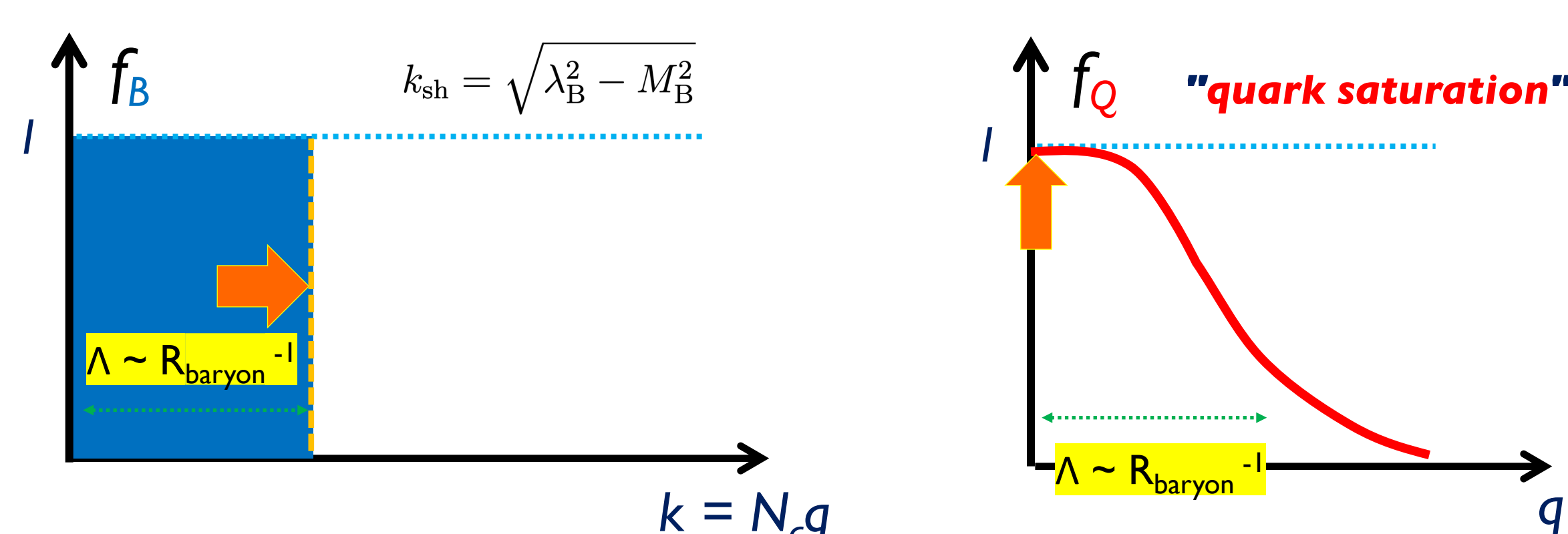
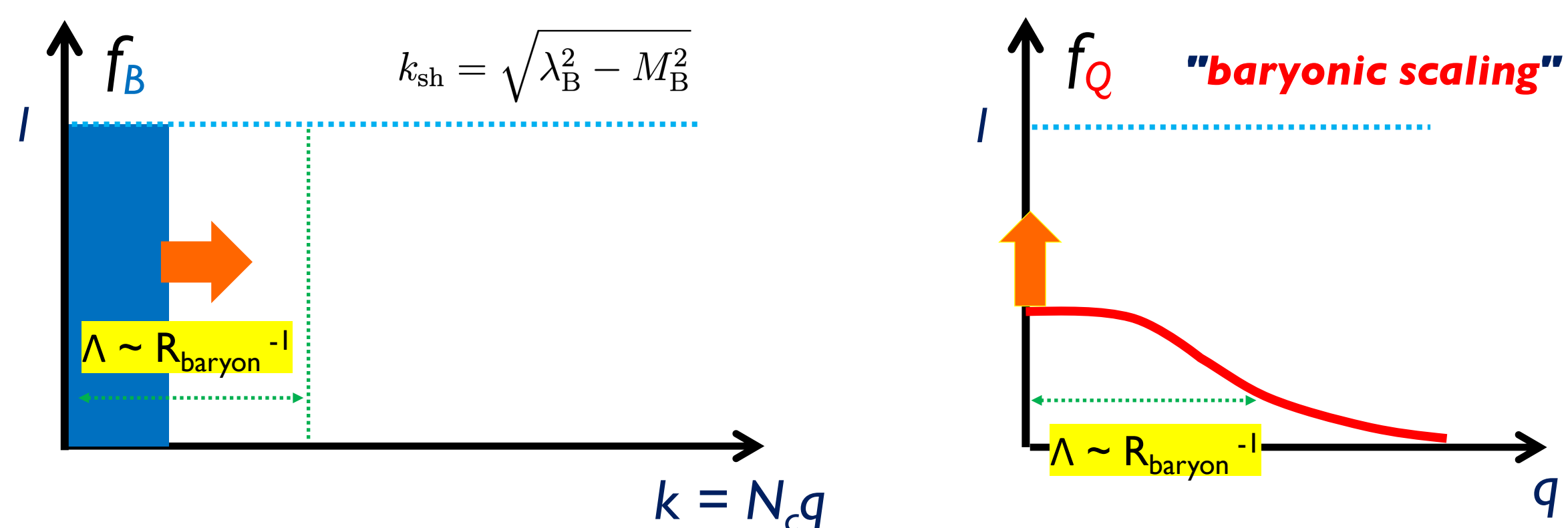
An ideal model (IdylliQ model)

- neglect interactions *except* confining forces
e.g.) 2-flavor hamiltonian: $\varepsilon_B[f_B] = 4 \int_k E_B(k) f_B(k)$
 $E_B(k) = \sqrt{M_B^2 + k^2}$
- keep using the same φ_Q (quarkyonic; quarks always confined)
- use a special quark distribution $\rightarrow$ sum rules analytically *invertible*
 $\varphi_{3d}(q) = \frac{2\pi^2}{\Lambda^3} \frac{e^{-q/\Lambda}}{q/\Lambda}$
 $\hat{L} = -\nabla^2 + \frac{1}{\Lambda^2}$
 $\hat{L}[\varphi(p-q)] = \frac{(2\pi)^3}{\Lambda^2} \delta(p-q)$
 $f_Q(q) = \int_{P_B} f_B(P_B) \varphi_Q^B(q - P_B/N_c)$
 $f_B(N_c q) = \frac{\Lambda^2}{N_c^3} \hat{L}[f_Q(q)]$

Quark saturation, quark matter formation, and stiffening

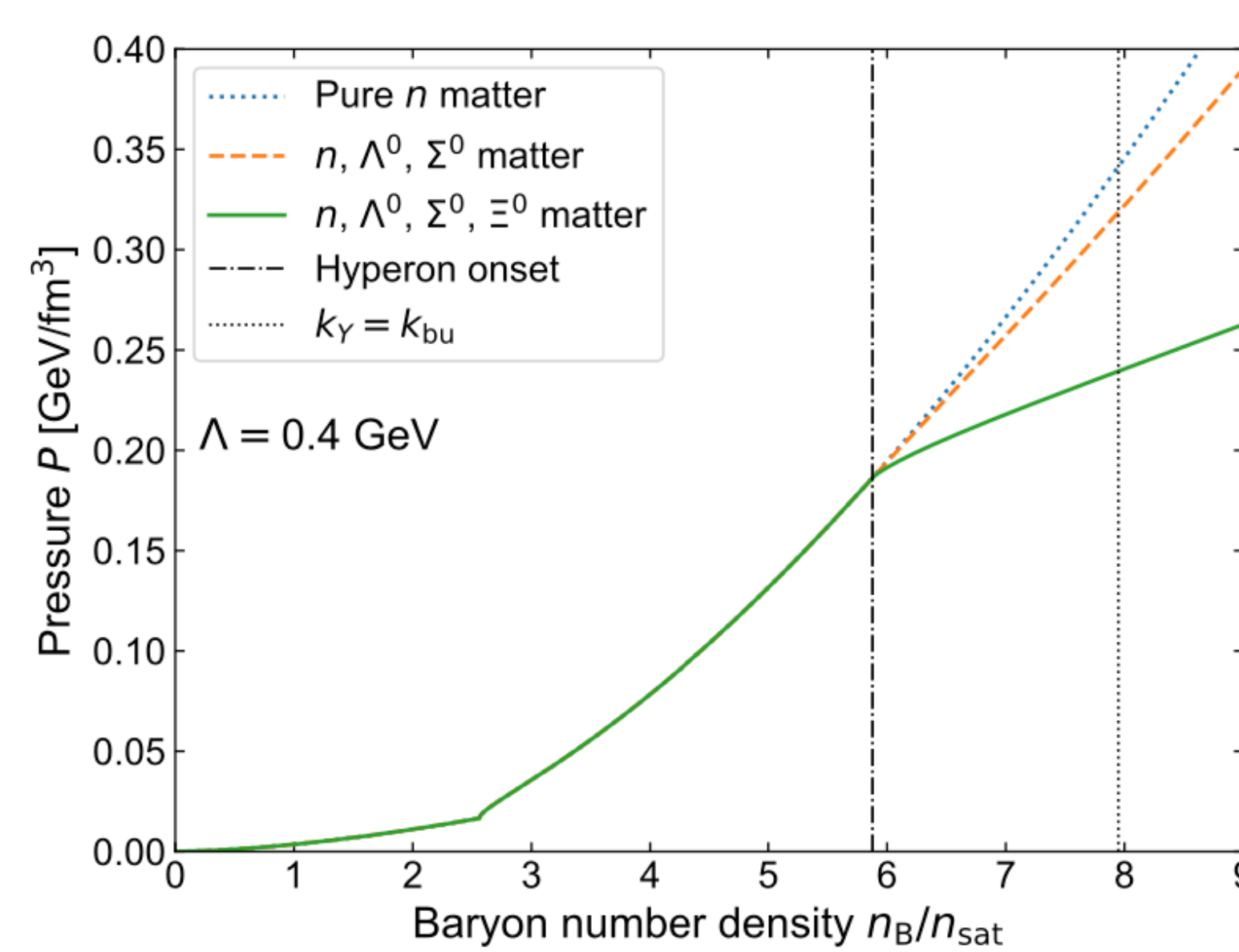
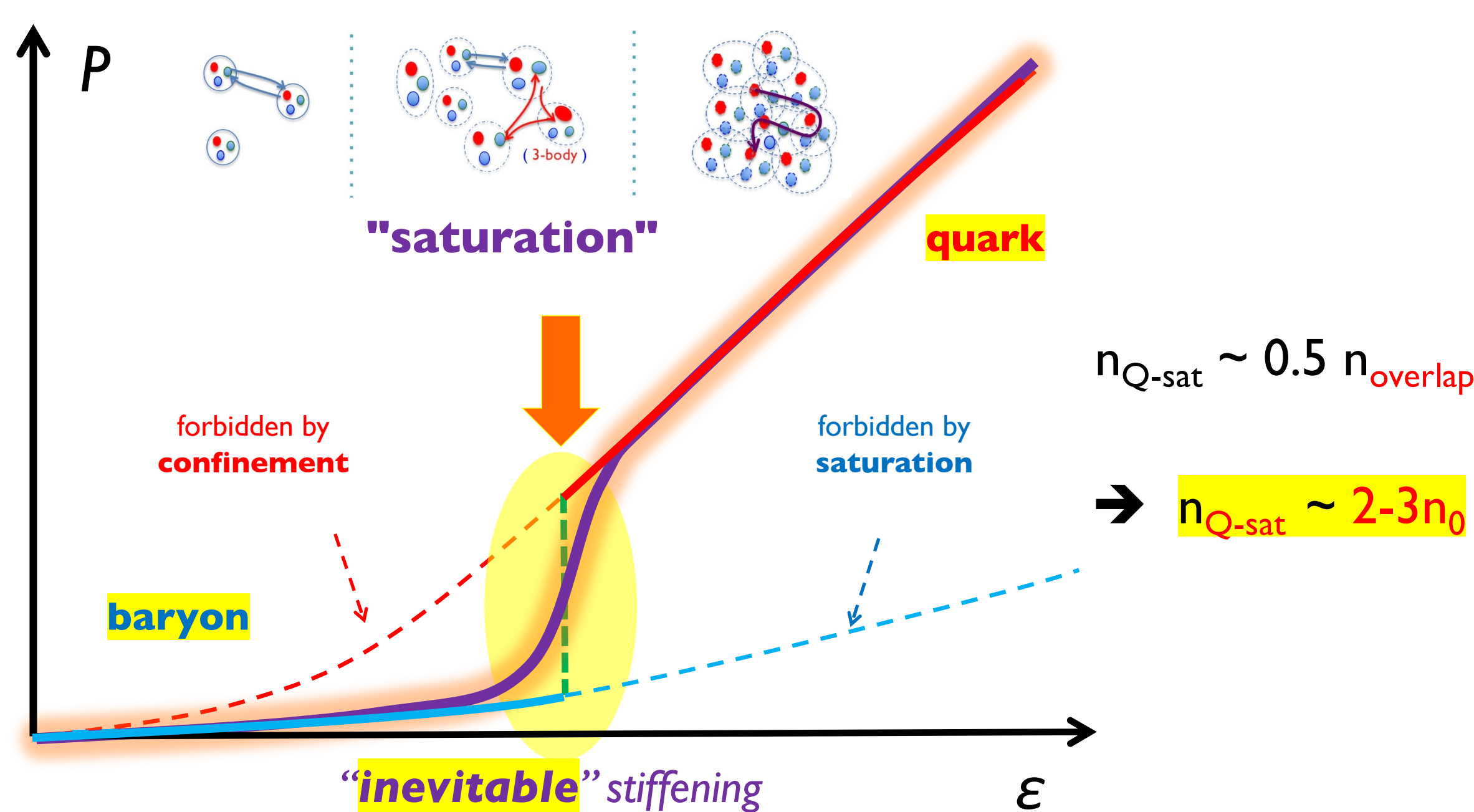
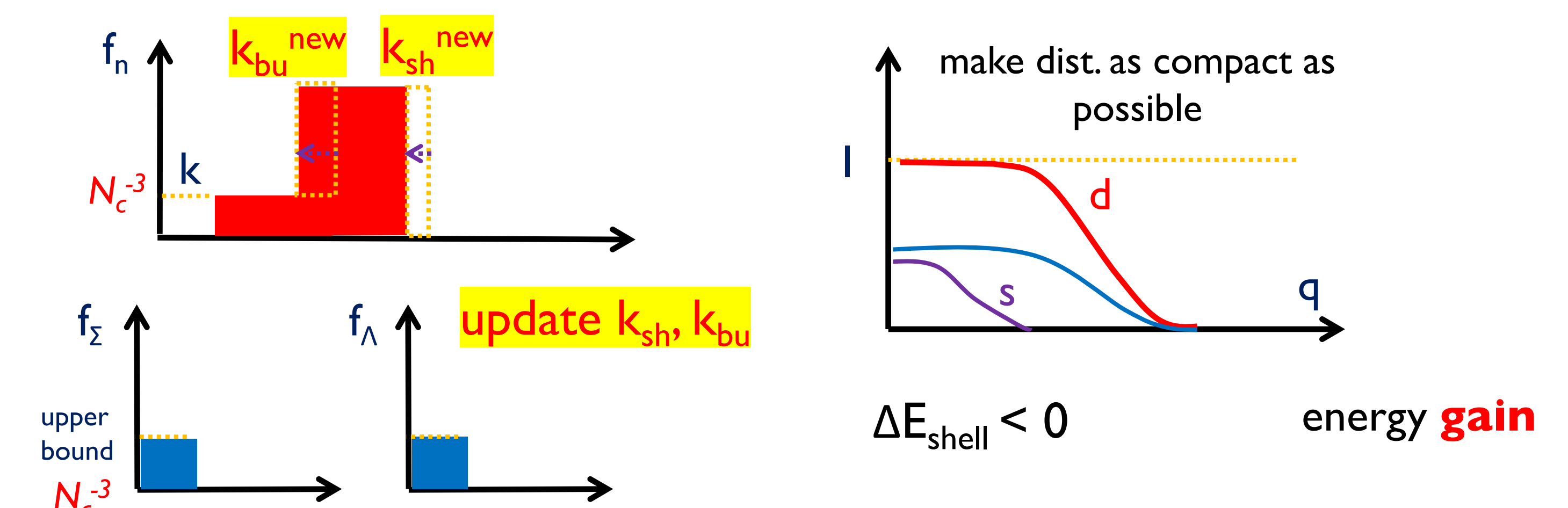
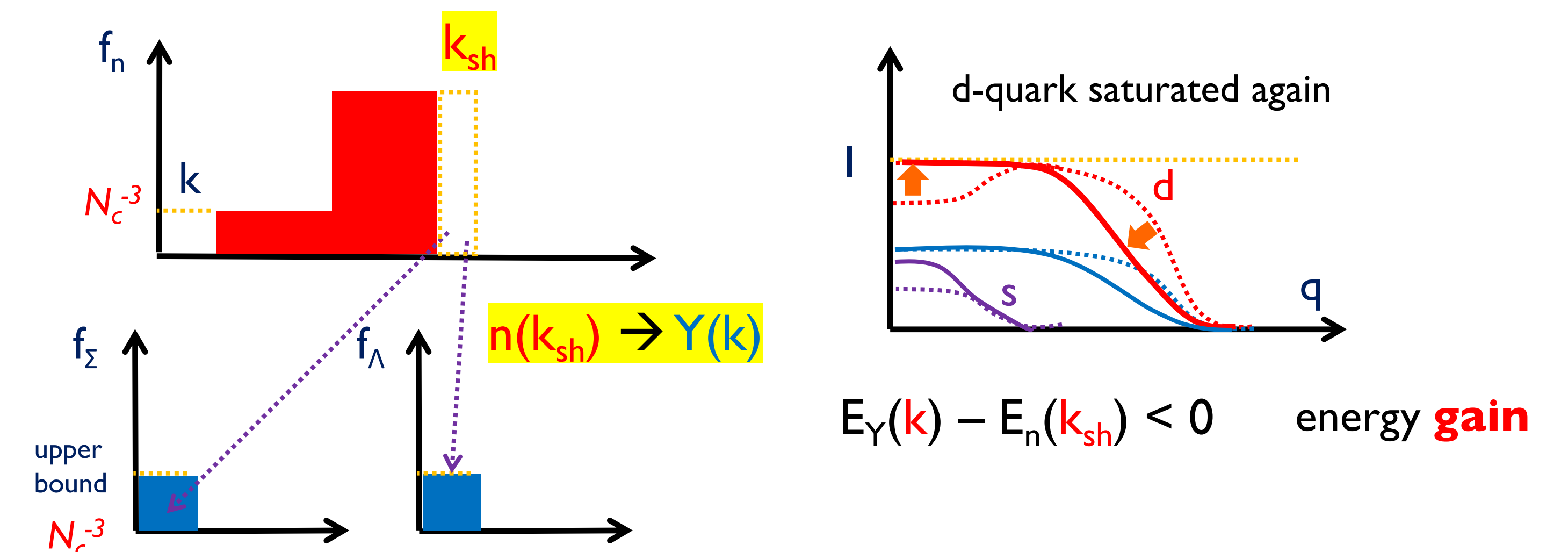
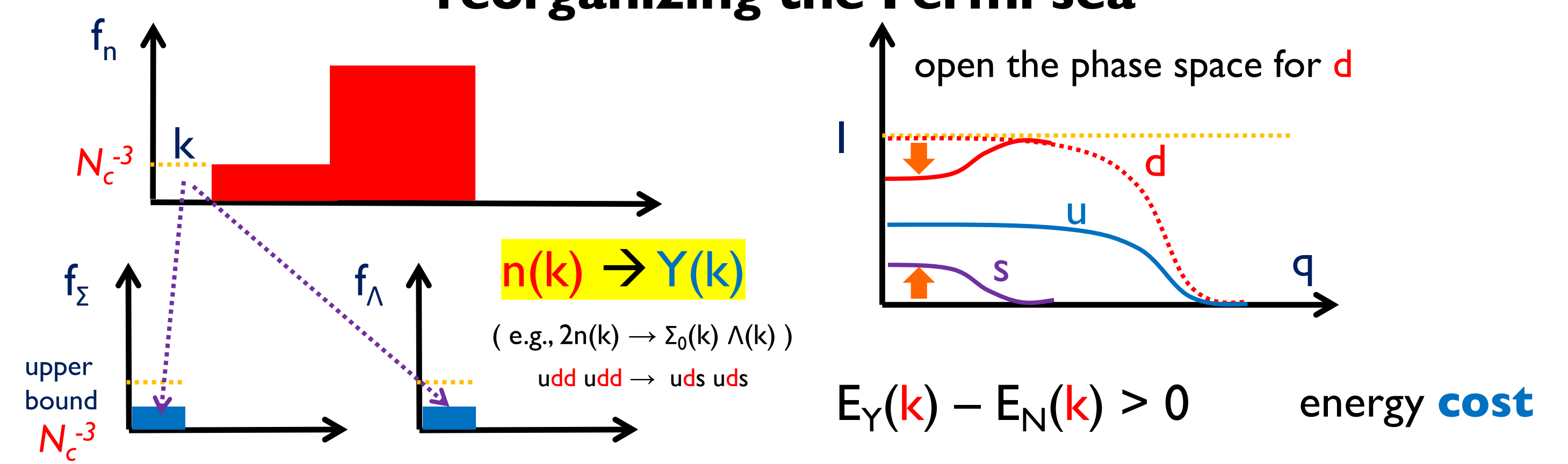
isospin symmetric matter

Density evolution



with hyperons (S=-1)

reorganizing the Fermi sea



Total change in energy :

$$\Delta E_{\text{tot}}(k) = 2E_Y(k) - E_N(k) - E_N(k_{\text{sh}}) + \Delta E_{\text{shell}}$$

($= -\mu_n = -\mu_B$)

$$\Delta E_{\text{tot}}(k=0) = 0$$

$$\mu_B^{\text{onset}} = 2M_Y - M_N$$

$$n_B^{Y, \Xi\text{-onset}} \sim 6n_0$$