

EXPLORING THE STRANGENESS FRONTIERS: KAONIC ATOMS X-RAY SPECTROSCOPY AT DAFNE WITH SIDDHARTA-2



HYP2025
TOKYO

Francesco Sgaramella INFN-LNF
on behalf of the SIDDHARTA-2 collaboration

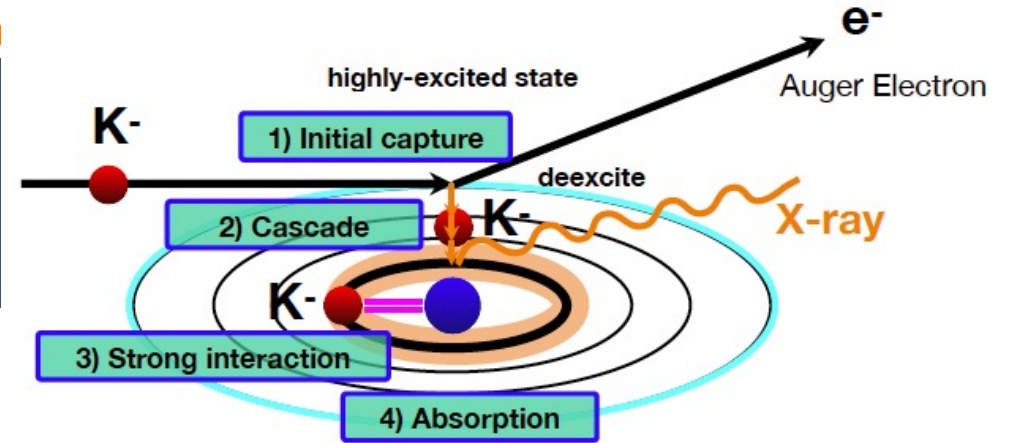


Istituto Nazionale di Fisica Nucleare
LABORATORI NAZIONALI DI FRASCATI



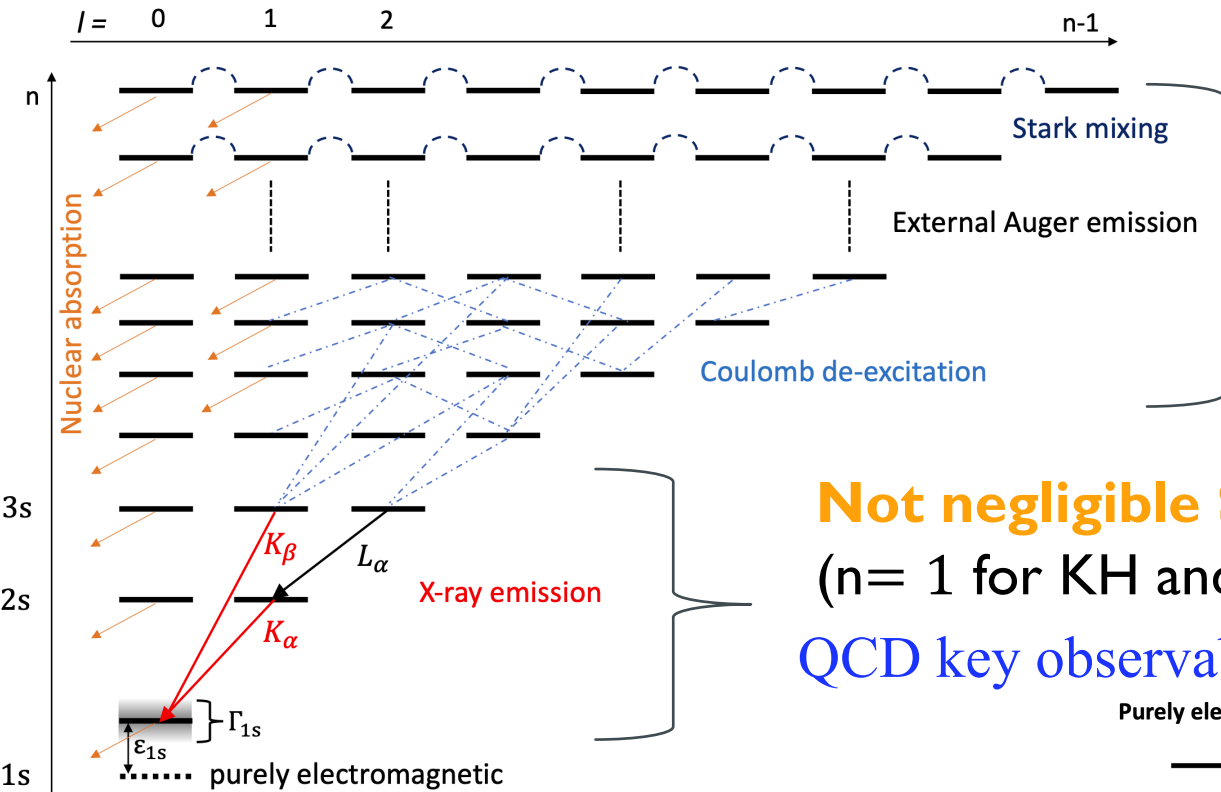
Kaonic Atoms X-ray Spectroscopy

Kaonic atoms X-ray spectroscopy to investigate the kaon nucleus interaction: from QED to QCD



Pure QED transition ($n > 4$)

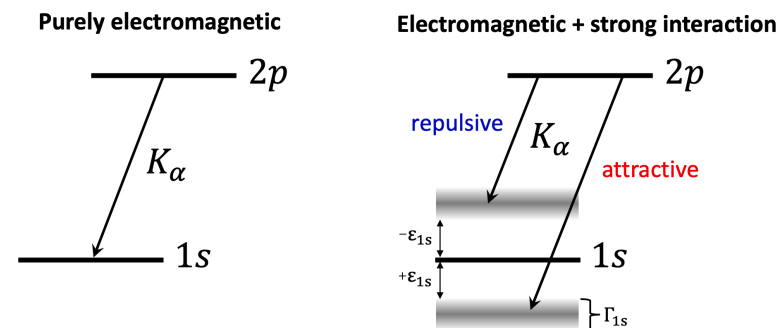
QED key observables: **Energy** and **X-ray Yield**
High precision test of BSQED



Not negligible Strong interaction effect

($n = 1$ for KH and Kd)

QCD key observables: **Shift (ϵ)**, **Width (Γ)** and **X-ray Yield**



Why Kaonic Atom?

Part. and Nuclear physics
QCD @ low-energy limit
Chiral symmetry, Lattice

Kaonic Atoms to Investigate Global Symmetry
Breaking Symmetry 12 (2020) 4, 547

Fundamental physics
New Physics

The modern era of light kaonic atom experiments
Rev.Mod.Phys. 91 (2019) 2, 025006

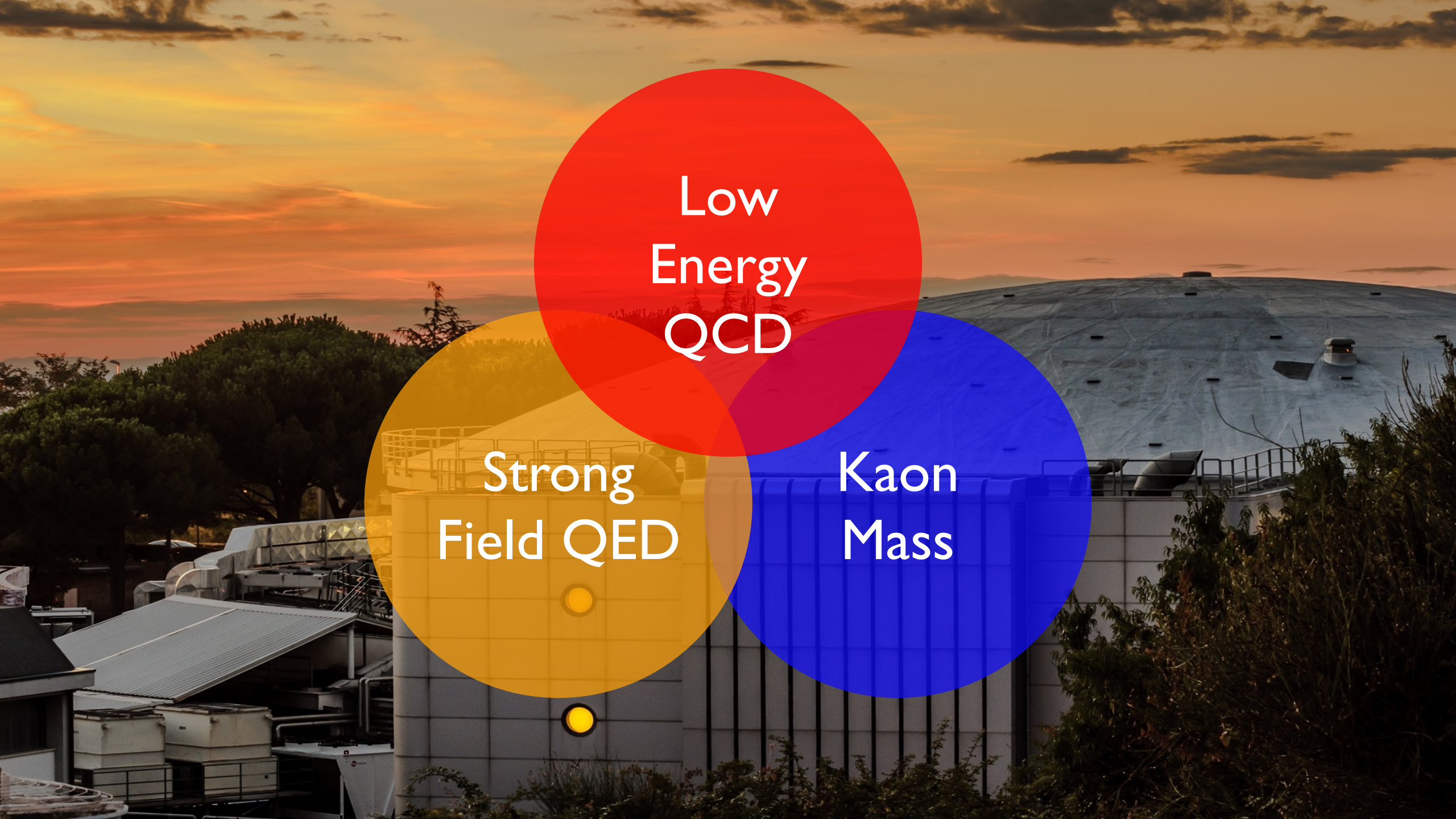
Kaonic atoms
Kaon-nuclei interactions (scattering and
nuclear interactions)

Dark Matter and BSM

Probing new hadronic forces with heavy exotic atoms
(2025) arXiv:2502.03537v1
On self-gravitating strange dark matter halos around
galaxies Phys.Rev.D 102 (2020) 8, 083015

The equation of state of dense matter: Stiff, soft, or
both? Astron.Nachr. 340 (2019) 1-3, 189

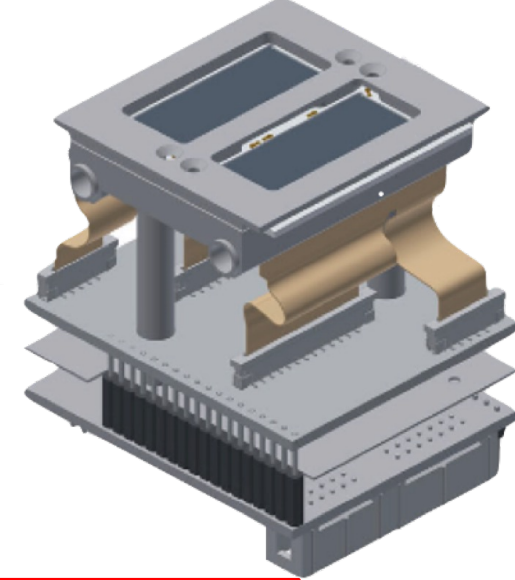
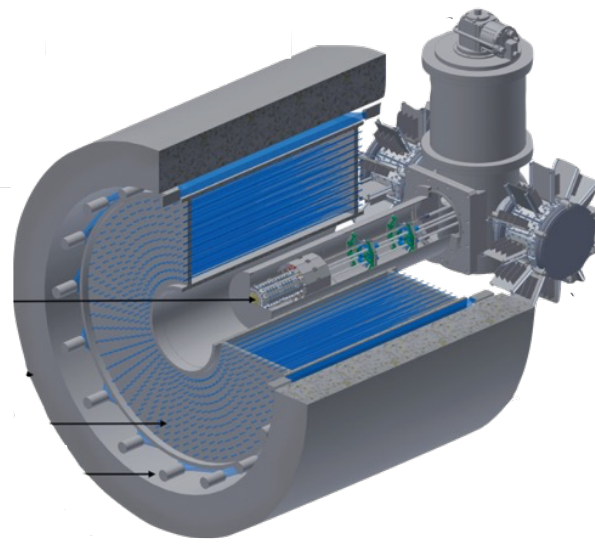
Astrophysics
EOS Neutron Stars



Low
Energy
QCD

Strong
Field QED

Kaon
Mass



The modern era of light kaonic atom experiments

Catalina Curceanu, Carlo Guaraldo, Mihail Iliescu, Michael Cargnelli, Ryugo Hayano, Johann Marton, Johann Zmeskal, Tomoichi Ishiwatari, Masa Iwasaki, Shinji Okada, Diana Laura Sirghi, and Hideyuki Tatsuno

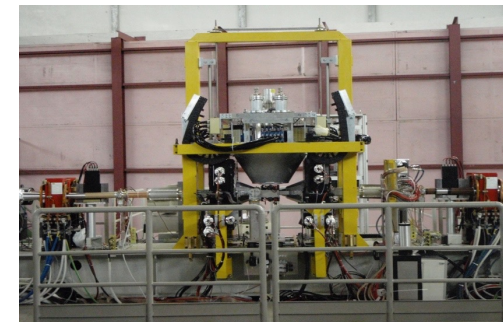
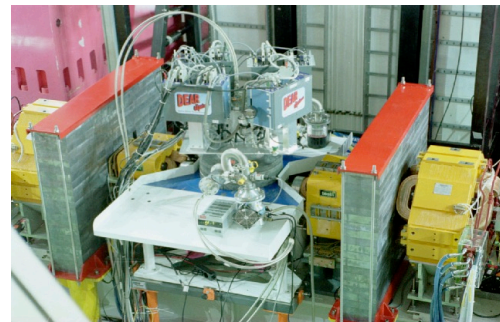
Rev. Mod. Phys. **91**, 025006 – Published 20 June 2019



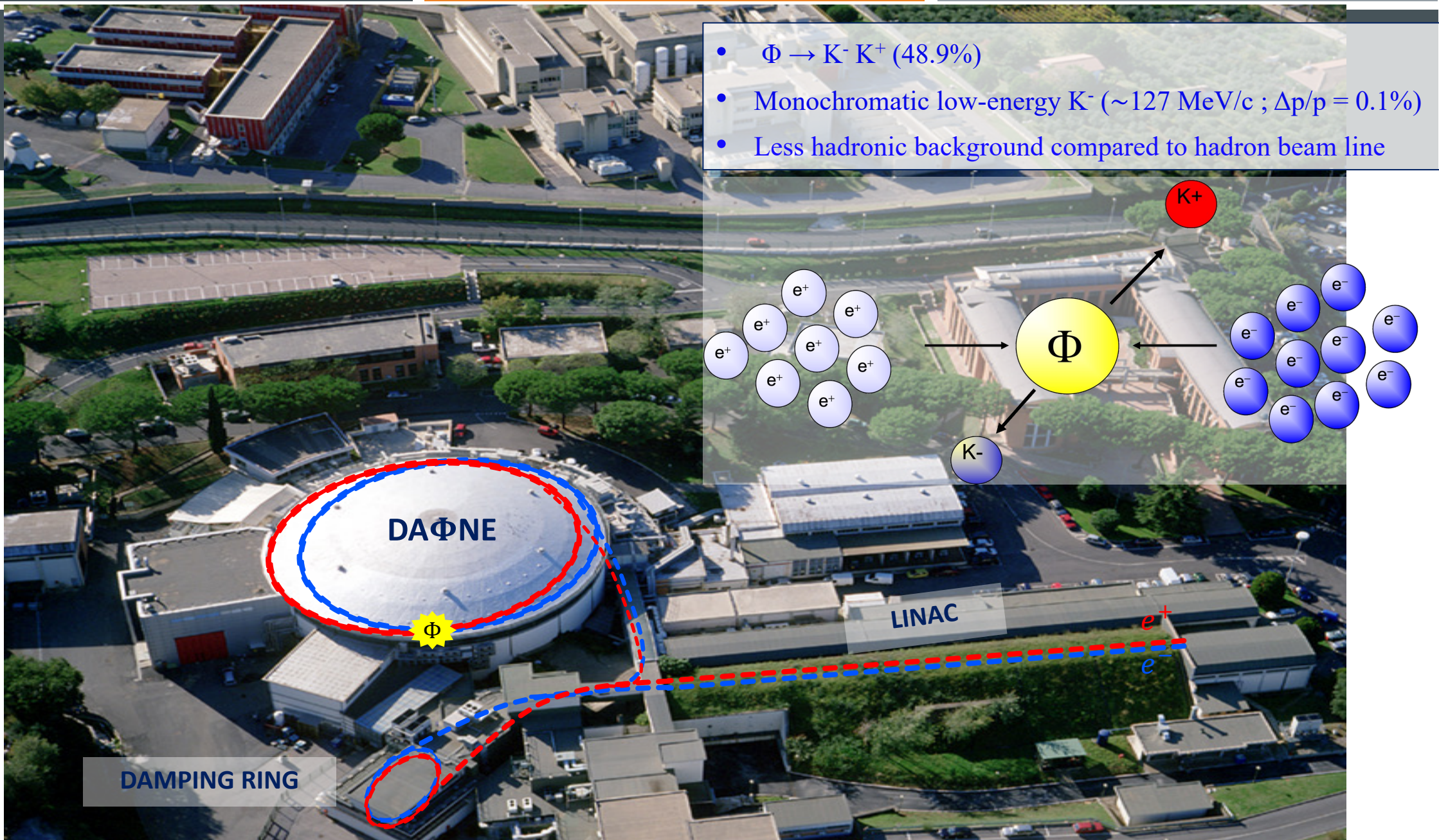
DEAR
2002

SIDDHARTA
2009

SIDDHARTA-2
2022

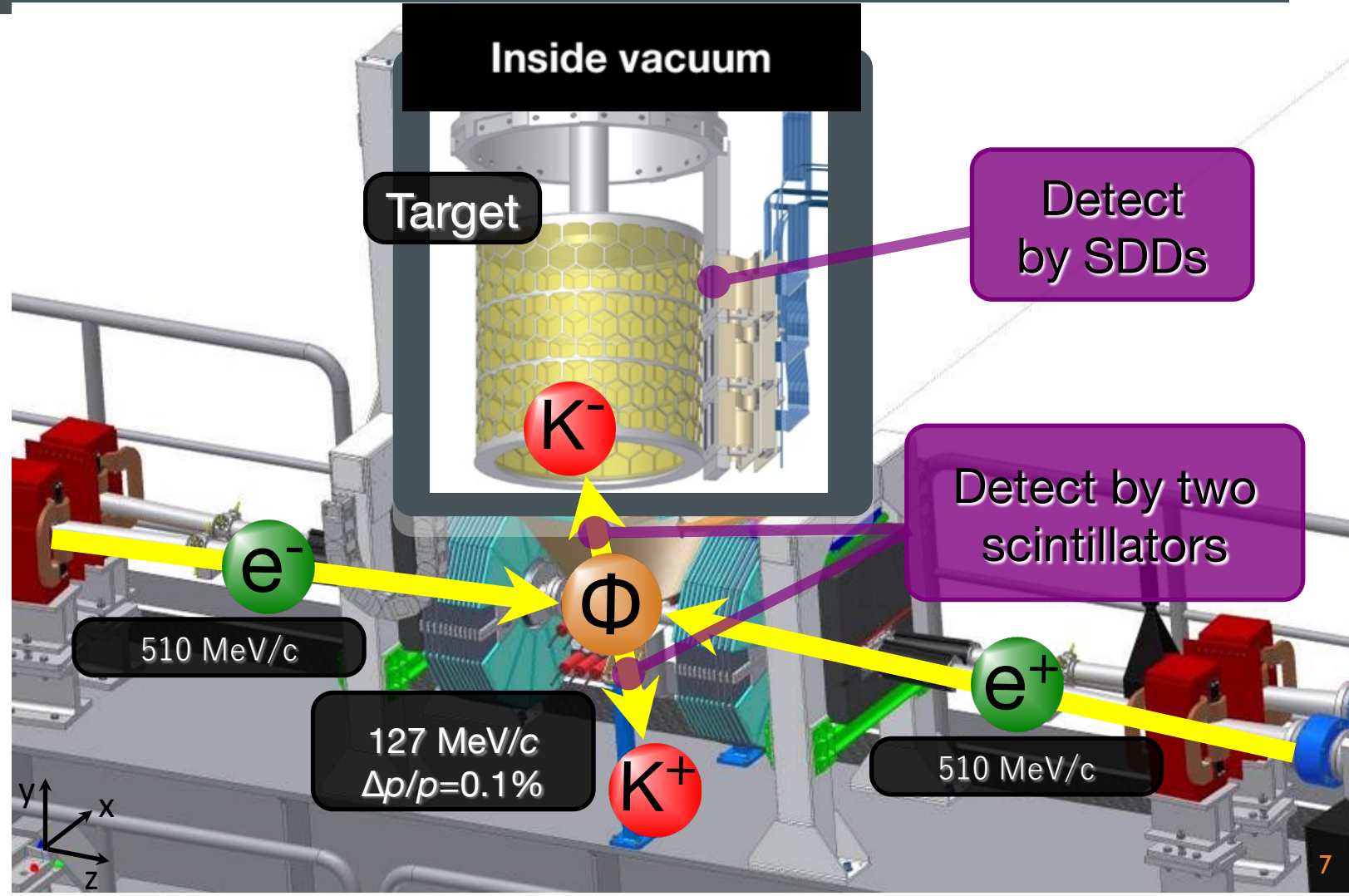
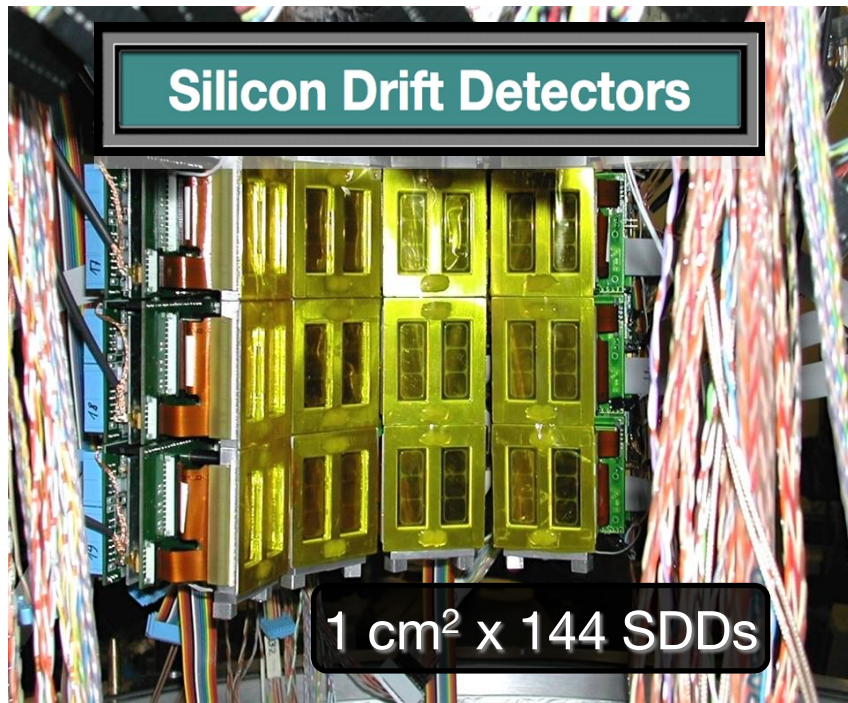


The DAΦNE collider of INFN-LNF



The SIDDHARTA Experiment (2009)

A cryogenic gaseous target and Silicon Drift Detectors to perform the kaonic hydrogen measurement

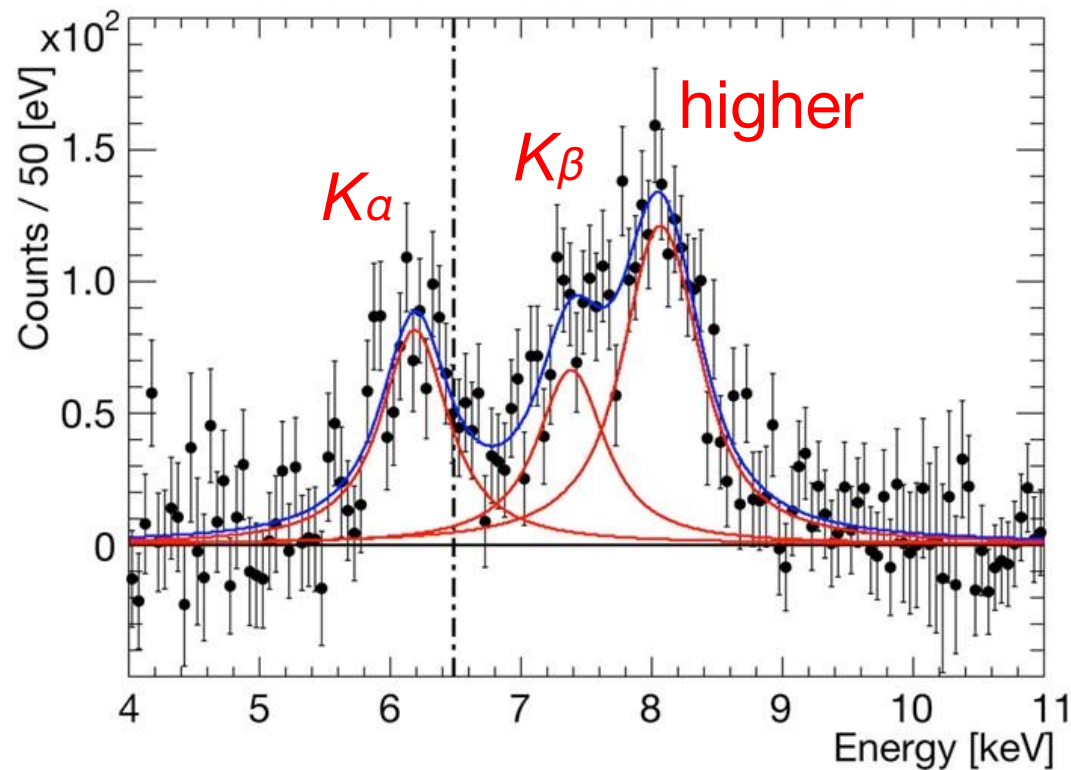


The SIDDHARTA Experiment (2009)

The most precise measurement of kaonic hydrogen 1s shift and width performed by SIDDHARTA was fundamental to constrain the description of the K-p interaction at threshold

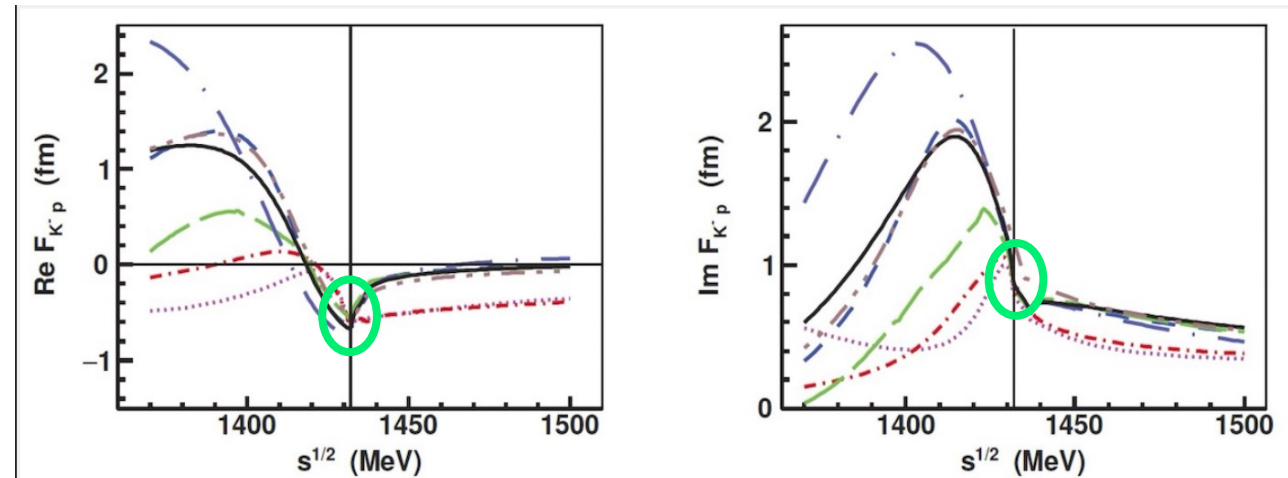
$$\varepsilon_{1s} = -283 \pm 36(\text{stat}) \pm 6(\text{syst}) \text{ eV}$$

$$\Gamma_{1s} = 541 \pm 89(\text{stat}) \pm 22(\text{syst}) \text{ eV}$$



C. Curceanu et al., *Phys. Lett. B* **704** (2011) 113

K-p scattering amplitudes generated by recent chirally motivated approaches. The vertical lines mark the threshold energy

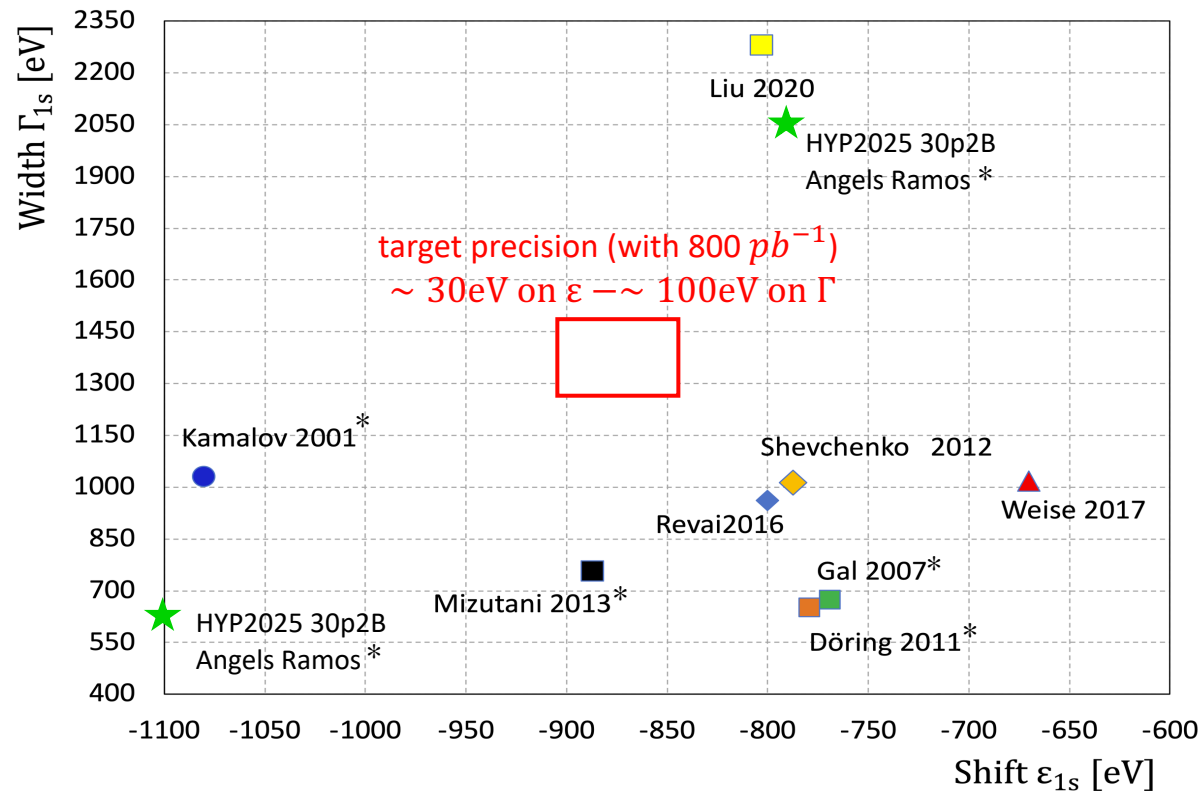


Ciepl y, A. et al. *From KN interactions to K-nuclear quasi-bound states. AIP Conf. Proc.* 2249, 030014 (2020).

The SIDDHARTA-2's experiment main aim

Main scientific goal: **first measurement ever of kaonic deuterium X-ray transition** to the ground state (1s-level) such as to determine its **shift and width induced by the strong interaction**, providing unique data to investigate the **QCD in the non-perturbative regime** with strangeness.

Kd theoretical prediction



“The most important experiment to be carried out in low energy K-meson physics today is the **definitive determination of the energy level shifts in the K-p and K-d atoms**, because of their direct connection with the physics of $\bar{K}N$ interaction and their complete independence from all other kinds of measurements which bear on this interaction”.

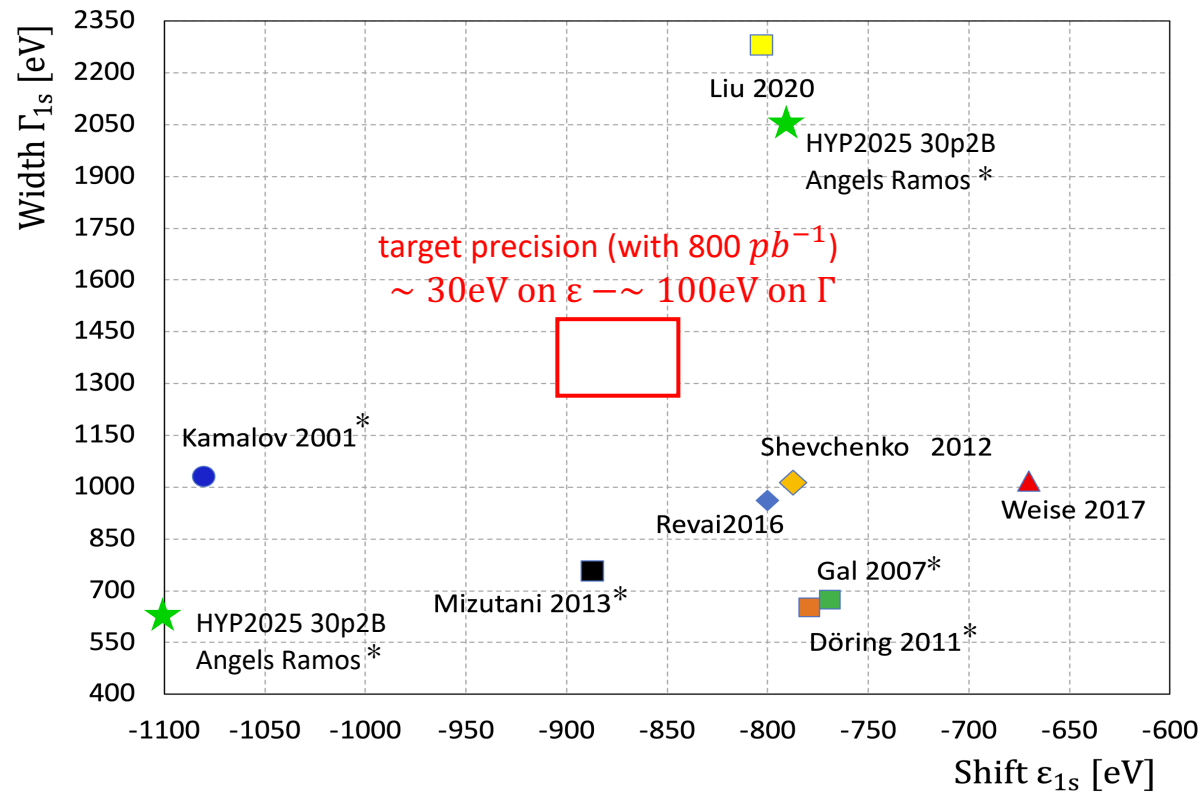
R.H. Dalitz (1982)

* The energy shift and width are extracted from scattering length using the Deser-Trueeman formula

The SIDDHARTA-2's experiment main aim

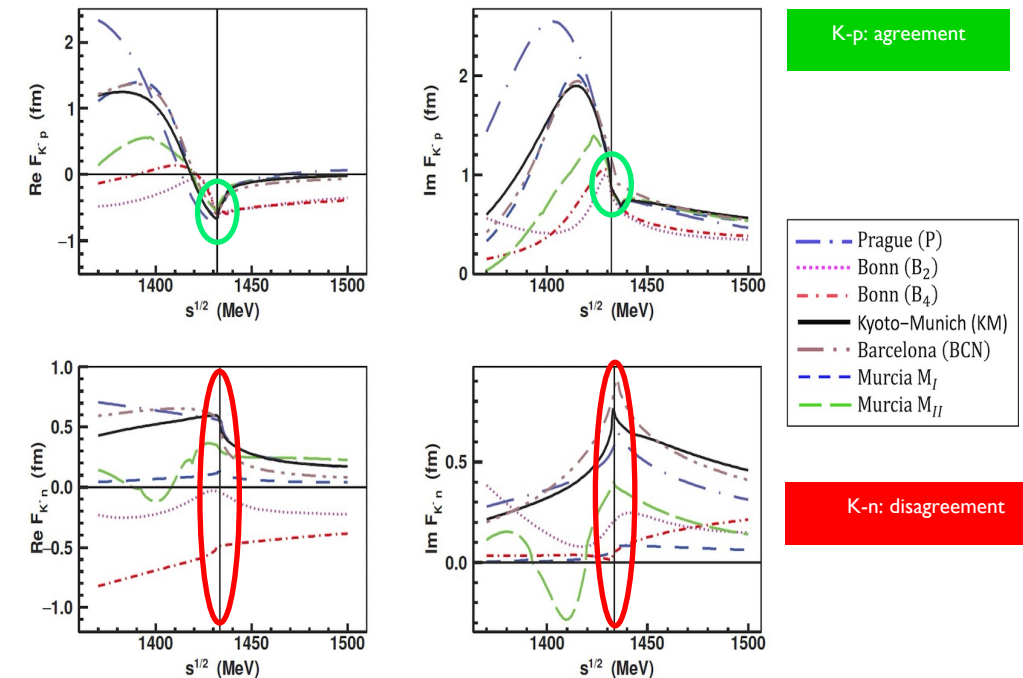
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Kd theoretical prediction



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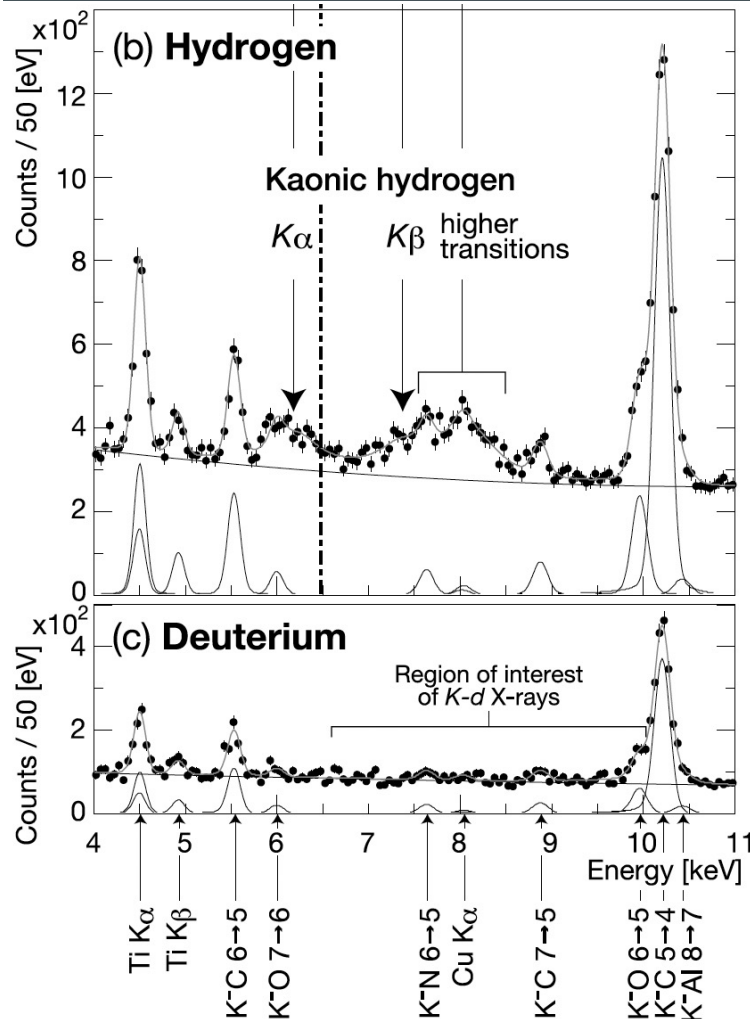
Combined analysis of the kaonic deuterium and kaonic hydrogen measurements to determine the isospin-dependent $\bar{K}N$ scattering lengths



Ciepl y, A. et al. From $\bar{K}N$ interactions to $\bar{K}$ -nuclear quasi-bound states. AIP Conf. Proc. 2249, 030014 (2020).

The kaonic deuterium challenge

The measurement of Kaonic deuterium $2p \rightarrow 1s$ is a true challenge: SIDDHARTA performed an exploratory run in 2009, collecting 100 pb^{-1} of data but without observing a visible signal.



Physics factors:

- Low X-ray yield (~ 10 times lower than KH) See K.Toho Poster
- Transition width broader than KH $2p \rightarrow 1s$

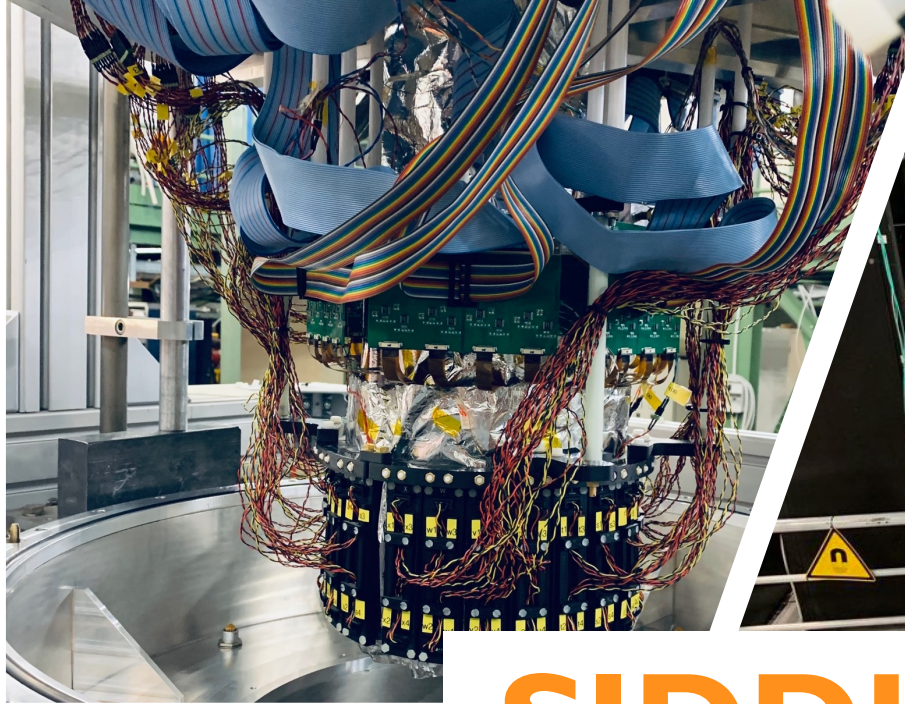
➤ Kd requires a higher integrate luminosity $\sim 800 \text{ pb}^{-1}$

Background:

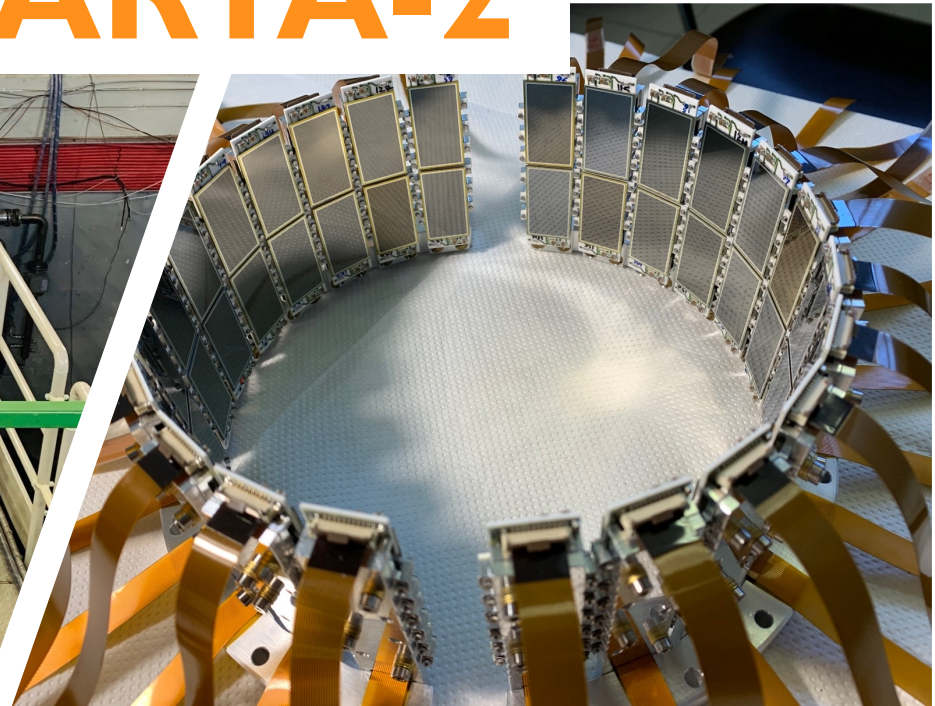
- The Kd measurement requires an improvement in signal/background ratio by a factor 10

➤ New experimental apparatus with improved SDDs, trigger and Veto systems

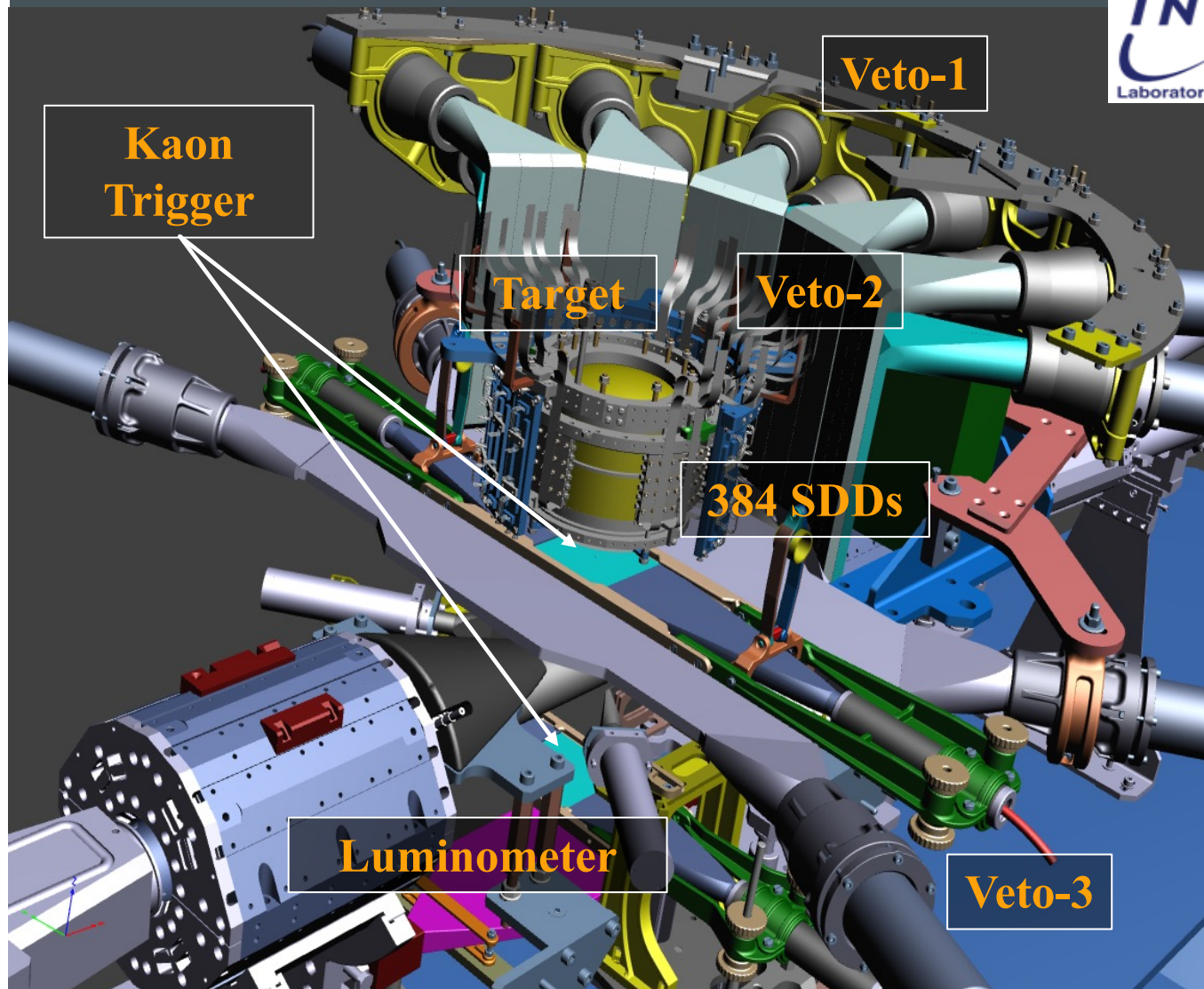
➤ Larger detection area



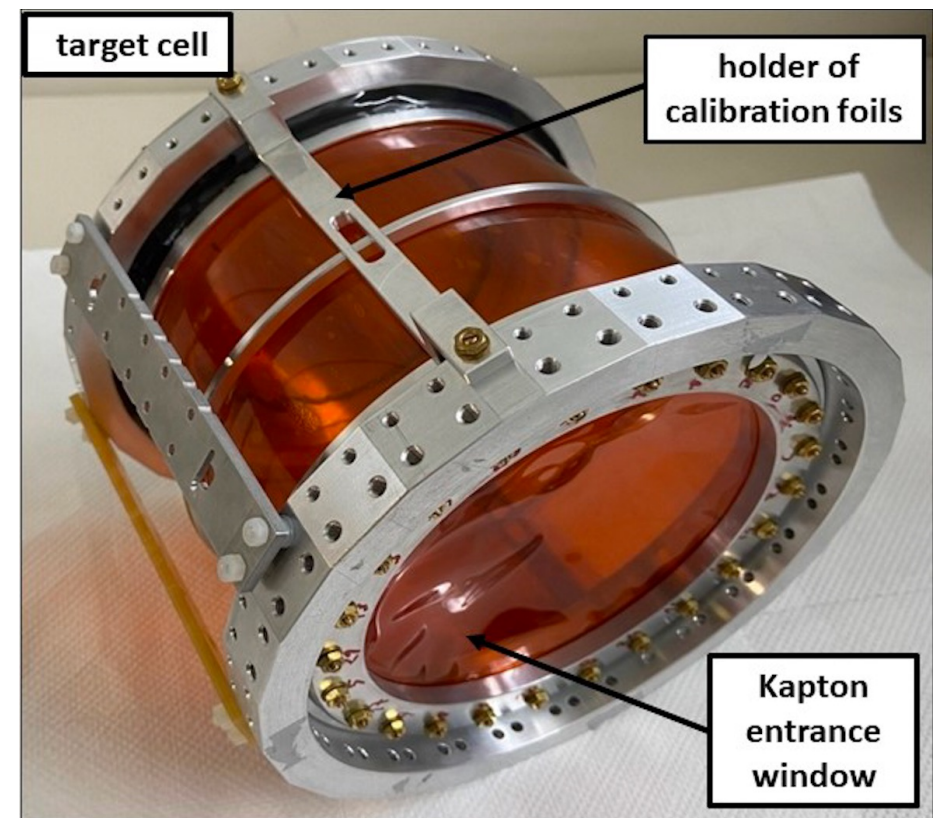
SIDDHARTA-2



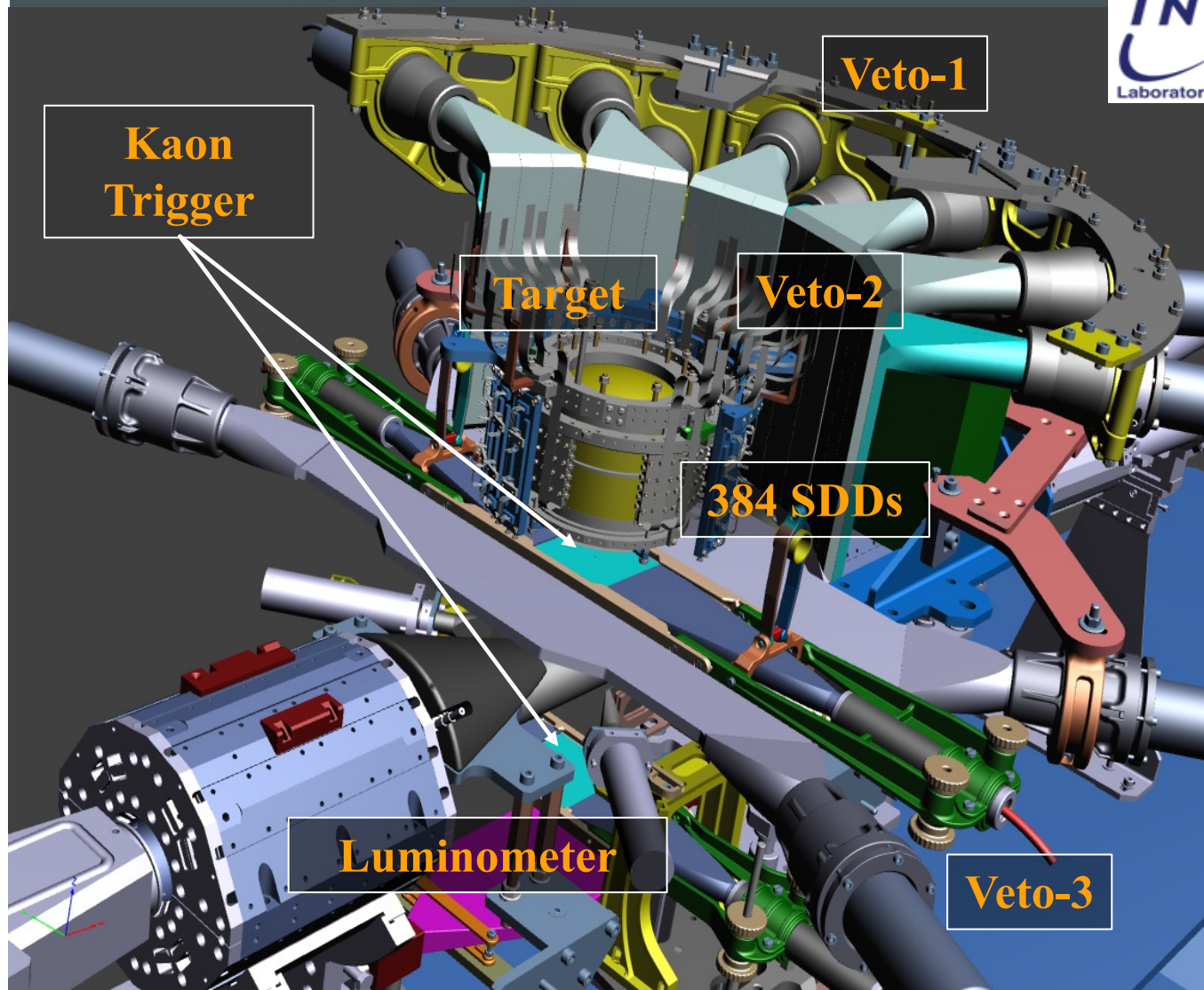
The SIDDHARTA-2 setup



Cryogenic gaseous target cell

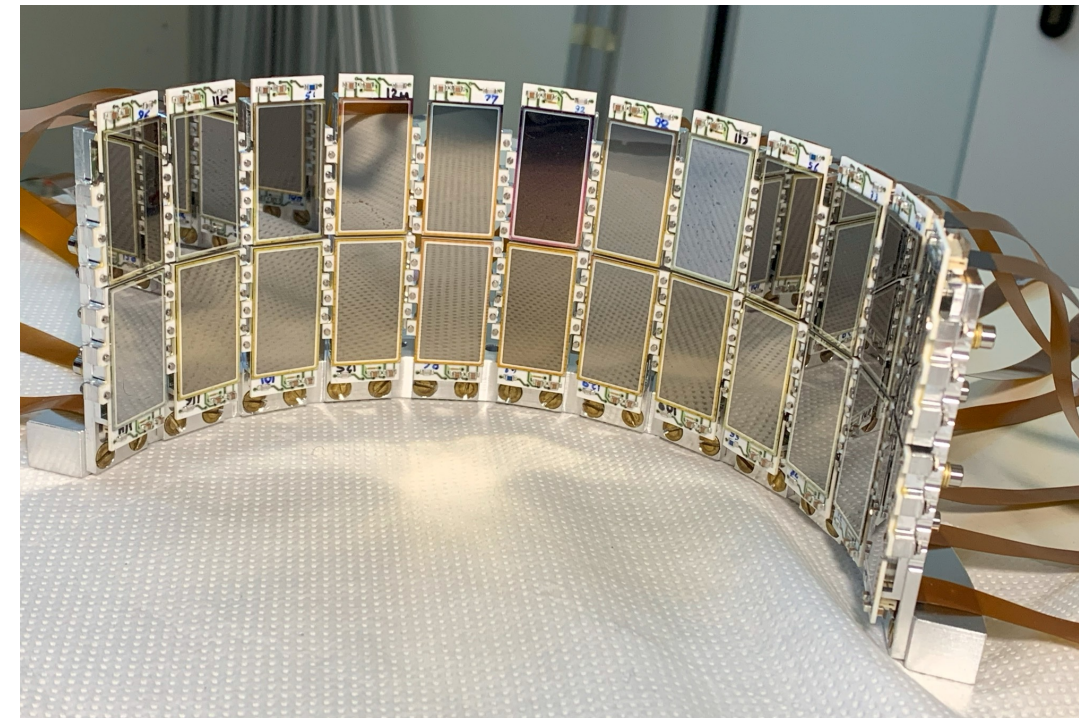


The SIDDHARTA-2 setup



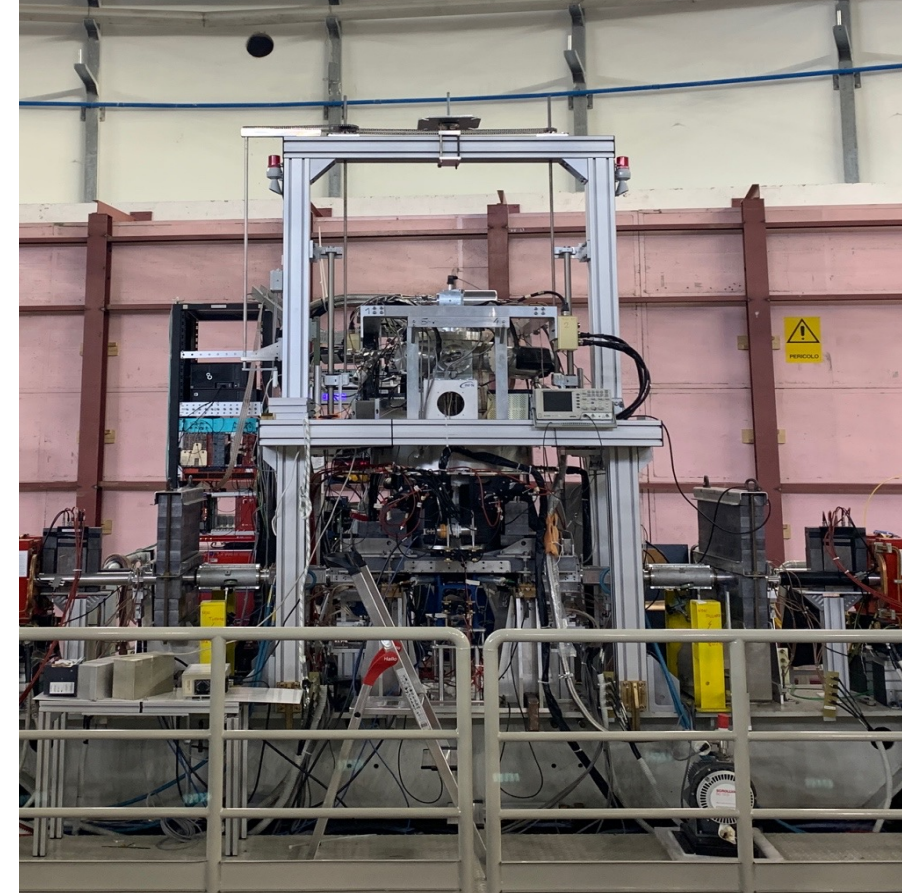
384 Silicon Drift Detector
for a total active area of 246 cm^2

The thickness of $450 \text{ }\mu\text{m}$ ensures a high collection efficiency for X-rays of energy in the 5-12 keV range



Kaonic Deuterium Data Taking Summary

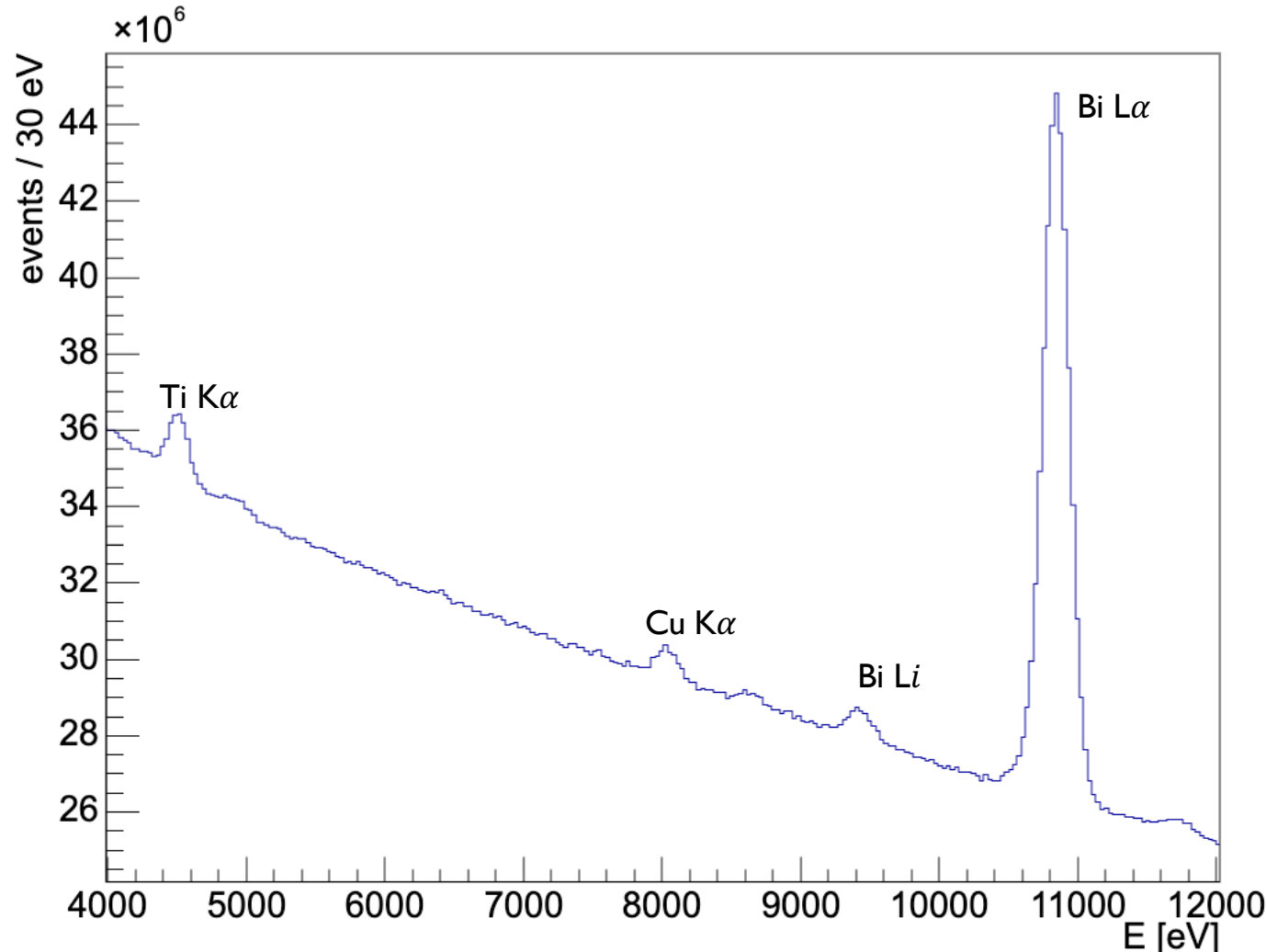
- *Kaonic neon*: initial calibration and optimization of the setup
- *K-d Run-1: int. luminosity 196 pb^{-1} (May – July 2023)*
- *Kaonic helium-4*: final calibration of the setup
- *Kaonic neon*: initial calibration of the setup
- *K-d Run-2: int. luminosity 338 pb^{-1} (October – December 2023)*
- *Kaonic hydrogen*: final calibration of the setup
- *Kaonic hydrogen*: initial calibration of the setup
- *K-d Run-3: int. luminosity 425 pb^{-1} (February – April 2024)*
- *Kaonic hydrogen*: final calibration of the setup
- *Low density run: int. luminosity 184 pb^{-1} (May – July 2024)*
- *Post Kd calibration run (July 2024) – 20 pb^{-1} with solid targets (B and F)*



K-d total integrated luminosity: 1143 pb^{-1}

Kaonic deuterium analysis

Inclusive energy spectrum: the continuous background and the fluorescence peaks are due to the electromagnetic (asynchronous) and hadronic (synchronous) background



- Electromagnetic (asynchronous)

background: the electromagnetic shower produced in the accelerator pipe (and other setup materials) invested by e-/e+ lost from the beam overlaps the signal; the loss rate in the interaction region reaches few MHz. The main contribution comes from Touschek effect → **Kaon Trigger** and **SDDs drift time**

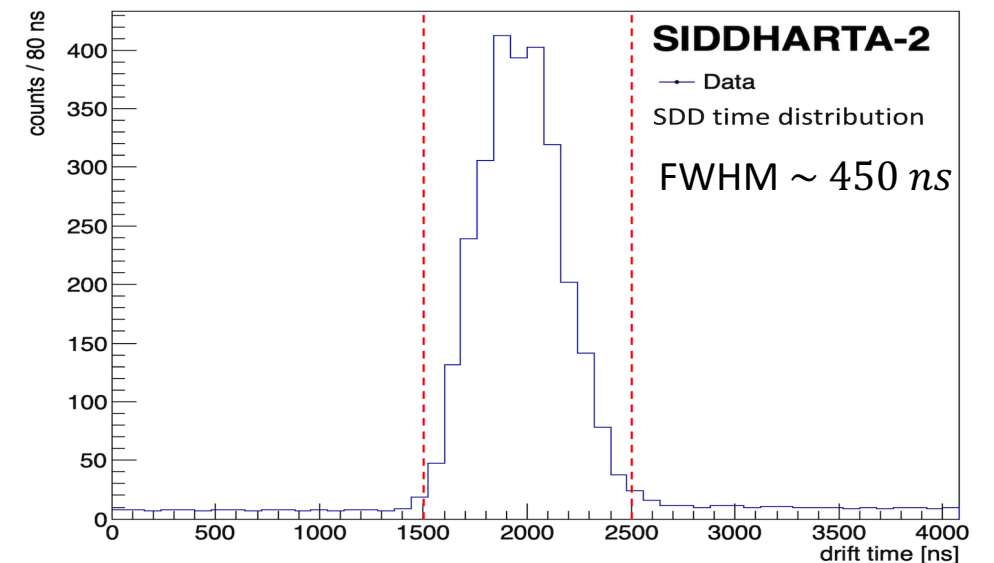
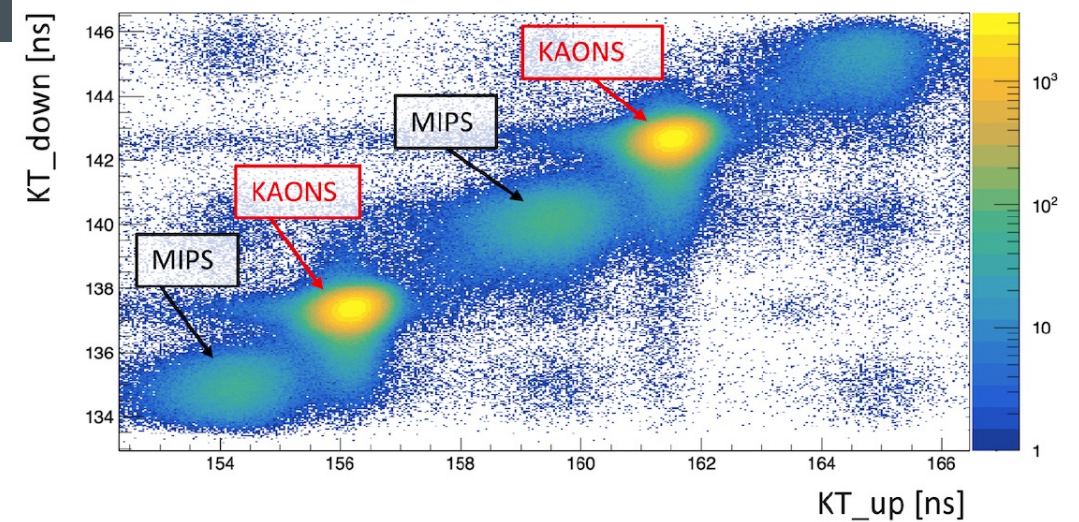
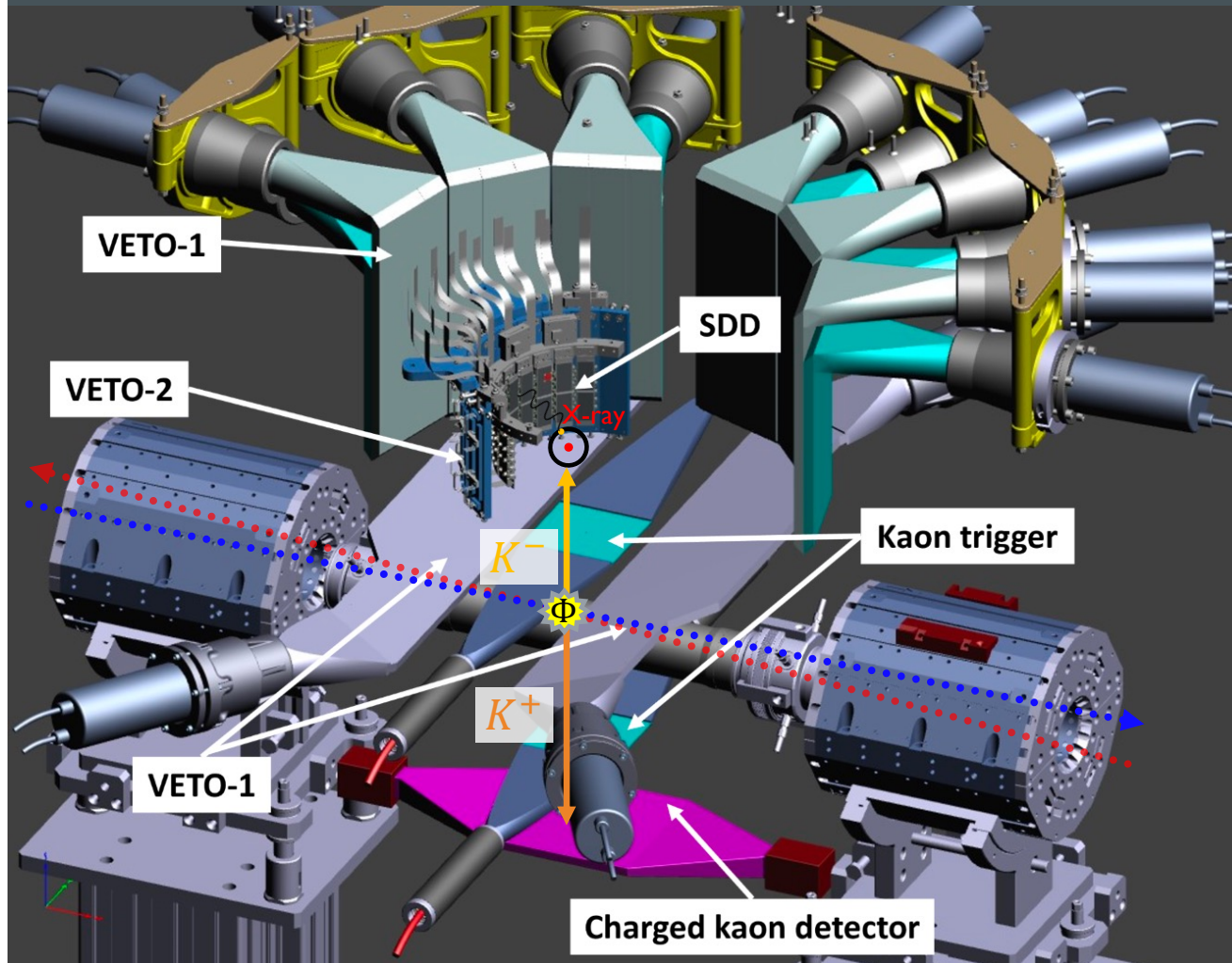
-Hadronic (synchronous) background:

associated to kaon absorption on materials nuclei, or to other Φ decay channels. It can be considered a hadronic background.

-Spectra contamination by Xray fluorescence or by X-rays produced in higher transitions of other kaonic atoms, formed in the setup materials → **Veto systems**

Kaonic deuterium analysis

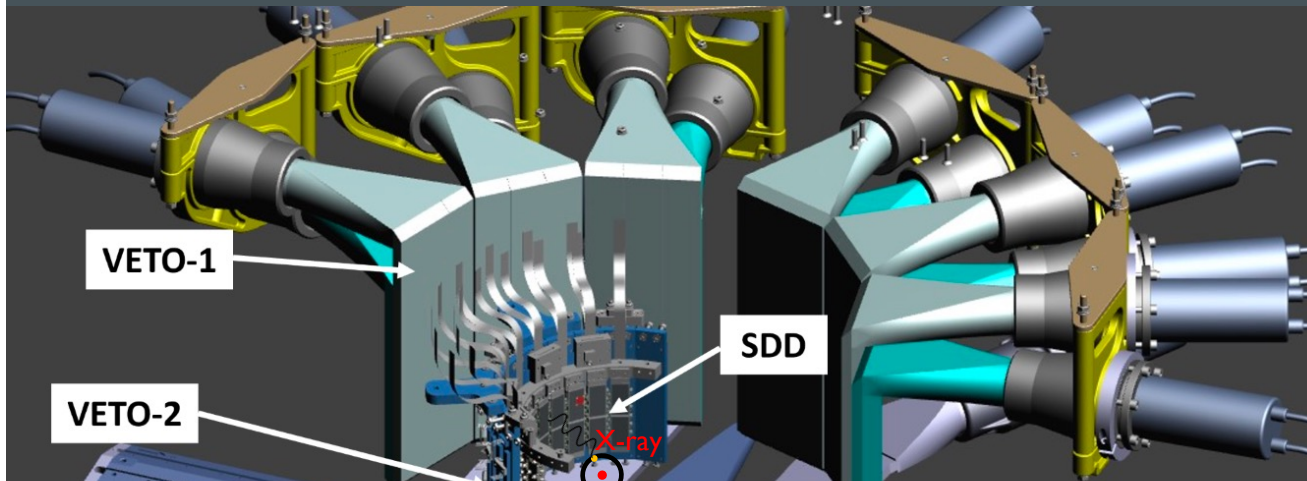
The trigger is generated by the coincidence of 2 back-to-back scintillators. The ToF is different (~ 2 ns) for Kaons, $m(K) \sim 500$ MeV/ c^2 and light particles originating from beam-beam and beam-environment interaction (MIPs).
ToF analysis can efficiently discriminate the Kaons from MIPs!



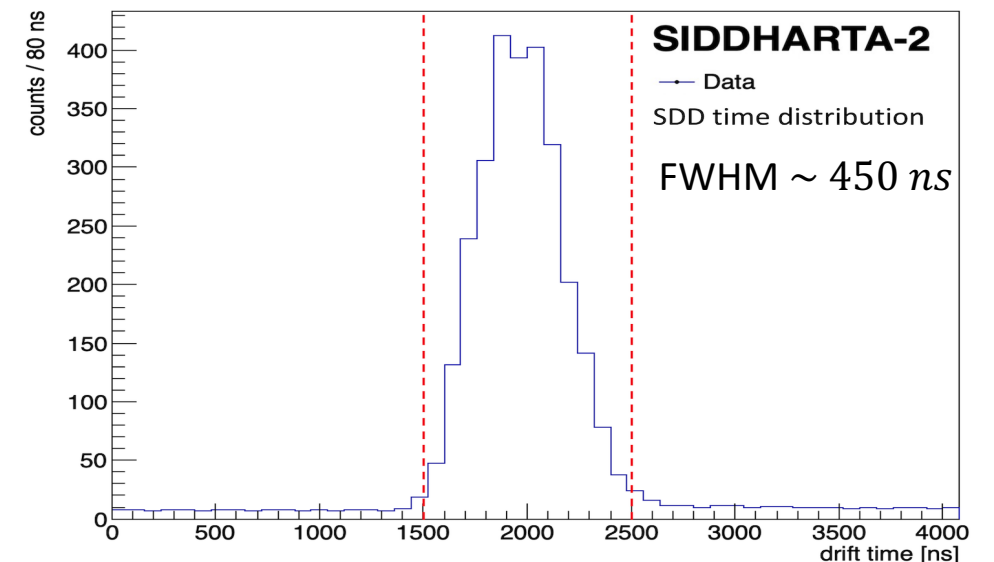
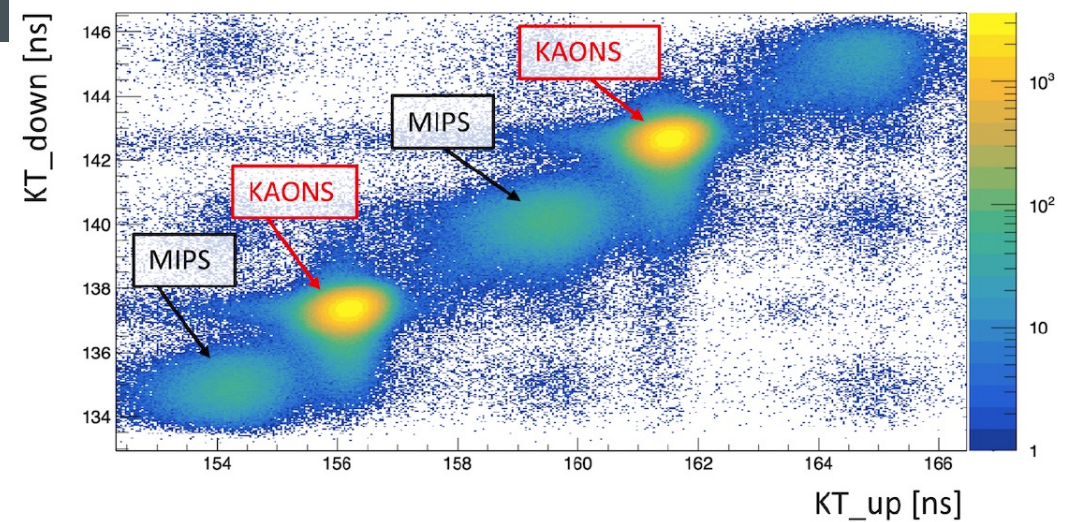
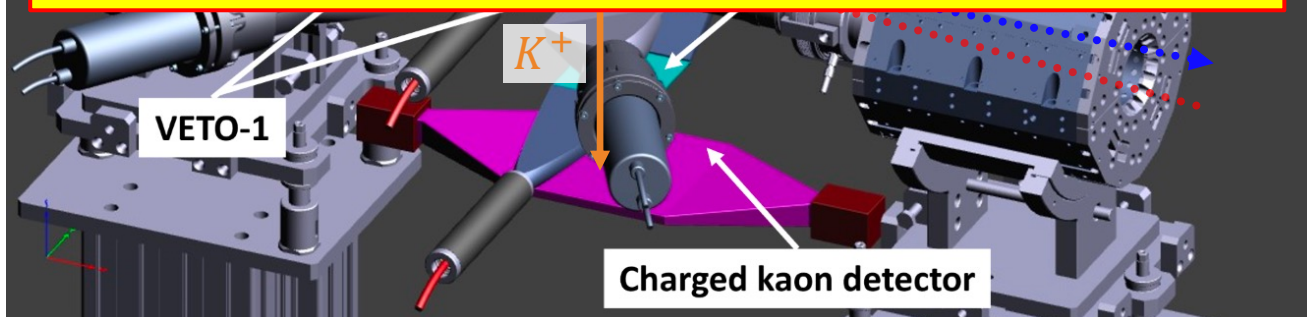
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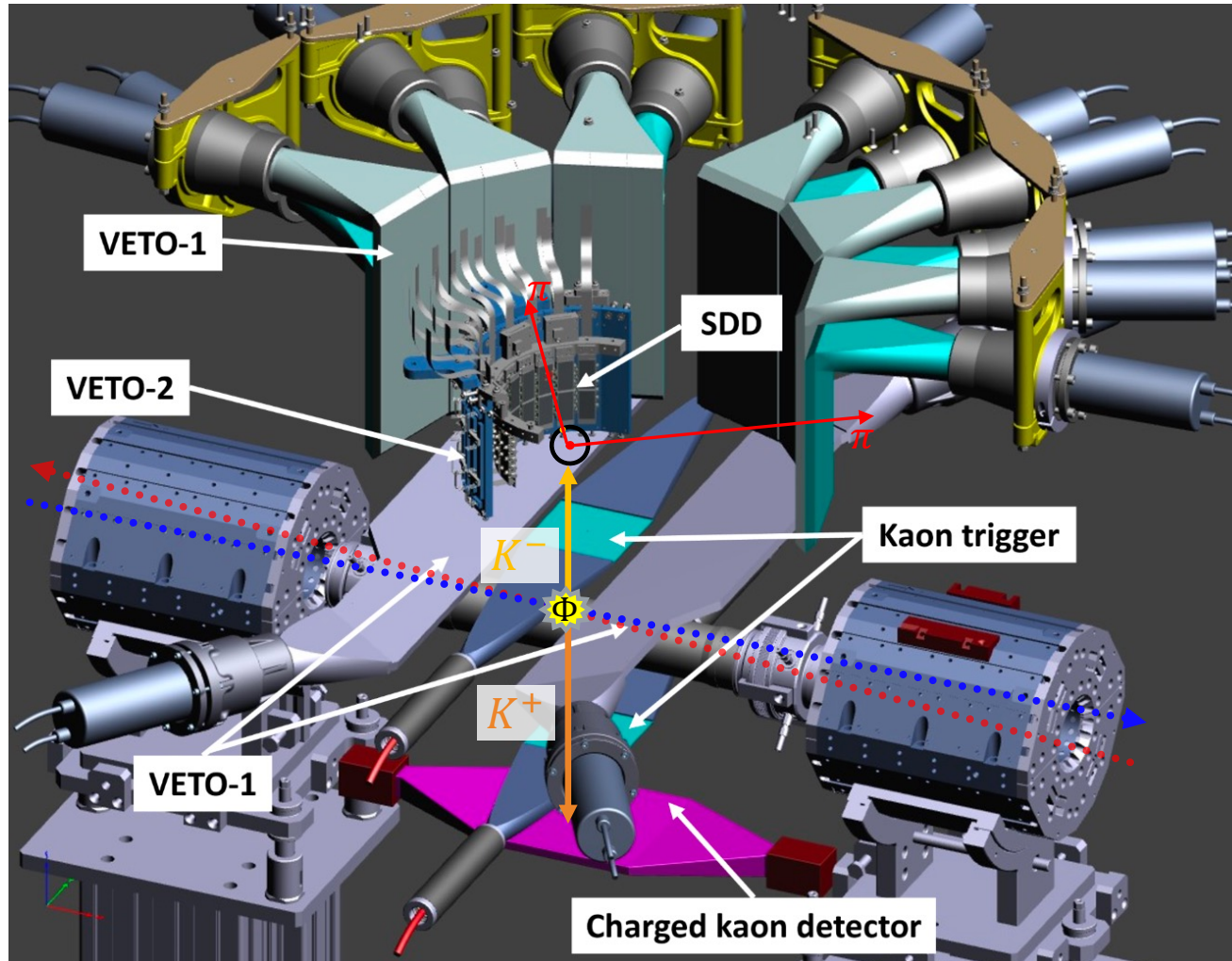


The combined use of Kaon Trigger and SDDs drift time reduces the background by a factor $\sim 10^4$



Kaonic deuterium analysis

The combined use of the veto systems further reduces the background by a factor of 3



Veto-1 (fiducial volume selection): to select the events occurring in the gas target, rejecting the X-ray background corresponding to K^- stopped in the solid elements of the setup.

Veto-2 (charged particle rejection): 48 plastic scintillator read by SiPMs to distinguish X-ray from charged particle by topological correlation with the SDD.

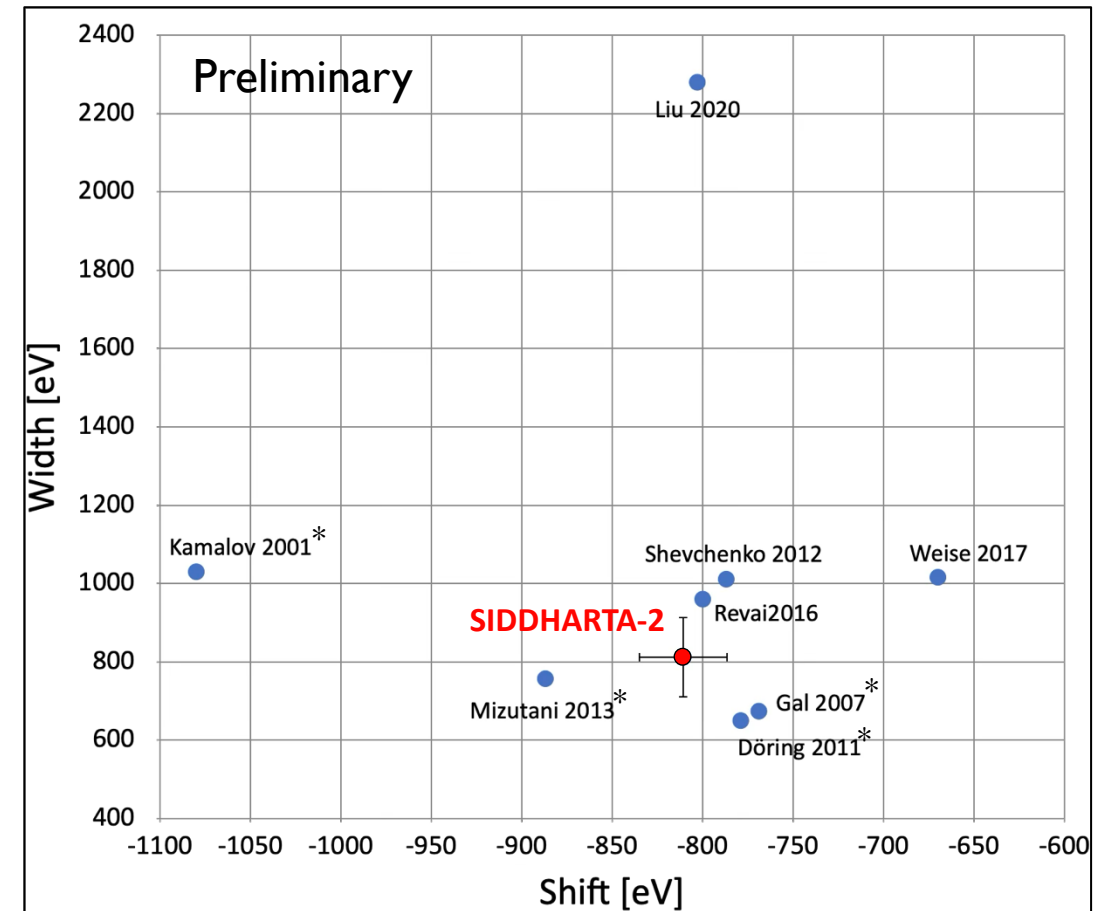
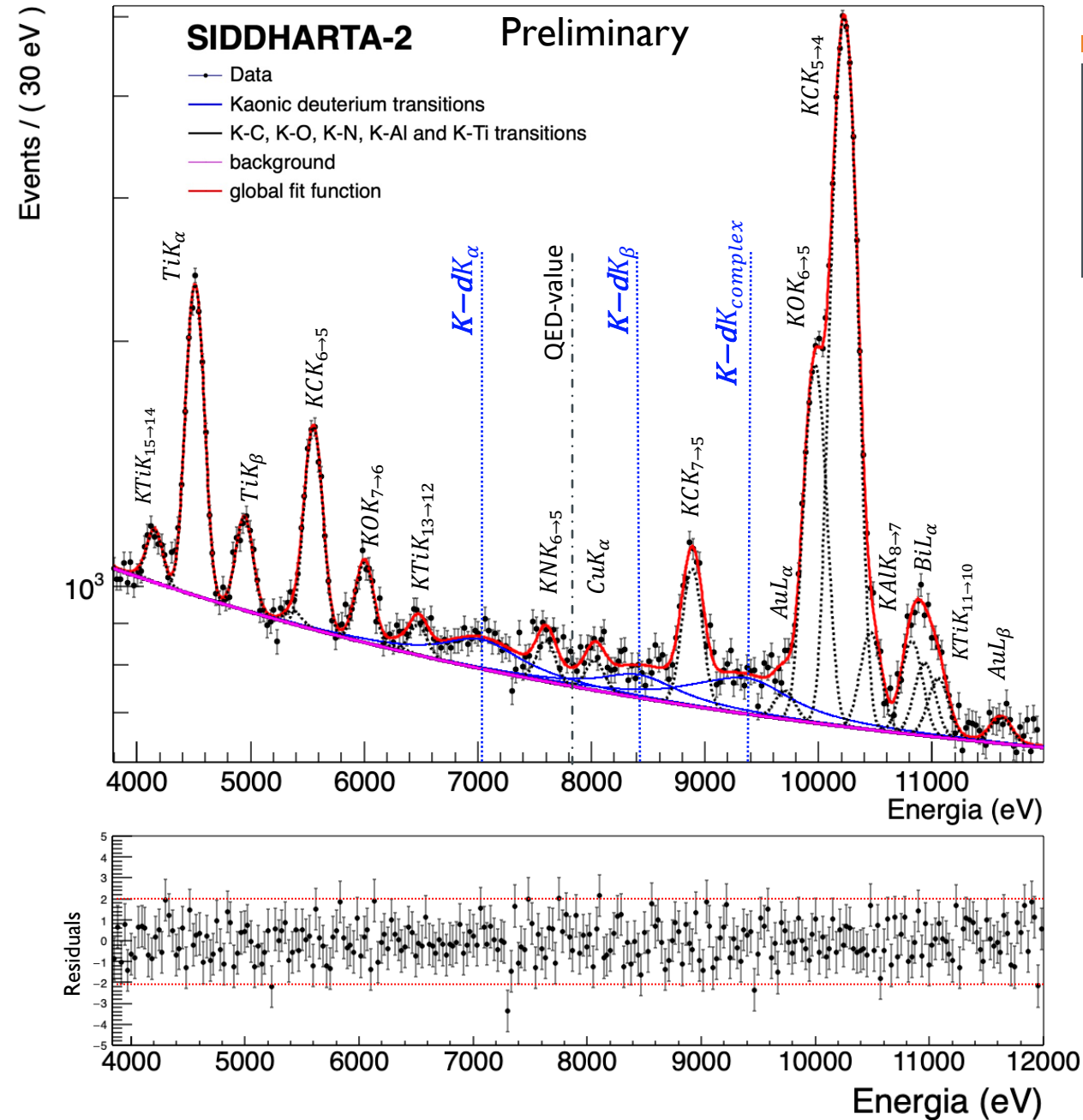
Charged kaon detector: distinguishes K^+ from K^- by exploiting their different interaction mechanisms with matter.

Kaonic Deuterium Results

$$\varepsilon_{1s} = -810.9 \pm 24.5(\text{stat}) \pm 2.1(\text{syst}) \text{ eV}$$

$$\Gamma_{1s} = 812 \pm 97(\text{stat}) \pm 33(\text{syst}) \text{ eV}$$

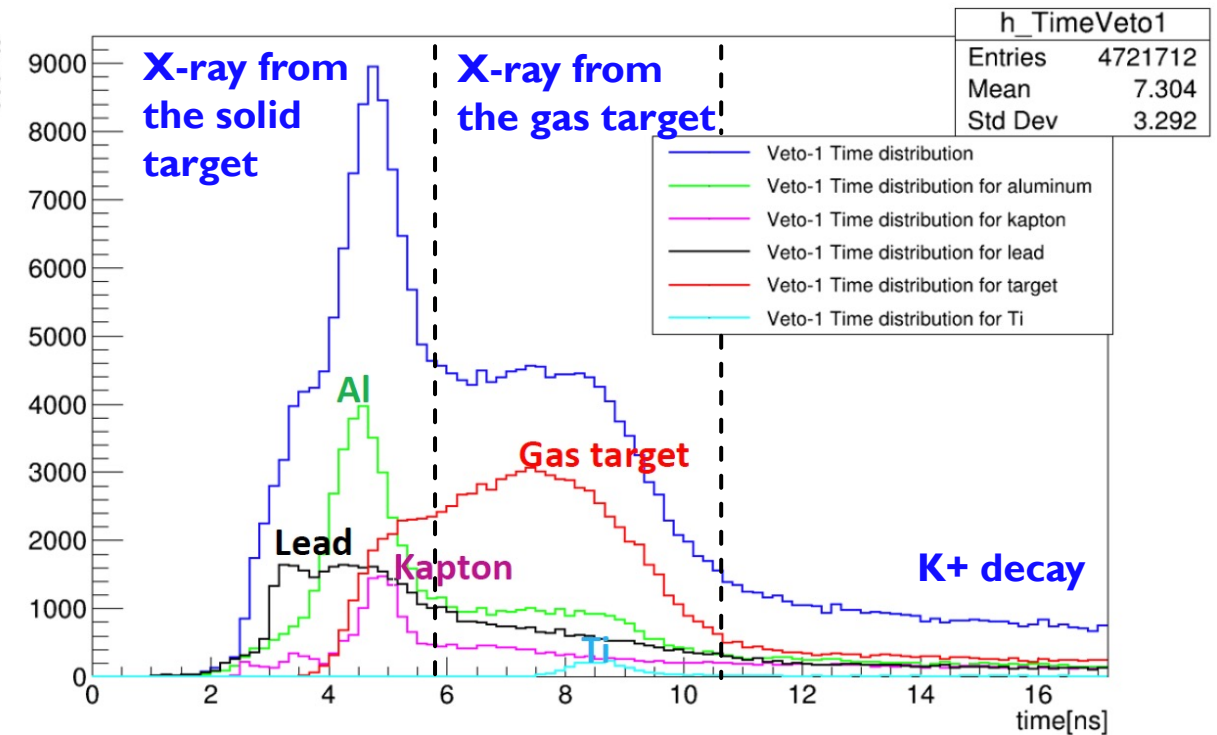
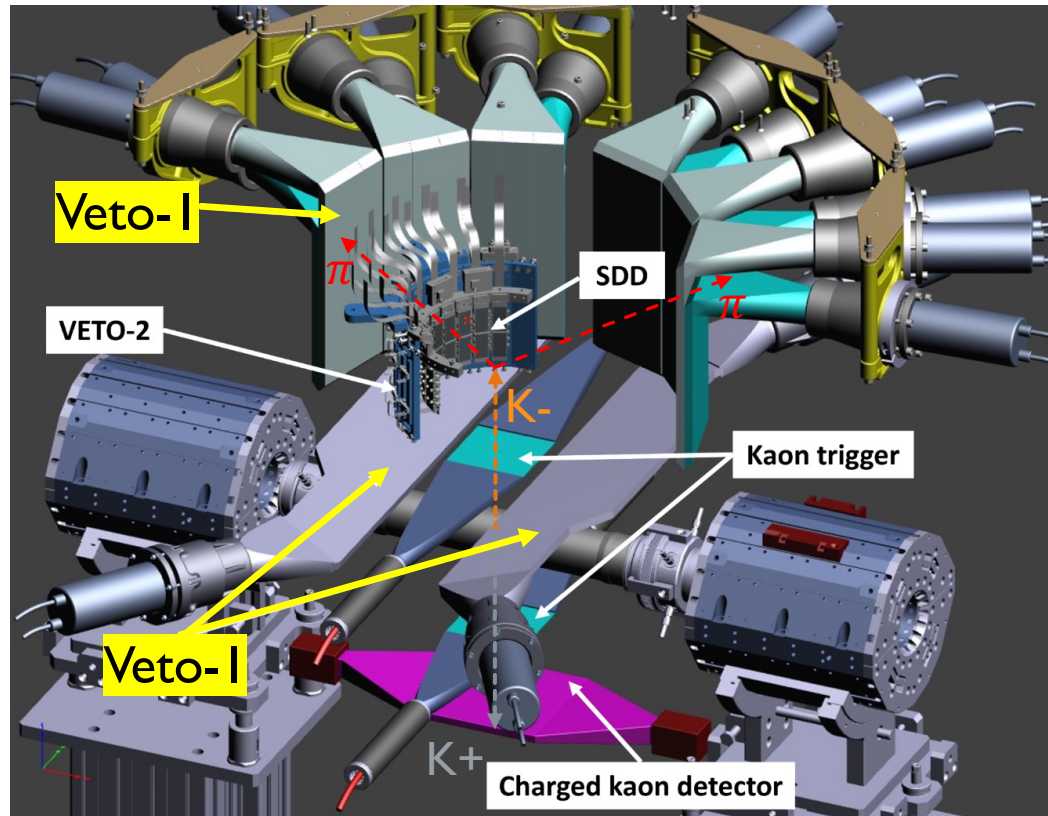
Targeted precision achieved!



* The energy shift and width are extracted from scattering length using the Deser-Trueman formula

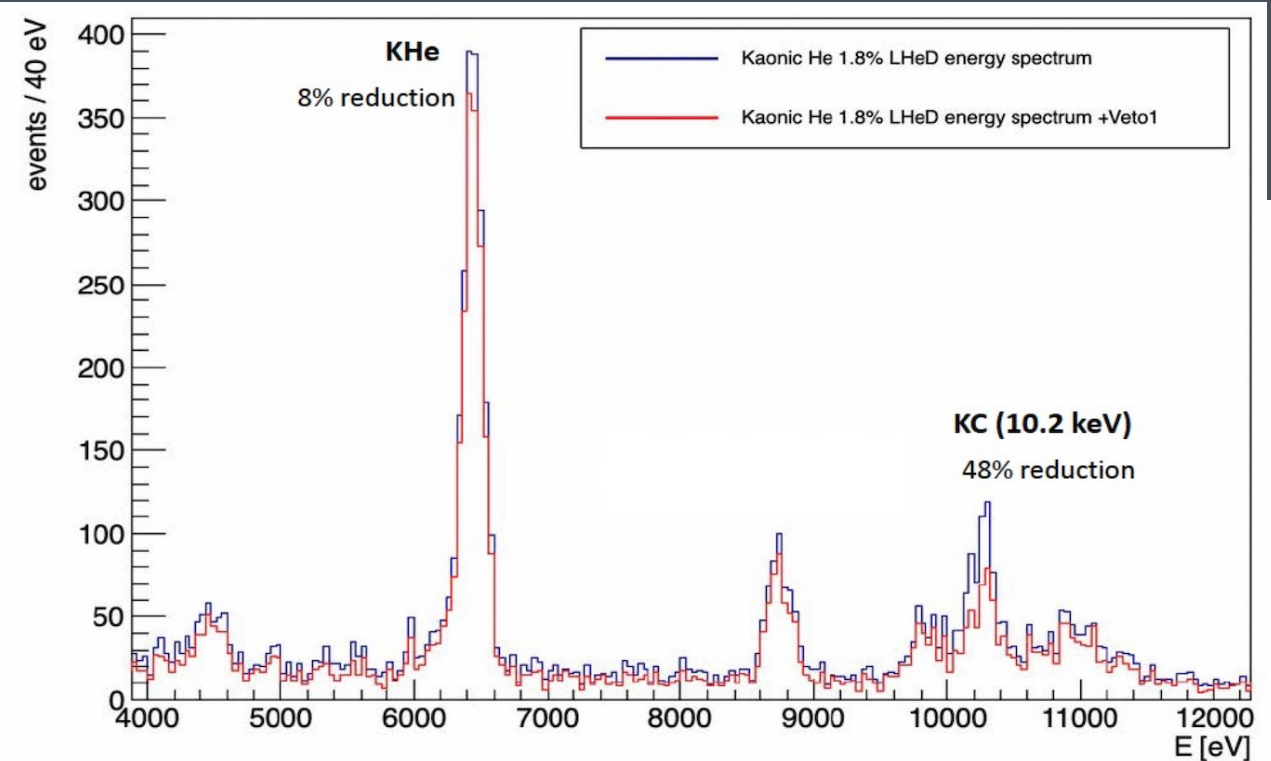
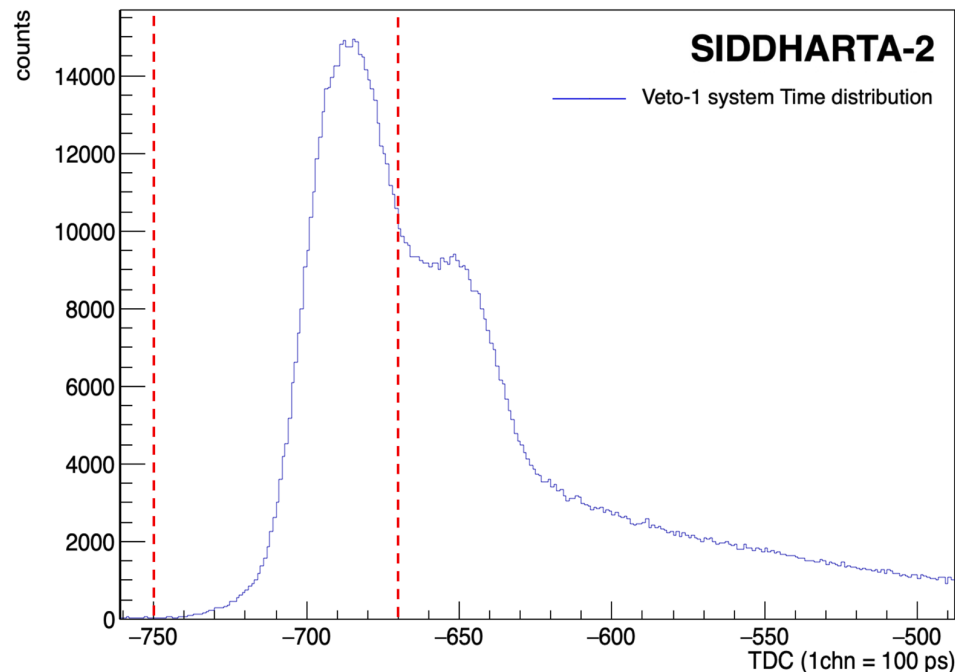
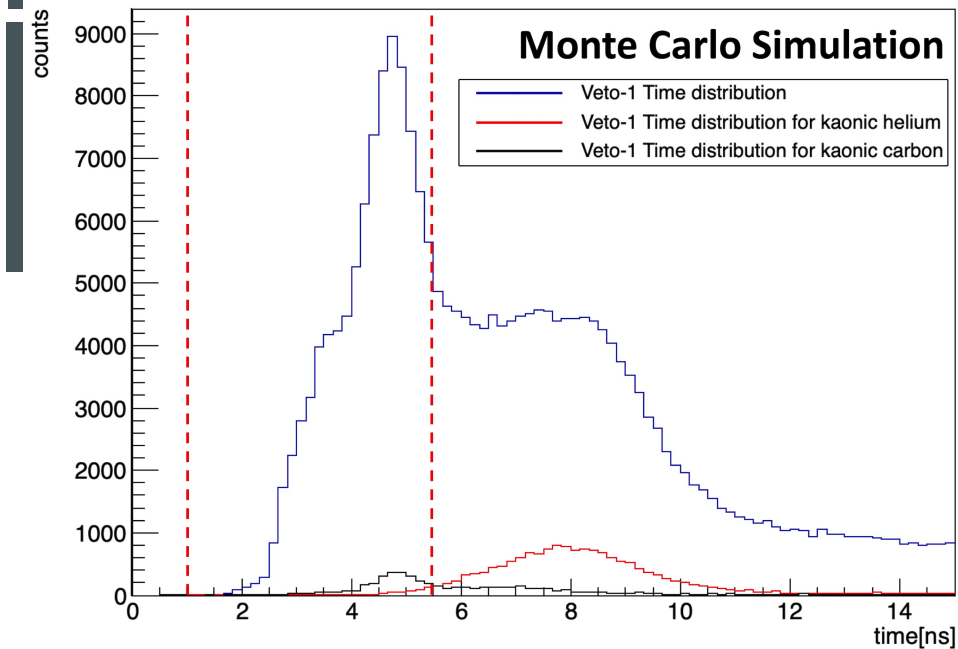
Kaonic deuterium analysis – Veto1 system

Veto-1 for hadronic background reduction: it measures the arrival time of charged particles emitted by the kaon-nucleus absorption to determine the origin of the signal — whether it comes from the deuterium gas target or from surrounding solid materials.



Veto-1: 14 plastic scintillators placed around and below the vacuum chamber

Veto-1 system optimization with kaonic He

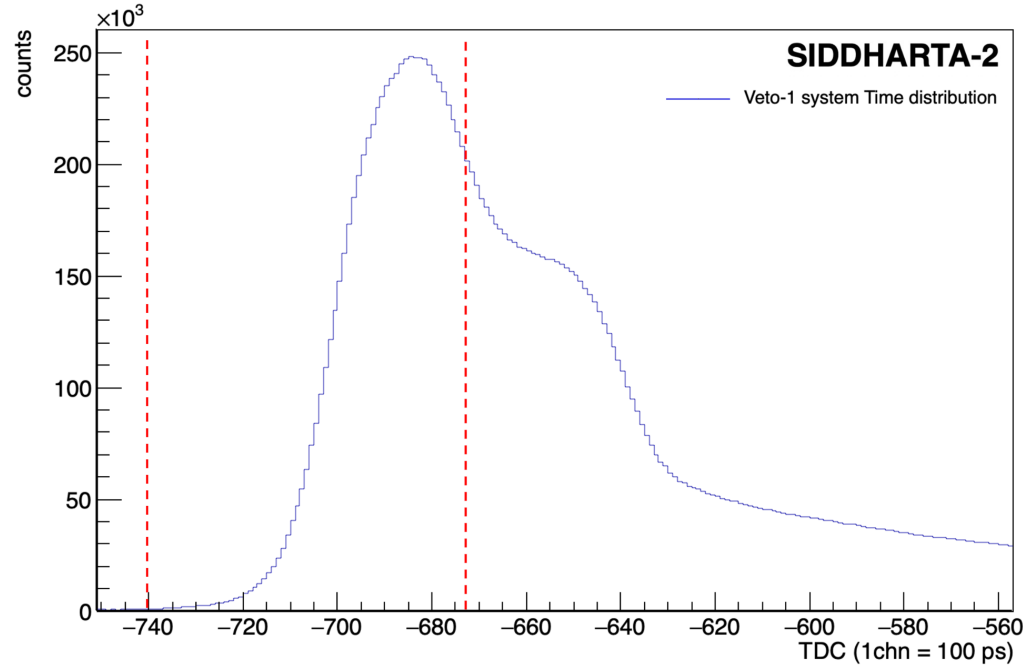


	SIDDHARTA-2 Veto-1 reduction factor	Monte Carlo Veto-1 reduction factor
K- ⁴ He L _α	(8 ± 1)%	4%
K-C _{5→4}	(48 ± 4)%	44%

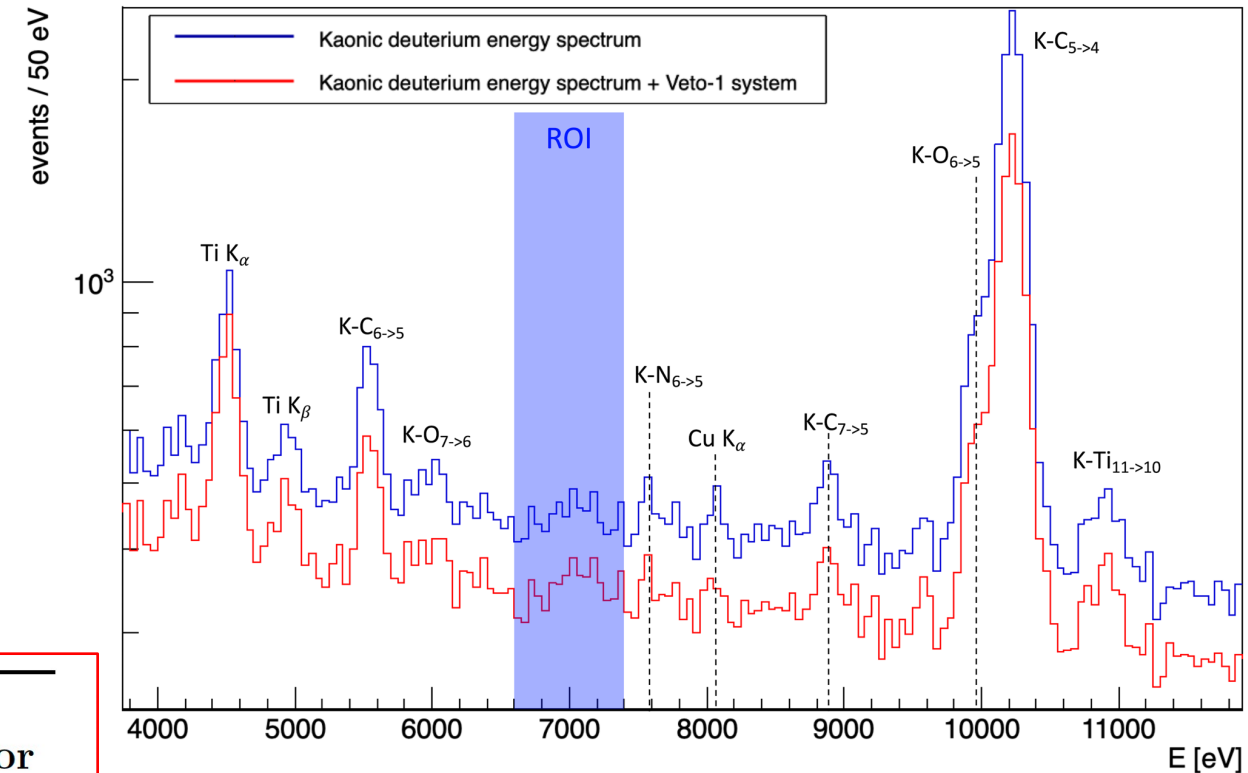
$$\text{reduction factor} = 1 - \frac{\text{\#events(with veto1)}}{\text{\#events}}$$

Kaonic Deuterium (run1): veto-1 system analysis

Veto-1 time distribution and time window used to reduce the background



SDDs X-ray energy spectra with and without the veto-1 (be aware logarithmic scale)

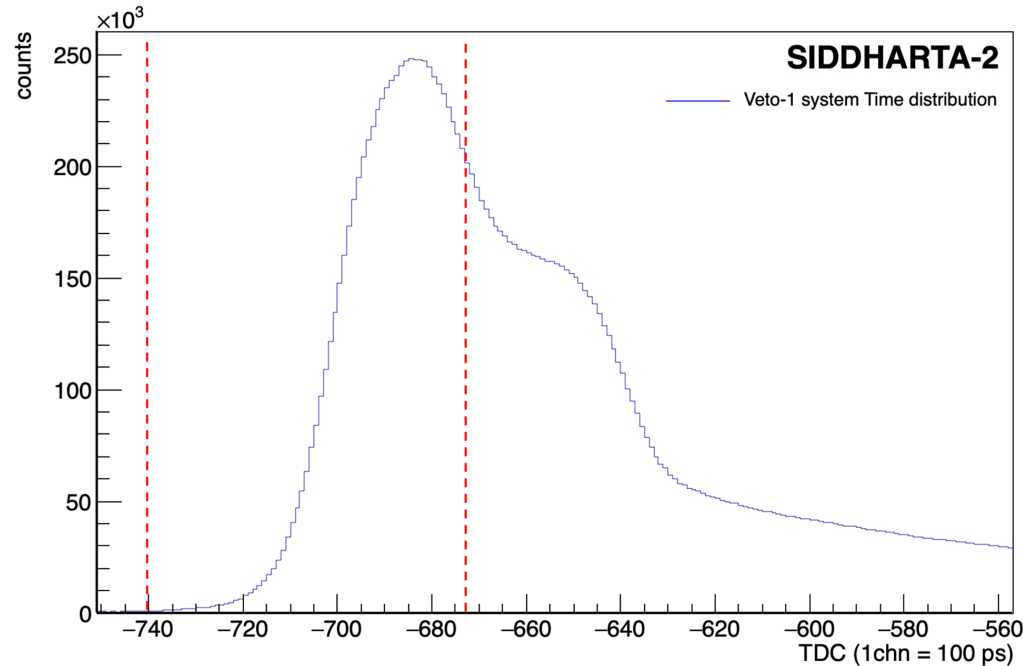


SIDDHARTA-2 Veto-1 reduction factor	Monte Carlo Veto-1 reduction factor
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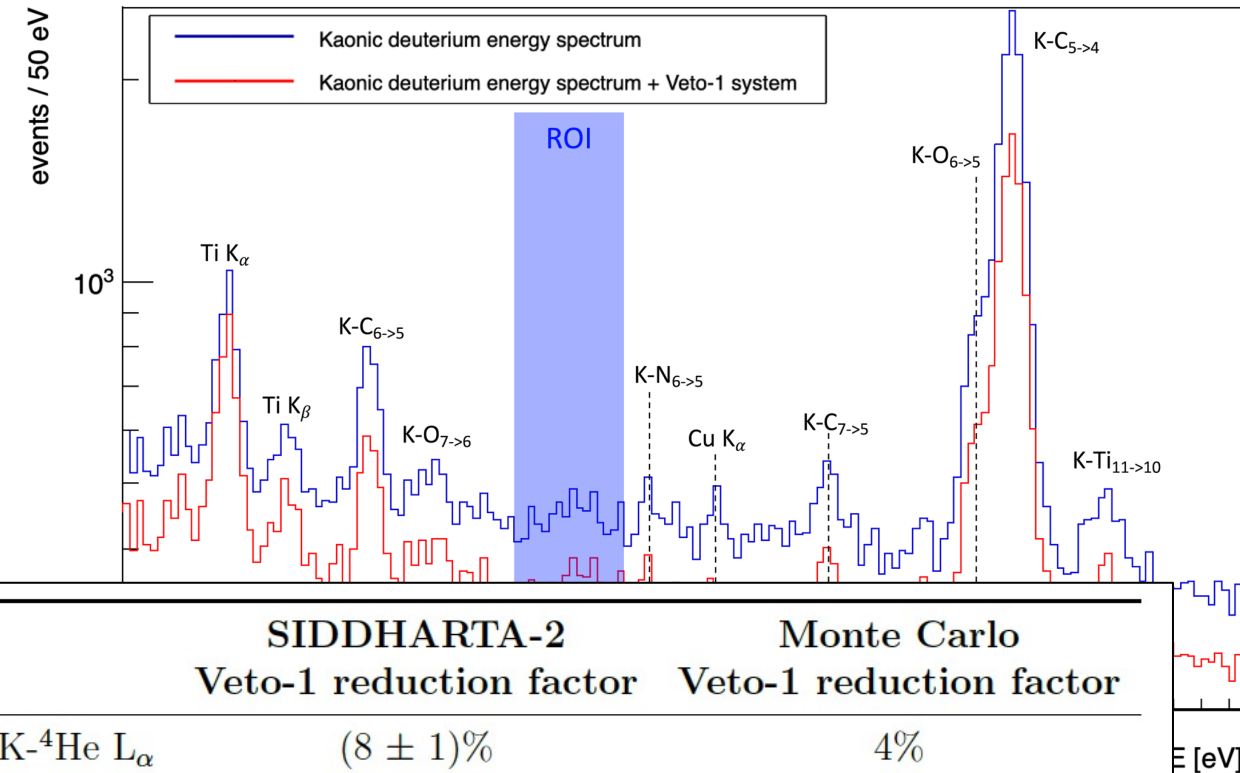
K-d K_α	$(11 \pm 3)\%$	4%
K-C ₅ →4	$(44 \pm 4)\%$	46%
K-C ₆ →5	$(39 \pm 5)\%$	45%
K-C ₇ →5	$(48 \pm 4)\%$	46%

Kaonic Deuterium: veto-1 system analysis

Veto-1 time distribution and time window used to reduce the background



SDDs X-ray energy spectra with and without the veto-1 (be aware logarithmic scale)



SIDDHARTA-2 Veto-1 reduction factor	Monte Carlo Veto-1 reduction factor
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K-d K_α	$(11 \pm 3)\%$	4%
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K-C ₇ → ₅	$(48 \pm 4)\%$	46%

	SIDDHARTA-2 Veto-1 reduction factor	Monte Carlo Veto-1 reduction factor
K- ⁴ He L_α	$(8 \pm 1)\%$	4%
K-C ₅ → ₄	$(48 \pm 4)\%$	44%

Low
Energy
QCD

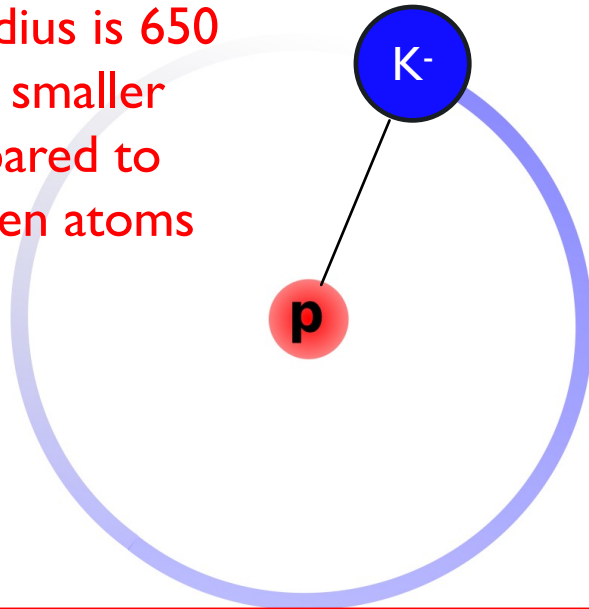
Strong
Field QED

Kaon
Mass

Probing strong field QED with kaonic atoms

Kaonic atoms are extremely compact, resulting in very strong electric fields that allow access to the strong-field regime of QED

Bohr radius is 650
times smaller
compared to
hydrogen atoms



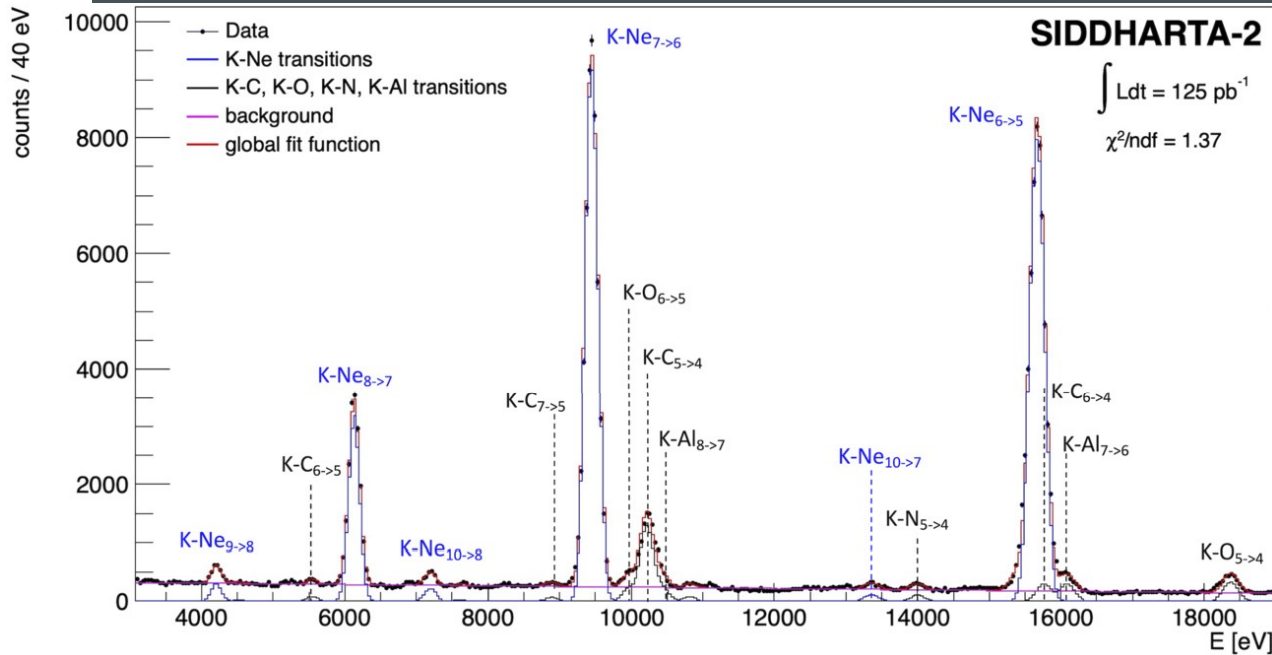
Particle	m [MeV/c ²]	B_{1s} [keV]	r_B [10^{-15} m]	Accessible interactions
$e p$	0.511	13.6×10^{-3}	0.53×10^5	Electro-weak
μp	105.7	2.53	279	Electro-weak
πp	139.6	3.24	216	Electro-weak + strong
$K p$	493.7	8.61	81	Electro-weak + strong
$\bar{p} p$	938.3	12.5	58	Electro-weak + strong

The electric field between the kaon and the nucleus is 430 000 times higher than that of hydrogen atom



Study of QED under strong field

Kaonic neon X-ray spectroscopy



Kaonic neon energy transitions and absolute yields at the density of 3.60 ± 0.18 g/l. The first error is statistical, the second systematic.

Transition	Energy [eV]	Yield
K-Ne (10 → 8)	$7191.21 \pm 4.91 \pm 2.00$	$0.010 \pm 0.001 \pm 0.001$
K-Ne (10 → 7)	$13352.20 \pm 10.07 \pm 3.00$	$0.004 \pm 0.002 \pm 0.001$
K-Ne (9 → 8)	$4206.35 \pm 3.75 \pm 2.20$	$0.137 \pm 0.012 \pm 0.010$
K-Ne (8 → 7)	$6130.86 \pm 0.71 \pm 1.50$	$0.228 \pm 0.004 \pm 0.011$
K-Ne (7 → 6)	$9450.08 \pm 0.41 \pm 1.50$	$0.277 \pm 0.002 \pm 0.014$
K-Ne (6 → 5)	$15673.30 \pm 0.52 \pm 9.00$	$0.308 \pm 0.003 \pm 0.015$



Letter

High precision X-ray spectroscopy of kaonic neon [☆]

F. Sgaramella^{a, b, *}, D. Sirghi^{b, a, c, **}, K. Toho^d, F. Clozza^a, L. Abbene^e, C. Amsler^f, F. Artibani^{a, g}, M. Bazzi^a, G. Borghi^{h, i}, D. Bosnar^j, M. Bragadireanu^c, A. Buttacavoli^e, M. Cargnelli^f, M. Carminati^{h, i}, A. Clozza^a, R. Del Grande^{k, a}, L. De Paolis^a, K. Dulski^{a, l, m}, L. Fabbietti^k, C. Fiorini^{h, i}, I. Friščić^j, C. Guaraldo^{a, l}, M. Iliescu^a, M. Iwasakiⁿ, A. Khreptak^{l, m}, S. Manti^a, J. Marton^f, P. Moskal^{l, m}, F. Napolitano^a, S. Niedźwiecki^{l, m}, H. Ohnishi^d, K. Piscicchia^{b, a}, F. Principato^e, A. Scordo^a, M. Silarski^l, F. Sirghi^{a, c}, M. Skurzok^{l, m}, A. Spallone^a, L.G. Toscano^{h, i}, M. Tüchler^f, O. Vazquez Doce^a, E. Widmann^f, J. Zmeskal^{f, l}, C. Curceanu^a

The kaonic neon measurement demonstrates the feasibility of precision studies in the field of QED with kaonic atoms (BSQED)

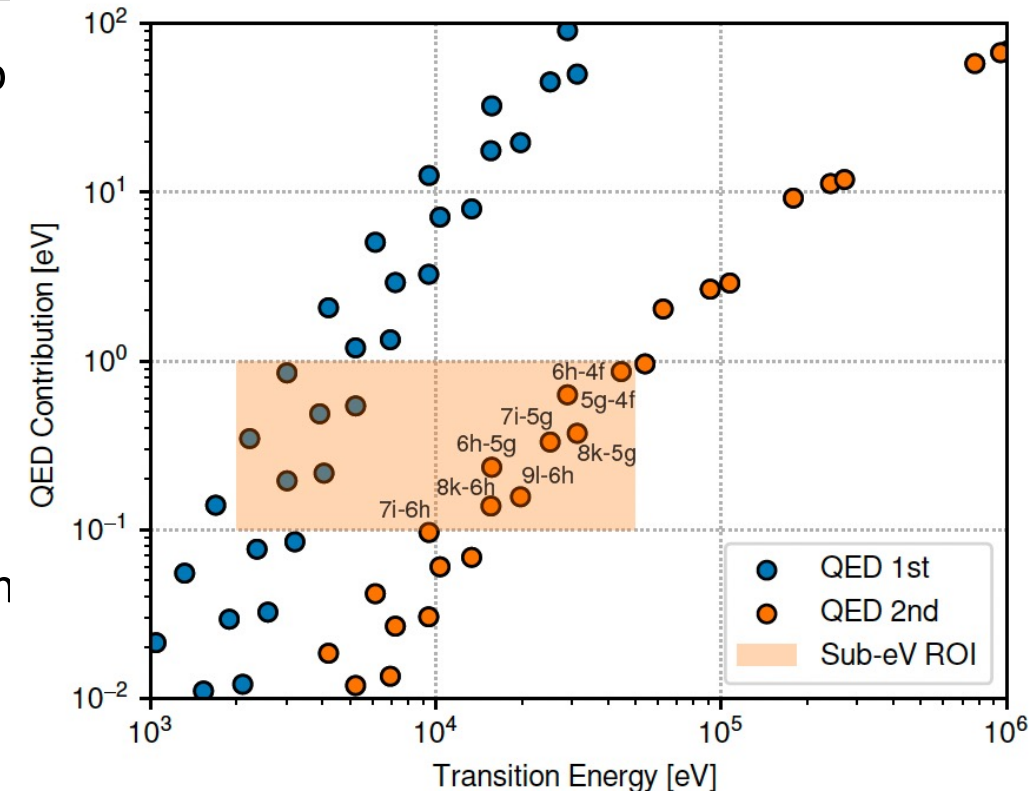
Kaonic atoms offer high sensitivity to BSQED corrections thanks to their large reduced mass compared to muonic atoms. Unlike antiprotonic and muonic atoms, they lack hyperfine structure effects, simplifying measurement and analysis.

Bound State QED with kaonic atoms

We (S. Manti) **Run, Test, and Contribute** to the **MCDF** code for Ab Initio Atomic Calculations, including Relativistic and QED effects, in collaboration with **Prof. Paul Indelicato (Paris–CNRS)**.

Transition	$E_{if}^{(\text{exp.})}$	$\delta E_{if}^{(\text{stat.})}$	$\delta E_{if}^{(\text{syst.})}$	$E_{if}^{(\text{calc.})}$	$E_{if}^{(\text{QED})}$	$E_{if}^{(\text{QED1})}$	$E_{if}^{(\text{QED2})}$	$\Delta E_{if}^{(\text{isot.})}$	$\Delta E_{if}^{(\text{screen.})}$	$\Delta E_{if}^{(\text{PDG})}$
9l-8k	4206.97	3.43	2.00	4201.45	2.09	2.07	0.02	9.90	-0.38	0.11
8k-7i	6130.57	0.65	1.50	6130.31	5.09	5.05	0.04	14.45	-0.27	0.16
7i-6h	9450.23	0.37	1.50	9450.28	12.66	12.56	0.10	22.28	-0.18	0.24
6h-5g ^a	15673.30	0.52	9.00	15685.39	32.75	32.51	0.24	37.01	-0.11	0.40

- **MCDF Calculations** include transition energies, QED effects, Electro Screening, and Isotopic Shifts ($\text{K}^{20}\text{Ne}, \text{K}^{22}\text{Ne}$)
- **Sub-eV Statistical Uncertainties** match the scale of **Second-Order QED corrections** (e.g., Vacuum Polarization).
- **K^{20}Ne 7i $\rightarrow$ 6h transition:**
 - **Exp:** 9450.23 ± 0.37 (stat) ± 1.50 (syst) eV
 - **Theory:** 9450.28 (with the QED contribution to the transition energy is 12.66 eV)
 $\rightarrow$ **Excellent agreement between Theory and Experiment**
- **Kaonic Neon** offers a clean system to test **QED calculations**

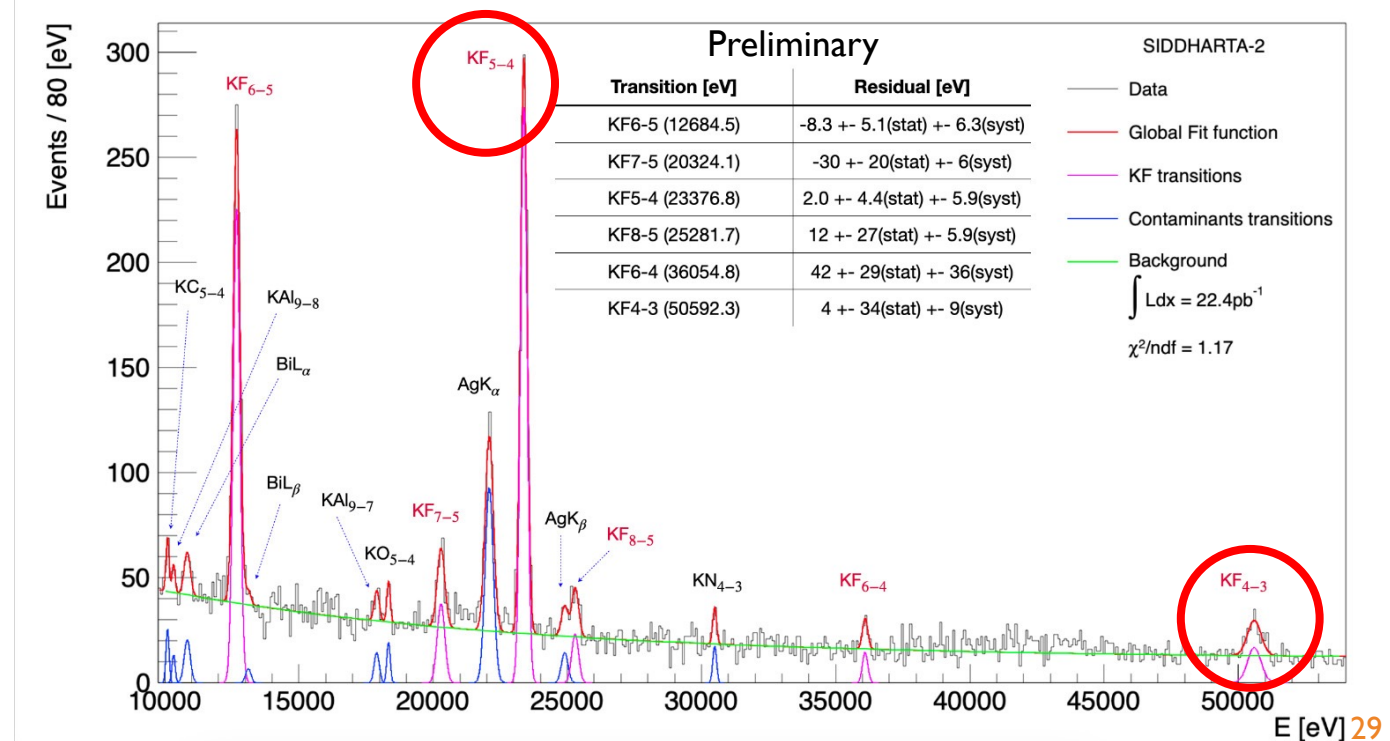
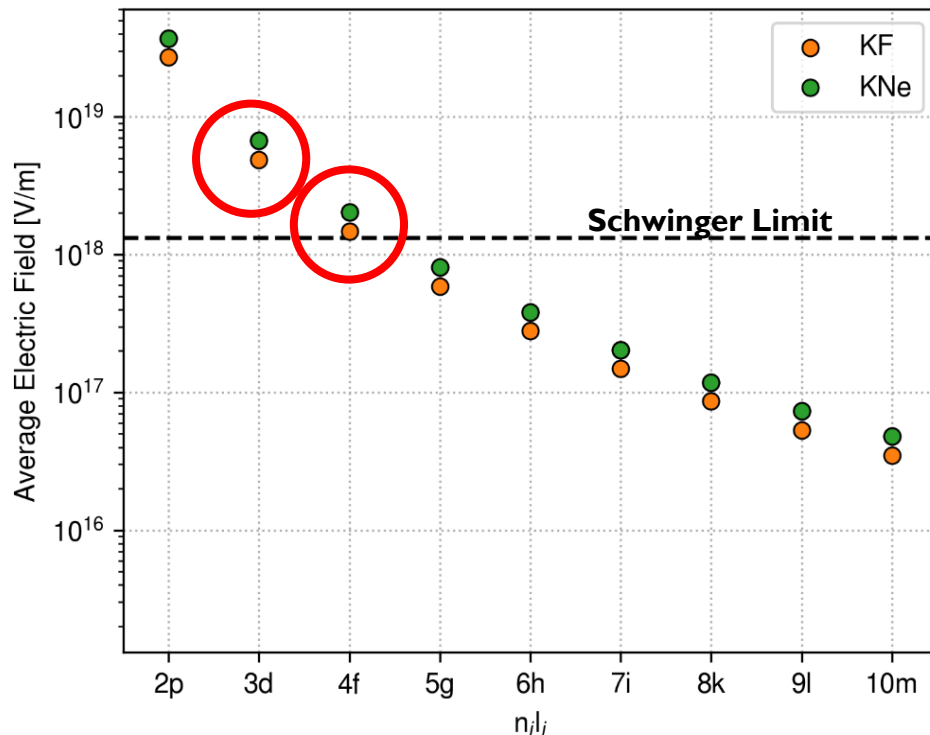


Kaonic Atoms as Probes of Strong Fields QED: Schwinger Limit for KF

- Exotic atoms (like KNe and KF) enable experimental access to Strong Electric Fields [Paul *et al.*, PRL 126, 173001 (2021)]
- The Schwinger Limit for spontaneous e^+e^- Pair Creation is:

$$E_c = \frac{m_e^2 c^3}{q_e \hbar} \approx 1.32 \times 10^{18} \text{ V/m} \quad \langle E \rangle_{nl} = \int d^3r |\psi_{nl}(\mathbf{r})|^2 E(\mathbf{r})$$

For the transition $\text{KF } 5 \rightarrow 4$, and $4 \rightarrow 3$, the Average Electric Field in kaonic orbitals approaches the Schwinger limit and the BSQED contribution are enhanced by the Nonlinear Electric-Field dependence of Vacuum Polarization



Low
Energy
QCD

Strong
Field QED

Kaon
Mass

The charged kaon mass puzzle

Charged kaon mass (K^+ , K^-)

$$493.677 \pm 0.013 \text{ MeV}$$

P.a. Zyla et al. (Particle Data Group), Prog.Theor. Exp. Phys. 2020, 083C01 (2020)

The charged kaon mass puzzle

Charged kaon mass (K^+ , K^-)

~~$493.677 \pm 0.013 \text{ MeV}$~~

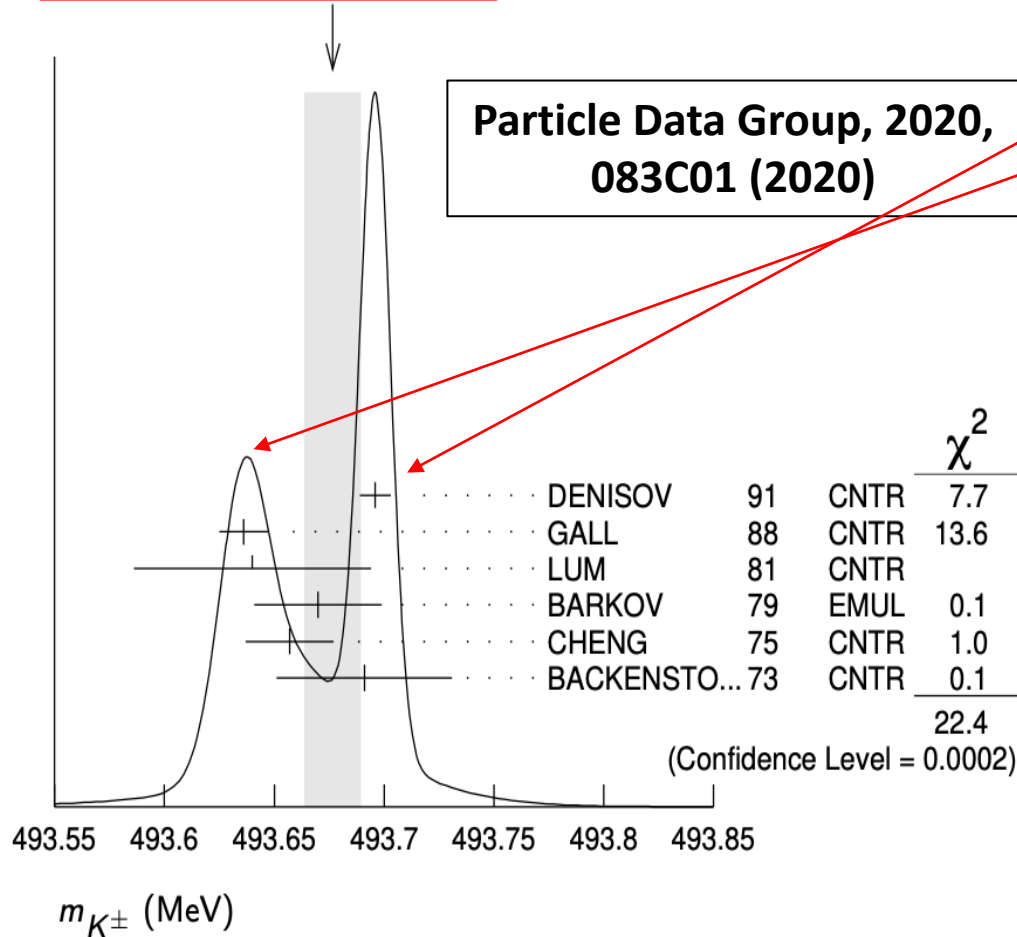
P.a. Zyla et al. (Particle Data Group), Prog.Theor. Exp. Phys. 2020, 083C01 (2020)

The charged kaon mass puzzle

60 keV discrepancy between the two most accurate measurement
Large uncertainty $\rightarrow$ 26 p.p.m, compared to charged pion (1.6 p.p.m)

WEIGHTED AVERAGE
 493.677 ± 0.013 (Error scaled by 2.4)

Particle Data Group, 2020,
083C01 (2020)



VALUE (MeV)	DOCUMENT ID	TECN	CHG	COMMENT	
493.677±0.016 OUR FIT	Error includes scale factor of 2.8.				
493.677±0.013 OUR AVERAGE	Error includes scale factor of 2.4. See the ideogram below.				
493.696±0.007	¹ DENISOV	91	CNTR	—	Kaonic atoms
493.636±0.011	² GALL	88	CNTR	—	Kaonic atoms
493.640±0.054	LUM	81	CNTR	—	Kaonic atoms
493.670±0.029	BARKOV	79	EMUL	±	e ⁺ e [−] → K ⁺ K [−]
493.657±0.020	² CHENG	75	CNTR	—	Kaonic atoms
493.691±0.040	BACKENSTO...	73	CNTR	—	Kaonic atoms

- The uncertainty on the charged kaon mass leads to an **error of 50 keV (σ) on the D^0 mass**
- **Large uncertainty on the charmonium spectrum**, in particular on precise values of charm-anticharm meson thresholds
- A particular case is that of $D^0 \bar{D}^{*0}$ which lies within the measured width of the best-known candidate for a **hadron-hadron molecule, the X(3872)**, an improved K-mass measurement would lead to a better interpretation of the X(3872), and of its radius.

Kaon Mass Determination: Application to KNe

Kaonic neon is the ideal system to precisely determine the kaon mass:

- **fully ionized** :no electron screening as in kaonic lead
- **low-density gas target**: negligible electron refill compared to solid targets like Pb and C

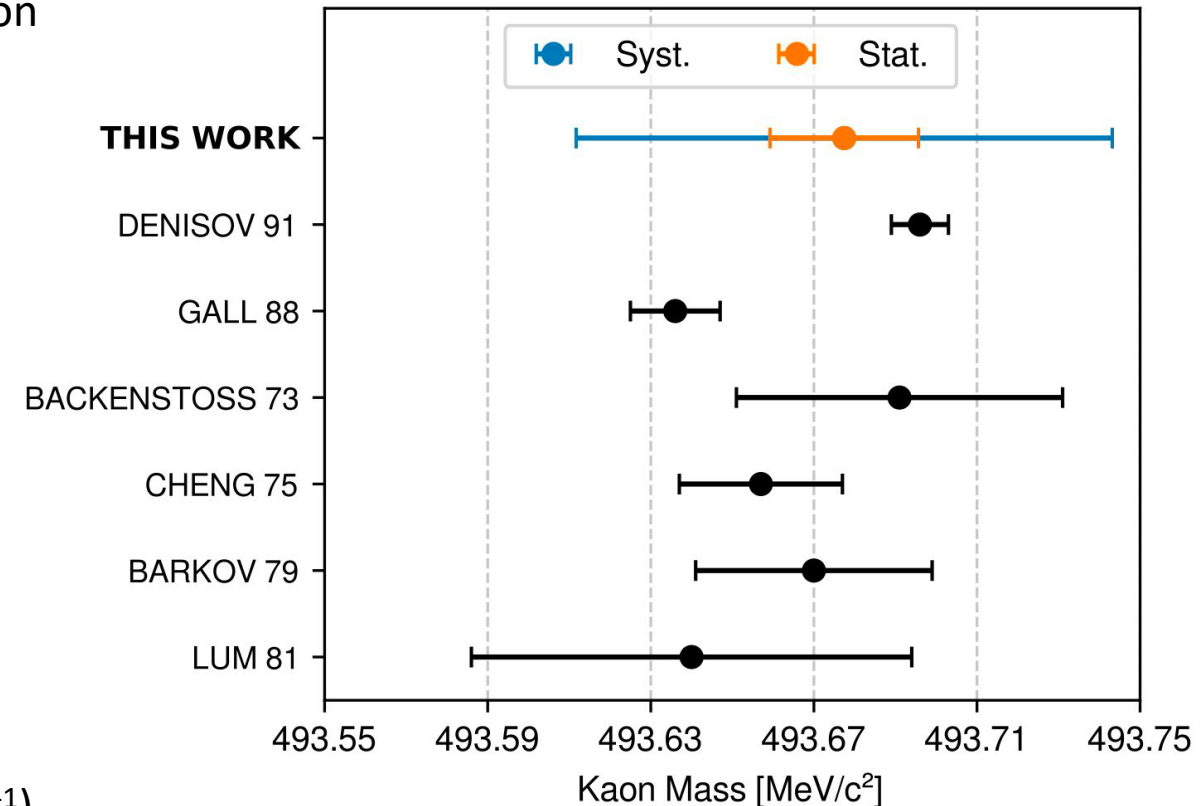
- **Theoretical Uncertainties** stem from **Electron Screening** and depend on occupation of a **K-shell** electron in the 1s during the transition

- $M_{K^-} = 493.677 + 0.018 \text{ (stat)} + 0.066 \text{ (syst)} \text{ MeV}$

Transition	M_{K^-} [MeV]	$\delta M_{K^-}^{\text{stat.}}$ [keV]	$\delta M_{K^-}^{\text{syst.}}$ [keV]
7i-6h	493.674	19	78
8k-7i	493.699	52	121
7i-6h + 8k-7i	493.677	18	66

- Transitions with **Sub-eV Statistical Precision** already provide a **Statistical Uncertainty of 18 keV** on the Kaon Mass

- **A KNe Dedicated Optimized Run can reach < 10 keV via:**
 - × 2 improvement by Doubling Statistics ($150 \text{ pb}^{-1} \rightarrow 300 \text{ pb}^{-1}$)
 - × 2 improvement by Optimized Gas Target Design and calibration system (lower statistical and systematic $\rightarrow 0$)



I.1 - High precision kaonic neon measurement

To extract the charged kaon mass with a precision of about 5 keV

BSQED and Physics beyond Schwinger limit

I.2 - Light kaonic atoms (LHKA)

- solid target **Li, Be, B**
- integration of Imm SDD

**Low energy QCD –
kaon multi-nucleon interaction**

EXKALIBUR

C. Curceanu et al., Front.in Phys. 11 (2023) 1240250

EXtensive

Kaonic

Atoms research: from

Lithium and

Beryllium to

URanium

Intermediate kaonic atoms (IMKA)

In parallel we plan dedicated runs for kaonic atoms (O, Al, S) with **CdZnTe detectors**
- 200 -300 pb⁻¹ of integrated luminosity/target

- **Feasibility** with minimal modifications/addings of the already existent SIDDHARTA-2
 - New calibration system (0.2 eV accuracy)
 - New Imm thick SDDs
 - New and improved CZT setup
- **Impact**: i.e. the maximal scientific outcome: KN and KNN interaction at threshold; Nuclear density distributions; Kaonic atoms cascade models; kaon mass; BSQED; Physics beyond Schwinger limit

Conclusion

- First measurement of kaonic deuterium 1s level shift and width

$$\varepsilon_{1s} = -810.9 \pm 24.5(\text{stat}) \pm 2.1(\text{syst}) \text{ eV}$$

$$\Gamma_{1s} = 812 \pm 97(\text{stat}) \pm 33(\text{syst}) \text{ eV}$$

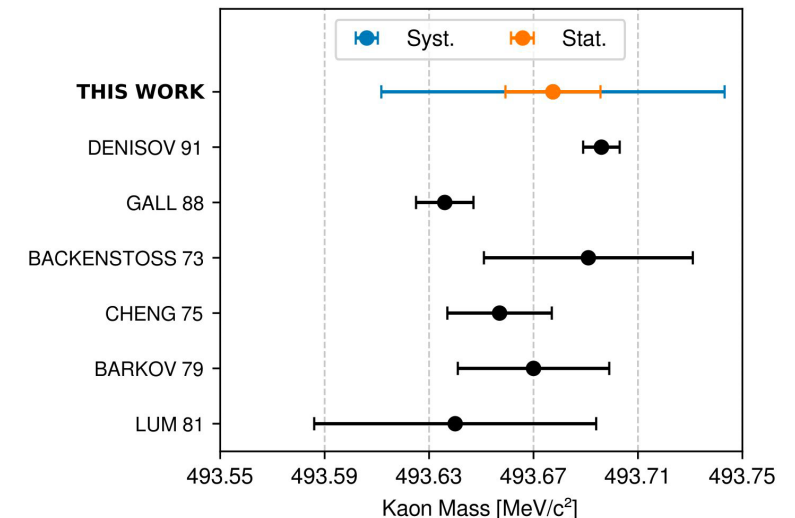
Preliminary

- First proof-of-concept for QED studies with kaonic atoms from an experimental and theoretical perspective.

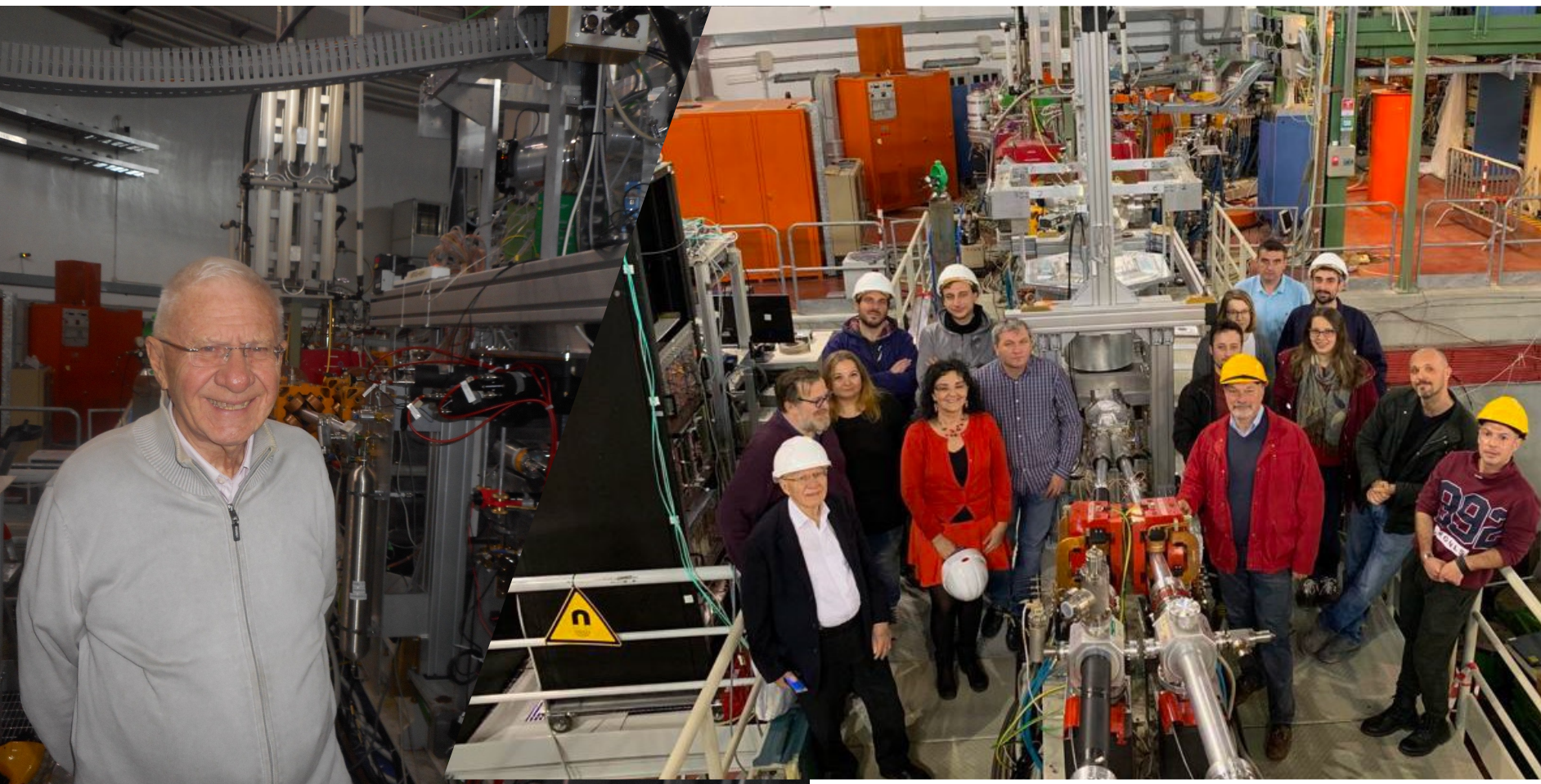
Transition	$E_{if}^{(\text{exp.})}$	$\delta E_{if}^{(\text{stat.})}$	$\delta E_{if}^{(\text{syst.})}$	$E_{if}^{(\text{calc.})}$	$E_{if}^{(\text{QED})}$	$E_{if}^{(\text{QED1})}$	$E_{if}^{(\text{QED2})}$	$\Delta E_{if}^{(\text{isot.})}$	$\Delta E_{if}^{(\text{screen.})}$	$\Delta E_{if}^{(\text{PDG})}$
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6h-5g ^a	15673.30	0.52	9.00	15685.39	32.75	32.51	0.24	37.01	-0.11	0.40

- Kaonic neon X-ray spectroscopy ideal candidate to solve the kaon mass discrepancy

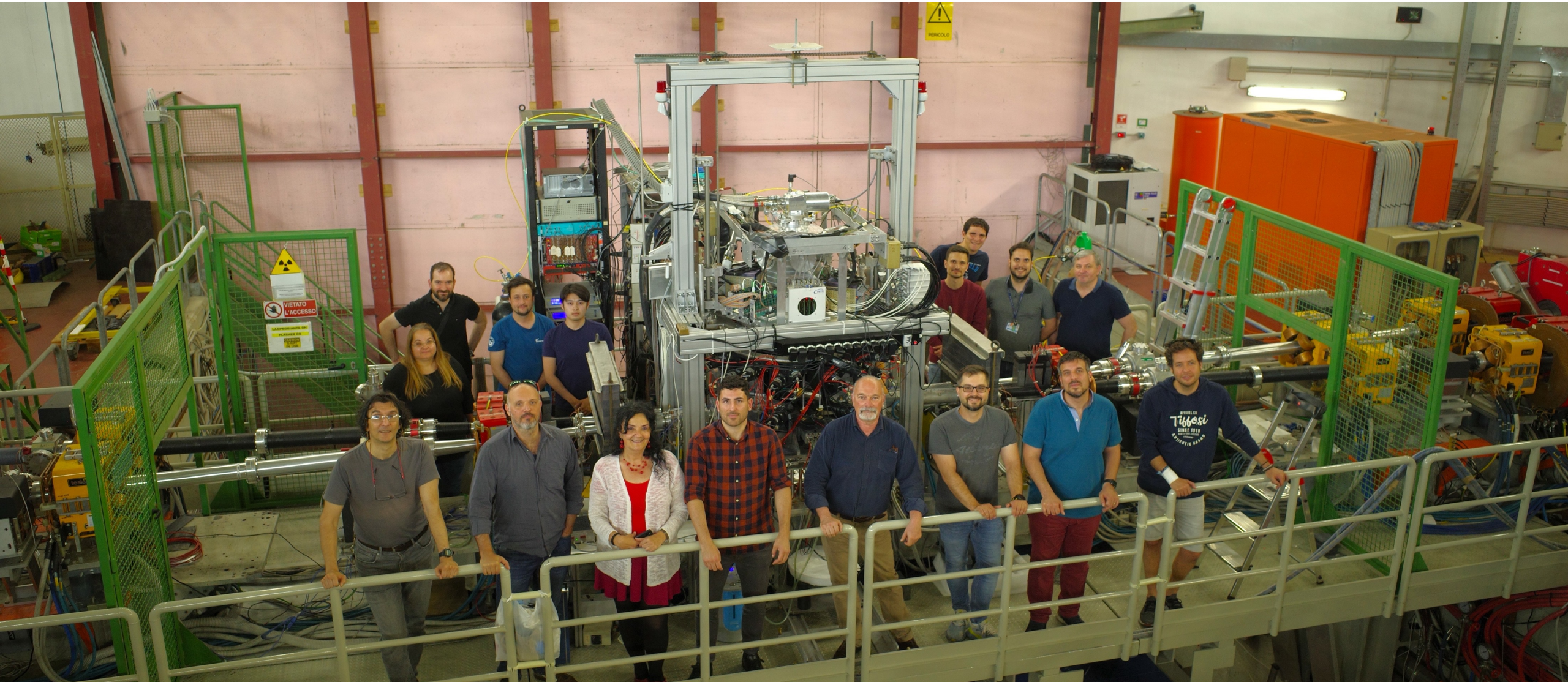
- **EXKALIBUR**: new X-ray detectors (SDDs – CZT - HPGe) have been developed/tested to perform kaonic atoms measurements along the periodic table providing new experimental data to probe the kaon-nucleus interaction – kaon mass - BSQED



We dedicate these results to our dear friends and colleagues
Prof Carlo Guaraldo and Dr. Johann Zmeskal
who passed away in 2024. You are very much missed!



THANK YOU





SPARE

EXPLORING THE STRANGENESS FRONTIERS: KAONIC ATOMS X-RAY SPECTROSCOPY AT DAFNE WITH SIDDHARTA-2



HYP2025
TOKYO

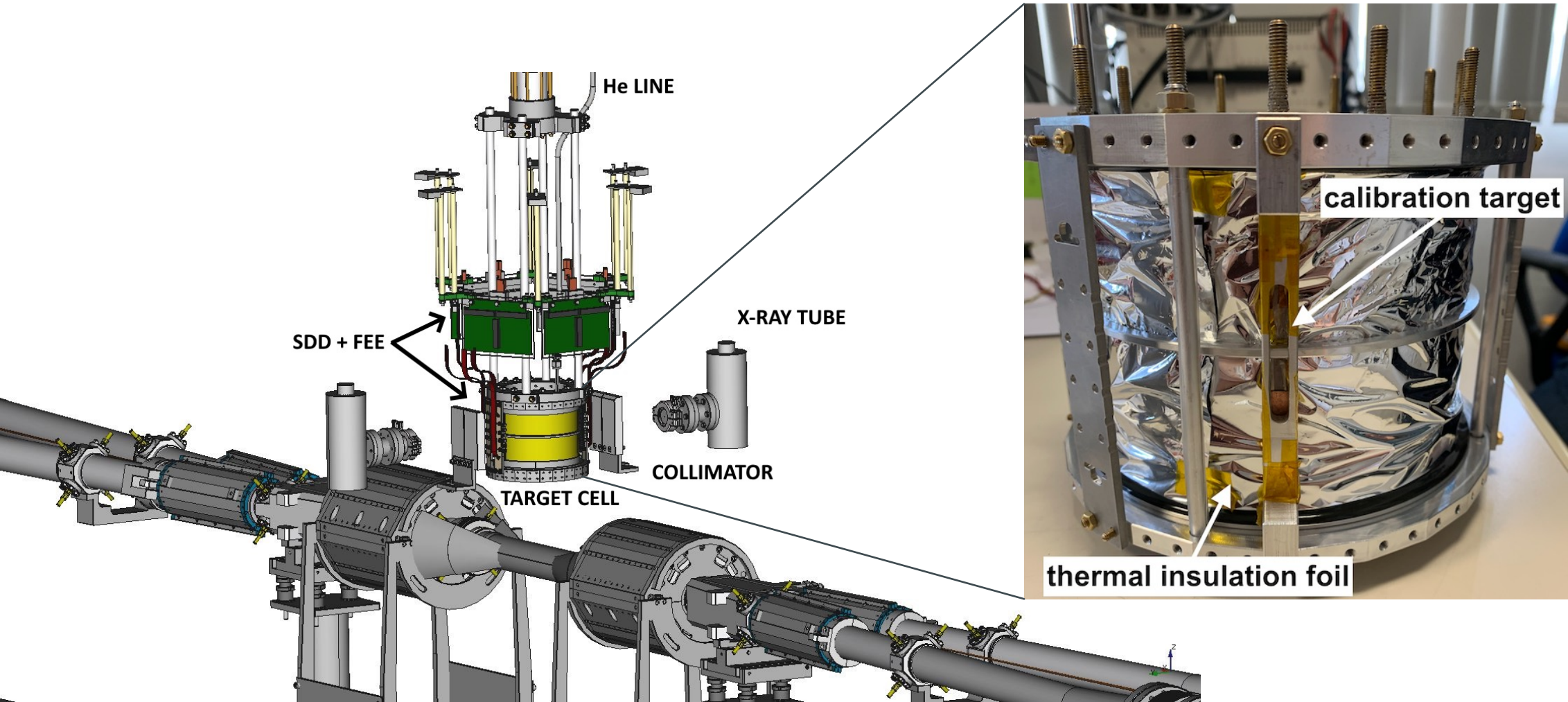
Francesco Sgaramella
on behalf of the SIDDHARTA-2 collaboration



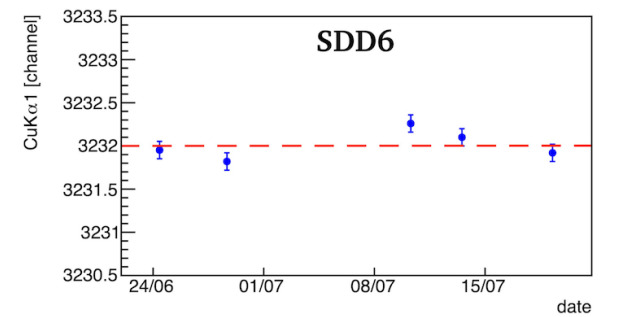
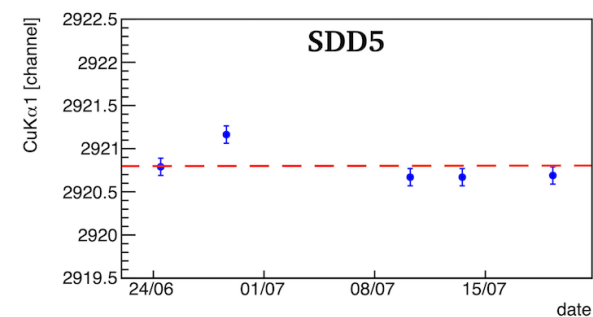
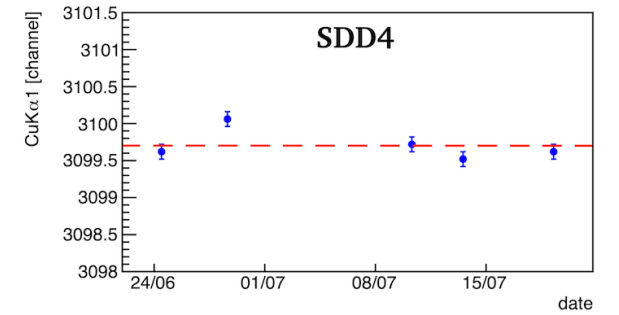
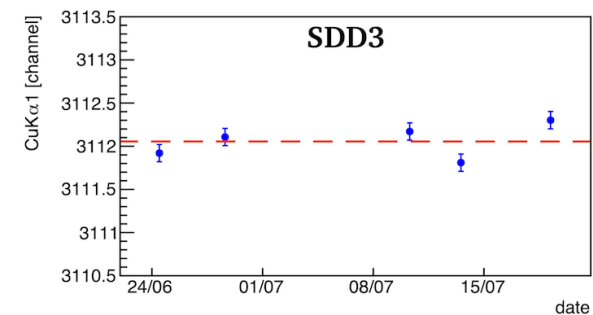
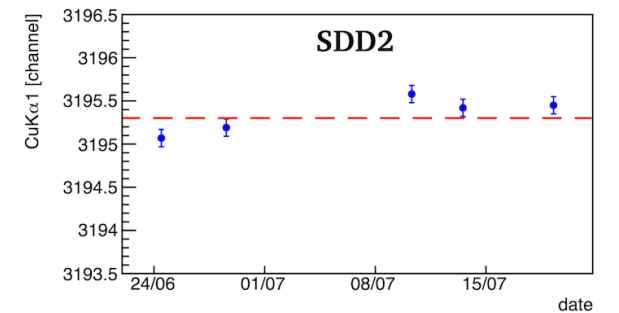
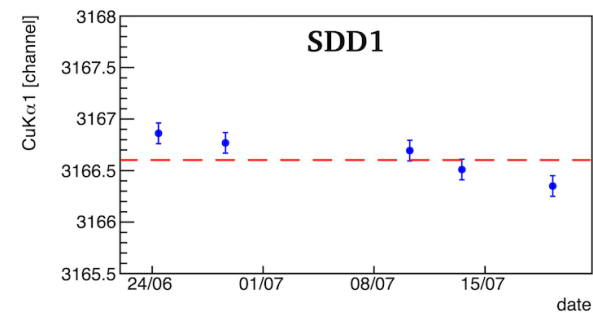
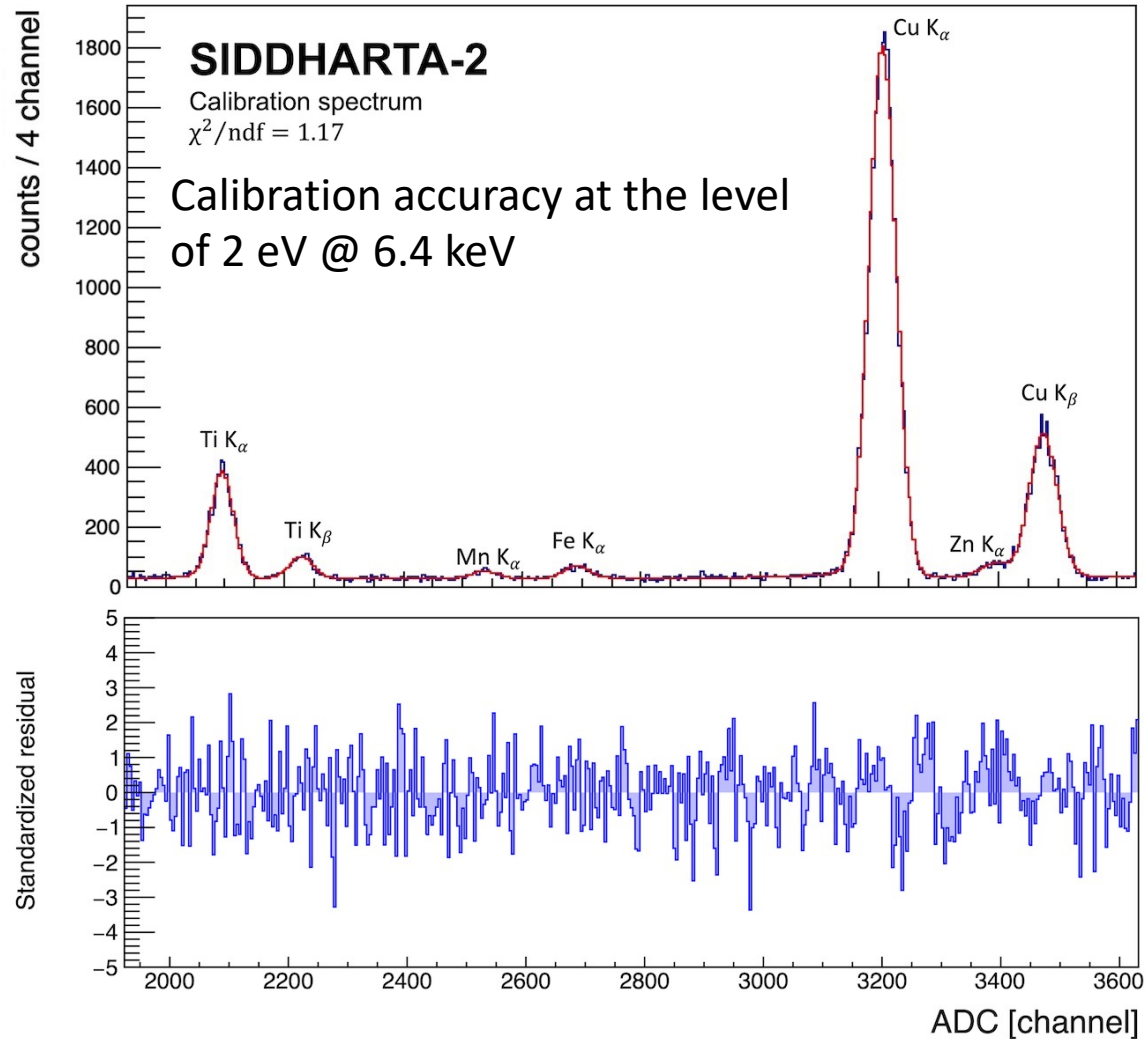
Istituto Nazionale di Fisica Nucleare
LABORATORI NAZIONALI DI FRASCATI



SDDs Calibration Procedure in DAΦNE



SDDs Calibration Procedure in DAΦNE



SIDDHARTINO - The kaonic ${}^4\text{He}$ $3d \rightarrow 2p$ measurement

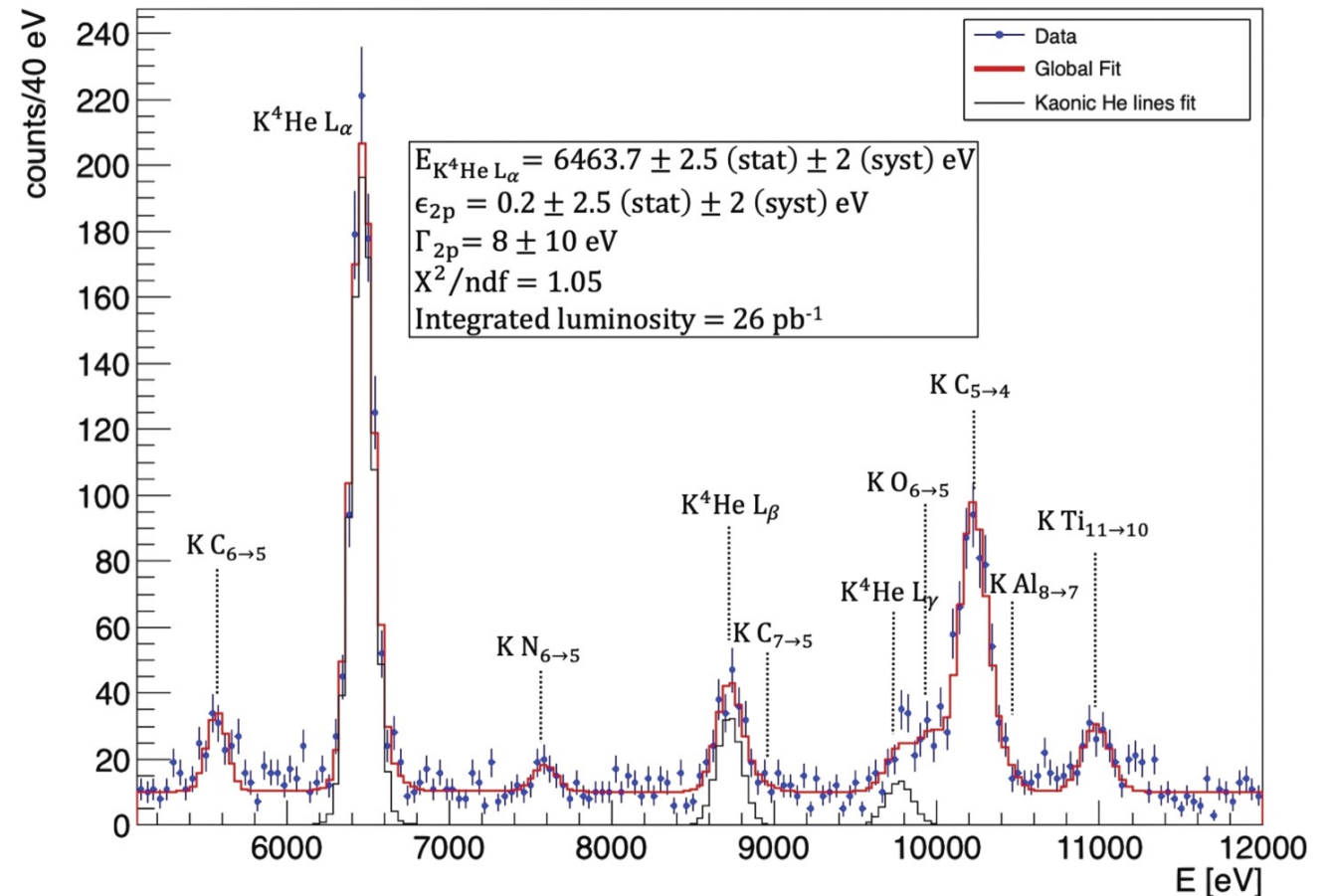
Characterization of the SIDDHARTA-2 apparatus and optimization of DAΦNE background through the kaonic helium measurement

Kaonic Helium puzzle:

SIDDHARTA and KEK results confirmed, improving accuracy

SIDDHARTINO:

reduced version of the SIDDHARTA-2 apparatus
It was used to optimize the DAFNE background and characterize the SDDs



SIDDHARTINO

The kaonic ^4He $3d \rightarrow 2p$ (L_α) measurement

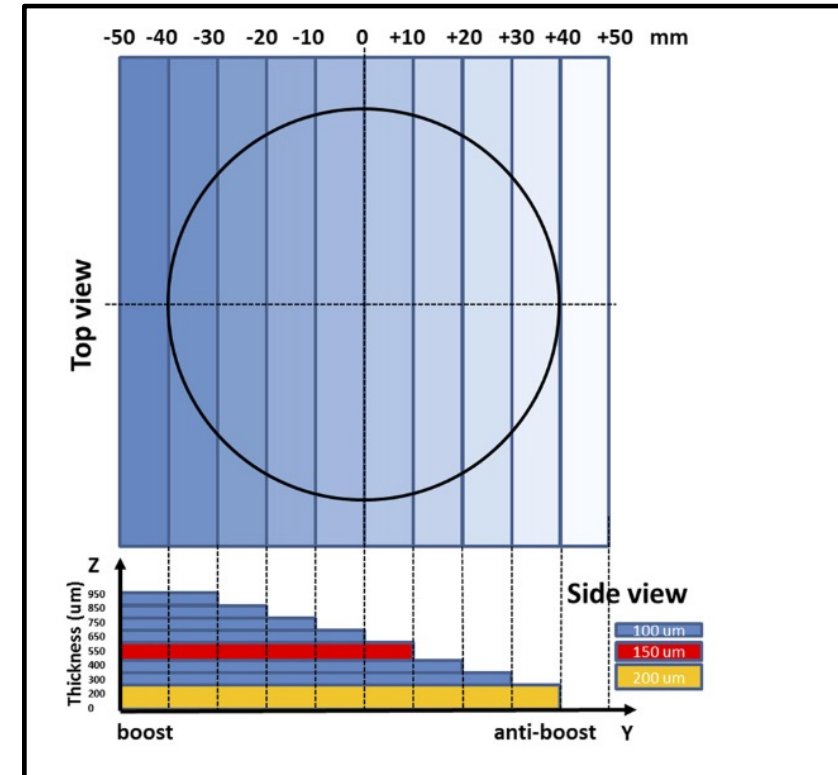
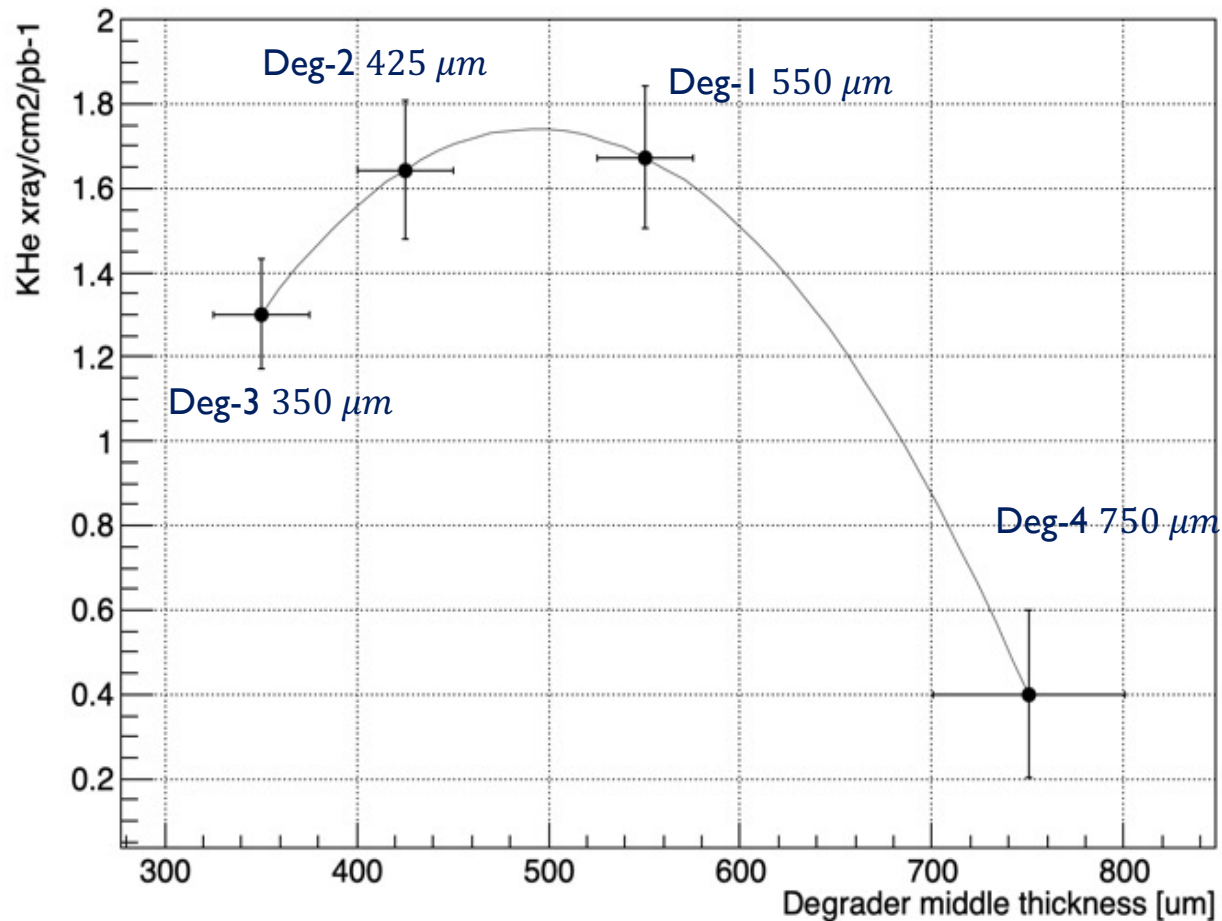


Figure 6. Nearest to optimal configuration of the size of the entrance window of the vacuum chamber side of the DAΦNE ring, corresponding to the ring has eight steps to compensate for the boost effect part of the figure.

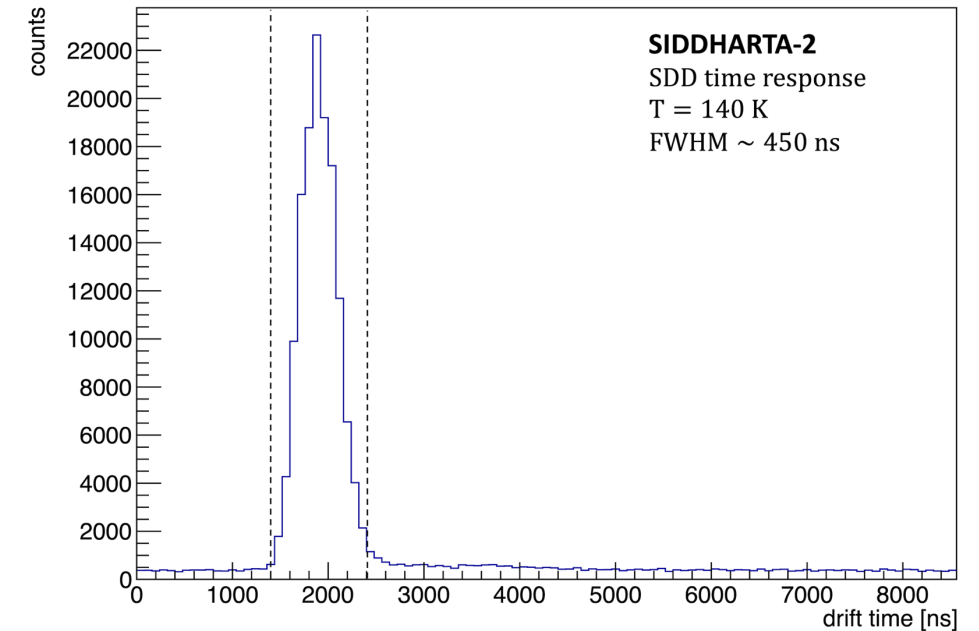
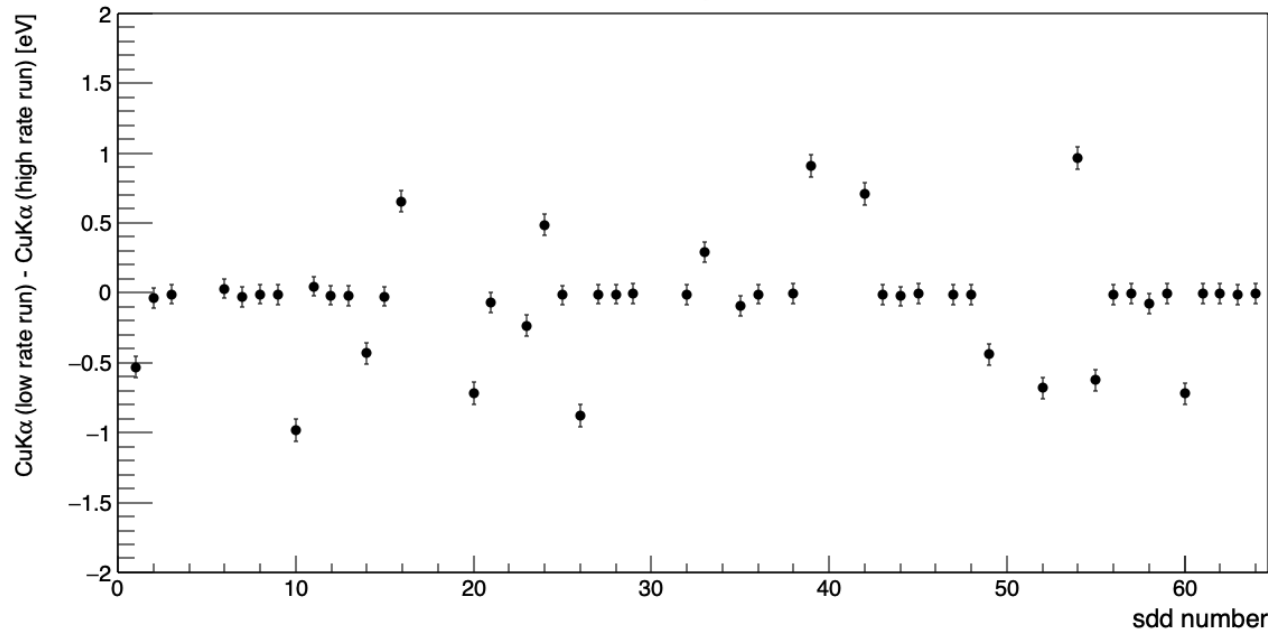
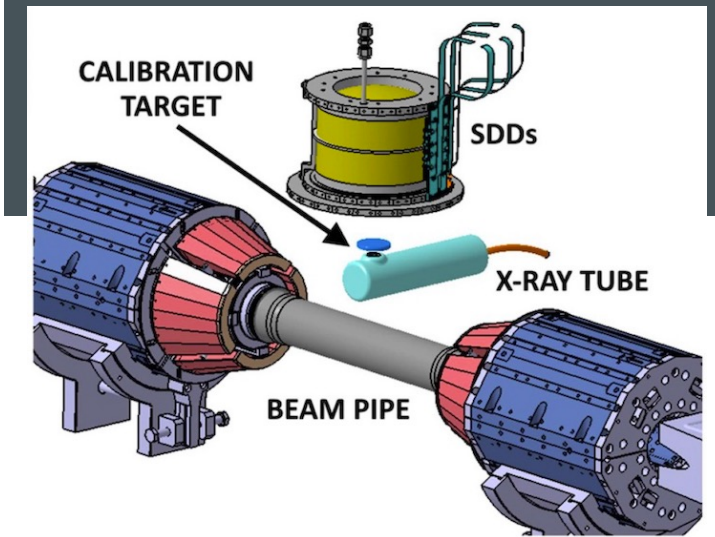
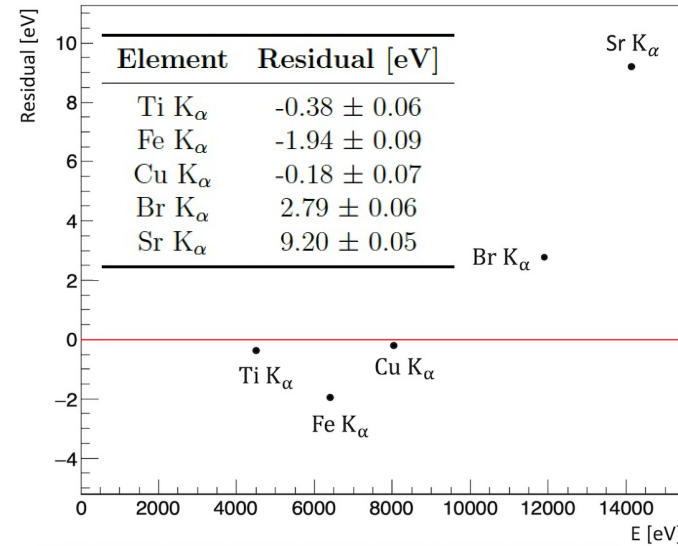
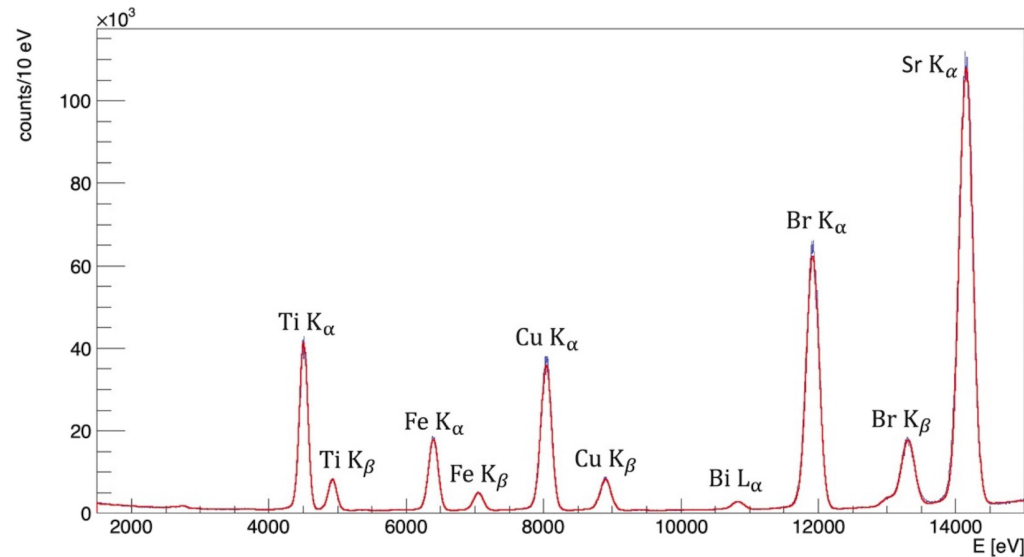
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IOP Publishing

Journal of Physics G: Nuclear and Particle Physics
J. Phys. G: Nucl. Part. Phys. 49 (2022) 055106 (14pp)

<https://doi.org/10.1088/1361-6471/ac5dac>

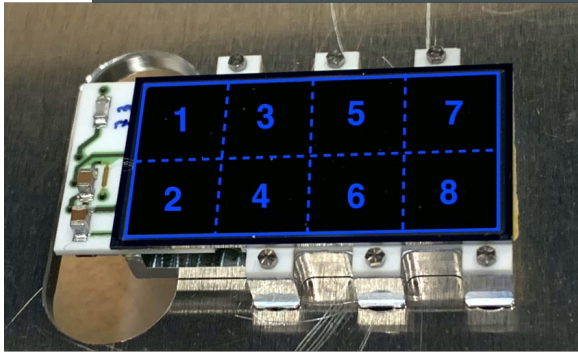
A new kaonic helium measurement in gas
by SIDDHARTINO at the DAΦNE collider*

Spectroscopy Response in a High Background Environment

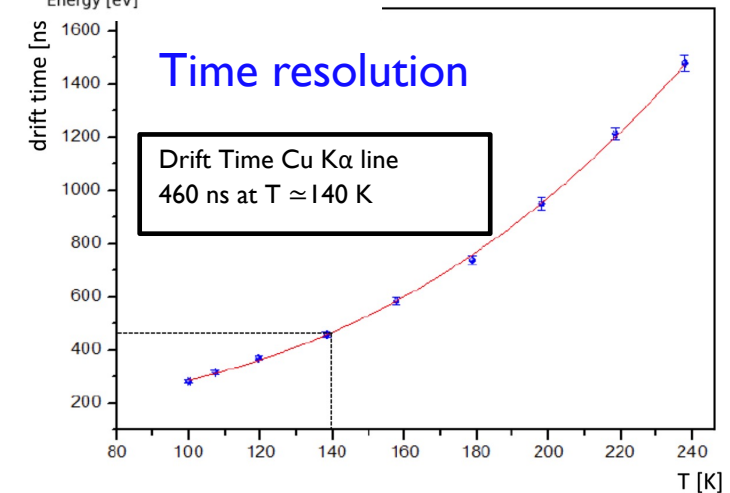
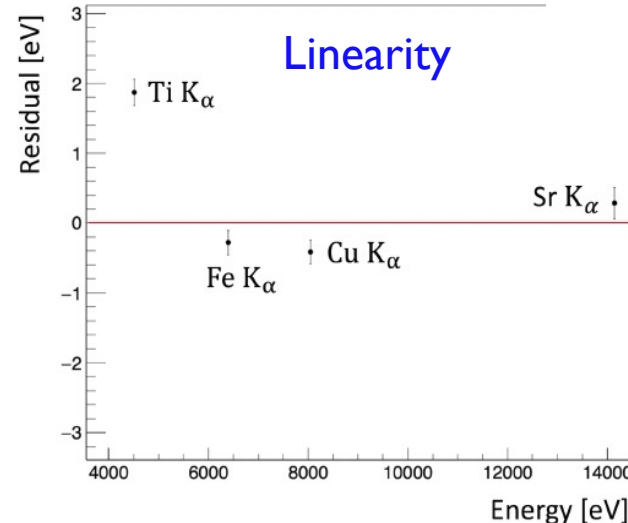
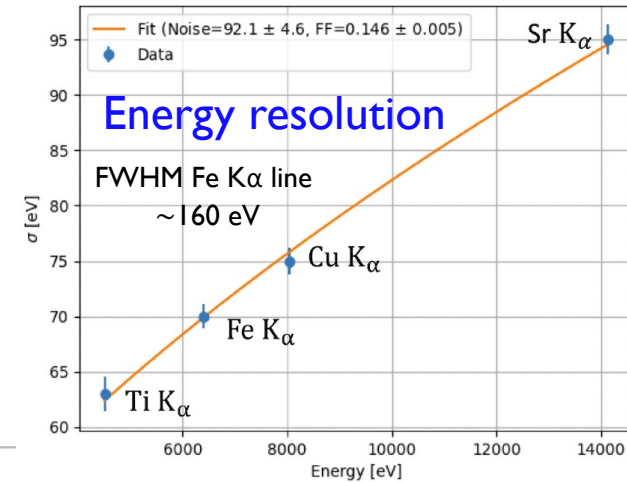
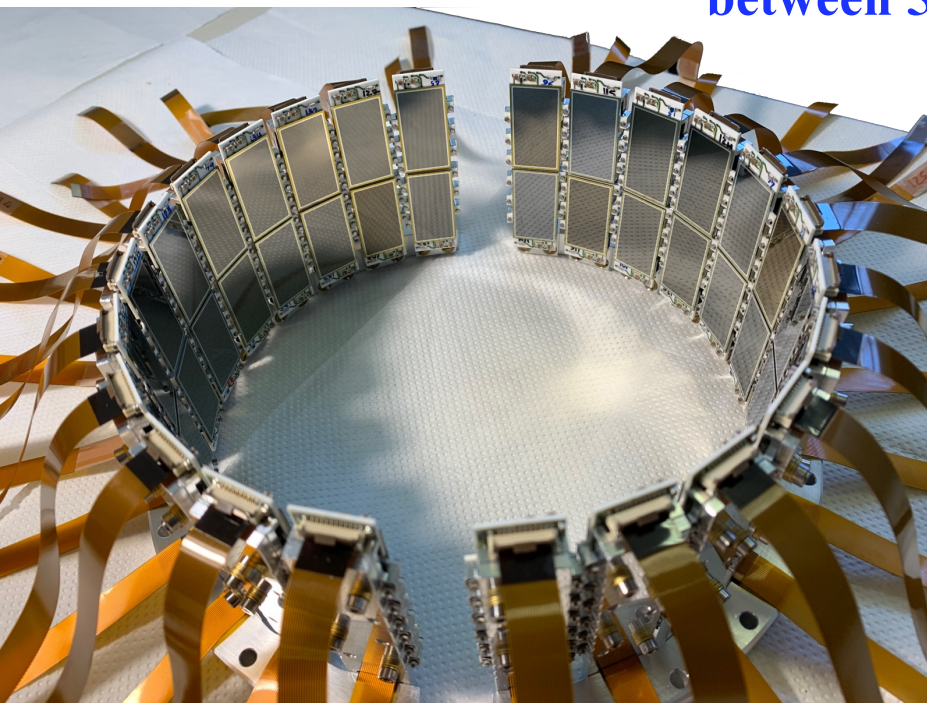


Silicon Drift Detectors

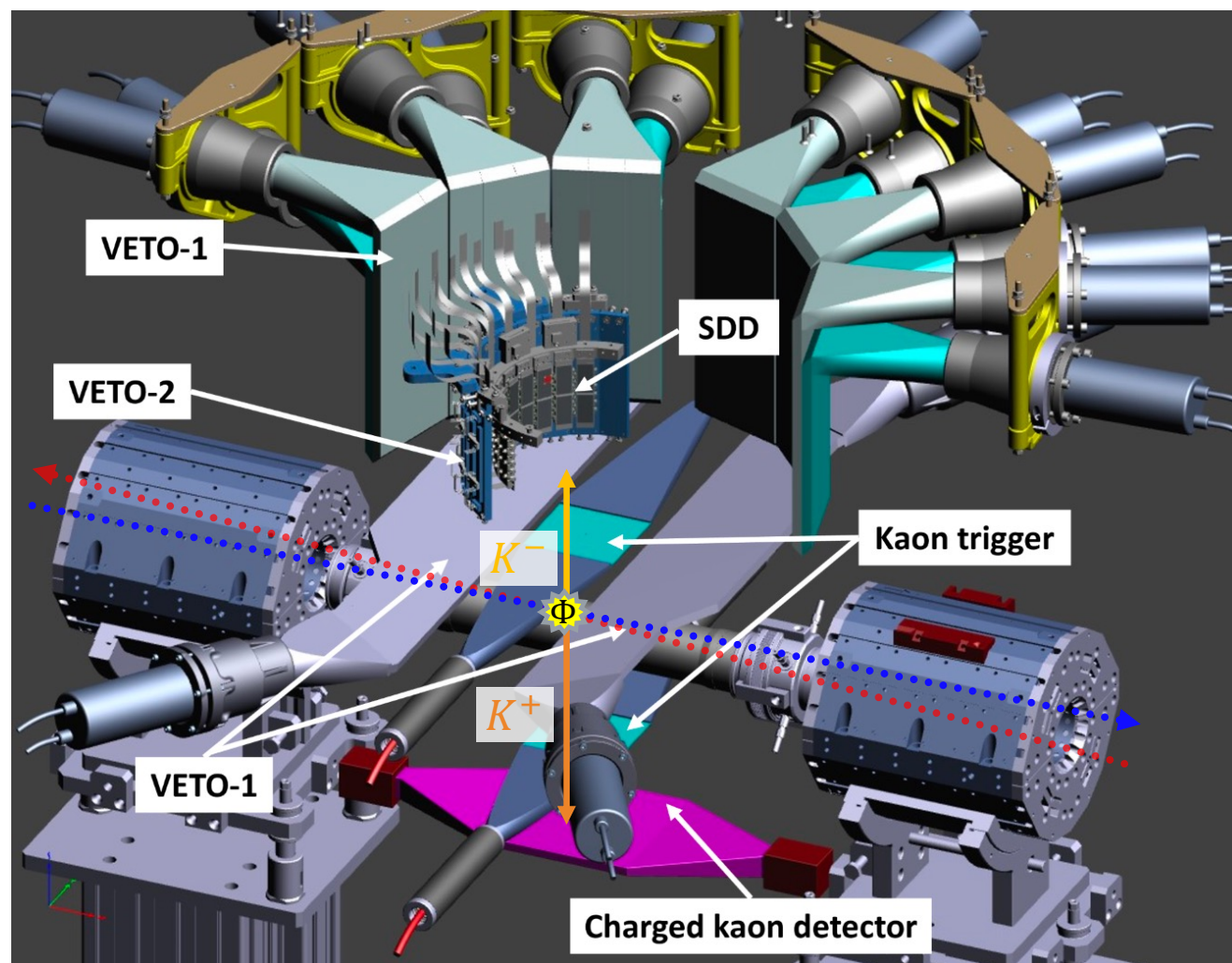
Large area Silicon Drift Detectors (SDDs) have been developed to perform high precision kaonic atoms X-ray spectroscopy



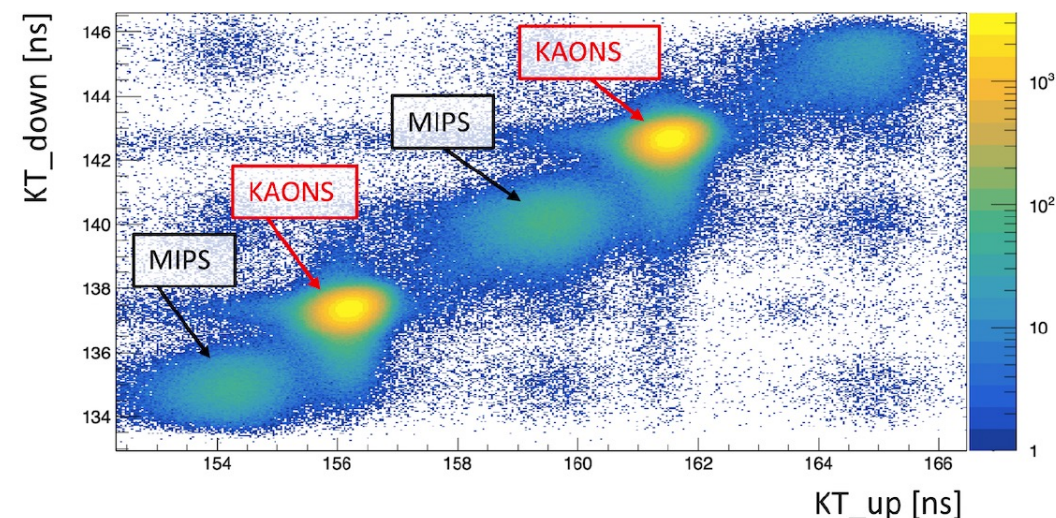
8 SDD units (0.64 cm^2)
for a total active area of 5.12 cm^2
Thickness of $450 \mu\text{m}$ ensures a high
collection efficiency for X-rays of energy
between 5 keV and 12 keV



The SIDDHARTA-2 apparatus



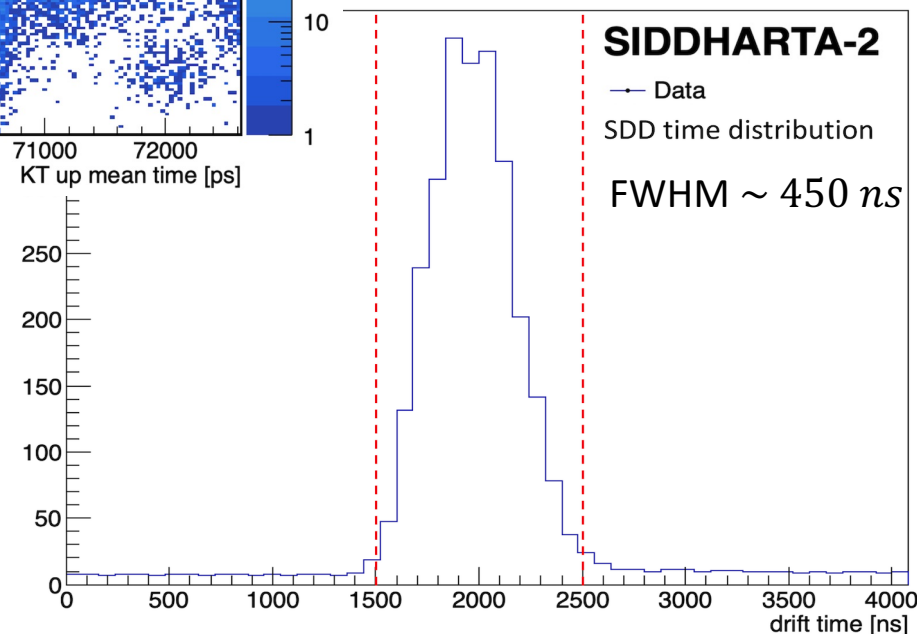
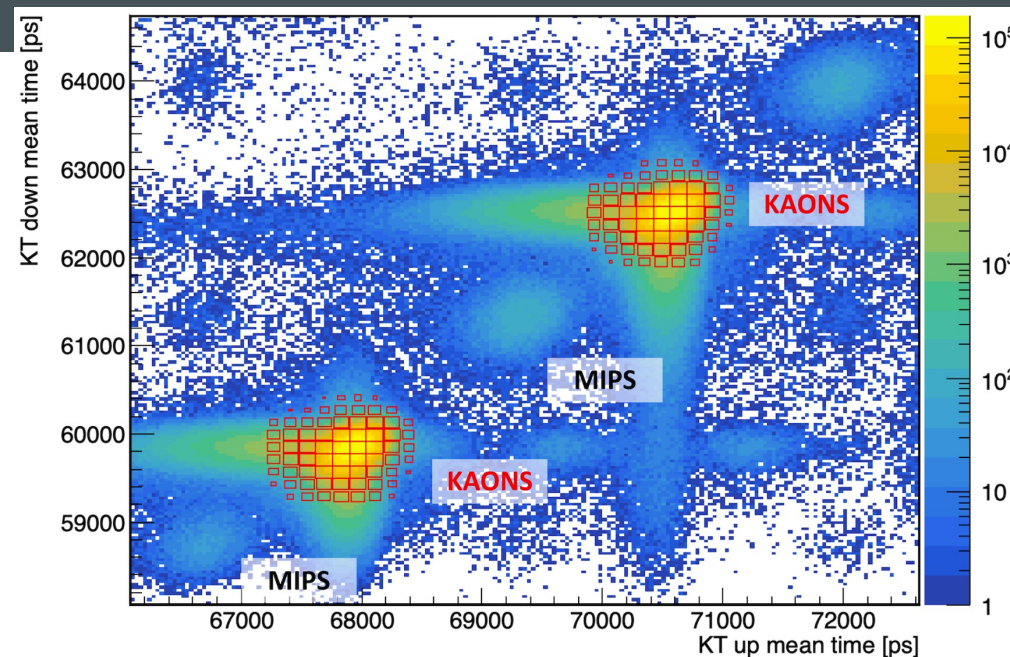
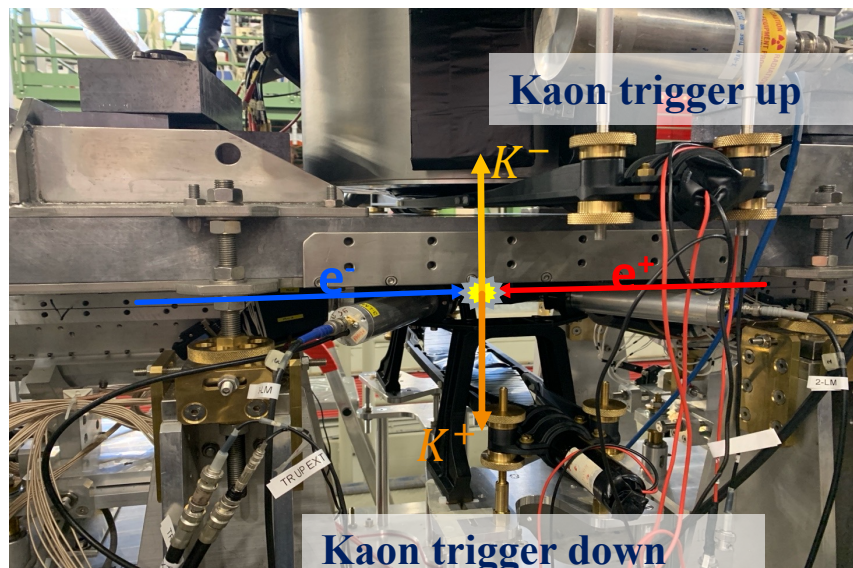
Kaon Trigger: two plastic scintillators read by photomultipliers placed above and below the interaction region.



The ToF is different for Kaons, $m(K) \sim 500 \text{ MeV}/c^2$ and light particles originating from beam-beam and beam-environment interaction (MIPs).
Can efficiently discriminate by ToF Kaons and MIPs!

The first kaonic deuterium measurement

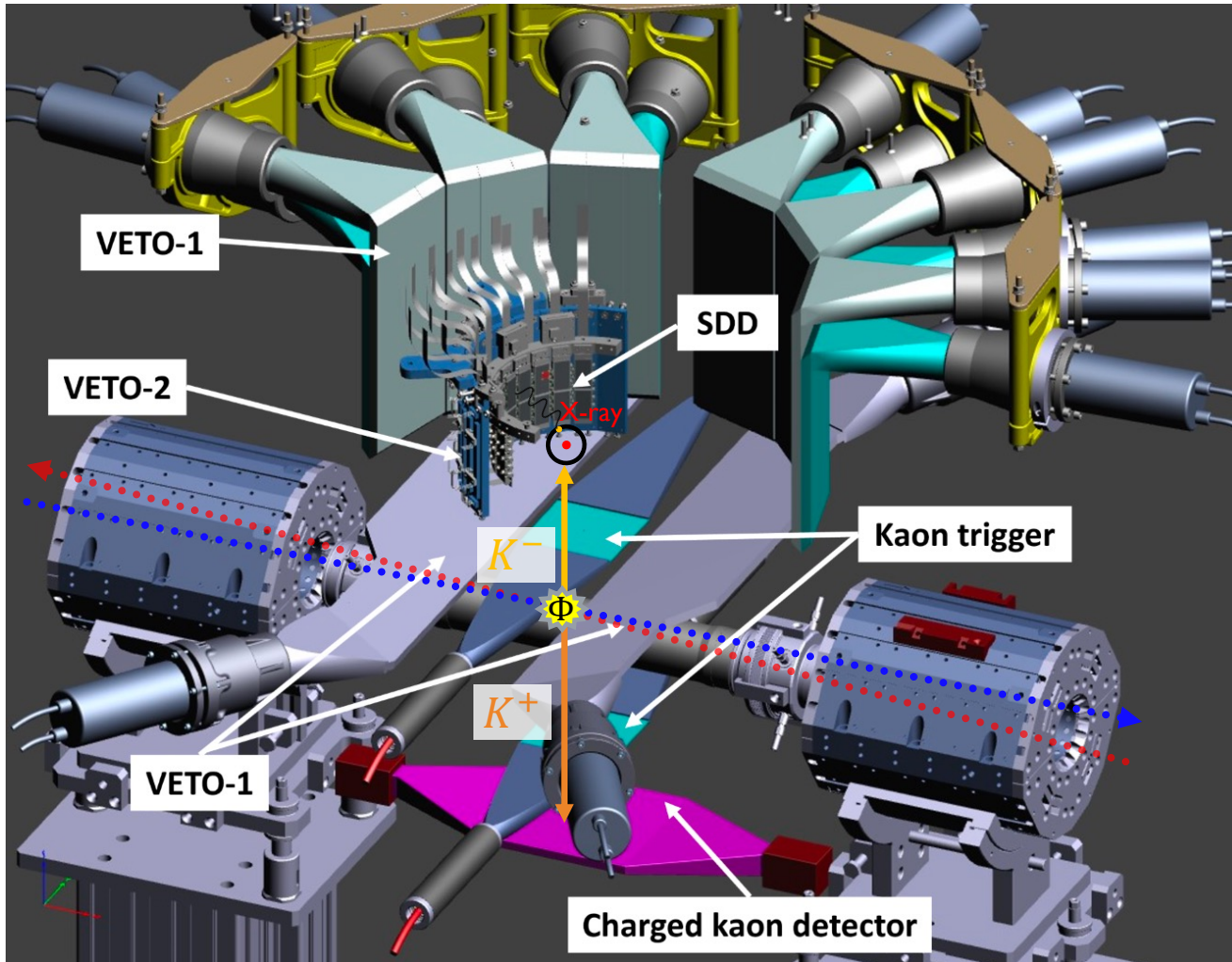
The combined use of Kaon Trigger and SDDs drift time allows to reduce the asynchronous background by a factor $\sim 2 \cdot 10^4$



The trigger is generated by the coincidence of 2 back-to-back scintillators

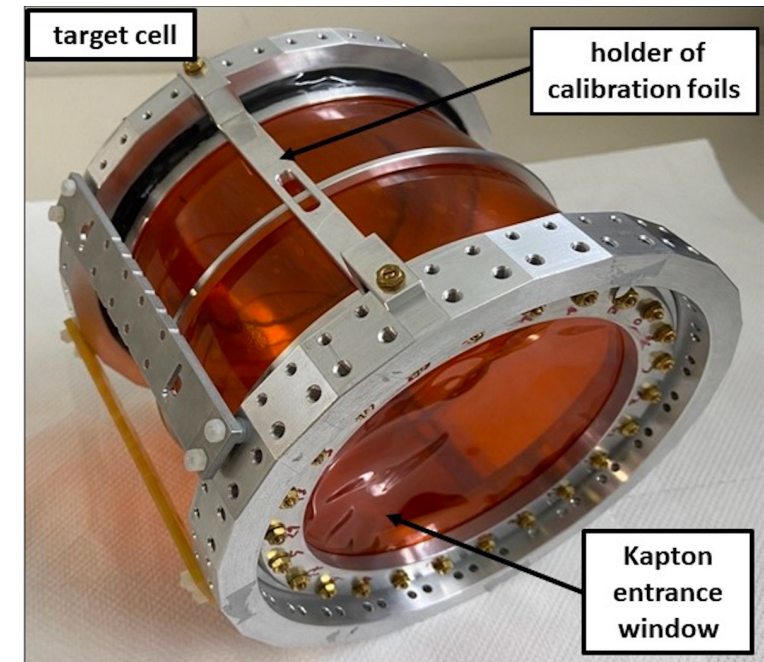
The ToF is different (~ 2 ns) for Kaons, $m(K) \sim 500$ MeV/ c^2 and light particles originating from beam-beam and beam-environment interaction (MIPs).
Can efficiently discriminate by ToF Kaons and MIPs!

The SIDDHARTA-2 apparatus

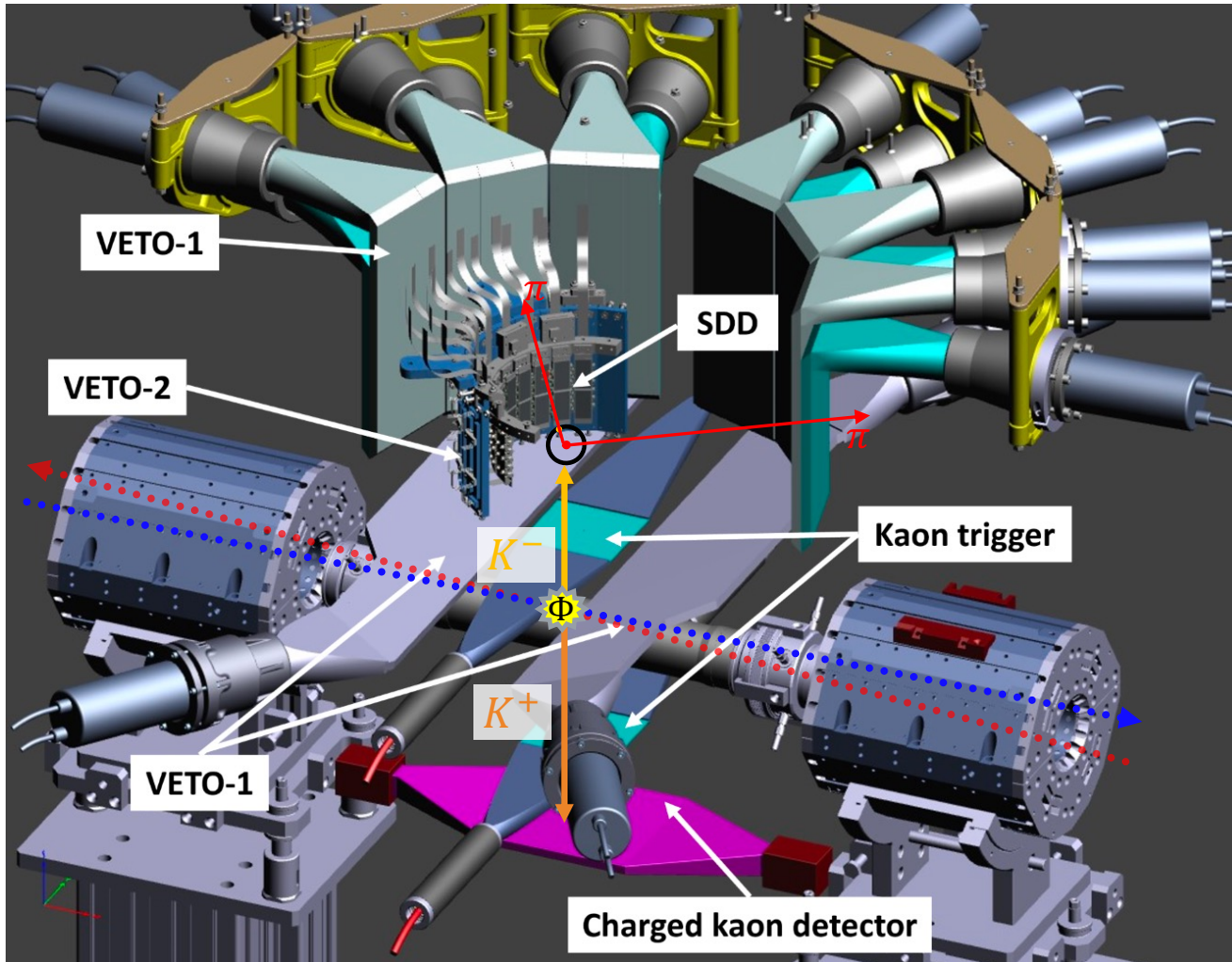


Kaon Trigger: two plastic scintillators read by photomultipliers placed above and below the interaction region.

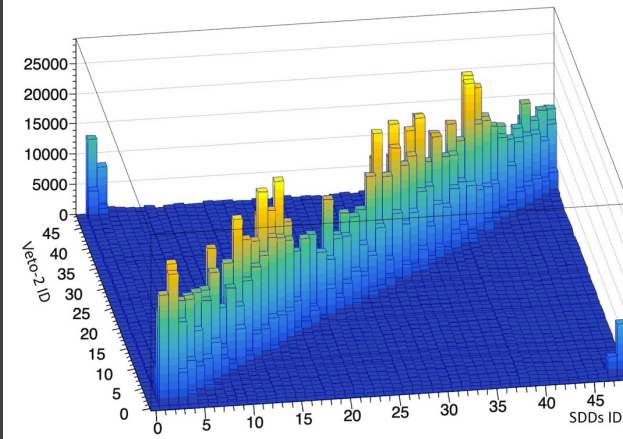
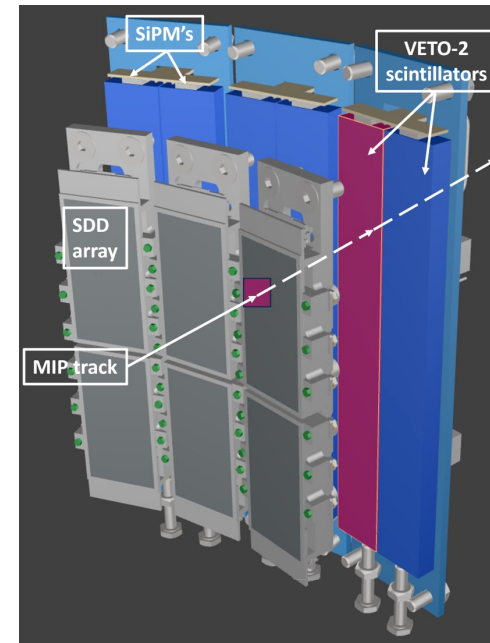
Cryogenic gaseous target cell surrounded by **384 SDDs**



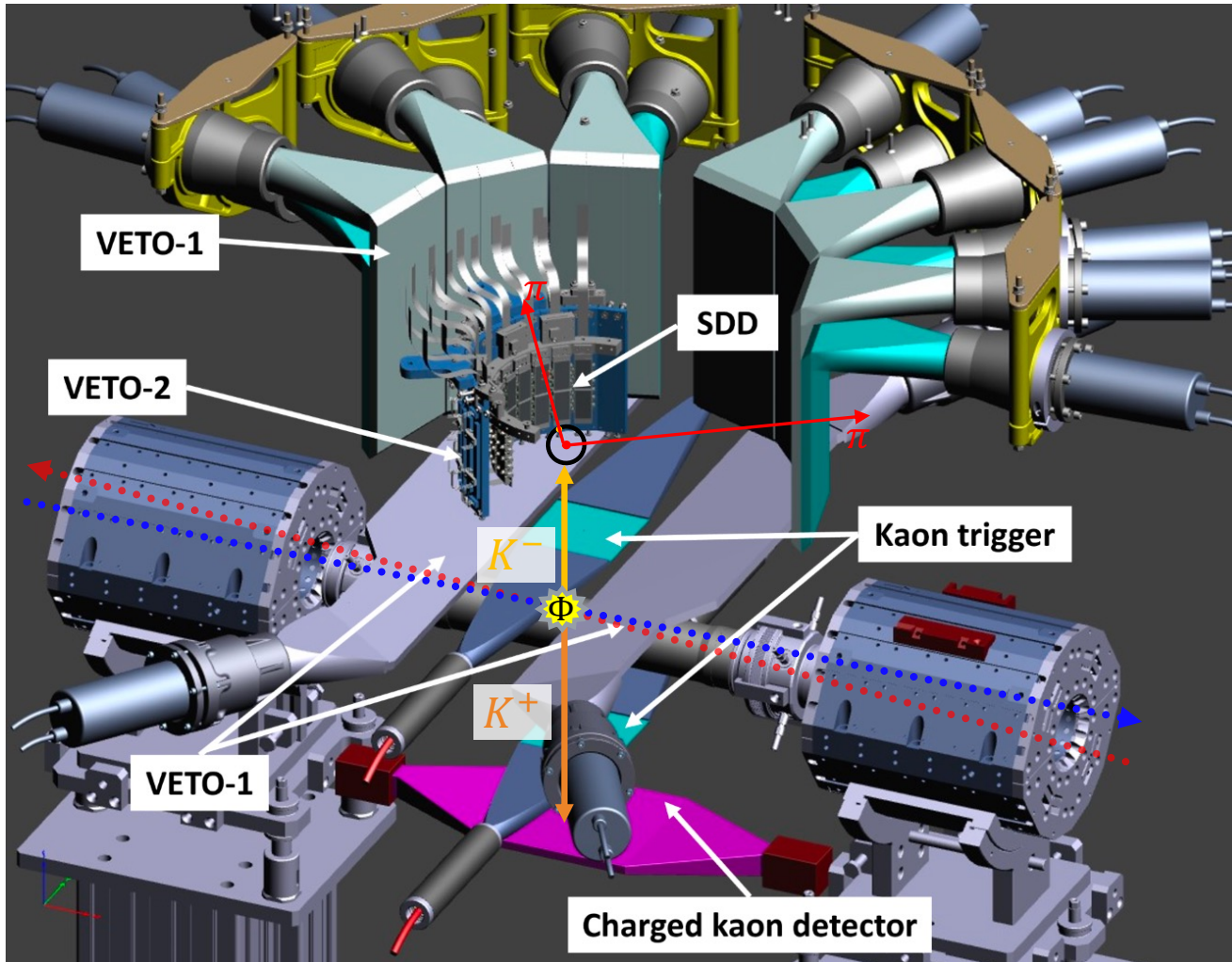
The SIDDHARTA-2 apparatus



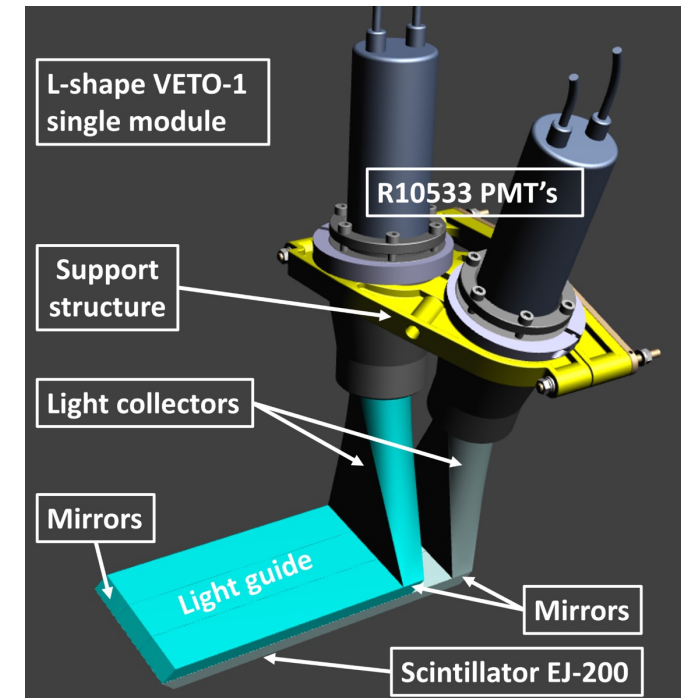
Veto-2 48 plastic scintillator read by SiPMs to suppress the background induced by particles produced by kaon absorption, passing through the SDDs.



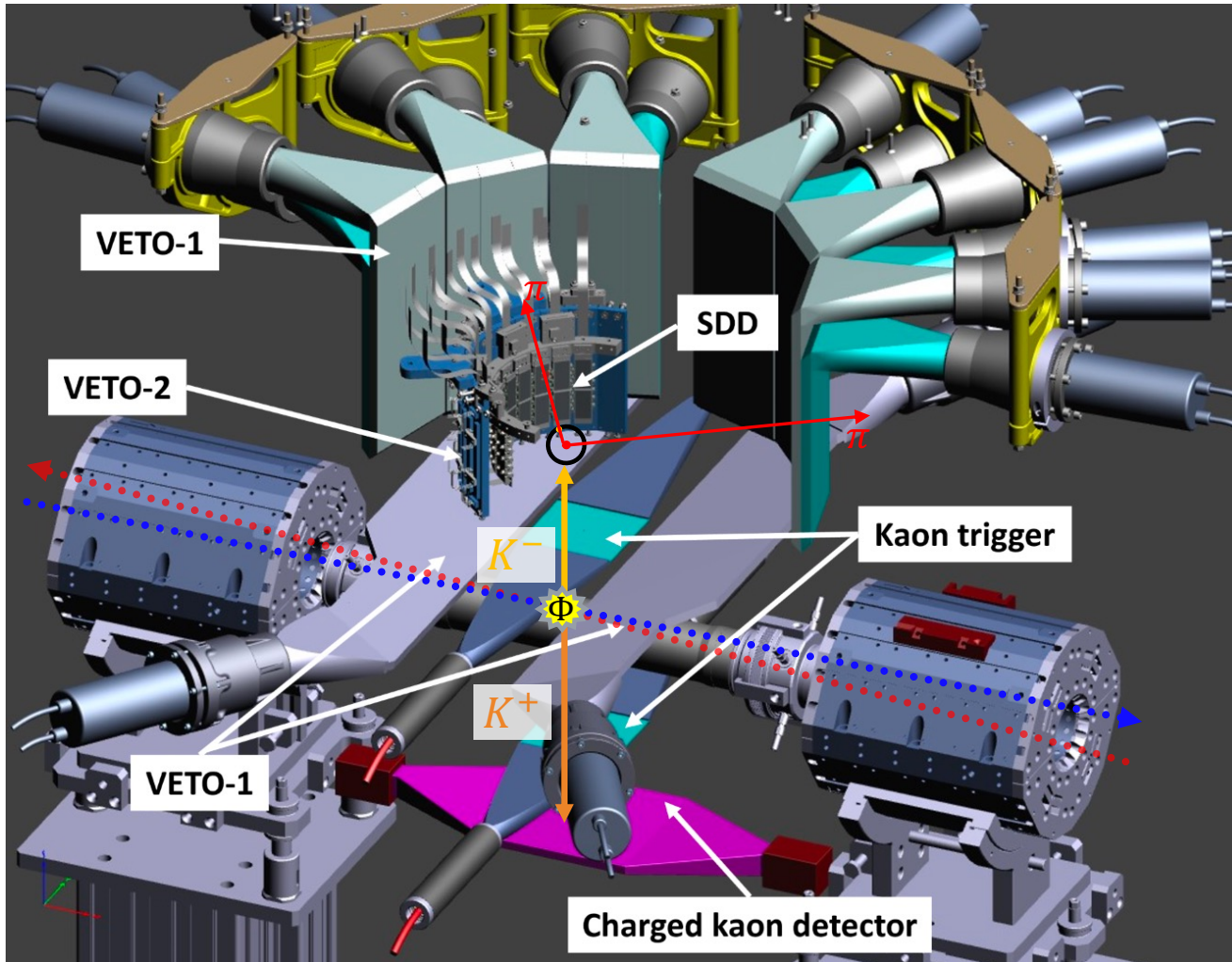
The SIDDHARTA-2 apparatus



Veto-1 14 plastic scintillator read by PMTs to select the events occurring in the gas target, rejecting the X-ray background corresponding to K- stopped in the solid elements of the setup

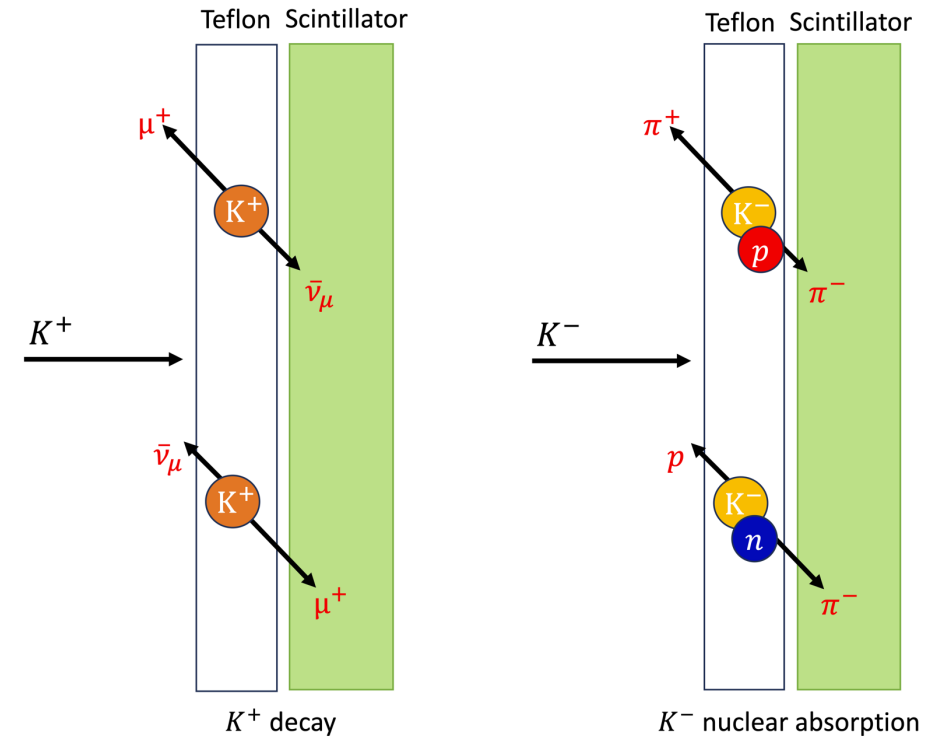


The SIDDHARTA-2 apparatus

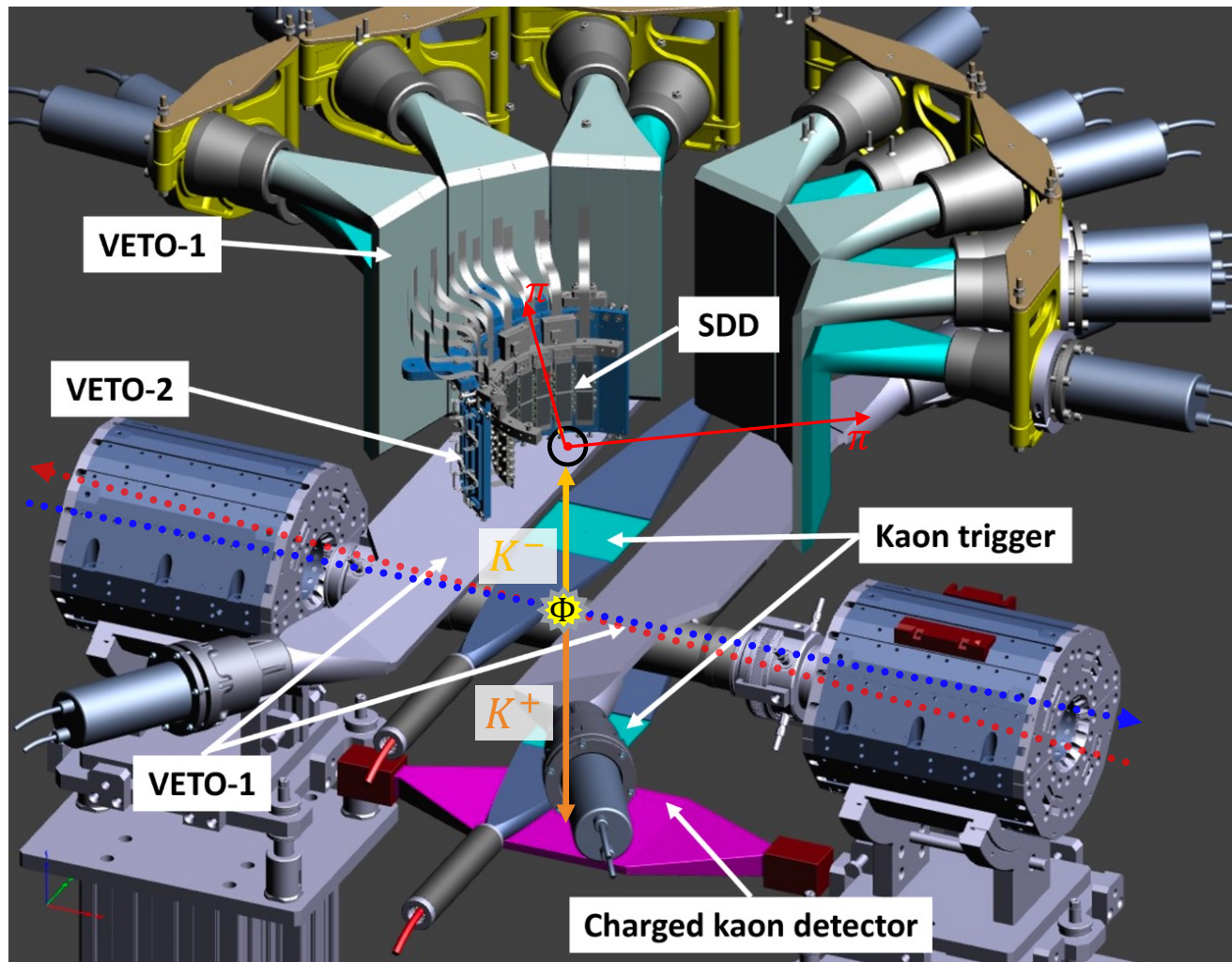


Charged Kaons Veto

Stop both K^+ and K^- in a passive layer (Teflon 3 mm) and detect secondaries charged particles using a plastic scintillator

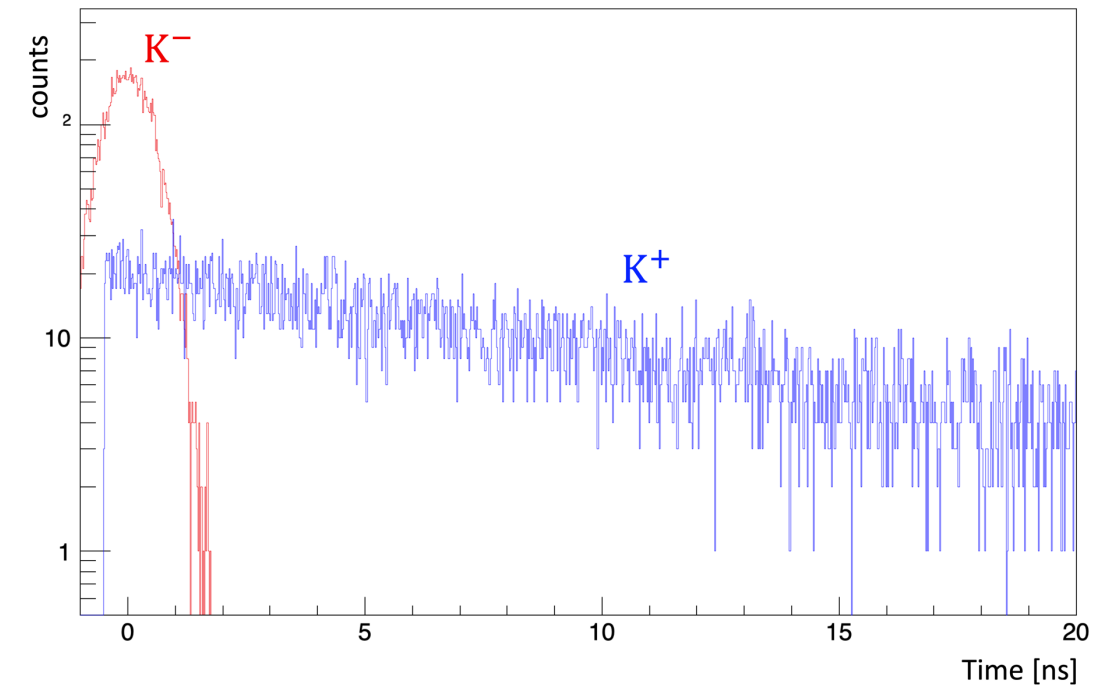


The SIDDHARTA-2 apparatus



Charged Kaons Veto

Stop both K^+ and K^- in a passive layer (Teflon 3 mm) and detect secondaries charged particles using a plastic scintillator



1. The available data on “lower levels” have big uncertainties
2. Many of them are actually UNmeasured
3. Many of them are hardly compatible among each other
4. Relative yields with upper levels are not always measured
5. Absolute yields are basically unknown (except for few transitions)
6. The REmeasured ones have been proved WRONG

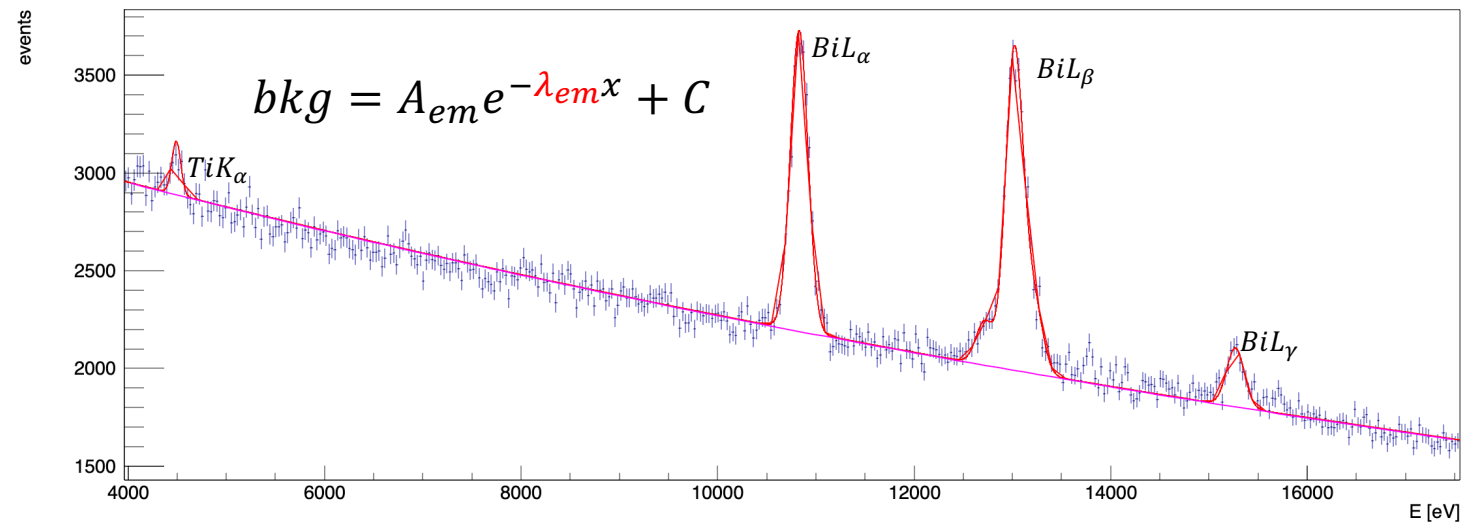
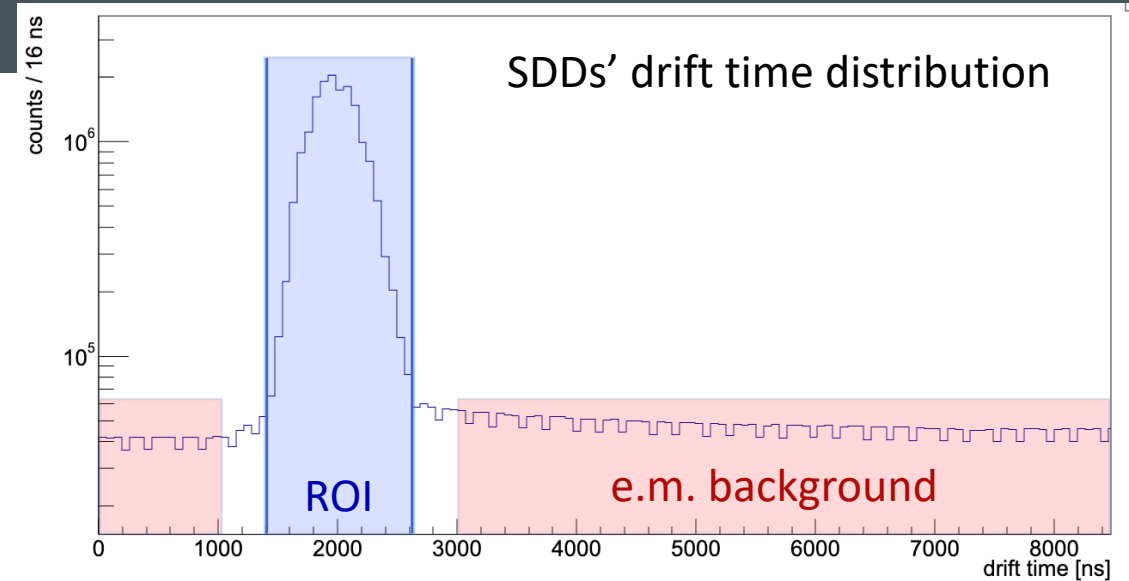
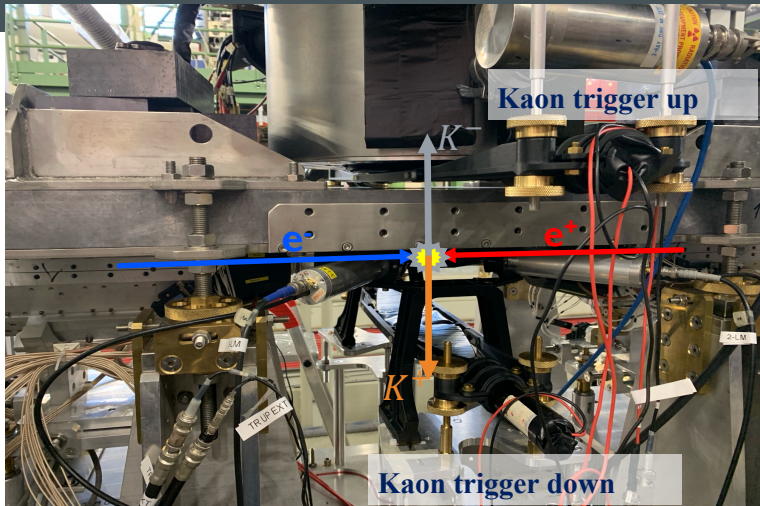
This situation would already be a
**proper justification for new
measurements**

Table 1
Compilation of K^- atomic data

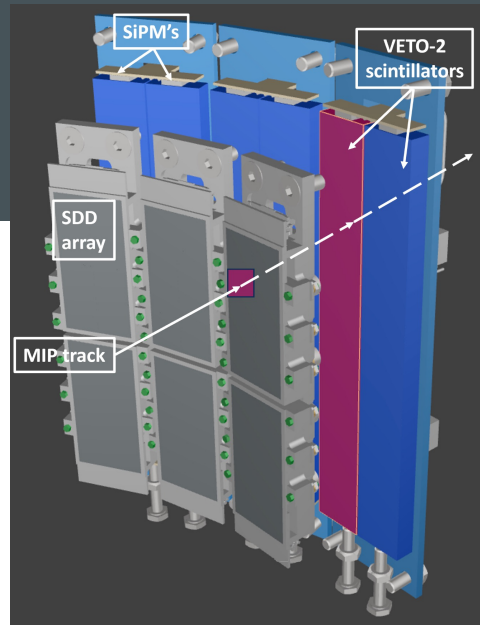
Nucleus	Transition	ϵ (keV)	Γ (keV)	Y	Γ_u (eV)	Ref.
He	$3 \rightarrow 2$	-0.04 ± 0.03	–	–	–	[15]
		-0.035 ± 0.012	0.03 ± 0.03	–	–	[16]
Li	$3 \rightarrow 2$	0.002 ± 0.026	0.055 ± 0.029	0.95 ± 0.30	–	[17]
Be	$3 \rightarrow 2$	-0.079 ± 0.021	0.172 ± 0.58	0.25 ± 0.09	0.04 ± 0.02	[17]
^{10}B	$3 \rightarrow 2$	-0.208 ± 0.035	0.810 ± 0.100	–	–	[18]
^{11}B	$3 \rightarrow 2$	-0.167 ± 0.035	0.700 ± 0.080	–	–	[18]
C	$3 \rightarrow 2$	-0.590 ± 0.080	1.730 ± 0.150	0.07 ± 0.013	0.99 ± 0.20	[18]
O	$4 \rightarrow 3$	-0.025 ± 0.018	0.017 ± 0.014	–	–	[19]
Mg	$4 \rightarrow 3$	-0.027 ± 0.015	0.214 ± 0.015	0.78 ± 0.06	0.08 ± 0.03	[19]
Al	$4 \rightarrow 3$	-0.130 ± 0.050	0.490 ± 0.160	–	–	[20]
		-0.076 ± 0.014	0.442 ± 0.022	0.55 ± 0.03	0.30 ± 0.04	[19]
Si	$4 \rightarrow 3$	-0.240 ± 0.050	0.810 ± 0.120	–	–	[20]
		-0.130 ± 0.015	0.800 ± 0.033	0.49 ± 0.03	0.53 ± 0.06	[19]
P	$4 \rightarrow 3$	-0.330 ± 0.08	1.440 ± 0.120	0.26 ± 0.03	1.89 ± 0.30	[18]
S	$4 \rightarrow 3$	-0.550 ± 0.06	2.330 ± 0.200	0.22 ± 0.02	3.10 ± 0.36	[18]
		-0.43 ± 0.12	2.310 ± 0.170	–	–	[21]
Cl	$4 \rightarrow 3$	-0.462 ± 0.054	1.96 ± 0.17	0.23 ± 0.03	2.9 ± 0.5	[19]
		-0.770 ± 0.40	3.80 ± 1.0	0.16 ± 0.04	5.8 ± 1.7	[18]
		-0.94 ± 0.40	3.92 ± 0.99	–	–	[22]
		-1.08 ± 0.22	2.79 ± 0.25	–	–	[21]
Co	$5 \rightarrow 4$	-0.099 ± 0.106	0.64 ± 0.25	–	–	[19]
Ni	$5 \rightarrow 4$	-0.180 ± 0.070	0.59 ± 0.21	0.30 ± 0.08	5.9 ± 2.3	[20]
		-0.246 ± 0.052	1.23 ± 0.14	–	–	[19]
Cu	$5 \rightarrow 4$	-0.240 ± 0.220	1.650 ± 0.72	0.29 ± 0.11	7.0 ± 3.8	[20]
		-0.377 ± 0.048	1.35 ± 0.17	0.36 ± 0.05	5.1 ± 1.1	[19]
Ag	$6 \rightarrow 5$	-0.18 ± 0.12	1.54 ± 0.58	0.51 ± 0.16	7.3 ± 4.7	[19]
Cd	$6 \rightarrow 5$	-0.40 ± 0.10	2.01 ± 0.44	0.57 ± 0.11	6.2 ± 2.8	[19]
In	$6 \rightarrow 5$	-0.53 ± 0.15	2.38 ± 0.57	0.44 ± 0.08	11.4 ± 3.7	[19]
Sn	$6 \rightarrow 5$	-0.41 ± 0.18	3.18 ± 0.64	0.39 ± 0.07	15.1 ± 4.4	[19]
Ho	$7 \rightarrow 6$	-0.30 ± 0.13	2.14 ± 0.31	–	–	[23]
Yb	$7 \rightarrow 6$	-0.12 ± 0.10	2.39 ± 0.30	–	–	[23]
Ta	$7 \rightarrow 6$	-0.27 ± 0.50	3.76 ± 1.15	–	–	[23]
Pb	$8 \rightarrow 7$	–	0.37 ± 0.15	0.79 ± 0.08	4.1 ± 2.0	[24]
		-0.020 ± 0.012	–	–	–	[25]
U	$8 \rightarrow 7$	-0.26 ± 0.4	1.50 ± 0.75	0.35 ± 0.12	45 ± 24	[24]

Electromagnetic background (asynchronous wrt kaons)

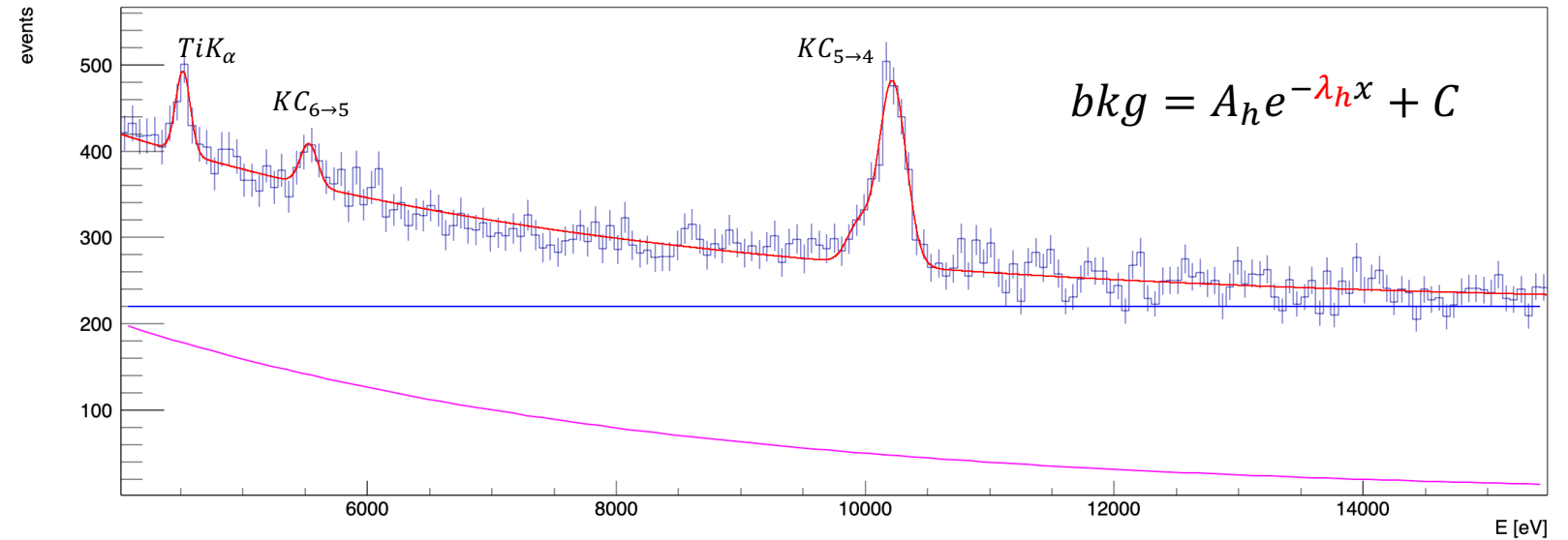
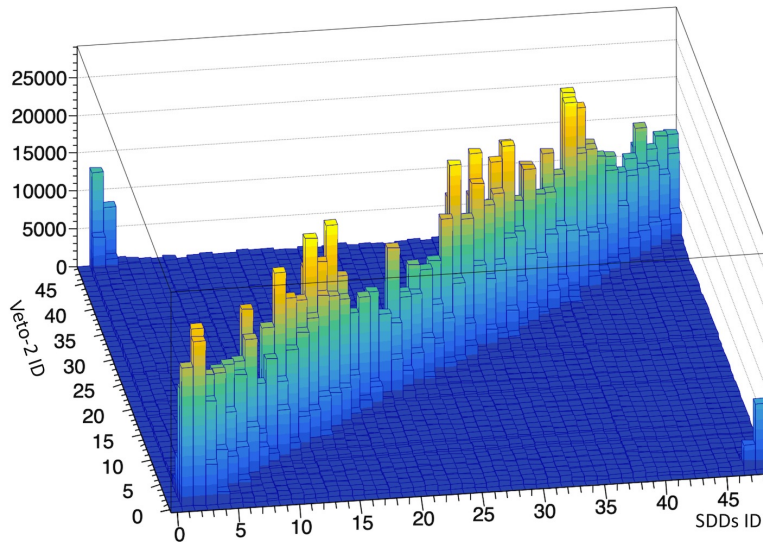
By exploiting the trigger and drift time of the SDDs, we can isolate a pure electromagnetic background spectrum by selecting events uncorrelated with kaon production.



Hadronic background (synchronous wrt kaons)



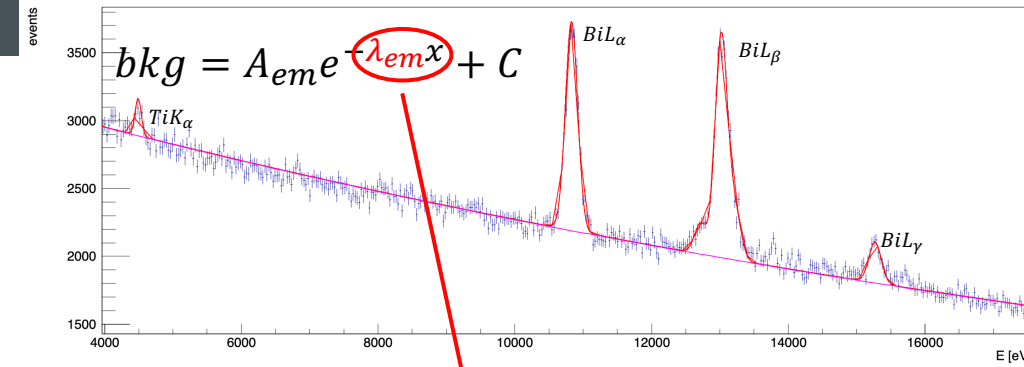
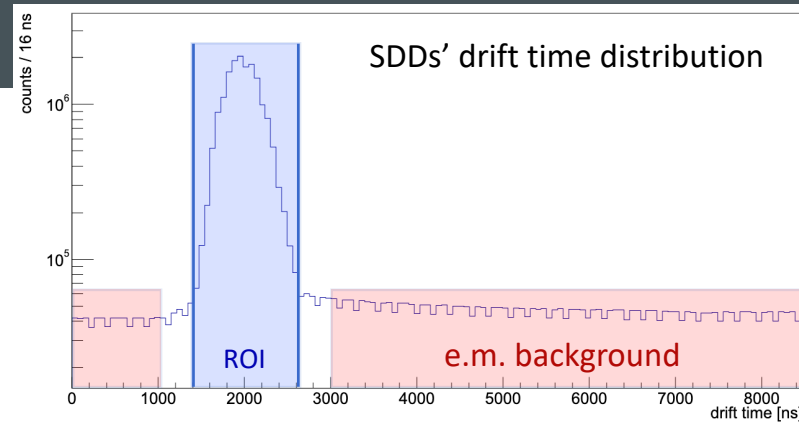
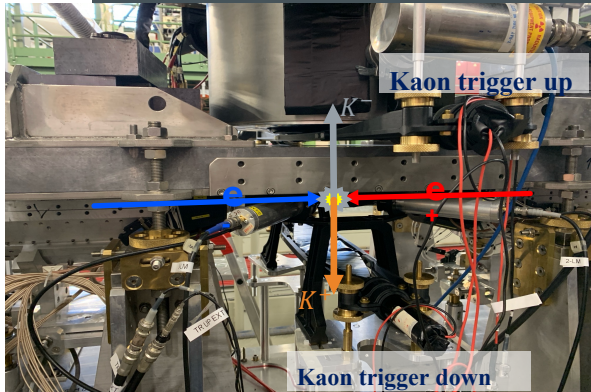
Hadronic background from MIPs (mainly pions) produced by kaon nuclear absorption can be isolated using the spatial correlation between Veto-2 and the SDDs, yielding a pure hadronic background spectrum.



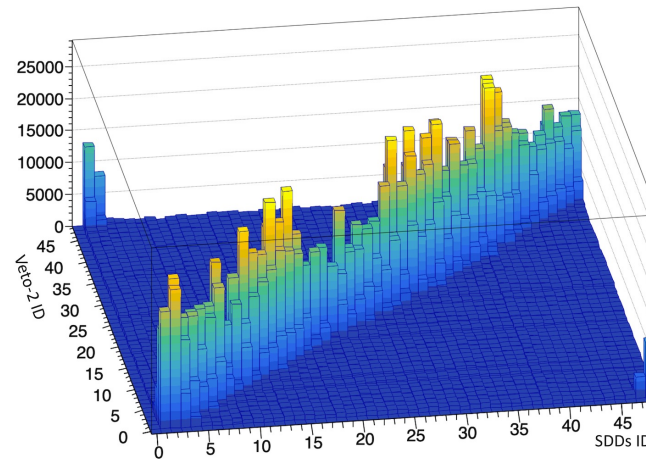
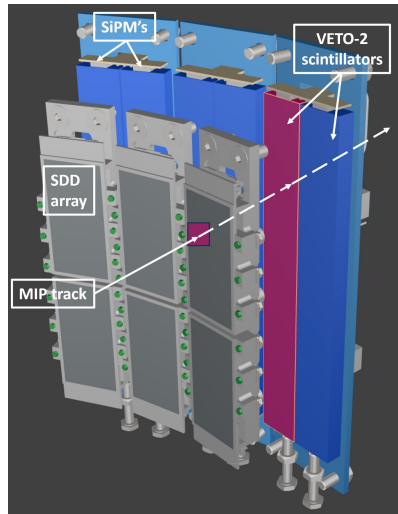
Plot of the topological correlation between Veto2's scintillators and SDDs

Study of background events (outside the signal window)

Electromagnetic (asynchronous) background

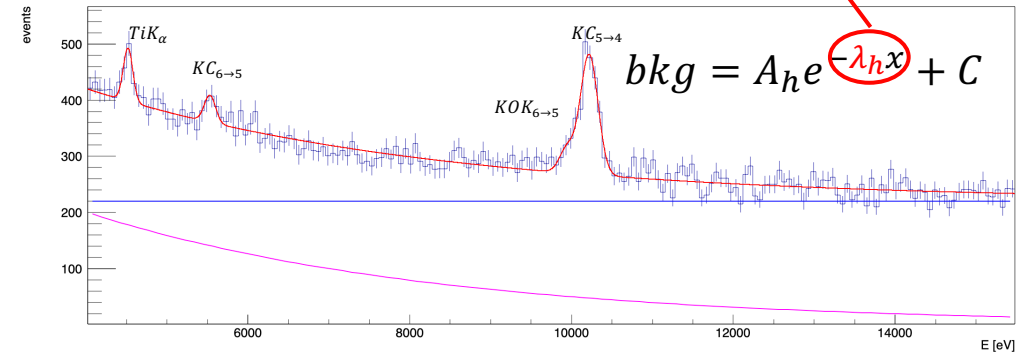


Hadronic (synchronous) background

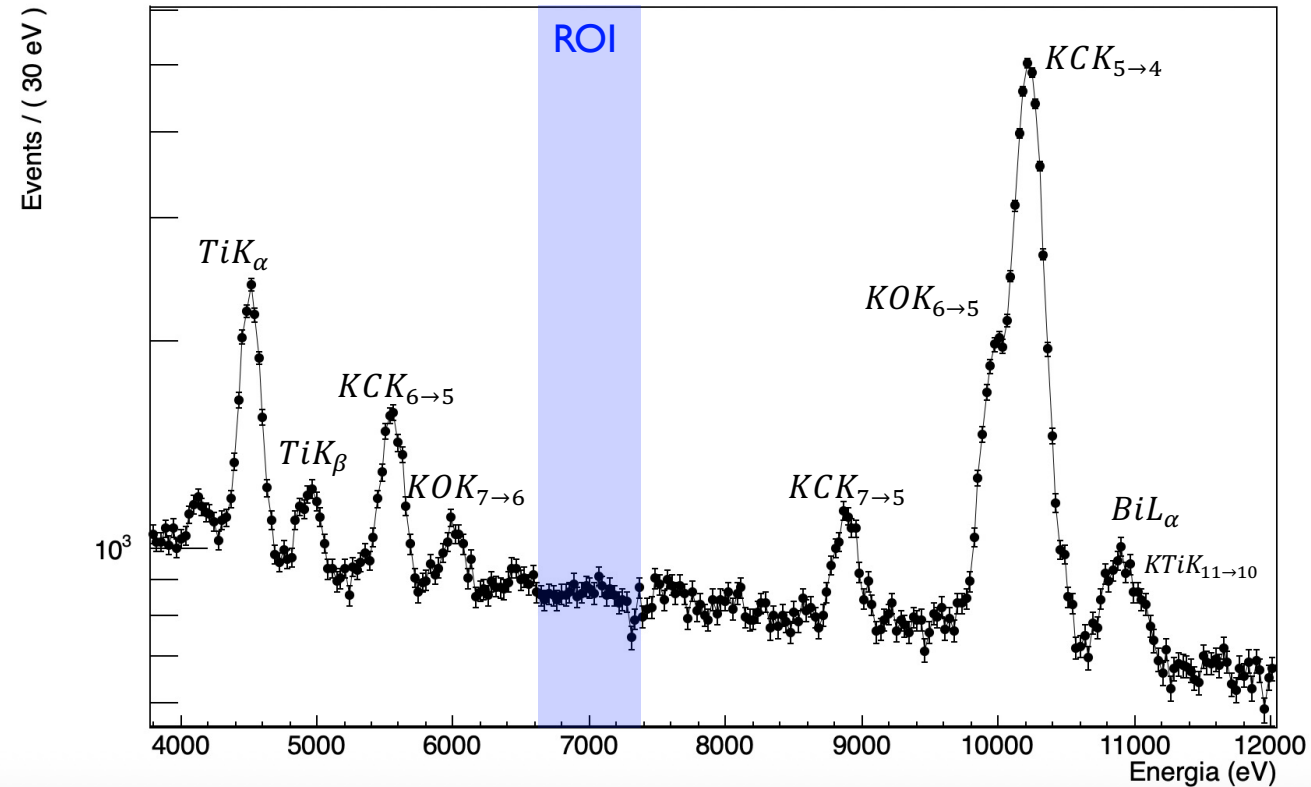
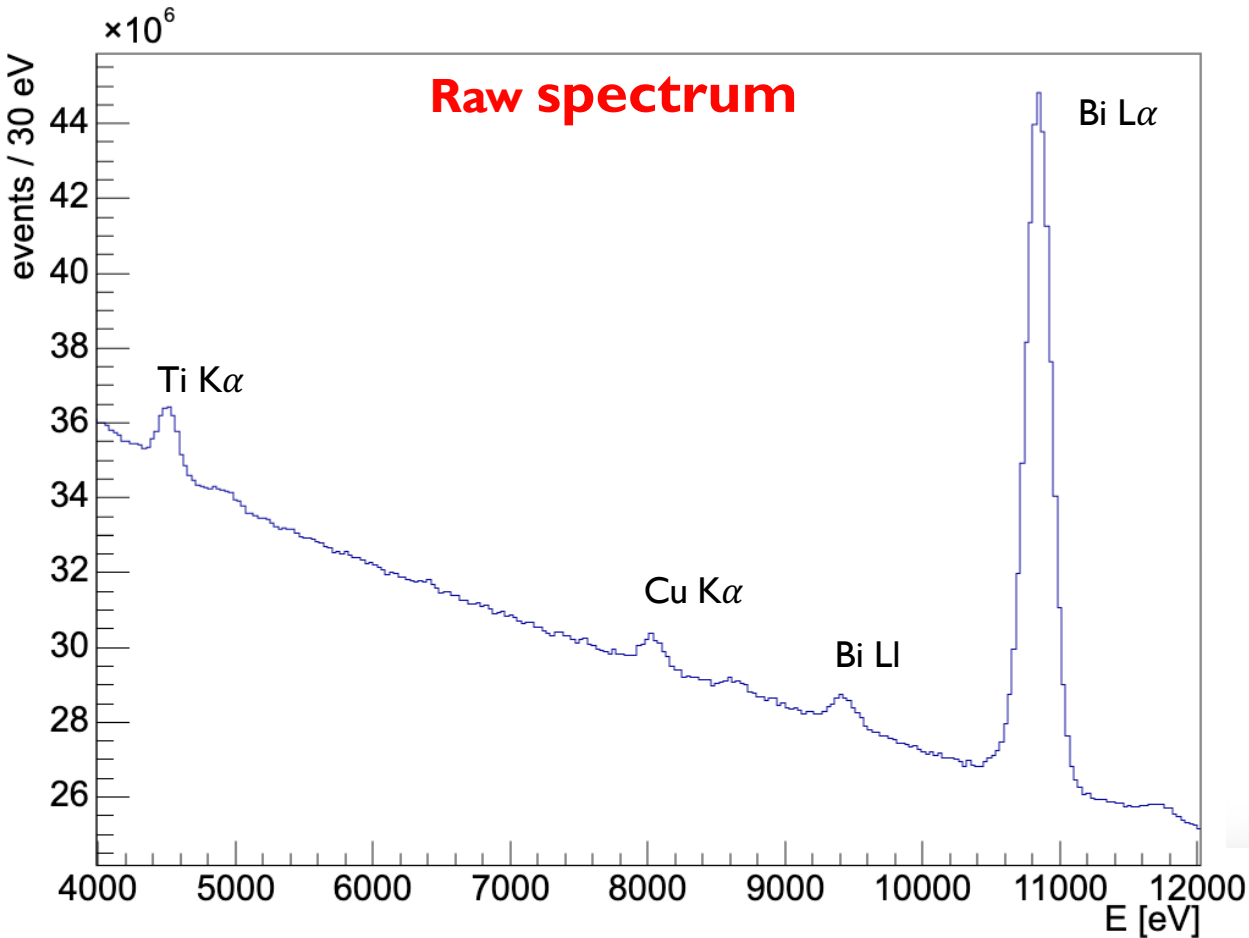


Plot of the topological correlation between Veto2's scintillators and SDDs

Input values for background description



Kaonic Deuterium Energy Spectrum



Background reduced by a factor 3×10^4

Kaonic Deuterium Energy Spectrum – Fit procedure

- We perform an extended maximum-likelihood fit to the binned spectrum, including systematic uncertainties as nuisance parameters
- The full model is implemented in ROOT/RooFit

$$f(E; \theta) = \underbrace{\sum_{n=0}^{18} y_n G_n(E; \mu, \sigma)}_{\text{X-ray transitions from known sources}} + \underbrace{y_{2p \rightarrow 1s} V_{2p \rightarrow 1s}(E; \mu, \sigma, \Gamma) + \sum_{n=3}^7 y_{n \rightarrow 1s} V_{n \rightarrow 1s}(E; \mu, \sigma, \Gamma^*)}_{\text{Kaonic deuterium lines}} + \underbrace{y_{em} e^{\lambda_{em} E} + y_{had} e^{\lambda_{had} E}}_{\text{Background E.M and hadronic}} + C$$

X-ray transitions from known sources
(*Eur. Phys. J. A* (2023) 59 3, 56)

Kaonic deuterium lines

Background
E.M and hadronic

$$-2\ln\mathcal{L}_{tot} = -2\ln\mathcal{L}_{ext} - 2\ln\mathcal{L}_{nuisance}$$

$$-2\ln\mathcal{L}_{nuisance} = \underbrace{\frac{(\delta\mu_{cal})^2}{\sigma_{cal}^2} + \frac{(\delta\mu_{stab})^2}{\sigma_{stab}^2}}_{\text{Systematic uncertainty on SDDs energy calibration and stability}} + \underbrace{\frac{(FF - FF_0)^2}{\sigma_{FF}^2} + \frac{(noise - noise_0)^2}{\sigma_{noise}^2}}_{\text{Systematic uncertainty on SDDs energy resolution}} + \underbrace{\frac{(\lambda_{em} - \lambda_{em}^0)^2}{\sigma_{\lambda_{em}}^2} + \frac{(\lambda_{had} - \lambda_{had}^0)^2}{\sigma_{\lambda_{had}}^2}}_{\text{Systematic uncertainty on e.m. and hadronic background}}$$

Systematic uncertainty on SDDs energy calibration and stability

Systematic uncertainty on SDDs energy resolution

Systematic uncertainty on e.m. and hadronic background

(*Phys.Scripta* 97 (2022) 11, 114002; *Measur.Sci.Tech.* 32 (2021) 9, 095501; *Measur.Sci.Tech.* 33 (2022) 9, 095502)

Beyond SIDDHARTA-2

Except for the most recent measurements at DAFNE and JPARC on KHe and KH, the database on kaonic atoms dates back to 1970s and 1980s. These measurements are the basis for all the theoretical model used for: KN, KNN interaction at threshold, Kaon mass, Kaonic atoms cascade models

Present status:

1. The available data on “lower levels” have big uncertainties
2. Many of them are hardly compatible among each other
3. Many atoms are actually Unmeasured
4. Absolute yields are basically unknown (except for few transitions)

We propose to do precision measurements along the periodic table at DAFNE for:

- Selected light kaonic atoms
- Selected intermediate mass kaonic atoms
- Selected heavy kaonic atoms

charting the periodic table

E. Friedman et al. / Nuclear Physics A579 (1994) 518–538

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Table 1
Compilation of K^- atomic data

Nucleus	Transition	ϵ (keV)	Γ (keV)	Y	Γ_a (eV)	Ref.
He	3 $\rightarrow$ 2	-0.04 ± 0.03	–	–	–	[15]
		-0.035 ± 0.012	0.03 ± 0.03	–	–	[16]
Li	3 $\rightarrow$ 2	0.002 ± 0.026	0.055 ± 0.029	0.95 ± 0.30	–	[17]
Be	3 $\rightarrow$ 2	-0.079 ± 0.021	0.172 ± 0.58	0.25 ± 0.09	0.04 ± 0.02	[17]
^{10}B	3 $\rightarrow$ 2	-0.208 ± 0.035	0.810 ± 0.100	–	–	[18]
^{11}B	3 $\rightarrow$ 2	-0.167 ± 0.035	0.700 ± 0.080	–	–	[18]
C	3 $\rightarrow$ 2	-0.590 ± 0.080	1.730 ± 0.150	0.07 ± 0.013	0.99 ± 0.20	[18]
O	4 $\rightarrow$ 3	-0.025 ± 0.018	0.017 ± 0.014	–	–	[19]
Mg	4 $\rightarrow$ 3	-0.027 ± 0.015	0.214 ± 0.015	0.78 ± 0.06	0.08 ± 0.03	[19]
Al	4 $\rightarrow$ 3	-0.130 ± 0.050	0.490 ± 0.160	–	–	[20]
		-0.076 ± 0.014	0.442 ± 0.022	0.55 ± 0.03	0.30 ± 0.04	[19]
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Beyond SIDDHARTA-2

A new era of kaonic atoms measurements along the periodic table requires suitable detectors to match the different transitions energies

*Journal of Instrumentation, 19
(2024) P07039*

- SDDs**
- 100 eV max resolution
 - 4-40 keV range
 - High efficiency

*Nuclear Instruments and Methods
in Physics Research A 1060 (2024)
169060*

- Cd(Zn)Te**
- 20-300 keV range
 - FWHM / E ~ %
 - High efficiency
 - Room Temperature

*Nuclear Instruments and Methods
in Physics Research A 1069 (2024)
169966*

- HPGe**
- 100-1000 keV range
 - FWHM / E ~ %
 - High efficiency
 - Cooling needed

