

The phi meson in dense matter from a theoretical perspective

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P. Gubler, M. Ichikawa, T. Song and E. Bratkovskaya, Phys. Rev. C **111**, 034908 (2025).

M. Ichikawa, P. Gubler, et al., (KEK-PS E325 Collaboration), Prog. Theor. Exp. Phys. **2025**, 093D01 (2025).

G. Balassa, K. Aoki, P. Gubler, S.H. Lee, H. Sako and G. Wolf, arXiv:2508.11344 [hep-ph] (to be published in PTEP).

Talk at the 15th International Conference on Hypernuclear and
Strange Particle Physics (HYP2025),
University of Tokyo, Japan
October 1, 2025

Work done in
collaboration
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H. Sako (JAEA)
K. Aoki (KEK)
S.H. Lee (Yonsei U.)
G. Balassa (Yonsei U.)
G. Wolf (Wigner RCP)

R. Muto et al. (KEK E325),
Phys. Rev. Lett. **98**, 042501 (2007).

J. Adamczewski-Musch et al. (HADES),
Phys. Rev. Lett. **123**, 022002 (2019).

E16/E88 experiments at J-PARC

Dilepton/ K^+K^- spectra
from pA reactions

Chiral mixing?

C. Sasaki,
Phys. Rev. D **106**, 054034 (2022).


R. Ejima et al.,
Phys. Rev. C **111**, 055201 (2025).

Exotic dispersion
relations

H.J. Kim and P. Gubler,
Phys. Lett. B **805**, 135412 (2020).

Introduction

Partial restoration of
chiral symmetry

$$|\langle \bar{s}s \rangle_\rho|$$


T. Hatsuda and S.H. Lee,
Phys. Rev. C **46**, R34 (1992).

J. Kim, P. Gubler and S.H. Lee,
Phys. Rev. D **105**, 114053 (2022).

Pion induced ϕ
meson production

P95 experiment at J-PARC

S. Acharya et al. (ALICE Coll.),
Phys. Rev. Lett. **127**, 172301 (2021).

Femtoscscopy from LHC

ϕ N bound state?

E. Chizzali et al.,
Phys. Lett. B **848**, 138358 (2024).

B.-X. Sun et al.,
Commun. Theor. Phys.
75, 055301 (2023).

Photoproduction

I.I. Strakovsky et al.,
Phys. Rev. C **101**, 045201 (2020).

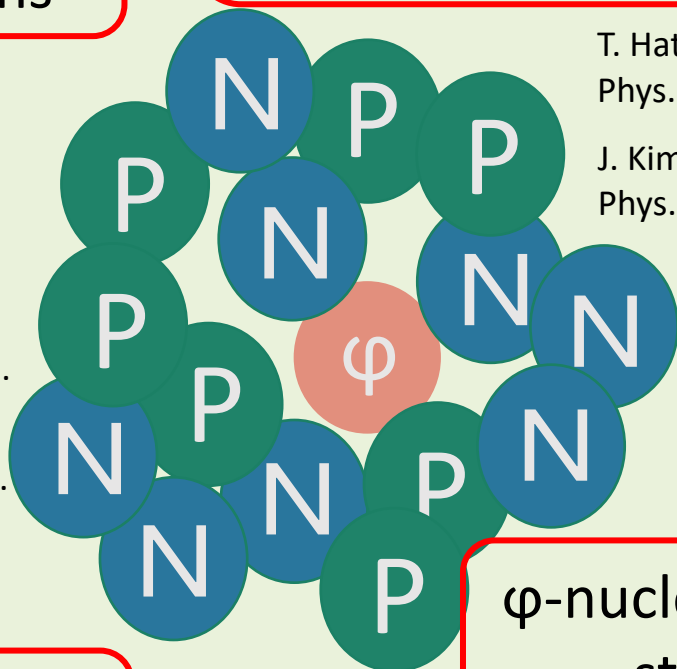
ϕ -nucleus bound
states?

J.J. Cobos-Martínez et al.,
Phys. Rev. C **96**, 035201 (2017).

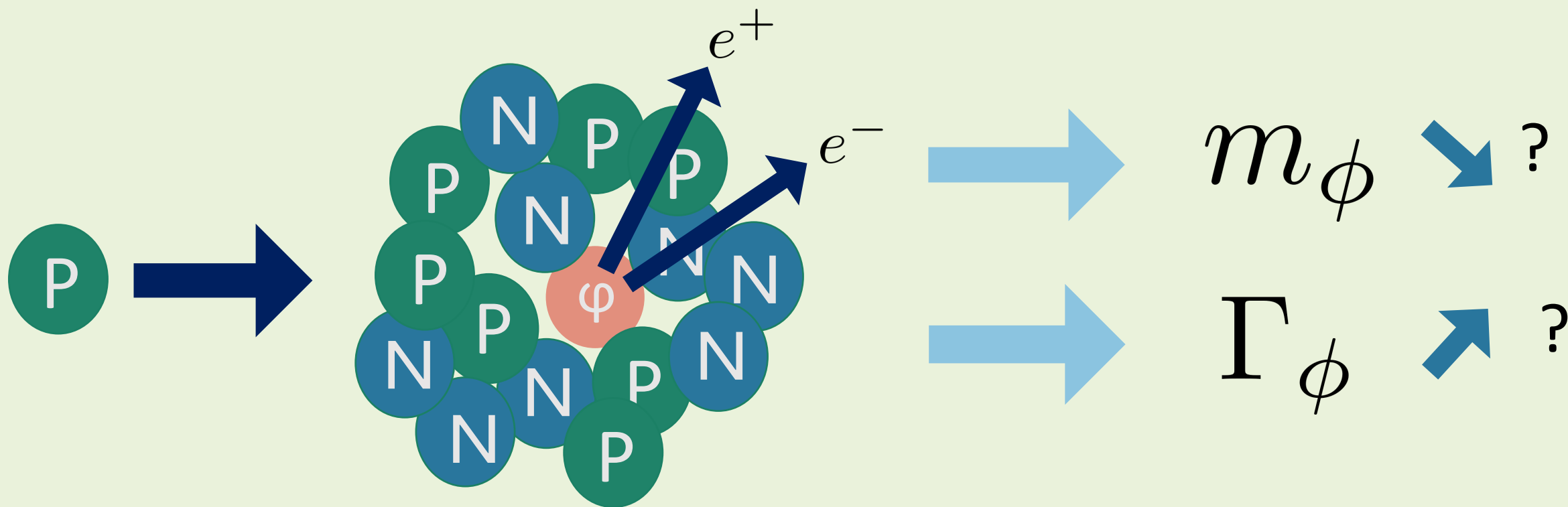
I. Filikhin et al.,
Phys. Rev. C **110**, 065202 (2024).

Lattice QCD (HAL)

Y. Liu et al.,
Phys. Rev. D **106**, 074507 (2022).

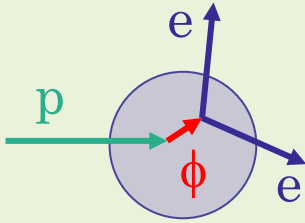


Dilepton spectrum from pA reactions (studied in a transport approach)



Previous experimental results

KEK
E325



12 GeV
pA-reaction

slow ϕ s

Pole mass:

$$\frac{m_\phi(\rho)}{m_\phi(0)} = 1 - k_1 \frac{\rho}{\rho_0}$$

0.034 ± 0.007

intermediate
 ϕ s

Pole width:

$$\frac{\Gamma_\phi(\rho)}{\Gamma_\phi(0)} = 1 + k_2 \frac{\rho}{\rho_0}$$

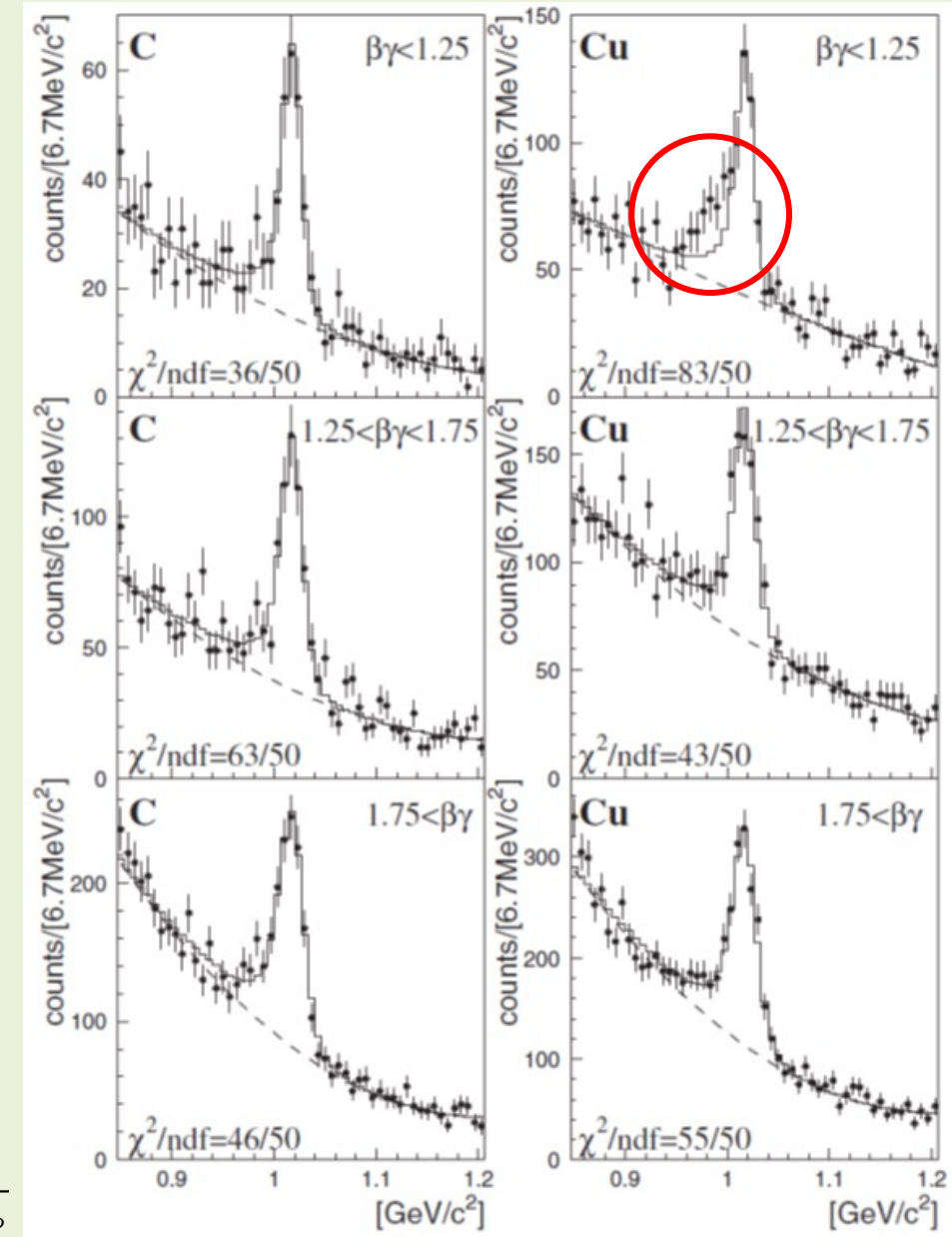
2.6 ± 1.5



Measurement is being repeated with
~100x increased statistics at the
J-PARC E16 experiment!

fast ϕ s

$$\beta\gamma = \frac{|\vec{p}|}{m_\phi}$$



Distortion of the target nucleus due to the incoming proton?

Not included in previous analysis!

p

density much below ρ_0

large probability of vector meson decay outside of the nucleus

Off-shell behavior of vector mesons at finite density

Not included in previous analysis!



A state-of-the-art transport simulation code: PHSD (Parton Hadron String Dynamics)

E.L. Bratkovskaya and W. Cassing, Nucl. Phys. A **807**, 214 (2008).

W. Cassing and E.L. Bratkovskaya, Phys. Rev. C **78**, 034919 (2008).



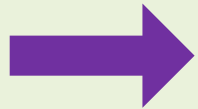
Equations of motion of all particles participating in the reaction are solved



Motion/deformation of target is taken into account!



Off-shell dynamics of vector mesons and kaons is included



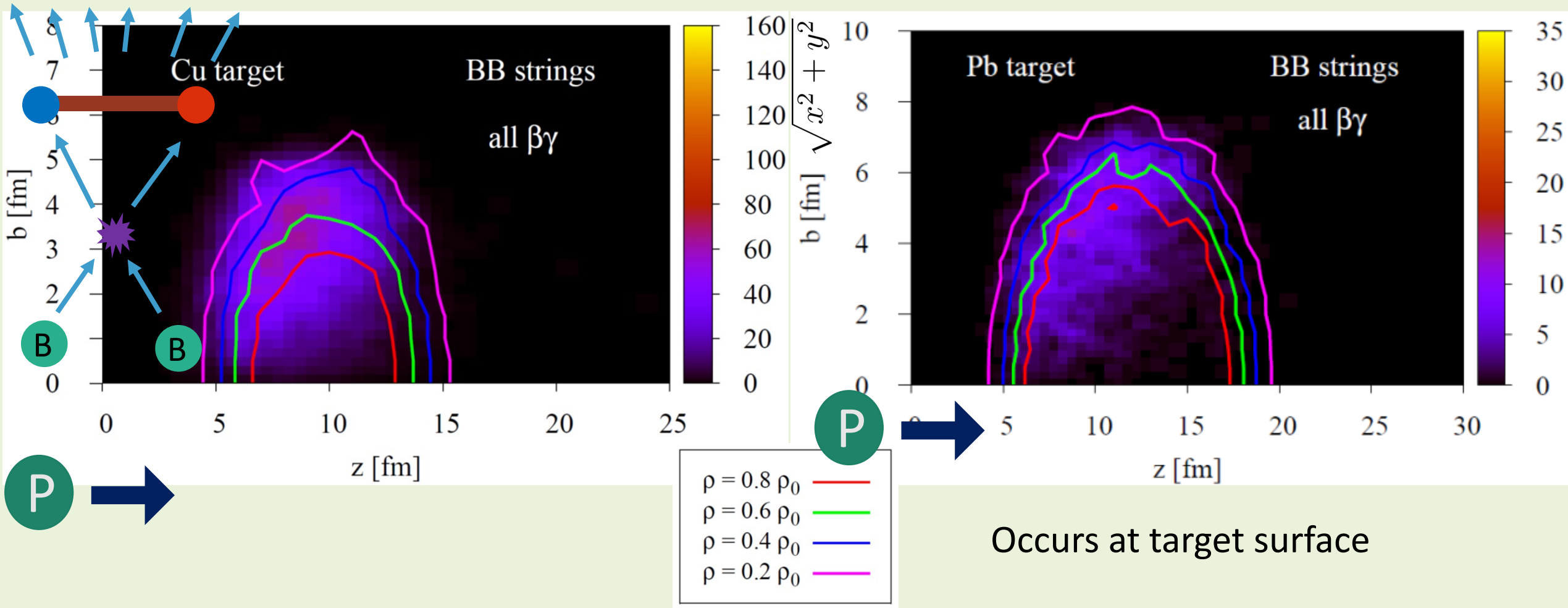
We use a relativistic Breit-Wigner with density dependent mass and width

$$C \frac{2}{\pi} \frac{M^2 \Gamma_{\phi}^*(M, \rho)}{[M^2 - M_{\phi}^{*2}(\rho)]^2 + M^2 \Gamma_{\phi}^{*2}(M, \rho)}$$

$$\text{with } \begin{cases} M_{\phi}^*(\rho) = M_{\phi}^{\text{vac}} \left(1 - \alpha^{\phi} \frac{\rho}{\rho_0}\right), \\ \Gamma_{\phi}^*(M, \rho) = \Gamma_{\phi}^{\text{vac}} + \alpha_{\text{coll}}^{\phi} \frac{\rho}{\rho_0} \end{cases}$$

How are φ mesons produced in 12 GeV pA collisions?

Production through initial high-energy collisions (via strings)

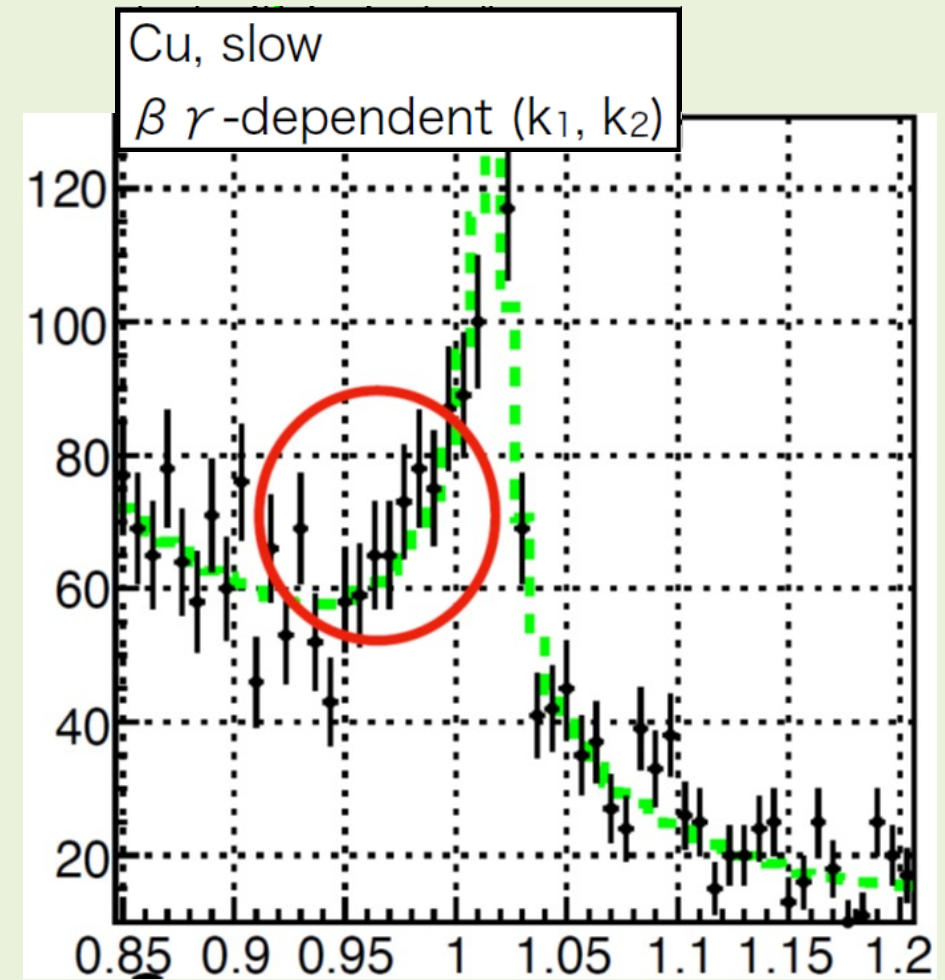


Comparison with KEK E325 dilepton data

Momentum-independent
mass shift and decay width



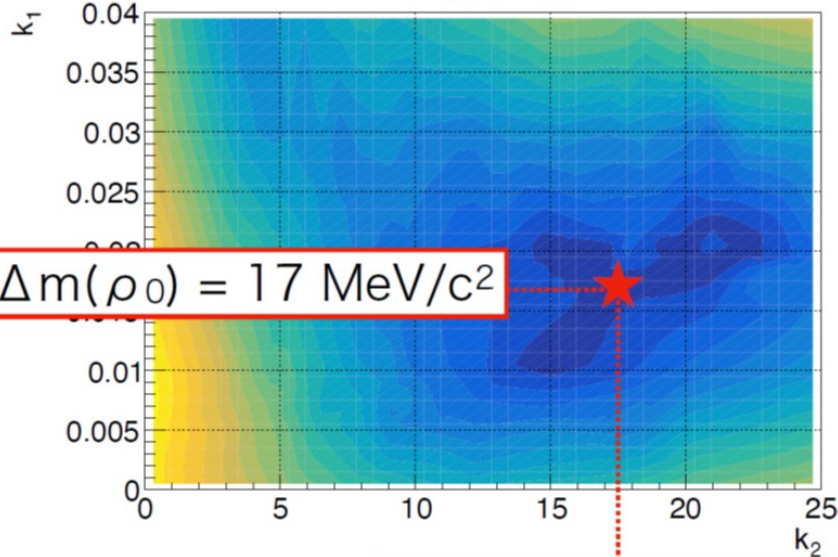
Momentum-dependent
mass shift and decay width



Comparison with KEK E325 dilepton data

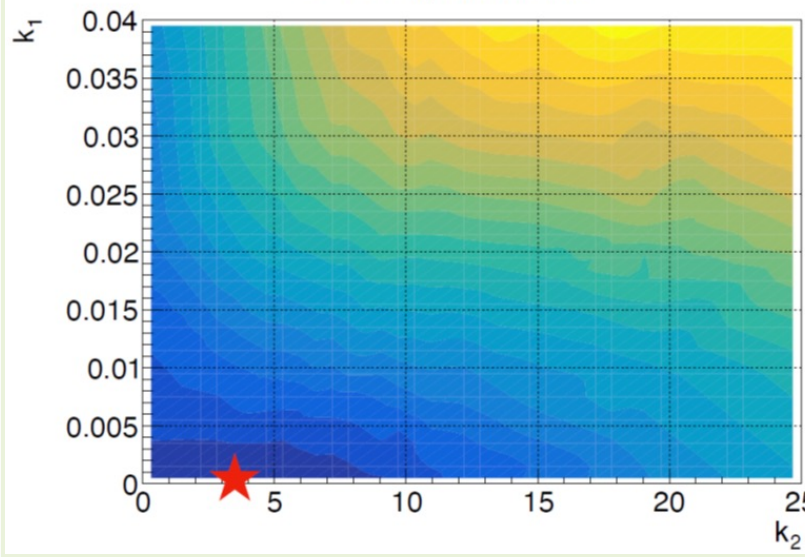
Momentum-dependent mass shift and decay width

C+Cu, slow

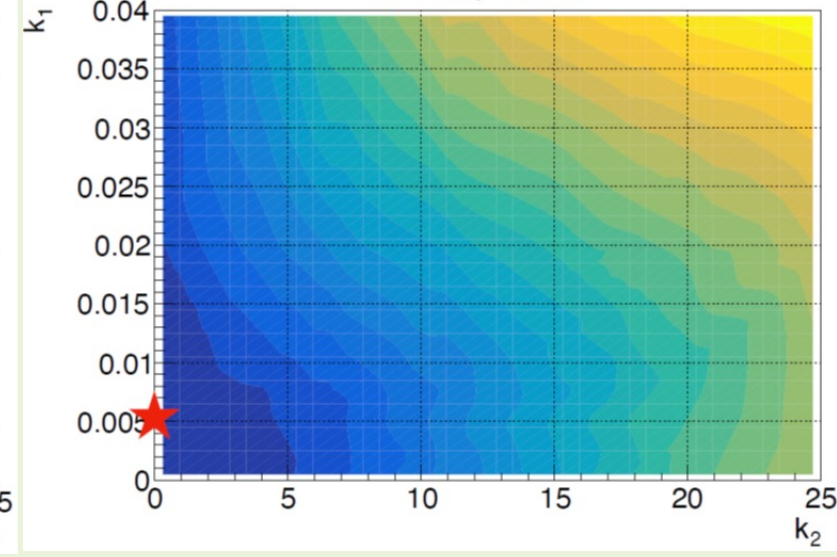


$$\Gamma(\rho_0) = 79 \text{ MeV}/c^2$$

C+Cu, middle



C+Cu, fast



Case of density dependent partial decay width Γ^{ee}

$$\frac{m_\phi(\rho)}{m_\phi(0)} = 1 - k_1 \frac{\rho}{\rho_0}$$

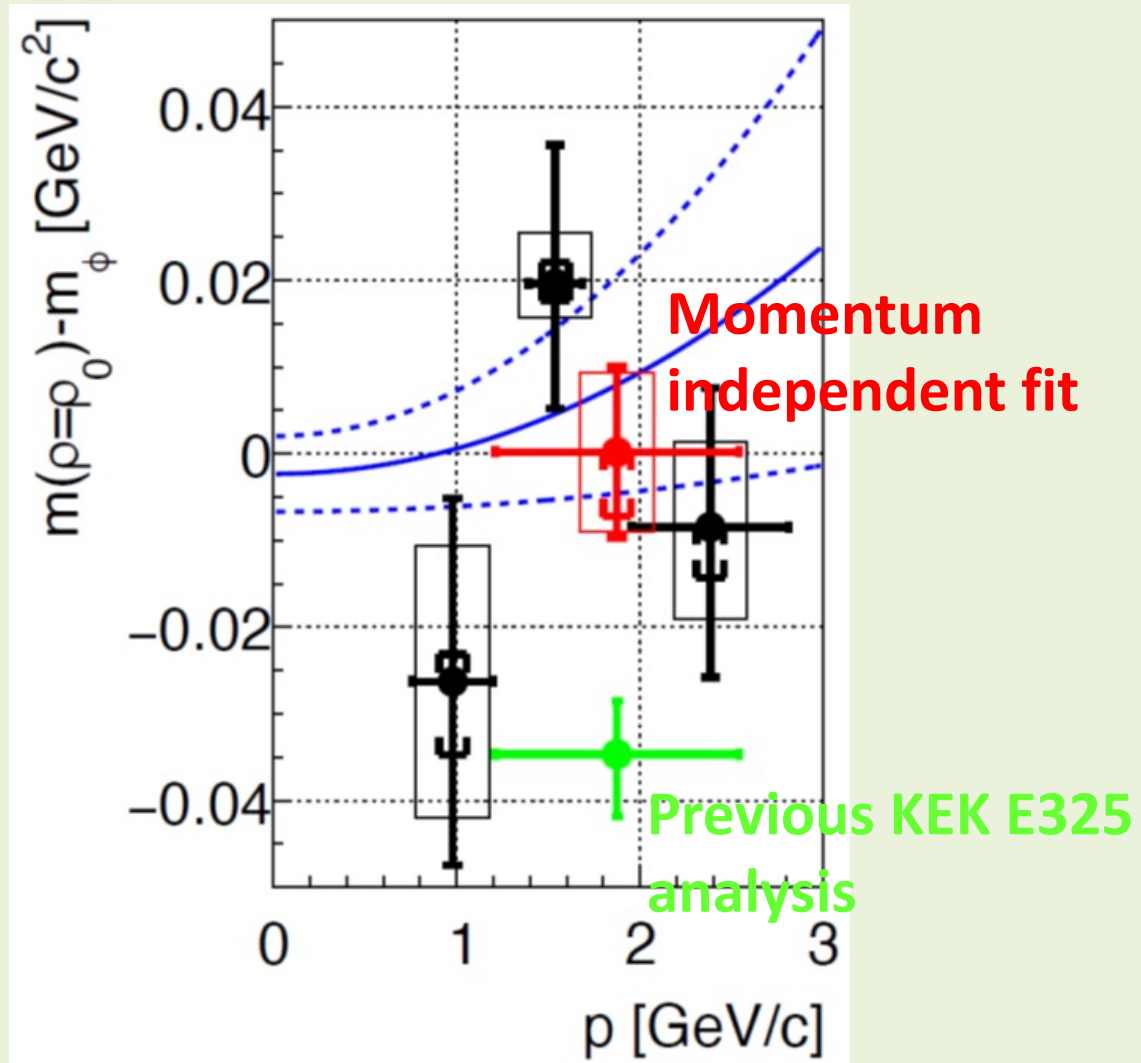
$$\frac{\Gamma_\phi(\rho)}{\Gamma_\phi(0)} = 1 + k_2 \frac{\rho}{\rho_0}$$

Momentum dependence is needed to explain the data!

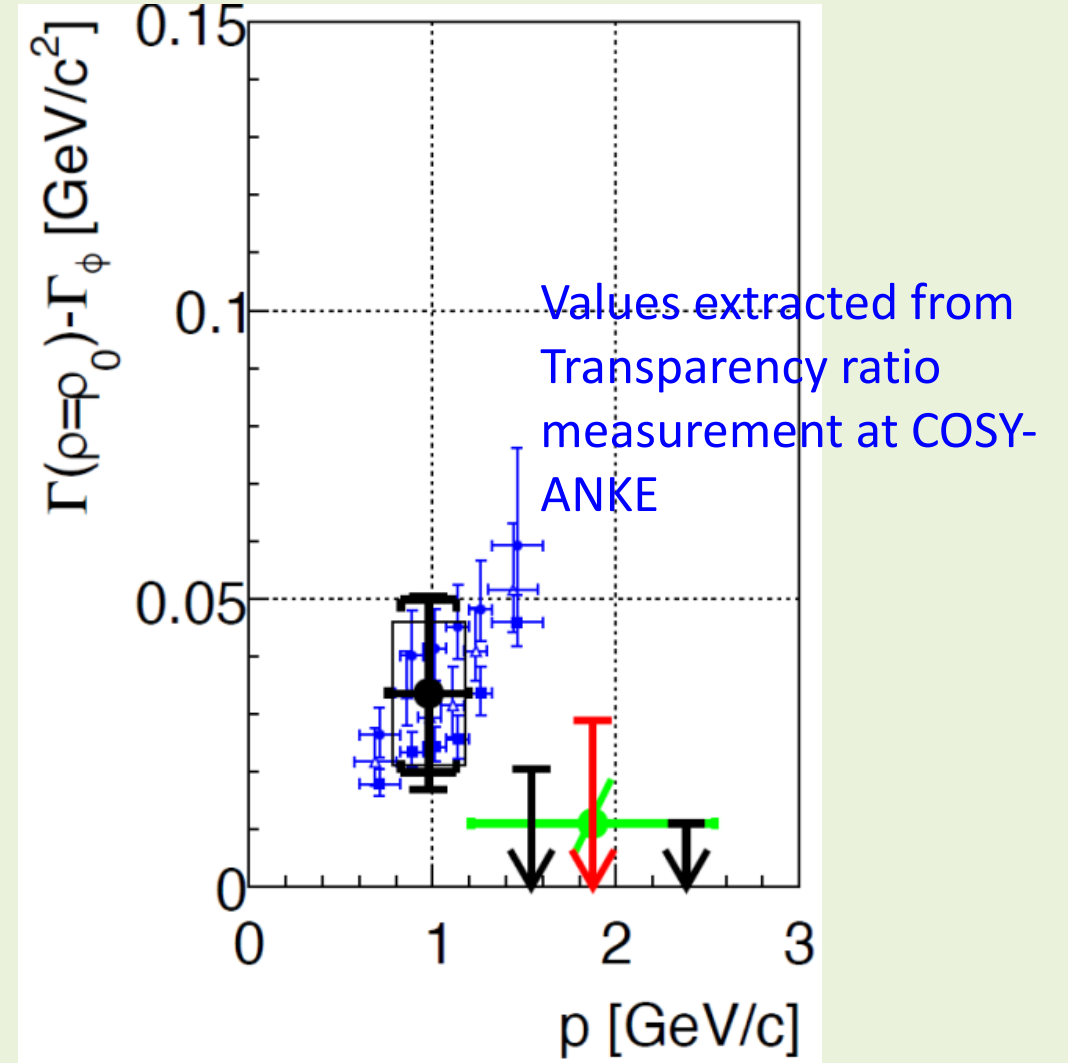
Smaller mass shift and stronger broadening than in the original E325 analysis!

Final analysis results

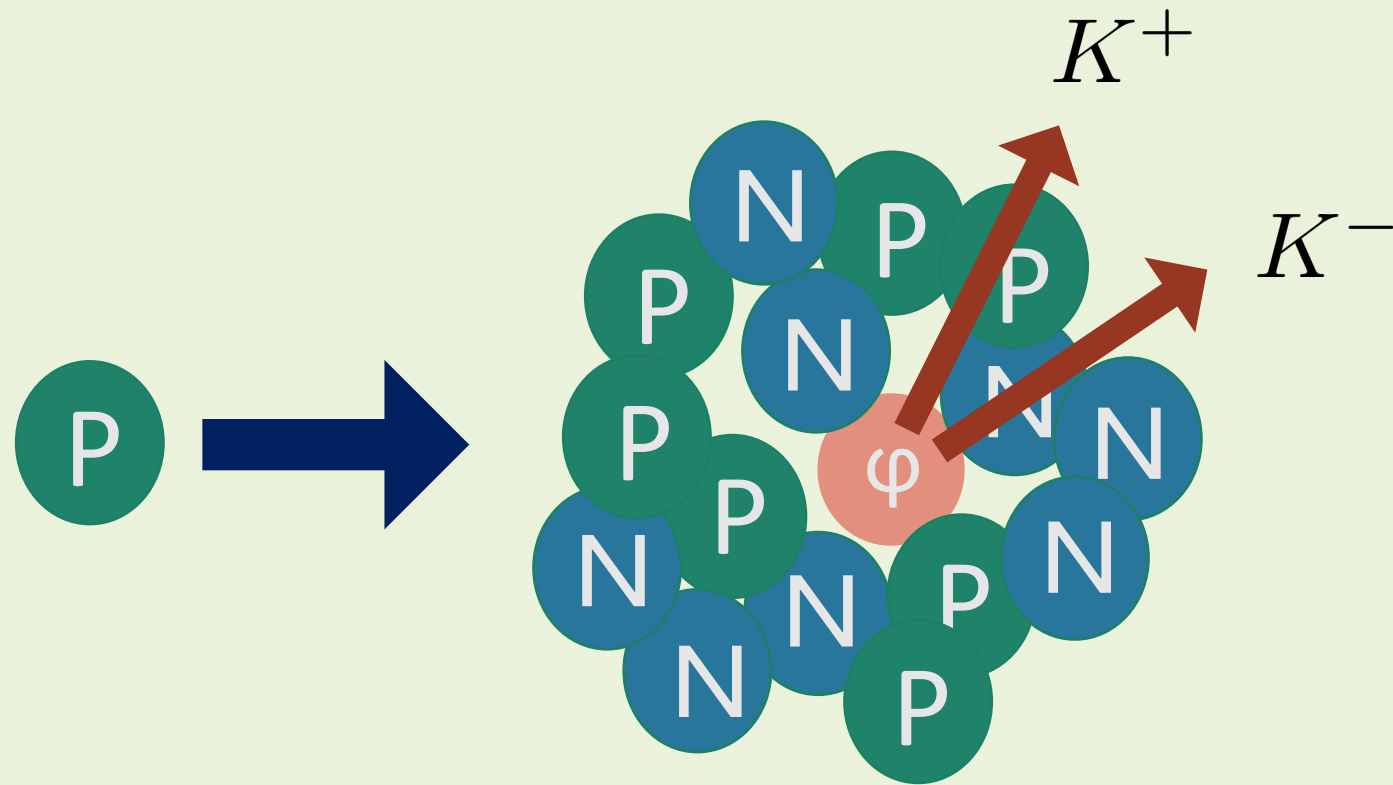
Mass shift at ρ_0



Width shift at ρ_0



What about the K^+K^- spectrum?

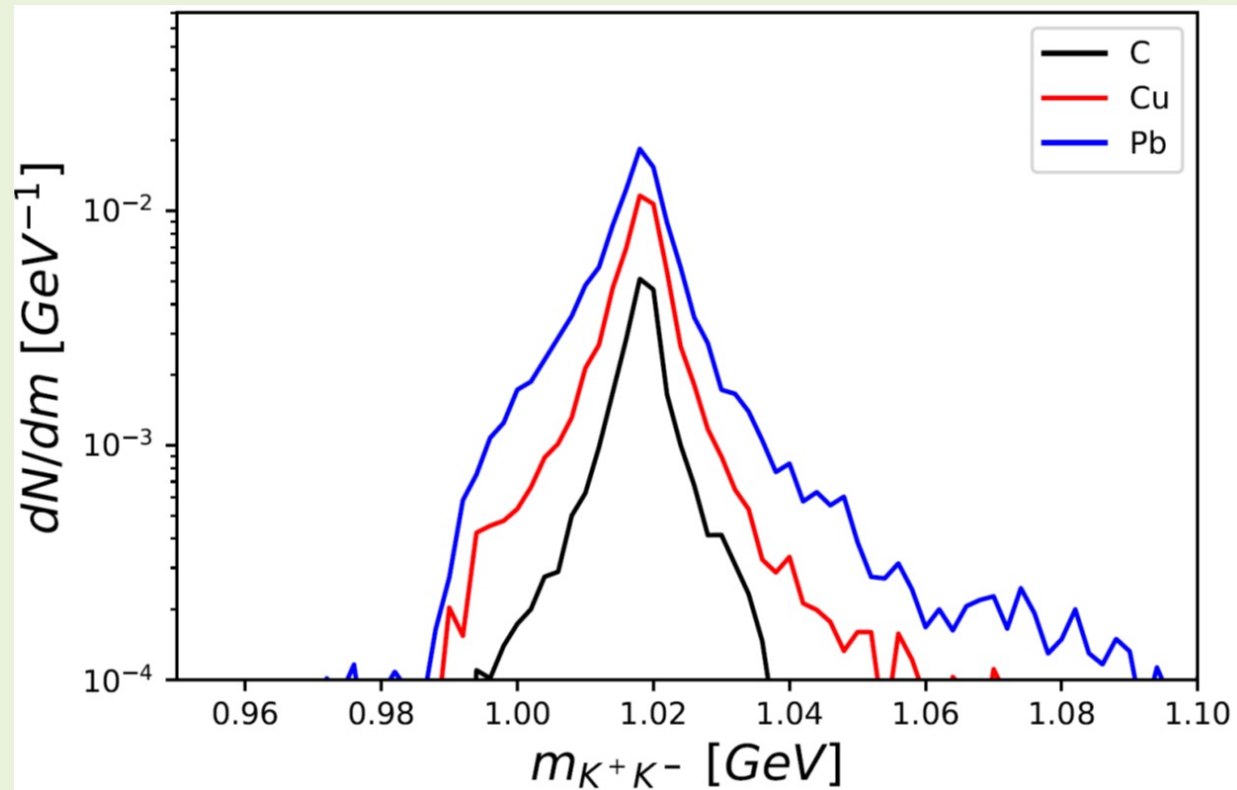


Will be studied at the J-PARC E88/SAPHIRE experiment!

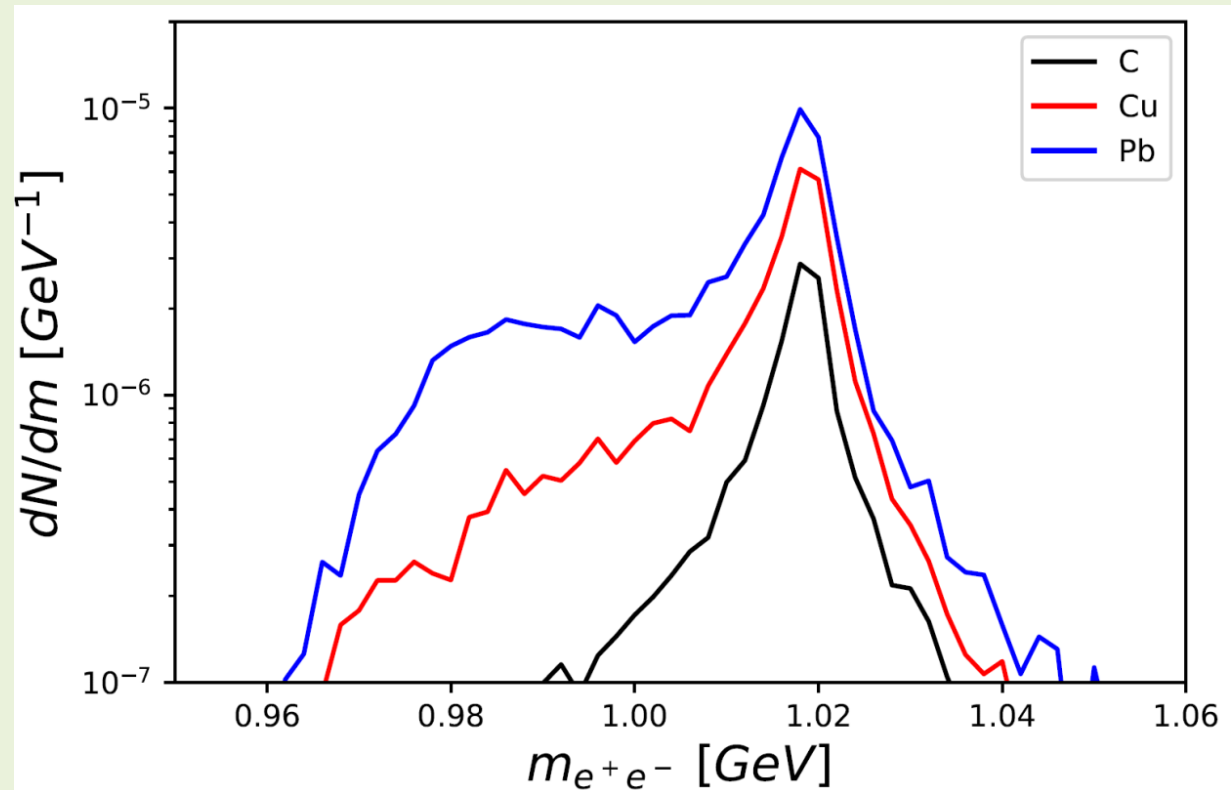
Comparative study of the K^+K^- and e^+e^- spectrum

(under the conditions of the J-PARC E16/E88 experiments:
30 GeV pA induced reactions)

K^+K^- invariant mass spectrum



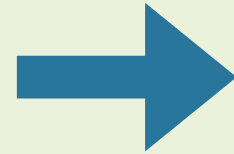
e^+e^- invariant mass spectrum



➡ Smaller effect for K^+K^- because of the closeness of its threshold to the ϕ meson mass

Summary and conclusions

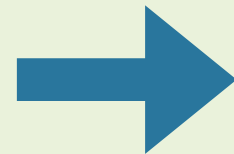
- ★ With the state-of-the-art PHSD transport approach, we can now study pA reactions more reliably



**Changed conclusions
from previous analysis of
experimental data**

- **Smaller mass shift**
- **Larger broadening**

- ★ Uncertainties are larger than in previous analysis

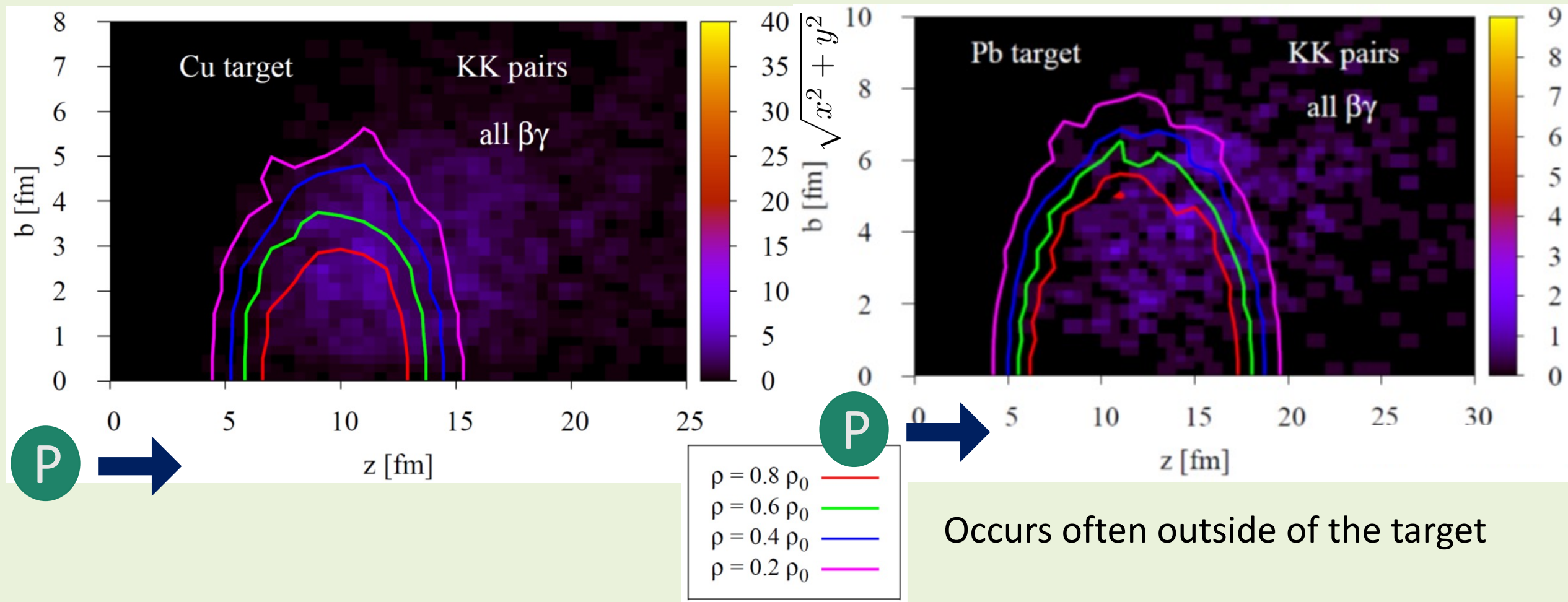


**More precise data from
the J-PARC E16 experiment
will be needed**

Backup slides

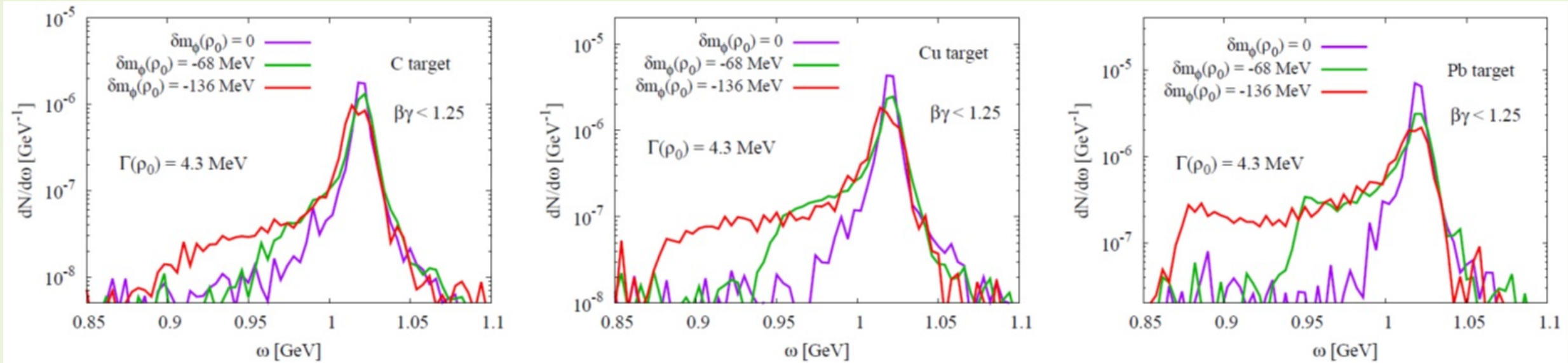
How are ϕ mesons produced?

Production through secondary low-energy hadron collisions



The obtained dilepton spectrum (without experimental effects)

Pure mass shift scenarios (no broadening)



No second peak, but only shoulder structure for mass shift scenarios (even before considering experimental resolution effects)

Secondary peak can be generated for sufficient large mass shift scenario if the target is large enough (Pb here)

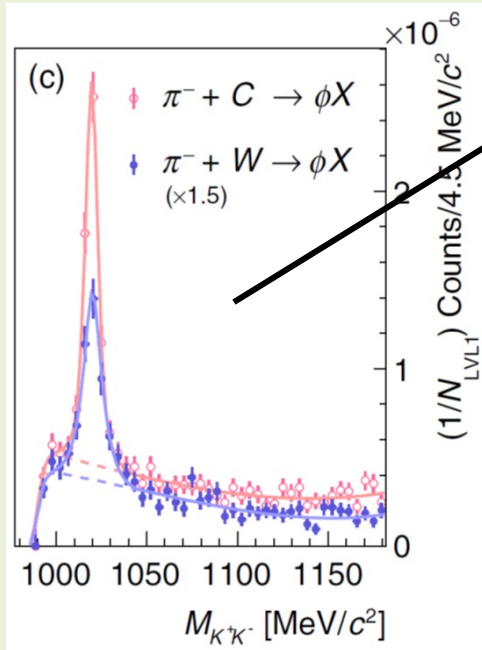
Issues of transport approaches

- ★ Reliability of transport simulations (dependence on different codes and approaches)
 - ➡ Need comparison of different codes (including QMC and BUU types)
- ★ Stability of the (undisturbed) target in the transport approach
 - ➡ Generally difficult to achieve in most transport models specialized in high-energy collisions, but is realized in some low-energy codes
 - ➡ Not so serious for KEK and J-PARC pA collisions, in which the phi mesons leave the target quickly??
- ★ Violation of energy or momentum conservation when mass shifted particles leave the dense medium
 - ➡ Relevant for E88, as kaons generated in dense matter are mass shifted

Many inconsistencies...

HADES: 1.7 GeV π^- A-reaction

K^+K^- - invariant mass spectrum



Theoretical analysis of the of the total ϕ meson production cross section:

★ **Attractive ϕ -nucleus potential:**
-(50 - 100) MeV

★ **Small imaginary part:**
20 – 25 MeV

E. Ya. Paryev, Nucl. Phys. A **1032**, 122624 (2023).

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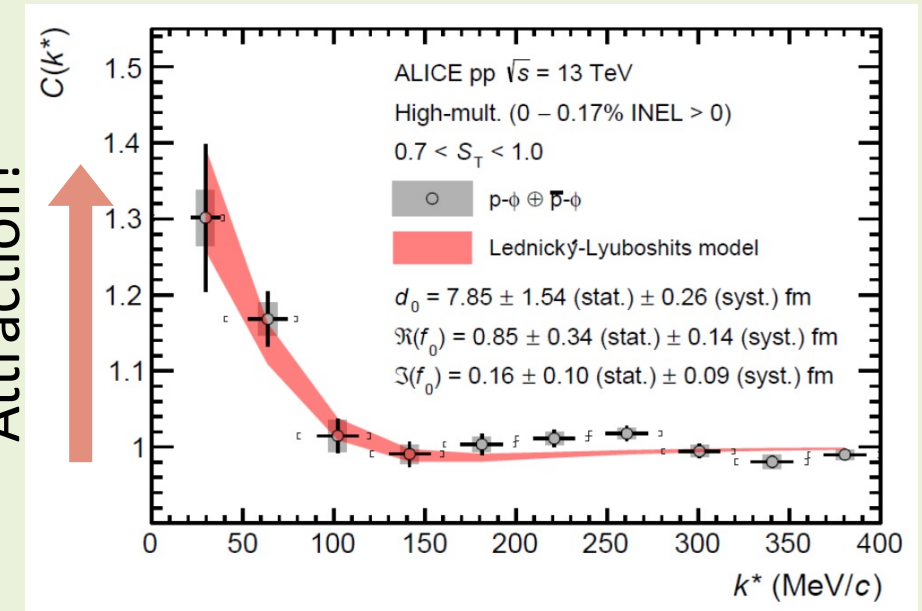
Large negative mass shift?
Small broadening?

★ Photoproduction measurements
I.I. Strakovsky et al. (CLAS),
Phys. Rev. C **101**, 045201 (2020).
 $|a_0| = 0.063 \pm 0.010$ fm

★ Hadronic Effective theory calculations

ALICE: pp

Measurement of ϕ N correlation



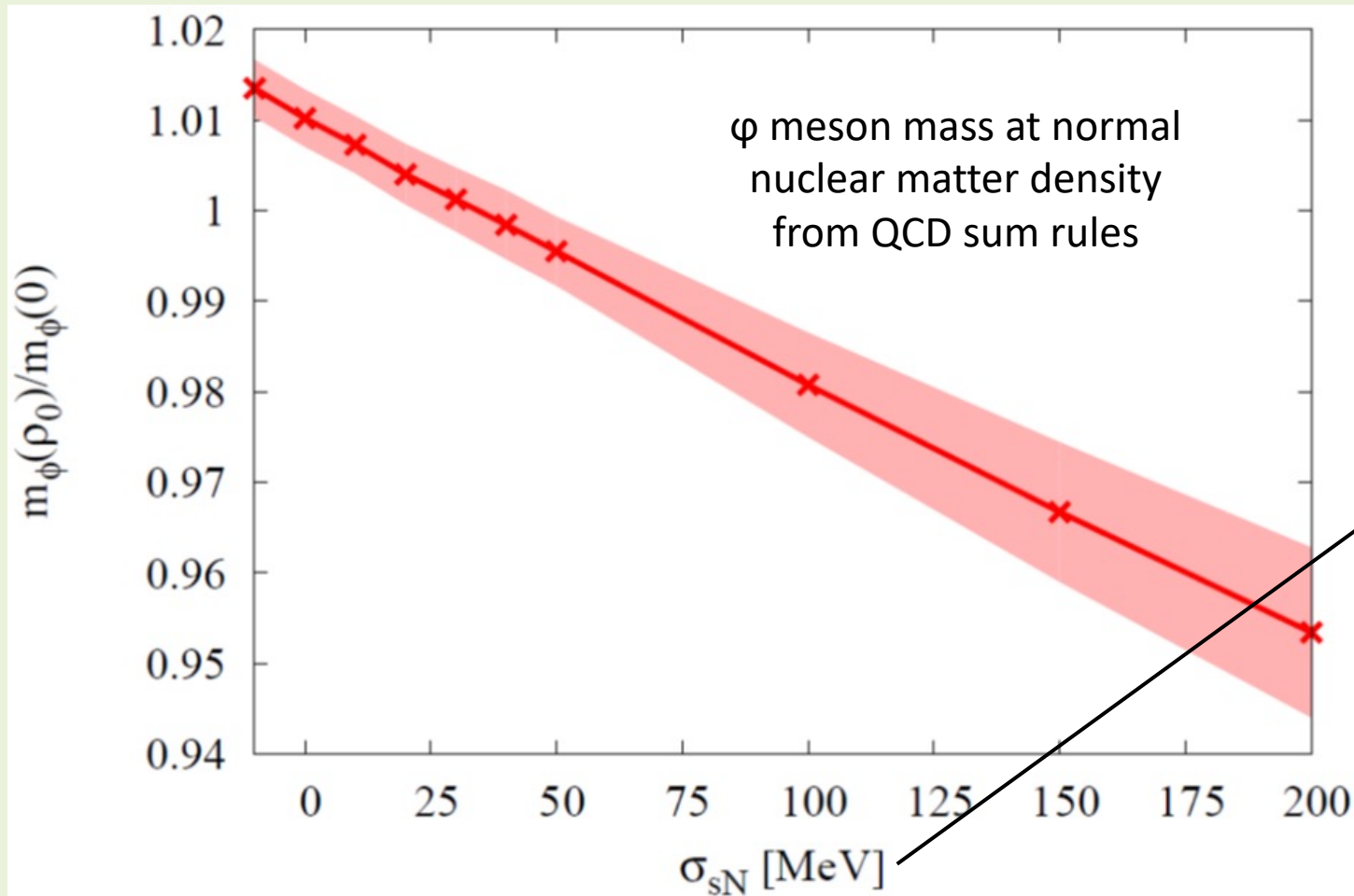
S. Acharya et al. (ALICE Coll.),
Phys. Rev. Lett. **127**, 172301 (2021).

Y. Lyu et al. (Lattice QCD, HAL QCD Collaboration),
Phys. Rev. D **106**, 074507 (2022).

$\rightarrow a_0^{3/2} = 1.43(23)_{\text{stat.}} \left({}^{+36}_{-06} \right)_{\text{syst.}} \text{ fm}$

Why should we be interested?

The ϕ meson mass in nuclear matter probes the strange quark condensate at finite density!



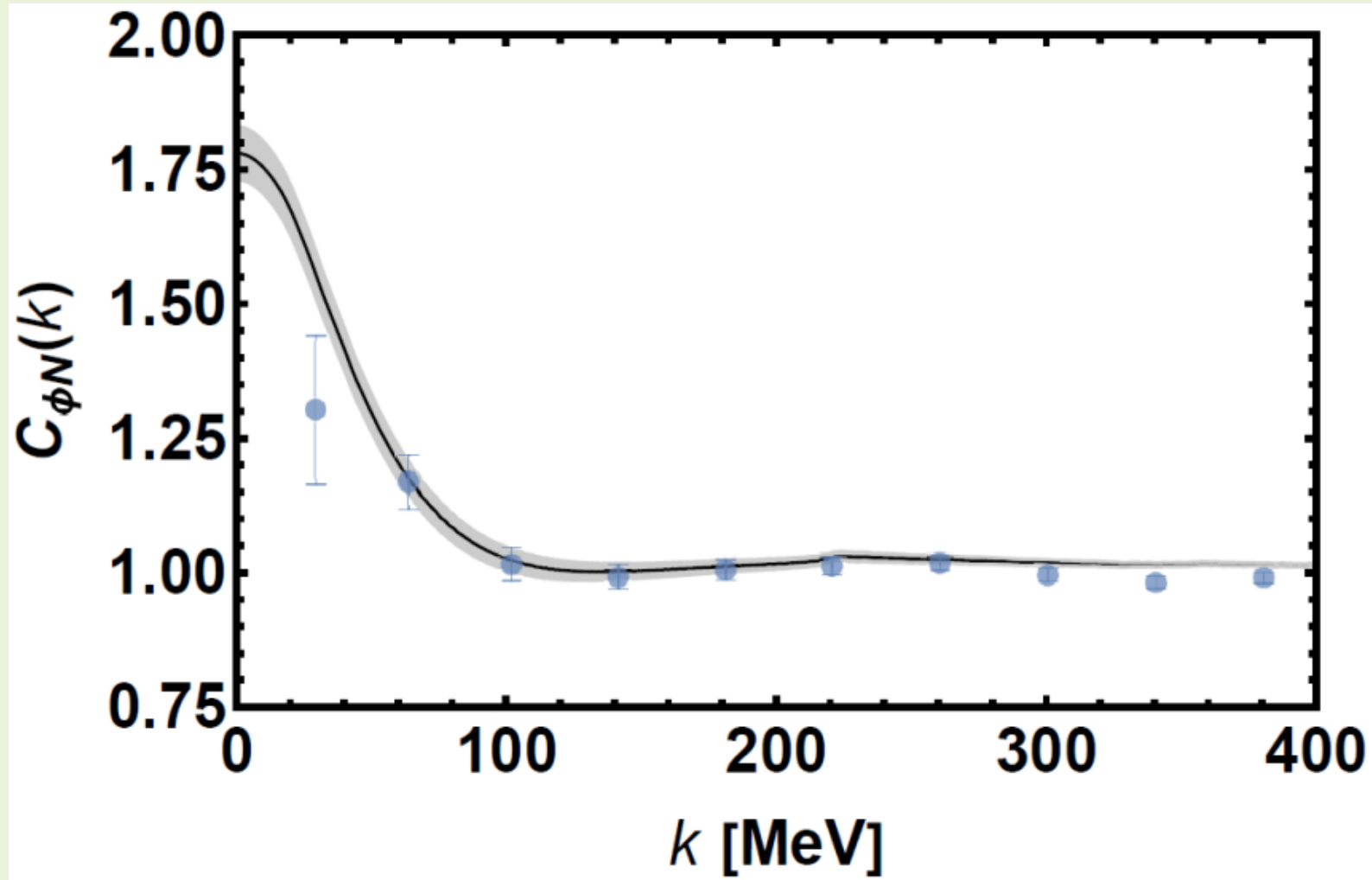
T. Hatsuda and S.H. Lee,
Phys. Rev. C **46**, R34 (1992).

$$\sigma_{sN} = m_s \langle N | \bar{s}s | N \rangle$$

$$|\langle \bar{s}s \rangle_\rho| = |\langle \bar{s}s \rangle_0| - \frac{\rho}{m_s} \sigma_{sN} + \dots$$

P. Gubler and K. Ohtani, Phys. Rev. D **90**, 094002 (2014).

The obtained correlation function (compared with ALICE data)



Reasonably good agreement
without any parameter fitting!

Simple relation between ϕ N scattering length and ϕ meson mass shift in nuclear matter

$$\begin{aligned} V_\phi(\rho) &= -\frac{2\pi}{m_\phi} \rho \left(1 + \frac{m_\phi}{m_N}\right) a_0 \\ &\simeq -85 \frac{\rho}{\rho_0} \left(\frac{a_0}{\text{fm}}\right) \text{MeV} \end{aligned} \quad \left. \vphantom{\begin{aligned} V_\phi(\rho) &= -\frac{2\pi}{m_\phi} \rho \left(1 + \frac{m_\phi}{m_N}\right) a_0 \\ &\simeq -85 \frac{\rho}{\rho_0} \left(\frac{a_0}{\text{fm}}\right) \text{MeV} \end{aligned}} \right\} \begin{array}{l} \text{Valid within the linear} \\ \text{density approximation} \end{array}$$

Larger than 100 MeV IF HAL QCD result is true for all spin configurations!

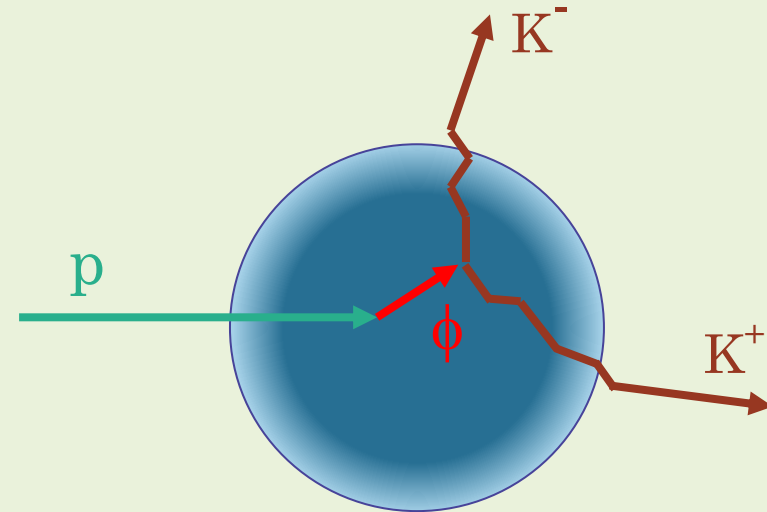
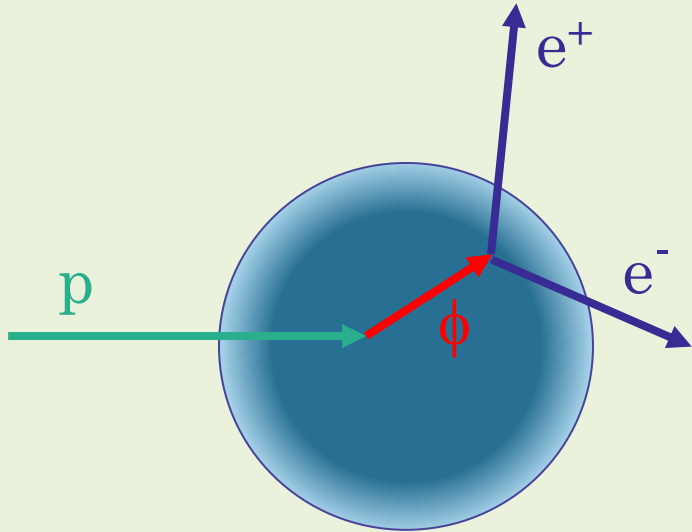
However, the above prescription seems problematic if a ϕ N bound state (or resonance) is formed.



Need better theoretical understanding!

Further tasks for theory

Have a good understanding of the production mechanisms of the ϕ mesons in nuclei from pA reactions.



- ★ Where (and at what densities) is the ϕ meson produced and where does it decay?
- ★ How do the final state interactions of the decay particles influence the decay spectrum (especially for K^+K^-)?



Realistic transport simulations using a transport approach (calculations using the PHSD code are ongoing)



See talk by L. Oliva

ϕ meson at rest in nuclear matter

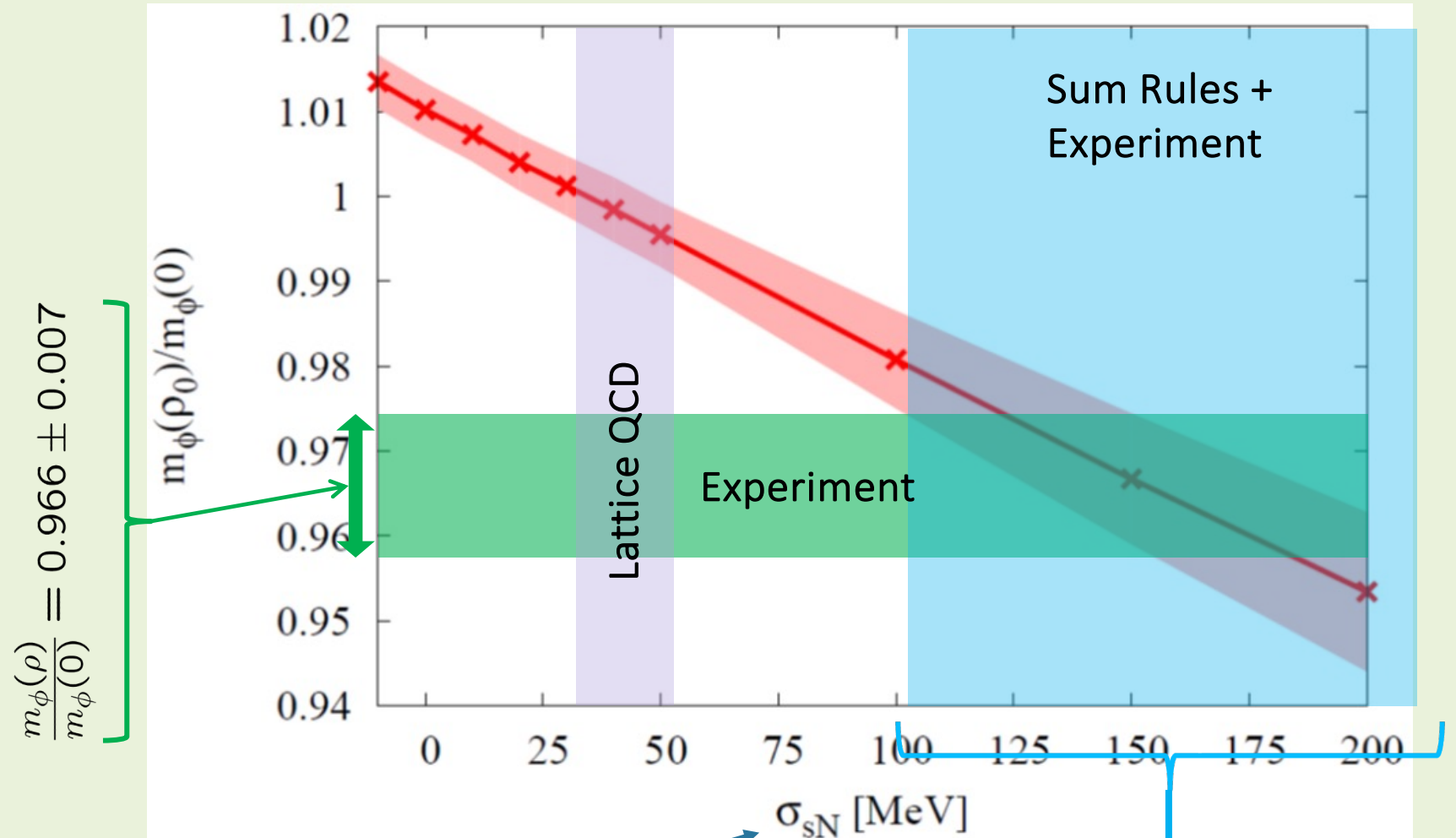
The ϕ meson mass in nuclear matter probes the strange quark condensate at finite density!

Not
consistent?

R. Muto et al.
(KEK, E325 Collaboration),
Phys. Rev. Lett. **98**,
042501 (2007).



Measurement will be
repeated at the
J-PARC E16 experiment
(with 100 times
increased statistics!)



$$\sigma_{sN} = m_s \langle N | \bar{s}s | N \rangle$$

$$\sigma_{sN} \sim 160 \pm 50 \text{ MeV}$$

Experimental di-lepton spectrum

