

Antikaon absorption in the nuclear medium: the role of hadron self-energies and implications for kaonic atoms

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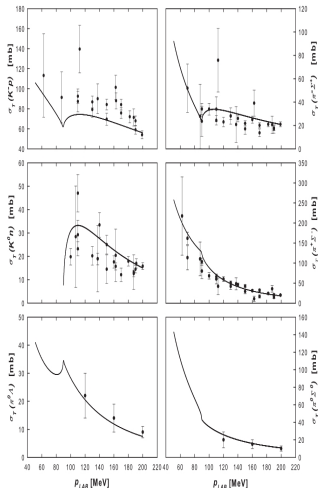
NPI, Řež

15th International Conference on Hypernuclear and Strange Particle Physics,
HYP2025,

29 September - 3 October 2025, Tokyo, Japan

K^-N interactions

K^-p scattering:



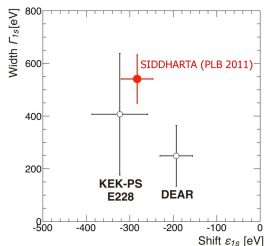
Threshold branching ratios:

$$\gamma = \frac{\Gamma(K^-p \rightarrow \pi^+\Sigma^-)}{\Gamma(K^-p \rightarrow \pi^-\Sigma^+)} = 2.36 \pm 0.04$$

$$R_c = \frac{\Gamma(K^-p \rightarrow \text{charged})}{\Gamma(K^-p \rightarrow \text{all})} = 0.664 \pm 0.011$$

$$R_n = \frac{\Gamma(K^-p \rightarrow \pi^0\Lambda)}{\Gamma(K^-p \rightarrow \text{neutral})} = 0.189 \pm 0.015$$

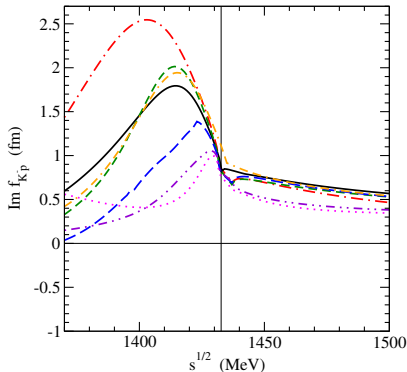
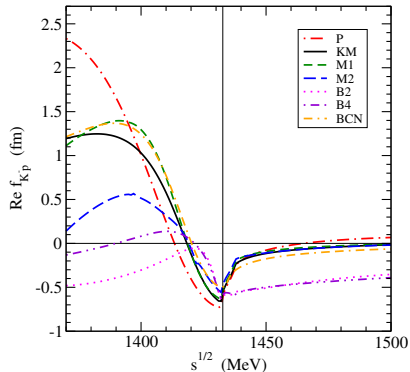
Kaonic hydrogen:



- K^-p data well described by several chiral models

A. Cieplý, M. Mai, U.-G. Meissner, J. Smejkal, NPA 945 (2016) 17

Free-space K^-p amplitudes in various chiral models



Prague (P)

Kyoto-Munich (KM)

Murcia (M1 and M2)

Bonn (B2 and B4)

Barcelona (BCN)

A. Cieplý, J. Smejkal, *Nucl. Phys. A* 881 (2012) 115

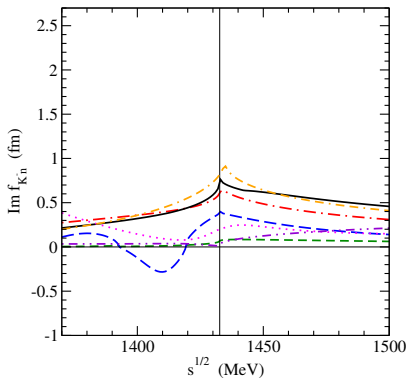
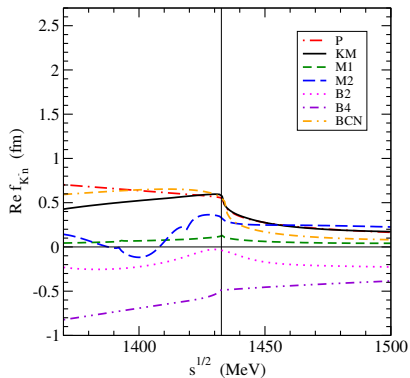
Y. Ikeda, T. Hyodo, W. Weise, *Nucl. Phys. A* 881 (2012) 98

Z. H. Guo, J. A. Oller, *Phys. Rev. C* 87 (2013) 035202

M. Mai, U.-G. Meißner, *Nucl. Phys. A* 900 (2013) 51

A. Feijoo, V. Magas, A. Ramos, *Phys. Rev. C* 99 (2019) 035211

Free-space K^-n amplitudes



- Measurement of kaonic deuterium by SIDDHARTA-2 Collaboration could provide new constrain on K^-N interaction in $I = 1$ channel
F. Sgaramella et al. [SIDDHARTA-2], NC C 47 (2024) 285

Kaonic atoms

- Info about K^-N interaction below threshold provided by kaonic atoms
65 data points (energy shifts, widths, yields=upper level widths)
from CERN, Argonne, RAL, BNL
- Global analysis of kaonic atom data using phenomenological potential
 $\rightarrow \text{Re} V_{K^-}(\rho_0) \sim -(150 - 200) \text{ MeV}$
E. Friedman, A. Gal, PR 452 (2007) 89
- K^- potential from chiral models $\rightarrow \text{Re} V_{K^-}(\rho_0) \sim -(40 - 60) \text{ MeV}$
A. Cieplý, E. Friedman, A. Gal, J. Mareš, NPA 696 (2001) 173
- Chirally motivated models fail to describe kaonic atom data
E. Friedman, A. Gal, NPA 959 (2017) 66

model	B2	B4	M1	M2	P	KM
$\chi^2(65)$	1174	2358	2544	3548	2300	1806

Multinucleon processes

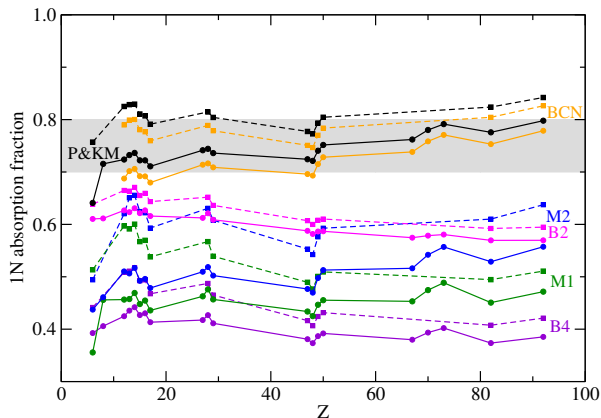
- Chiral models include only $K^- N \rightarrow \pi Y$ ($Y = \Lambda, \Sigma$) decay channel
- K^- interactions with two and more nucleons should be included, (e.g., $K^- + N + N \rightarrow Y + N$) $\leftarrow$ analysis of kaonic atom data
E. Friedman, A. Gal, NPA 959 (2017) 66

$$V_{K^- \text{ multiN}}^{\text{phen}} = -4\pi B \left(\frac{\rho}{\rho_0} \right)^\alpha \rho$$

B is a complex amplitude, ρ is nuclear density distribution, ρ_0 is saturation density and α is positive

- equally good description of data with $\chi^2/\text{dof} \leq 2$

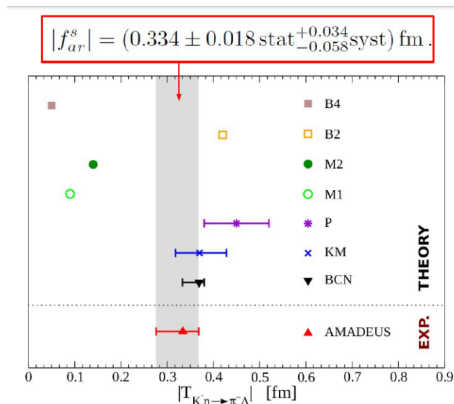
Single- vs. multi-nucleon processes



- Fraction of *single-nucleon* absorption 0.75 ± 0.05 (average value, gray band) used as an **additional constraint**.

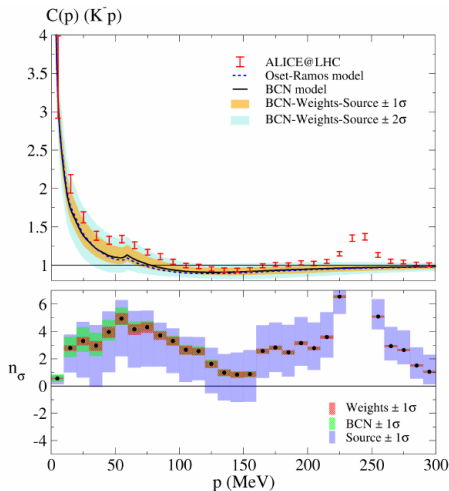
→ Only **P**, **KM** and **BCN** models found acceptable in kaonic atom analysis

E. Friedman, A. Gal, NPA 959 (2017) 66



K^-p correlation function from ALICE

P. Encarnación et al., PRD 111 (2025) 11



Multinucleon processes

- K^- multi-nucleon absorption in the surface region of atomic nuclei represents about 20%
NC 53 (1968) 313 (Berkeley), NPB 35 (1971) 332 (BNL), NC 39A (1977) 538 (CERN)
- K^- multi-nucleon absorption in atoms described by phenomenological optical potential
E. Friedman, A. Gal, NPA 959 (2017) 66
- Model for $K^- NN$ absorption in nuclear matter using free-space chiral amplitudes
T. Sekihara et al., PRC 86 (2012) 065205
- New experimental data on $K^- NN$ absorption (AMADEUS@DAΦNE)
K. Piscicchia et al., PLB 782 (2018) 339
R. Del Grande et al., EPJ C79 (2019) 190
- Solid microscopic model for $K^- NN$ absorption needed!

Microscopic model for $K^- NN$ absorption in nuclear matter

Microscopic model for K^- two-nucleon absorption in symmetric nuclear matter *J. Hrtánková, Á. Ramos, PRC 101 (2020) 035204*

- based on a meson-exchange approach
H. Nagahiro et al., PLB 709 (2012) 87
- **P** and **BCN** chiral $K^- N$ amplitudes employed
- **Pauli correlations** in the medium for $K^- N$ amplitudes considered
- **real part of the $K^- NN$ optical potential** evaluated as well
- $K^- N$ optical potential derived within the same approach

K^-N absorption in nuclear matter

$$K^-N \rightarrow \pi Y \quad (Y = \Lambda, \Sigma)$$

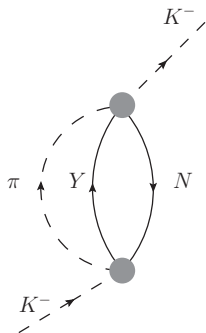


Fig.1: Feynman diagram for K^- absorption on a single nucleon in nuclear matter. The shaded circles denote the K^-N t-matrices derived from a chiral model.

$K^- NN$ absorption in nuclear matter

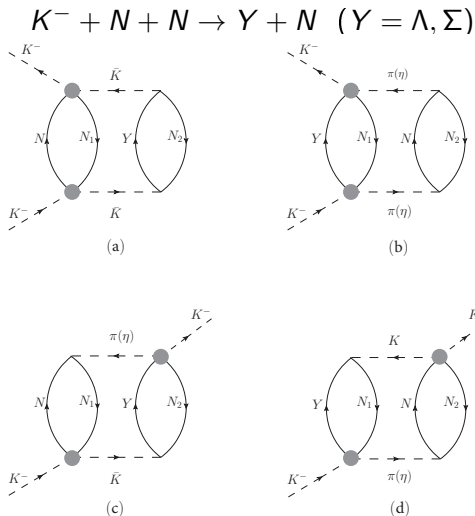


Fig.2: Two-fermion-loop (2FL) Feynman diagrams for non-mesonic K^- absorption on two nucleons N_1 , N_2 in nuclear matter. The shaded circles denote the $K^- N$ t-matrices derived from a chiral model.

K^-NN absorption in nuclear matter

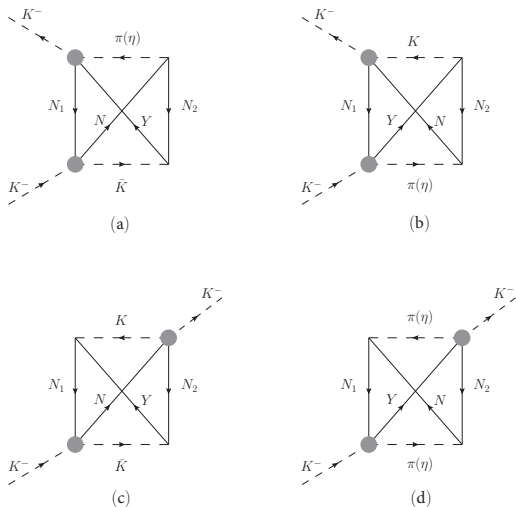


Fig.3: One-fermion-loop (1FL) Feynman diagrams for non-mesonic K^- absorption on two nucleons N_1 , N_2 in nuclear matter. The shaded circles denote the K^-N t-matrices derived from a chiral model.

$K^- NN$ absorption in nuclear matter

$$V_{K^-N} = \sum_{\text{channels}} V_{K^-N \rightarrow \pi Y} \text{ (Fig.1)}$$

$$V_{K^-NN} = \sum_{\text{channels}} V_{K^-NN}^{2\text{FL}} + V_{K^-NN}^{1\text{FL}} \text{ (Fig.2 and 3)}$$

→ contributions from 37 2FL and 28+33 1FL diagrams

Table 1: All considered channels for mesonic and non-mesonic K^- absorption in matter.

$K^- N$	$\rightarrow \pi Y$	$K^- N_1 N_2$	$\rightarrow YN$
$K^- p$	$\rightarrow \pi^0 \Lambda$	$K^- pp$	$\rightarrow \Lambda p$
	$\rightarrow \pi^0 \Sigma^0$		$\rightarrow \Sigma^0 p$
	$\rightarrow \pi^+ \Sigma^-$		$\rightarrow \Sigma^+ n$
	$\rightarrow \pi^- \Sigma^+$	$K^- pn(np)$	$\rightarrow \Lambda n$
$K^- n$	$\rightarrow \pi^- \Lambda$		$\rightarrow \Sigma^0 n$
	$\rightarrow \pi^- \Sigma^0$		$\rightarrow \Sigma^- p$
	$\rightarrow \pi^0 \Sigma^-$	$K^- nn$	$\rightarrow \Sigma^- n$

AMADEUS: Ratio for 2N absorption

R. Del Grande et al., EPJ C79 (2019) 190

$$R = \frac{\text{BR}(K^- pp \rightarrow \Lambda p)}{\text{BR}(K^- pp \rightarrow \Sigma^0 p)} = 0.7 \pm 0.2(\text{stat.})_{-0.3}^{+0.2}(\text{syst.})$$

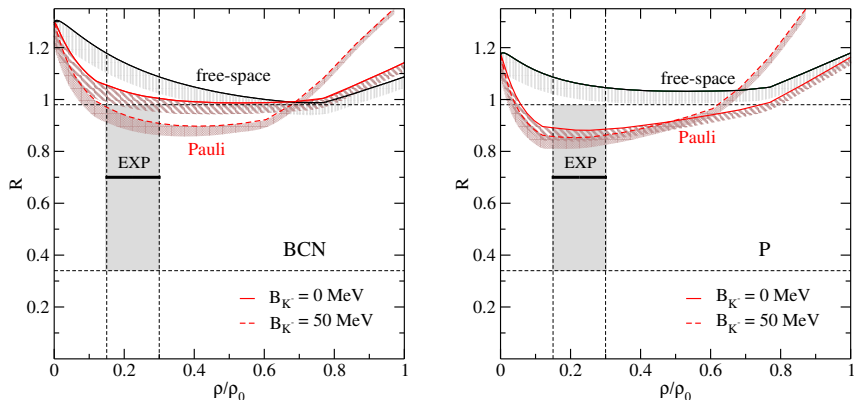


Fig.4: The ratio R as a function of relative density, calculated using the free-space and Pauli blocked amplitudes for $B_{K^-} = 0$ MeV and $B_{K^-} = 50$ MeV. Color bands denote the uncertainty due to different cut-off values $\Lambda_c = 800 - 1200$ MeV.

Application of K^-NN model to kaonic atoms

J. Óbertová, E. Friedman, J. Mareš, PRC 106 (2022) 065201

Table 2: Values of $\chi^2(65)$ resulting from comparisons of predictions with kaonic atom data using K^-N , $K^-N + K^-NN$, and $K^-N + K^-NN + \text{phen. multiN}$ potentials. Values of complex amplitude B and parameter α for potentials based on Pauli blocked and WRW (*T. Wass, M. Rho, W. Weise, NPA 617 (1997) 449*) modified BCN amplitudes.

	K^-N	$K^-N + K^-NN$	+ phen.	$\text{Re}B$ (fm)	$\text{Im}B$ (fm)	α
Pauli	825	565	105	-1.97(13)	-0.93(11)	1.4
WRW	2378	1123	116	-0.90(9)	0.72(10)	0.6

- best fit $K^-N + \text{phen. multiN}$ potential based on BCN amplitudes
 $\text{Re}B = -1.3 \text{ fm}$, $\text{Im}B = 1.9 \text{ fm}$, $\alpha = 1$, $\chi^2(65) = 112.3$

The role of hadron self-energies

- inclusion of Y , N , K^- and π self-energies into the K^-N chiral BCN amplitudes as well as into the K^-NN model.
- considered baryon potentials

$$V_N = -70 \frac{\rho}{\rho_0} \text{ MeV},$$

$$V_\Lambda = -58 \frac{\rho}{\rho_0} + 31 \left(\frac{\rho}{\rho_0} \right)^2 \text{ MeV} \rightarrow V_\Lambda(\rho_0) = -26.4 \text{ MeV},$$

$$V_\Sigma = 30 \frac{\rho}{\rho_0} \text{ MeV}.$$

- considered pion S- and P-wave self-energy

A. Ramos, E. Oset, NPA 671 (2000) 481

The role of hadron self-energies: K^-p amplitudes

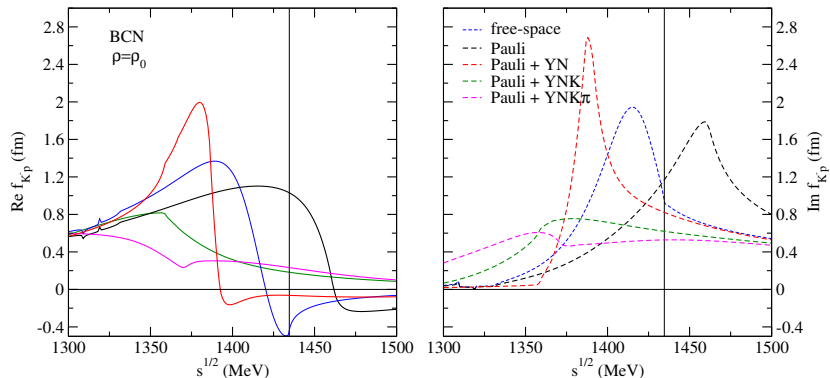


Fig.5: Comparison of $K^-p \rightarrow K^-p$ BCN amplitudes with Pauli blocking only (black), Pauli+ YN SE (red), Pauli + YNK^- SE (green), and Pauli + $YNK^- \pi$ SE (magenta) at saturation density $\rho_0 = 0.17 \text{ fm}^{-3}$.

The role of hadron self-energies: K^-n amplitudes

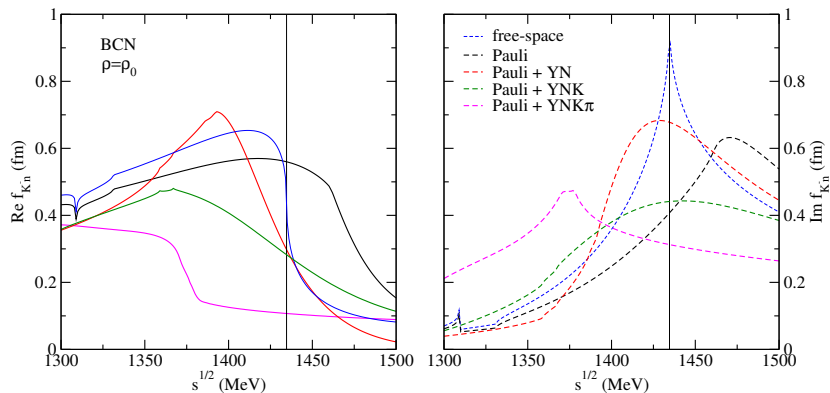


Fig.6: Comparison of $K^-n \rightarrow K^-n$ BCN amplitudes with Pauli blocking only (black), Pauli+YN SE (red), Pauli + YNK $^-$ SE (green), and Pauli + YNK $^-\pi$ SE (magenta) at saturation density $\rho_0 = 0.17 \text{ fm}^{-3}$.

Comparison with AMADEUS measurement

K. Piscicchia et al., PLB 782 (2018) 339 (AMADEUS); J.K. Kim, PRL 21 (1965) 29 (KIM)

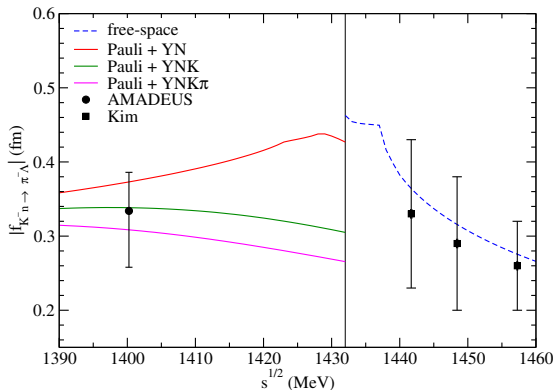


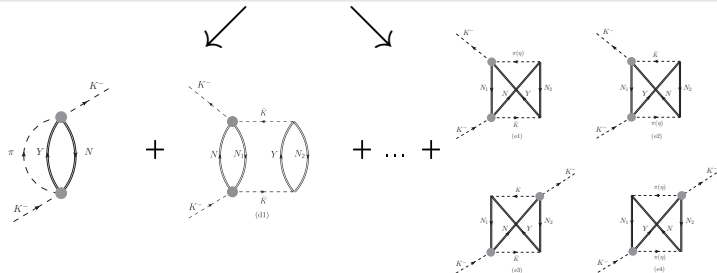
Fig.7: Comparison of $|f_{K^- n \rightarrow \pi^- \Lambda}|$ measured by AMADEUS with in-medium amplitudes: Pauli + YN SE (red), Pauli + YNK SE (green), and Pauli + YNK π SE (magenta) from the BCN model (taken at $\rho = 0.3\rho_0$). For comparison there is also the corresponding free-space amplitude (blue) above threshold (vertical line) compared with the amplitude extracted from cross section data.

The role of hadron self-energies: K^- potential

- Pauli + YN SE amplitudes $\Rightarrow V_{K^-} = V_{K^-N} + V_{K^-NN}$
- Pauli + $YNK^-(YNK^-\pi)$ SE amplitudes



$$V_{K^-} = t\rho + V_{K^-NN}^{\text{corr}}$$



The role of hadron self-energies: K^- potential

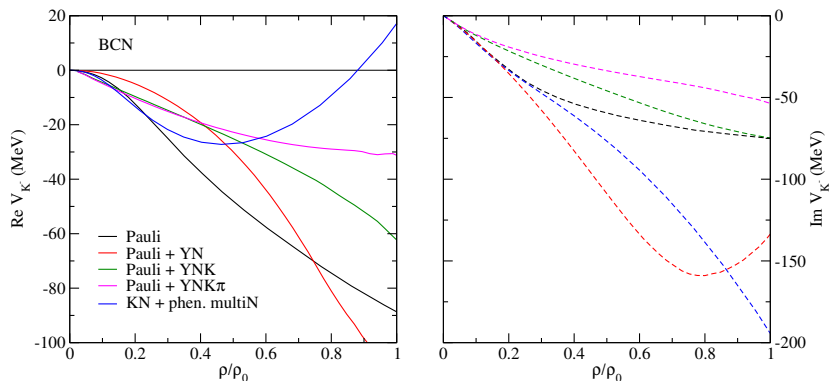


Fig.8: Real (left) and imaginary (right) parts of the total K^- potential as a function of relative density ρ/ρ_0 , calculated with Pauli (black), Pauli + YN SE (red), Pauli + YNK^- SE (green), and Pauli + $YNK^- \pi$ SE (magenta) BCN amplitudes. For comparison there is the best fit $K^-N + \text{phen. multiN}$ potential (blue) based on BCN amplitudes.

Calculations of energy shifts and widths in kaonic atoms

J. Óbertová, A. Ramos, J. Mareš, accepted in PRC (2025), arXiv:2508.07921

Table 3: Values of $\chi^2(64)$ obtained in calculations of kaonic atoms (24 nuclear species, 64 data points) using $K^-N + K^-NN$ potentials based on BCN Pauli, Pauli+YN, Pauli+YNK, and Pauli+YNK π amplitudes. For comparison, there are also results of calculations with K^-N +phen. multiN potential based on WRW modified BCN amplitudes (*T. Wass, M. Rho, W. Weise, NPA 617 (1997) 449*).

BCN	Pauli	Pauli+YN	Pauli+YNK	Pauli+YNK π	WRW
	$K^-N + K^-NN$	$K^-N + K^-NN$	$t\rho + V_{K^-NN}^{corr}$	$t\rho + V_{K^-NN}^{corr}$	K^-N +phen. multiN
$\chi^2(64)$	553.2	265.5	169.2	96.0	122.0
$\chi^2/\text{d.p.}$	8.6	4.2	2.6	1.5	1.9

- A. Baca et al. (*NPA 673 (2000)*) obtained value of $\chi^2/\text{dof} = 3.8$ with K^- potential based on chiral model of Oset and Ramos (*NPA 671 (2000)*).

Total K^- potential in $^{12}\text{C}+K^-$ and $^{208}\text{Pb}+K^-$ atoms

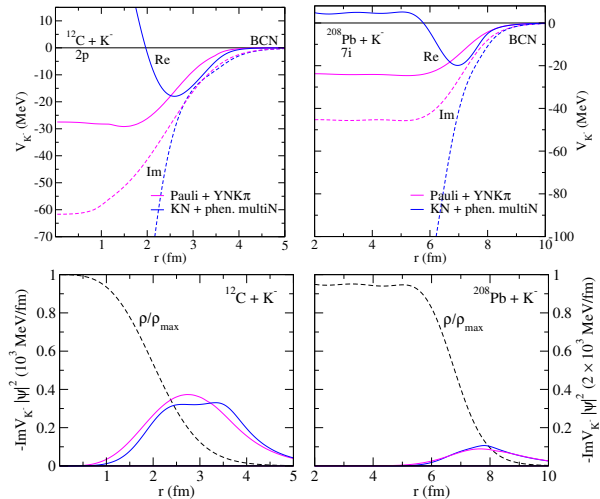


Fig.9: Total K^- potential in kaonic carbon (left) and kaonic lead (right) calculated with the Pauli+YNK π BCN amplitudes (magenta) compared with the K^-N +phen. multiN potential (blue) based on the BCN amplitudes. Lower panel shows the corresponding overlap of $\text{Im}V_{K^-}$ with the K^- wave function squared $|\psi|^2$.

Branching ratios in $^{12}\text{C}+K^-$ atom

Table 4: Primary-interaction branching ratios (in %) for mesonic ($K^-N \rightarrow Y\pi$, $Y = \Lambda, \Sigma$) and non-mesonic absorption ($K^-NN \rightarrow YN$) of K^- in $^{12}\text{C}+K^-$ atom (3d level), calculated with total K^- potential based on Pauli + YNK π BCN amplitudes. The experimental data corrected for primary interaction are shown for comparison.

$^{12}\text{C} + K^-$ (3d)	BCN		Exp.			
mesonic ratio	Pauli+YNK π	Sum	[1]	Sum	[2]	Sum
$\Sigma^+\pi^-$	14.4	} 33.7	29.4 ± 1.0	} 38.8 ± 1.2	14.4 ± 2.3	} 39.2 ± 3.0
$\Sigma^0\pi^-$	8.2		2.6 ± 0.6		1.2 ± 0.4	
$\Lambda\pi^-$	11.1	} 23.4	6.8 ± 0.3	} 26.0 ± 0.9	23.6 ± 1.9	} 24.8 ± 1.8
$\Sigma^-\pi^0$	8.1		2.6 ± 0.6		1.2 ± 0.4	
$\Sigma^0\pi^0$	10.0		20.0 ± 0.7		11.8 ± 1.4	
$\Lambda\pi^0$	5.3		3.4 ± 0.2		11.8 ± 1.0	
$\Sigma^-\pi^+$	12.8		13.1 ± 0.4		10.3 ± 1.7	
total 1N ratio (3d)	70.0		77.9 ± 1.6		74.3 ± 3.9	
total 1N ratio (3d + 4f)	71.8					
total multiN ratio (3d)	30.0		19.0 ± 2.0		25.7 ± 3.1	
total multiN ratio (3d + 4f)	28.2		$27 \pm 3(\text{stat.})^{+5}_{-6}(\text{syst.})$ [3]			

[1] C. Vander Velde-Wilquet et al., NC 39 A (1977) 538

[2] H. Davis et al., NC 53 A (1968) 313

[3] R. Del Grande et al., PS 95 (2020) 084012

Summary

- K^-N interaction described by chiral meson-baryon coupled channel interaction models
- Interactions of K^- with two and more nucleons important for realistic description of the K^- -nucleus interaction
 - ▶ only P, KM, and BCN models compatible with available data
- We have developed a microscopic model for K^-NN absorption in nuclear matter using amplitudes derived from the P and BCN chiral meson-baryon interaction models

J. Hrtánková, Ā. Ramos, PRC 101 (2020) 035204

 - ▶ Pauli blocked amplitudes included $\rightarrow$ medium effects non-negligible
 - ▶ Calculated ratios in a good agreement with experimental data

Summary

- Further extension of the K^-NN model and K^-N BCN amplitude model $\rightarrow$ inclusion of hadron (Y, N, K^-, π) self-energies
- Considering all hadron SE $\rightarrow \text{Re} V_{K^-}(\rho_0) \sim -30 \text{ MeV}$ and $\text{Im} V_{K^-}(\rho_0) \sim -50 \text{ MeV}$
- Calculations of energy shifts and widths in kaonic atoms:
 - ▶ inclusion of K^-NN absorption improves the description of data considerably
J. Óbertová, E. Friedman, J. Mareš, PRC 106 (2022) 065201
 - ▶ **full model Pauli+YNK π SE describes the data as good as the best fit K^-N +phen. multiN potential based on BCN amplitudes!**
J. Óbertová, A. Ramos, J. Mareš, accepted in PRC (2025), arXiv:2508.07921

BACK-UP SLIDES

Total K^- potential: version 2025 vs. 2000

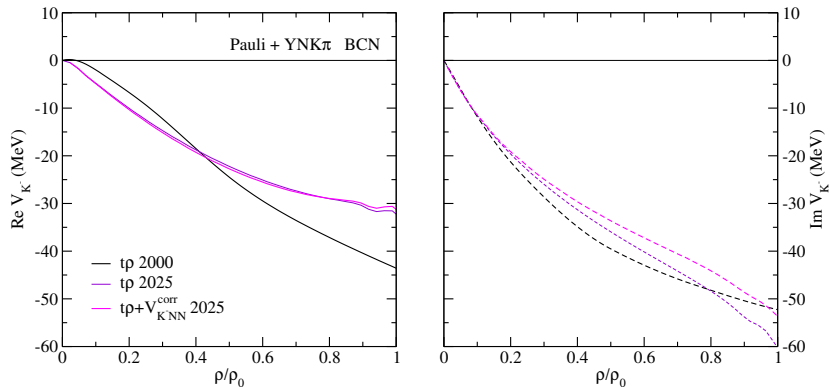


Fig.10: Real (left) and imaginary (right) parts of the total K^- potential calculated with Pauli + YNK π (magenta) BCN amplitudes compared with the K^- optical potential including hadron SE based on the model of Oset and Ramos (black). Thin violet lines represent the $t\rho$ part of the current K^- optical potential obtained with Pauli YNK π BCN amplitudes.

Branching ratios from FINUDA

Table 5: Primary-interaction ratios (in %) for the $\Sigma^-\pi^+$ and $\Sigma^+\pi^-$ channels in light kaonic atoms calculated using the Pauli+YNK π BCN amplitudes. Experimental data from the FINUDA collaboration [4] are corrected for pion attenuation and for $\Sigma N - \Lambda N$ conversion with probabilities (a) 60 % for Σ^+ and 22.5 % for Σ^- (b) 50 % for both Σ 's.

nucleus (<i>nl</i>)	BCN		Exp. [4]			
	Pauli+YNK π		(a)		(b)	
	$\Sigma^+\pi^-$	$\Sigma^-\pi^+$	$\Sigma^+\pi^-$	$\Sigma^-\pi^+$	$\Sigma^+\pi^-$	$\Sigma^-\pi^+$
${}^6\text{Li} (2p + 3d)$	16.1	17.5	45.8 ± 4.3	19.4 ± 1.5	36.7 ± 3.4	30.0 ± 2.3
${}^7\text{Li} (2p + 3d)$	12.6	11.0	20.5 ± 1.4	8.7 ± 1.0	16.4 ± 1.1	13.4 ± 1.5
${}^9\text{Be} (3d)$	13.1	12.6	14.6 ± 1.1	5.9 ± 0.7	11.7 ± 0.9	9.2 ± 1.2
${}^{12}\text{C}^a (3d)$	14.7	13.2	25.0 ± 0.8	10.2 ± 0.3	20.0 ± 0.7	15.9 ± 0.4
${}^{13}\text{C} (3d)$	12.7	10.3	18.6 ± 2.7	6.5 ± 0.8	14.9 ± 2.2	10.1 ± 1.2
${}^{16}\text{O} (3d)$	14.5	11.6	19.3 ± 1.9	8.4 ± 0.8	15.5 ± 1.5	13.0 ± 1.2

^a Branching ratios for ${}^{12}\text{C}$ were measured in bubble chamber experiment [1] but they served as reference values in Table 1 in Ref. [4].

[4] *M. Agnello et al., PLB 704 (2011) 474*

K^- absorption probabilities in p-shell nuclei

S. Manti, cascade calculations

