

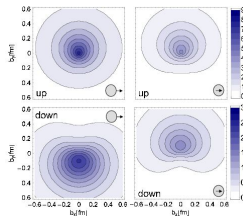
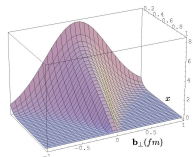
# Spin Physics at EIC

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October 20, 2011

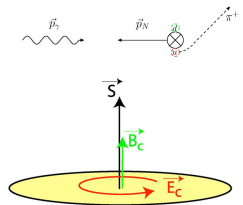
- Deeply virtual Compton scattering (DVCS)
- ↳ Generalized parton distributions (GPDs)
- ↳ 'transverse imaging'
- Probabilistic interpretation of GPDs as Fourier transforms of impact parameter dependent PDFs
  - $H(x, 0, -\Delta_{\perp}^2) \rightarrow q(x, \mathbf{b}_{\perp})$
  - $E(x, 0, -\Delta_{\perp}^2) \rightarrow \perp$  deformation of PDFs when the target is  $\perp$  polarized
- ↳ Ji relation
- Chromodynamik lensing and  $\perp$  single-spin asymmetries (SSA)



transverse distortion of PDFs  
+ final state interactions }  $\Rightarrow$

↳ SSA in  $\gamma N \rightarrow \pi + X$

- quark-gluon correlations  $\rightarrow \perp$  force on  $q$  in DIS
- Summary

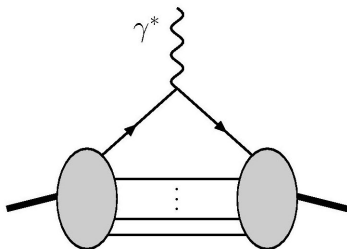


3 D imaging — join the experience!

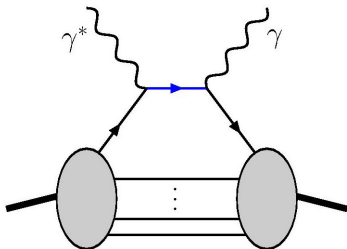


- virtual Compton scattering:  $\gamma^* p \rightarrow \gamma p$  (actually:  $e^- p \rightarrow e^- \gamma p$ )
  - ‘deeply’:  $-q_\gamma^2 \gg M_p^2, |t| \rightarrow$  Compton amplitude dominated by (coherent superposition of) Compton scattering off single quarks
- $\hookrightarrow$  only difference between form factor (a) and DVCS amplitude (b) is replacement of photon vertex by two photon vertices connected by **quark** (energy denominator depends on quark momentum fraction  $x$ )
- $\hookrightarrow$  DVCS amplitude provides access to momentum-decomposition of form factor = **Generalized Parton Distribution (GPDs)**.

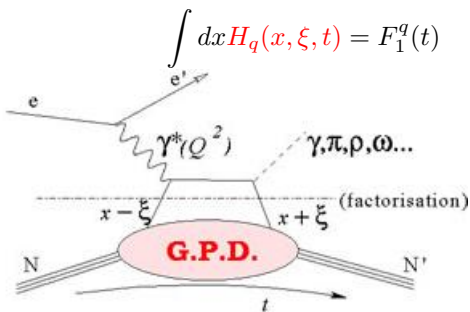
$$\int dx H_q(x, \xi, t) = F_1^q(t)$$



$$\int dx E_q(x, \xi, t) = F_2^q(t)$$



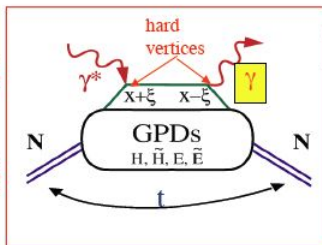
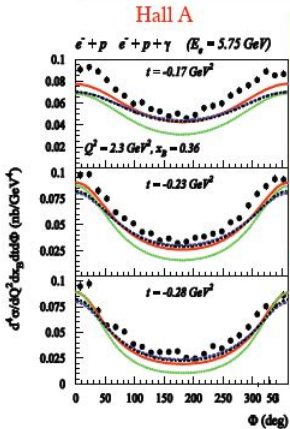
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$$\int dx E_q(x, \xi, t) = F_2^q(t)$$

- wide  $Q^2$  range
- $\hookrightarrow$   $Q^2$  dependence to disentangle convolution integral in  $\mathcal{A}_{DVCS}$
- high luminosity:
  - $\sigma_{DVCS} = \mathcal{O}(\alpha^3)$
  - study dependence on  $x, t, Q^2$

Unprecedented set of Deeply Virtual Compton Scattering data accumulated in **Hall A** and with **CLAS in Hall B at JLab**



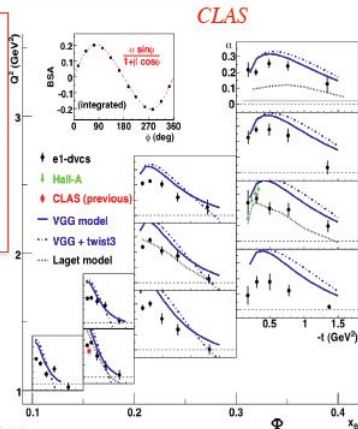
$$A = \frac{\sigma^+ - \sigma^-}{\sigma^+ + \sigma^-} = \frac{\Delta\sigma_{LU}}{2\sigma}$$

Polarized beam, unpolarized target:

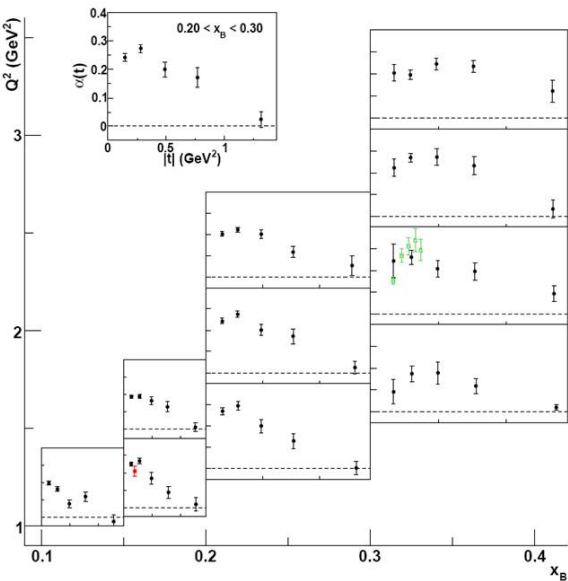
Kinematically suppressed

$$\Delta\sigma_{LU} \sim \sin\phi \{ F_1 H + \xi (F_1 + F_2) \bar{H} + k F_2 E \} U \phi$$

*Phys.Rev.Lett.*97:262002,2006



*Phys.Rev.Lett.*100:162002,2008



- form factors:  $\overleftarrow{FT} \rho(\vec{r})$
- $GPDs(x, \vec{\Delta})$ : form factor for quarks with momentum fraction  $x$
- ↪ suitable FT of  $GPDs$  should provide spatial distribution of quarks with momentum fraction  $x$
- careful: cannot measure longitudinal momentum ( $x$ ) and longitudinal position simultaneously (Heisenberg)
- ↪ consider purely transverse momentum transfer

### Impact Parameter Dependent Quark Distributions

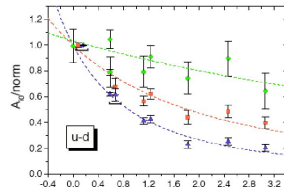
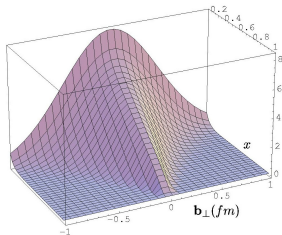
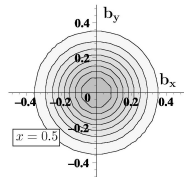
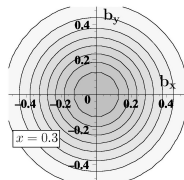
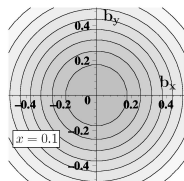
$$q(x, \mathbf{b}_\perp) = \int \frac{d^2 \Delta_\perp}{(2\pi)^2} H(x, \xi = 0, -\Delta_\perp^2) e^{-i\mathbf{b}_\perp \cdot \Delta_\perp}$$

$q(x, \mathbf{b}_\perp)$  = parton distribution as a function of the separation  $\mathbf{b}_\perp$  from the transverse center of momentum  $\mathbf{R}_\perp \equiv \sum_{i \in q, g} \mathbf{r}_{\perp, i} x_i$   
 MB, Phys. Rev. D62, 071503 (2000)

- No relativistic corrections (Galilean subgroup!)
- ↪ corollary: interpretation of 2d-FT of  $F_1(Q^2)$  as charge density in transverse plane also free of relativistic corrections
- probabilistic interpretation

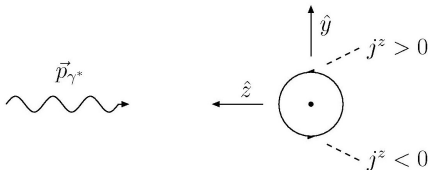
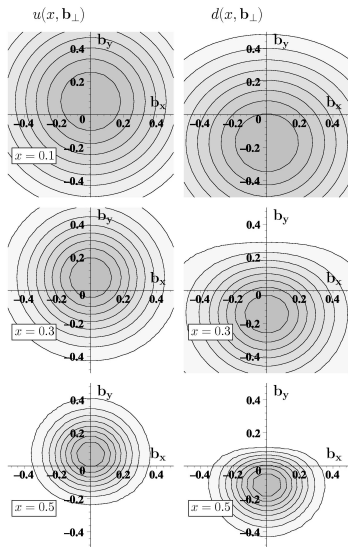


$q(x, \mathbf{b}_\perp)$  for unpol. p



### unpolarized proton

- $q(x, \mathbf{b}_\perp) = \int \frac{d^2\Delta_\perp}{(2\pi)^2} H(x, 0, -\Delta_\perp^2) e^{-i\mathbf{b}_\perp \cdot \Delta_\perp}$
  - $x$  = momentum fraction of the quark
  - $\vec{b}_\perp$  =  $\perp$  distance of quark from  $\perp$  center of momentum
  - small  $x$ : large 'meson cloud'
  - larger  $x$ : compact 'valence core'
  - $x \rightarrow 1$ : active quark becomes center of momentum
- $\hookrightarrow \vec{b}_\perp \rightarrow 0$  (narrow distribution) for  $x \rightarrow 1$

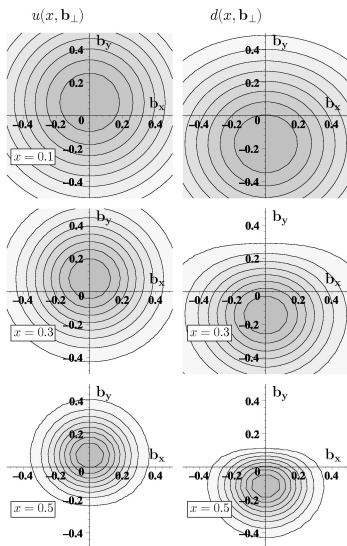


proton polarized in  $+\hat{x}$  direction

no axial symmetry!

$$q(x, \mathbf{b}_\perp) = \int \frac{d^2 \Delta_\perp}{(2\pi)^2} H_q(x, 0, -\Delta_\perp^2) e^{-i\mathbf{b}_\perp \cdot \Delta_\perp} \\ - \frac{1}{2M} \frac{\partial}{\partial b_y} \int \frac{d^2 \Delta_\perp}{(2\pi)^2} E_q(x, 0, -\Delta_\perp^2) e^{-i\mathbf{b}_\perp \cdot \Delta_\perp}$$

Physics: relevant density in DIS is  $j^+ \equiv j^0 + j^3$  and left-right asymmetry from  $j^3$



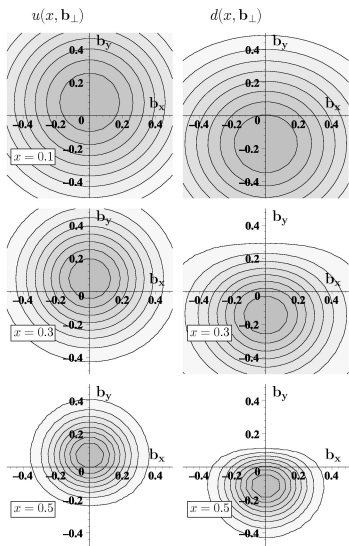
proton polarized in  $+\hat{x}$  direction

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sign & magnitude of the average shift

model-independently related to p/n  
anomalous magnetic moments:

$$\langle b_y^q \rangle \equiv \int dx \int d^2 b_\perp q(x, \mathbf{b}_\perp) b_y \\ = \frac{1}{2M} \int dx E_q(x, 0, 0) = \frac{\kappa_q}{2M}$$



sign & magnitude of the average shift  
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$$\kappa^P = 1.913 = \frac{2}{3} \kappa_u^P - \frac{1}{3} \kappa_d^P + \dots$$

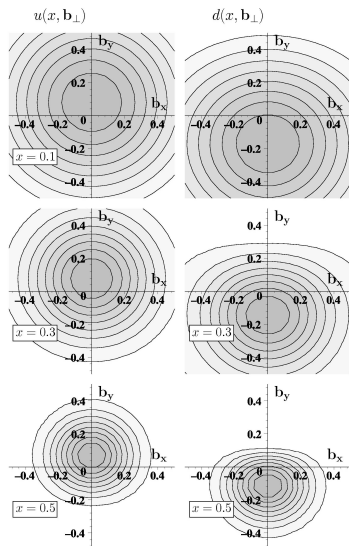
- $u$ -quarks:  $\kappa_u^P = 2\kappa_p + \kappa_n = 1.673$

↪ shift in  $+\hat{y}$  direction

- $d$ -quarks:  $\kappa_d^P = 2\kappa_n + \kappa_p = -2.033$

↪ shift in  $-\hat{y}$  direction

- $\langle b_y^q \rangle = \mathcal{O}(\pm 0.2 \text{ fm})$  !!!!

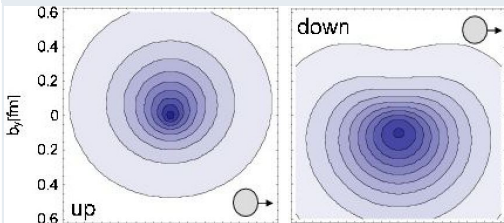


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lattice QCD (QCDSF): lowest moment



transverse images  $\leftrightarrow$  Ji relation for quark angular momentum:

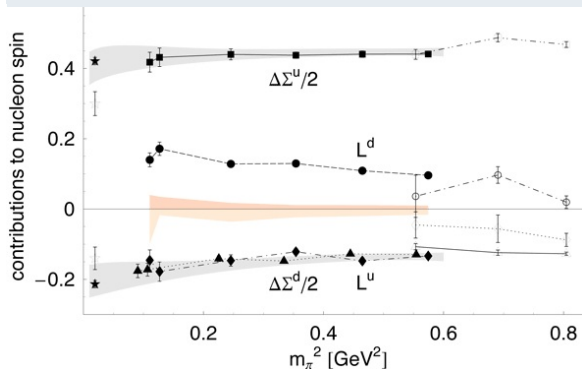
- $J_q^x = m_N \int dx x r^y q(x, \mathbf{r}_\perp)$  with  $b^y = r^y - \frac{1}{2m_N}$ , where  $q(x, \mathbf{r}_\perp)$  is distribution relative to CoM of whole nucleon
- recall:  $q(x, \mathbf{b}_\perp)$  for nucleon polarized in  $+\hat{x}$  direction

$$q(x, \mathbf{b}_\perp) = \int \frac{d^2 \Delta_\perp}{(2\pi)^2} H_q(x, 0, -\Delta_\perp^2) e^{-i\mathbf{b}_\perp \cdot \Delta_\perp} \\ - \frac{1}{2M_N} \frac{\partial}{\partial b_y} \int \frac{d^2 \Delta_\perp}{(2\pi)^2} E_q(x, 0, -\Delta_\perp^2) e^{-i\mathbf{b}_\perp \cdot \Delta_\perp}$$

$$\Rightarrow J_q^x = M_N \int dx x r^y q(x, \mathbf{r}_\perp) = \int dx x \left( m_N b^y + \frac{1}{2} \right) q(x, \mathbf{r}_\perp) \\ = \frac{1}{2} \int dx x [H(x, 0, 0) + E(x, 0, 0)]$$

- X.Ji(1996): rotational invariance  $\Rightarrow$  apply to all components of  $\vec{J}_q$
- partonic interpretation exists only for  $\perp$  components!

lattice: QCDSF



$$J^q = \frac{1}{2} \int dx x [H(x, 0, 0) + E(x, 0, 0)]$$

$$L^q = J^q - \frac{1}{2} \Delta\Sigma^q$$

## TMDs

- Transverse Momentum Dependent Parton Distributions
- 8 structures possible at leading twist (only 3 for PDFs)
- $f_{1T}^\perp$  and  $h_1^\perp$  require both **orbital angular momentum** and **final state interaction**
- can be measured in SIDIS and DY

“TMDs”

### Sivers function

correlation between the transverse spin of the nucleon and the transverse momentum of the quark

*sensitive to orbital angular momentum*
















### Boer-Mulders function

correlation between the transverse spin and the transverse momentum of the quark in unpol nucleons

*T-odd*

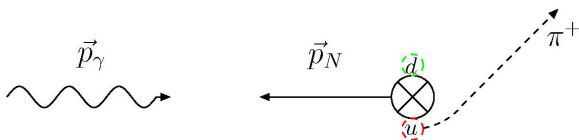
quark polarisation

nucleon polarisation

	U	L	T
U	$f_1$  number density $q$		$f_{1T}^\perp$  -  Sivers
L		$g_1$  -  helicity $\Delta q$	$g_{1T}$  - 
T	$h_1^\perp$  -  Boer Mulders	$h_{1L}^\perp$  - 	$h_1$  -  transversity $h_{1T}^\perp$  - 



example: semi-inclusive deep-inelastic scattering (SIDIS)  $\gamma p \rightarrow \pi X$



- $u, d$  distributions in  $\perp$  polarized proton have left-right asymmetry in  $\perp$  position space (T-even!); sign 'determined' by  $\kappa_u$  &  $\kappa_d$
  - attractive FSI deflects active quark towards the CoM
- $\hookrightarrow$  FSI translates position space distortion (before the quark is knocked out) in  $+\hat{y}$ -direction into momentum asymmetry that favors  $-\hat{y}$  direction  $\rightarrow$  'chromodynamic lensing'

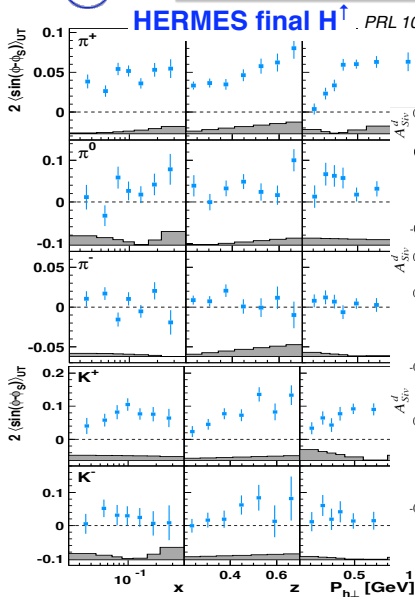
$\Rightarrow$

$\kappa_p, \kappa_n \longleftrightarrow$  sign of SSA!!!!!!!

- confirmed by HERMES (and recent COMPASS)  $p$  data; consistent with vanishing isoscalar Sivers (COMPASS)



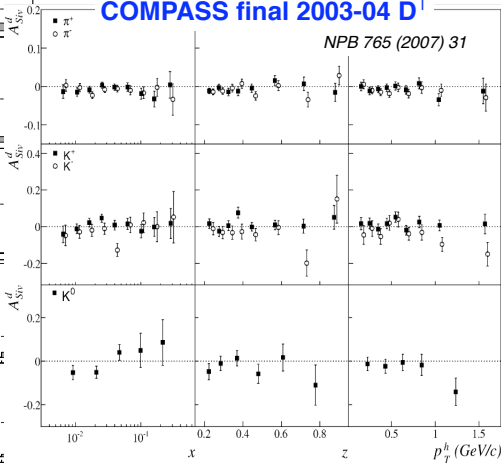
## Sivers Moments for $\pi$ and $K$ from $H^\uparrow$ & $D^\uparrow$

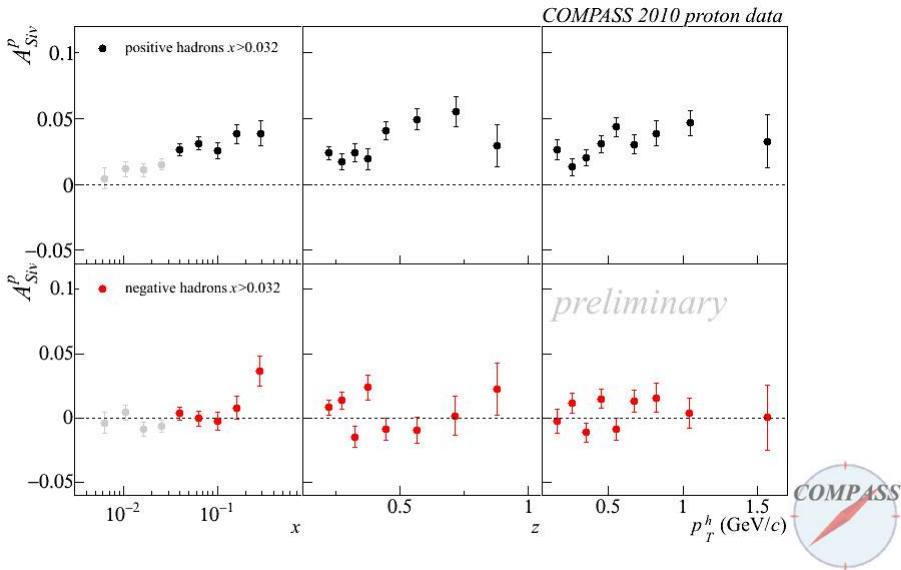


$$f_{1T}^\perp(x, k_T) \otimes D_1^\perp(z)$$

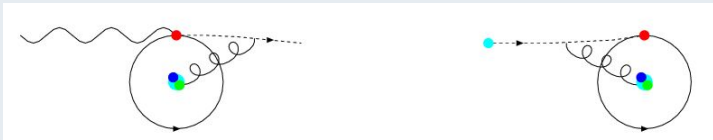


## COMPASS final 2003-04 $D^\uparrow$





compare FSI for 'red'  $q$  that is being knocked out of nucleon with ISI for 'anti-red'  $\bar{q}$  that is about to annihilate with a 'red' target  $q$



### FSI in SIDIS

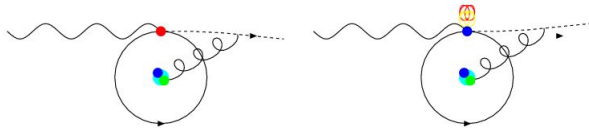
- knocked-out  $q$  'red'
- ↪ spectators 'anti-red'
- ↪ interaction between knocked-out quark and spectators **attractive**

### ISI in DY

- incoming  $\bar{q}$  'anti-red'
- ↪ struck target  $q$  'red'
- ↪ spectators also 'anti-red'
- ↪ interaction between incoming  $\bar{q}$  and spectators **repulsive**

test of  $f_{1T}^\perp(x, \mathbf{k}_\perp)_{DY} = -f_{1T}^\perp(x, \mathbf{k}_\perp)_{SIDIS}$  and  $h_1^\perp(x, \mathbf{k}_\perp)_{DY} = -h_1^\perp(x, \mathbf{k}_\perp)_{SIDIS}$   
**critical test** of TMD factorization approach

'Chromodynamic lensing' mechanism for  $\perp$  SSA requires long range coherence of color field!

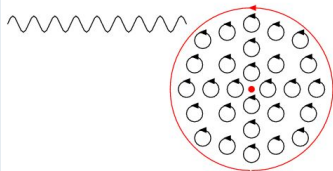


before 'dressing'      active quark 'dressed' with glue

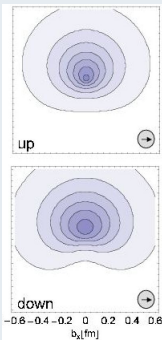
QCD-evolution: long-range color decoherence!

- after 'dressing' itself with a gluon, previously red quark more likely to be blue or green
- ↪ attraction to far-away spectators mostly gone
- only attracted to close-by (high  $Q^2$ )  $g$  from dressing
- high  $Q^2$ :  $q$  at low  $x$  likely to have dressed itself with perturbative gluon!
- ↪ 'Chromodynamic lensing' mechanism suppressed for high  $Q^2$  & small  $x$ ?
- ↪ interesting  $Q^2$  dependence?

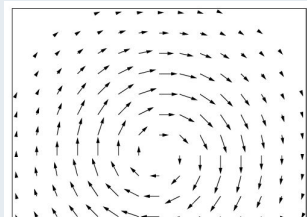
$q$  with polarization  $\odot$



lattice calculation (QCDSF)



unpolarized target



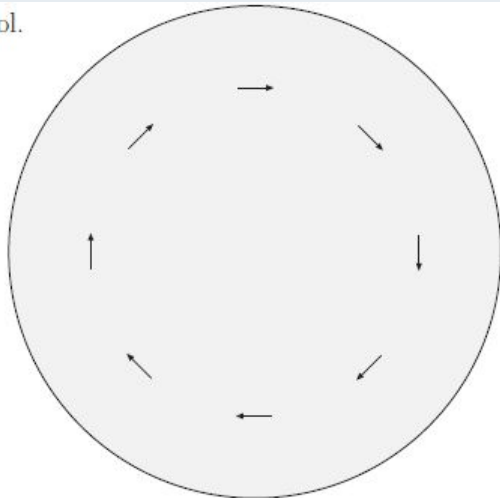
- transversity distribution in unpol. target described by chirally odd GPD  $\bar{E}_T$
- $\bar{E}_T > 0$  for  $u$  &  $d$  (QCDSF)
- connection  $h_1^\perp(x, \mathbf{k}_\perp) \leftrightarrow \bar{E}_T$  similar to  $f_{1T}^\perp(x, \mathbf{k}_\perp) \leftrightarrow E$
- $h_1^\perp(x, \mathbf{k}_\perp) < 0$  for  $u/p, d/p, u/\pi, \bar{d}/\pi$
- $h_{1\text{SIDIS}}^\perp = -h_{1\text{DY}}^\perp$

experimental access

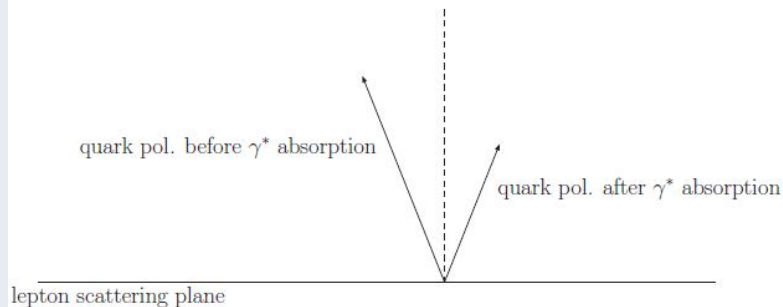
no polarization needed!

## Primordial Quark Transversity Distribution

→  $\perp$  quark pol.



## Primordial Quark Transversity Distribution



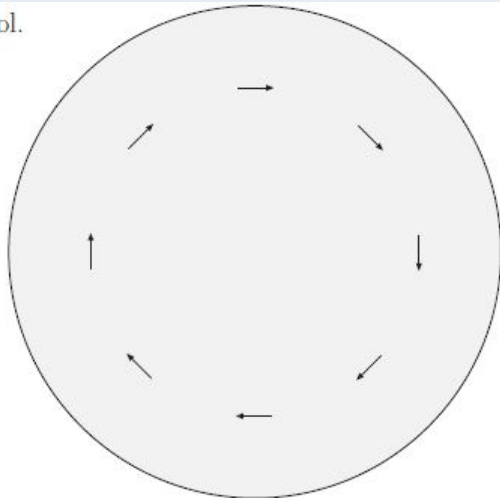
when  $\perp$  pol. quark absorbs  $\gamma^*$ ,  $\perp$  polarization

- gets reduced in size
- tilted symmetrically w.r.t. normal of lepton scattering plane



## Primordial Quark Transversity Distribution

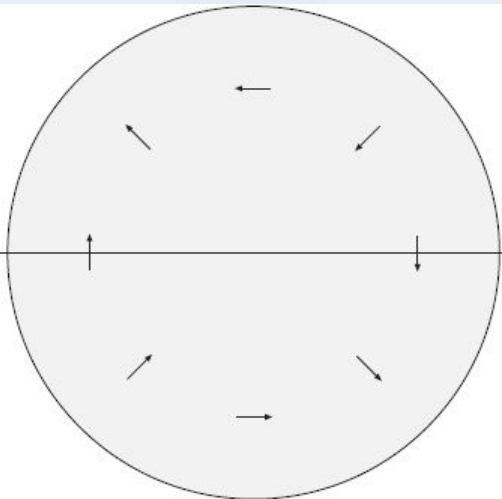
→  $\perp$  quark pol.



Quark Transversity Distribution after  $\gamma^*$  Absorption

→  $\perp$  quark pol.

lepton scattering plane

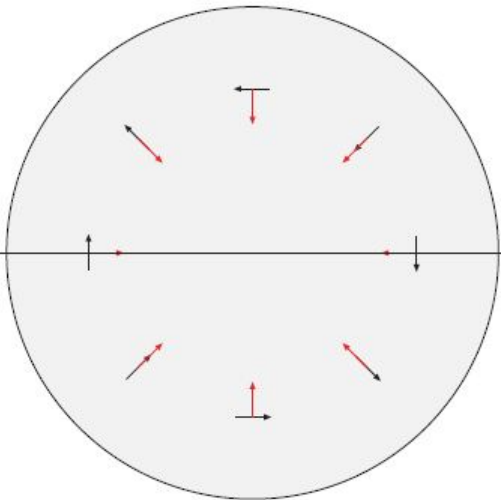


$\perp$  momentum (of  $q$ ) due to FSI

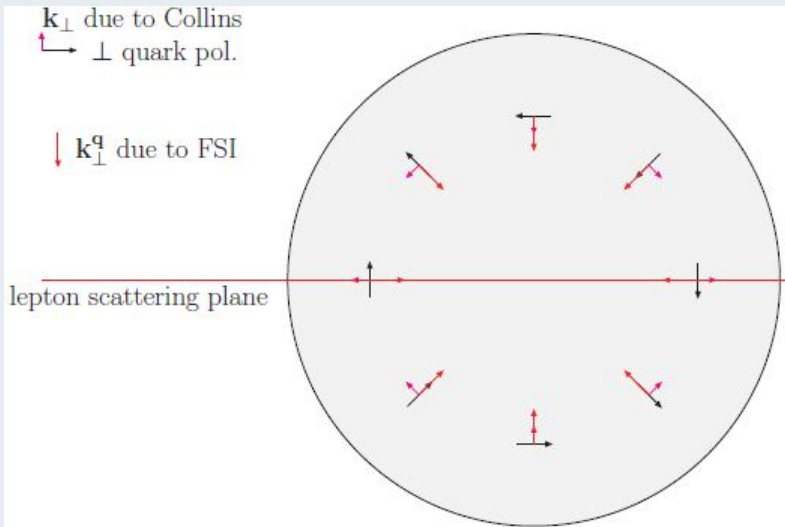
$\rightarrow$   $\perp$  quark pol.

$\downarrow$   $k_{\perp}^q$  due to FSI

lepton scattering plane



additional  $\perp$  momentum (of  $\pi$ ) due to Collins effect



Collins: favored  $\pi$  momentum preferentially to **left** (quark spin up)

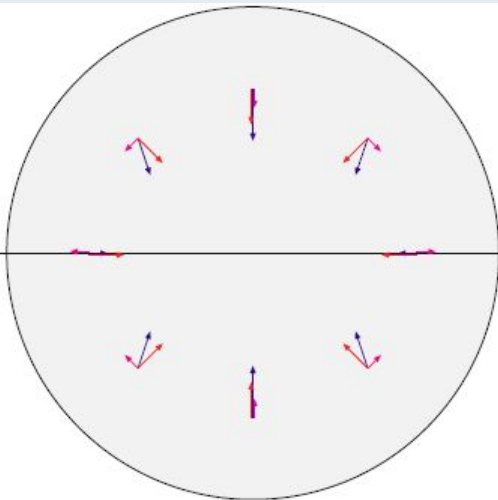
net  $k_{\perp}^{\pi}$  (FSI + Collins)

$\downarrow$   
 $k_{\perp}$  due to Collins

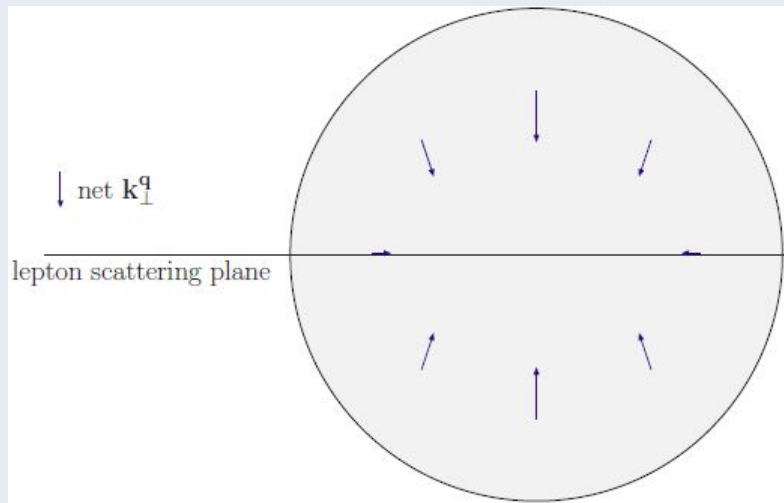
$\downarrow$   $k_{\perp}^q$  due to FSI

$\downarrow$  net  $k_{\perp}^q$

lepton scattering plane



net  $k_{\perp}^{\pi}$  (FSI + Collins)



$\cos 2\pi$  modulation of  $k_{\perp}^{\pi}$

## higher twist in polarized DIS

- $\sigma_{LL} \propto g_1 - \frac{2Mx}{\nu} g_2$
  - $g_1 = \frac{1}{2} \sum_q e_q^2 g_1^q$  with  $g_1^q = q^\uparrow(x) + \bar{q}^\uparrow(x) - q^\downarrow(x) - \bar{q}^\downarrow(x)$
  - $g_2$  involves quark-gluon correlations
- ↪ no parton interpret. as difference between number densities for  $g_2$
- for  $\perp$  pol. target,  $g_1$  &  $g_2$  contribute equally

$$\sigma_{LT} \propto g_T \equiv g_1 + g_2$$

- ↪ 'clean' separation between  $g_2$  and  $\frac{1}{Q^2}$  corrections to  $g_1$

What can we learn from  $g_2$ ?

- $g_2 = g_2^{WW} + \bar{g}_2$  with  $g_2^{WW}(x) \equiv -g_1(x) + \int_x^1 \frac{dy}{y} g_1(y)$

$$d_2 \equiv 3 \int dx x^2 \bar{g}_2(x) = \frac{1}{2MP^{+2}S^x} \langle P, S | \bar{q}(0) g G^{+y}(0) \gamma^+ q(0) | P, S \rangle$$

$$d_2 \equiv 3 \int dx x^2 \bar{g}_2(x) = \frac{1}{2MP^{+2}S^x} \langle P, S | \bar{q}(0) g G^{+y}(0) \gamma^+ q(0) | P, S \rangle$$

$$\sqrt{2}G^{+y} = G^{0y} + G^{zy} = -E^y + B^x$$



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color Lorentz force

$$\sqrt{2}G^{+y} = G^{0y} + G^{zy} = -E^y + B^x = -\left(\vec{E} + \vec{v} \times \vec{B}\right)^y \text{ for } \vec{v} = (0, 0, -1)$$

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$\leftrightarrow$   $d_2 \leftrightarrow$  average **color Lorentz force** acting on quark moving with  $v = c$  in  $-\hat{z}$  direction in the instant after being struck by  $\gamma^*$

$$\langle F^y \rangle = -2M^2 d_2 = -\frac{M}{P^{+2}S^x} \langle P, S | \bar{q}(0) g G^{+y}(0) \gamma^+ q(0) | P, S \rangle$$

cf. Qiu-Sterman matrix element  $\langle k_\perp^y \rangle \equiv \int_0^1 dx \int d^2 k_\perp k_\perp^2 f_{1T}^\perp(x, k_\perp^2)$

$$\langle k_\perp^y \rangle = -\frac{1}{2p^+} \left\langle P, S \left| \bar{q}(0) \int_0^\infty dx^- g G^{+y}(x^-) \gamma^+ q(0) \right| P, S \right\rangle$$

semi-classical interpretation: average  $k_\perp$  in SIDIS obtained by correlating the quark density with the transverse impulse acquired from (color) Lorentz force acting on struck quark along its trajectory to (light-cone) infinity

matrix element defining  $d_2$

$\leftrightarrow$

1<sup>st</sup> integration point in QS-integral

$$d_2 \equiv 3 \int dx x^2 \bar{g}_2(x) = \frac{1}{2MP^{+2}S_x} \langle P, S | \bar{q}(0) g G^{+y}(0) \gamma^+ q(0) | P, S \rangle$$

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sign of  $d_2 \leftrightarrow \perp$  imaging

- $\kappa_q/p \rightarrow$  sign of deformation
- $\hookrightarrow$  direction of average force
- $\hookrightarrow d_2^u > 0, d_2^d < 0$
- cf.  $f_{1T}^{\perp u} < 0, f_{1T}^{\perp d} < 0$

lattice (Göckeler et al., 2005)

$$d_2^u \approx 0.010, d_2^d \approx -0.0056$$

magnitude of  $d_2$

- $\langle F^y \rangle = -2M^2 d_2 = -10 \frac{\text{GeV}}{fm} d_2$
- expect partial cancellation of forces in SSA
- $\hookrightarrow |\langle F^y \rangle| \ll \sigma \approx 1 \frac{\text{GeV}}{fm}$
- $\hookrightarrow d_2 = \mathcal{O}(0.01)$

## color Lorentz force

$e_2 \leftrightarrow$  average **color Lorentz force** (in  $\hat{y}$ -direction) acting on quark (with transversity  $\hat{x}$ ) moving with  $v = c$  in  $-\hat{z}$  direction in the instant after being struck by  $\gamma^*$

$$\langle F^y \rangle = M^2 e_2 \equiv M^2 \int_0^1 dx \bar{e}_2(x) = \frac{M}{4P^{+2}} \sum_{i=1,2} \langle P | \bar{q}(0) g G^{+i}(0) \sigma^{+i} q(0) | P \rangle$$

## chirally even

- GPD  $E_q \Rightarrow \mathbf{b}_\perp$  deformation of unpol.  $q$  distr. in  $\perp$  pol. target

$$\hookrightarrow f_{1T}^\perp$$

$$\hookrightarrow d_2 \equiv \int dx x^2 \bar{g}_2 \text{ force}$$

## chirally odd

- GPD  $\bar{E}_T \Rightarrow \mathbf{b}_\perp$  deformation of quarks with transversity in unpol. target

$$\hookrightarrow h_1^\perp$$

$$\hookrightarrow e_2 \equiv \int dx x^2 \bar{e}_2 \text{ force}$$

## lattice (Göckeler et al., 2005)

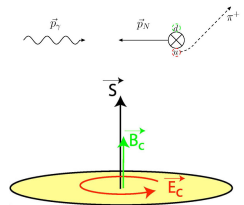
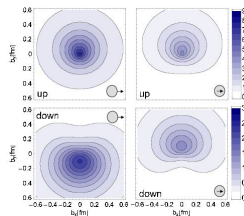
$$d_2^u \approx 0.010, d_2^d \approx -0.0056$$

please: up to date lattice calcs.

## lattice

B.Musch ...

- Deeply Virtual Compton Scattering  $\rightarrow$  GPDs
- $\hookrightarrow$  impact parameter dependent PDFs  $q(x, \mathbf{b}_\perp)$
- $E^q(x, 0, -\Delta_\perp^2) \leftrightarrow \kappa_{q/p}$  (contribution from quark flavor  $q$  to anomalous magnetic moment)
- $E^q(x, 0, -\Delta_\perp^2) \rightarrow \perp$  deformation of PDFs for  $\perp$  polarized target
- $\perp$  deformation  $\leftrightarrow$  (sign of) SSA (Sivers; Boer-Mulders)
- parton interpretation for Ji-relation
- higher-twist  $(\int dx x^2 \bar{g}_2(x), \int dx x^2 \bar{e}(x)) \leftrightarrow \perp$  force in DIS
- $\perp$  deformation  $\leftrightarrow$  (sign of) quark-gluon correlations  $(\int dx x^2 \bar{g}_2(x), \int dx x^2 \bar{e}(x))$



combine complementary information from deeply-virtual Compton scattering, semi-inclusive DIS & Drell-Yan to study orbital angular momentum and map the 3-d structure of hadrons