

Recent Results in Particle and Nuclear Physics from Lattice QCD



T. Hatsuda
(Univ. Tokyo / RIKEN)



THE UNIVERSITY
OF TOKYO



Recent **Selected** Results in Particle and Nuclear Physics from Lattice QCD



1. LQCD basics
2. Precision LQCD
3. Thermal LQCD
4. Nuclear LQCD
5. Summary

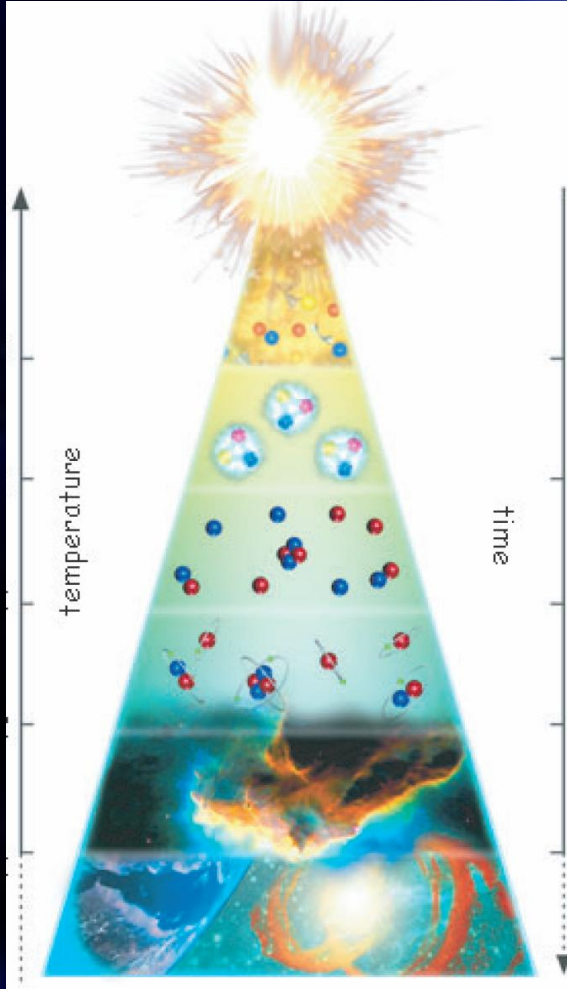
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Current challenges in QCD



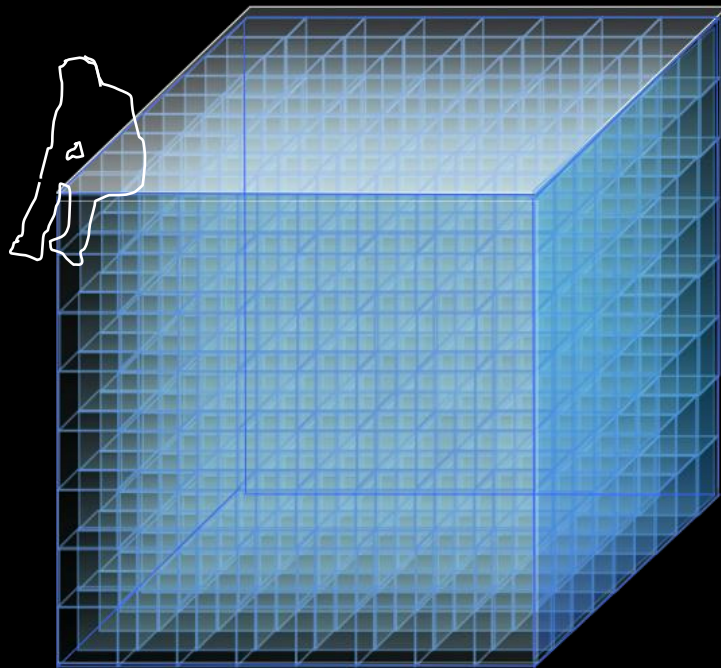
- **Primordial form of matter**
quark-gluon plasma, hadron structure
- **Origin of heavy elements**
in explosive astrophysical phenomena
- **Super dense matter**
neutron star, exotic matter, ...
- **Inputs for “new physics” search**
dark matter, ...



Lattice QCD provides

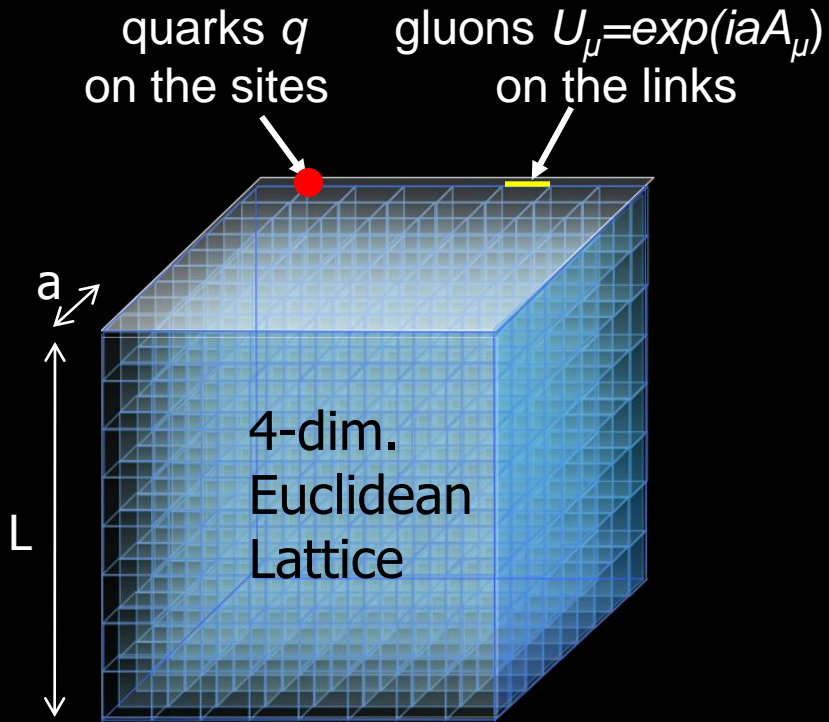
(1) precision calculations & (2) qualitative pictures

Lattice QCD basics

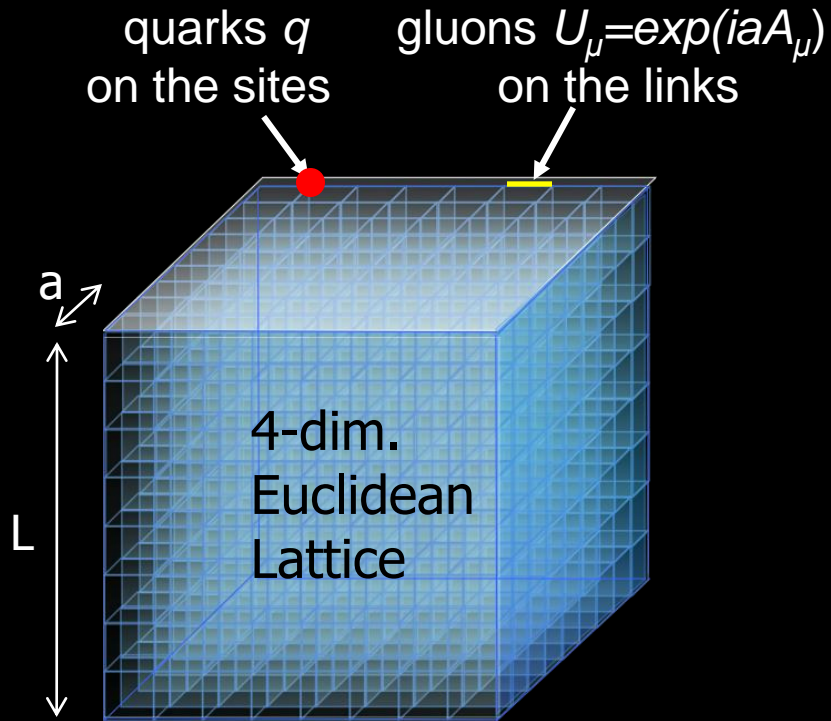


$$\mathcal{L} = -\frac{1}{4}G_{\mu\nu}^a G_a^{\mu\nu} + \bar{q}\gamma^\mu (i\partial_\mu - gt^a A_\mu^a)q - m\bar{q}q$$

$$\mathcal{L} = -\frac{1}{4}G_{\mu\nu}^a G_a^{\mu\nu} + \bar{q}\gamma^\mu (i\partial_\mu - gt^a A_\mu^a)q - m\bar{q}q$$



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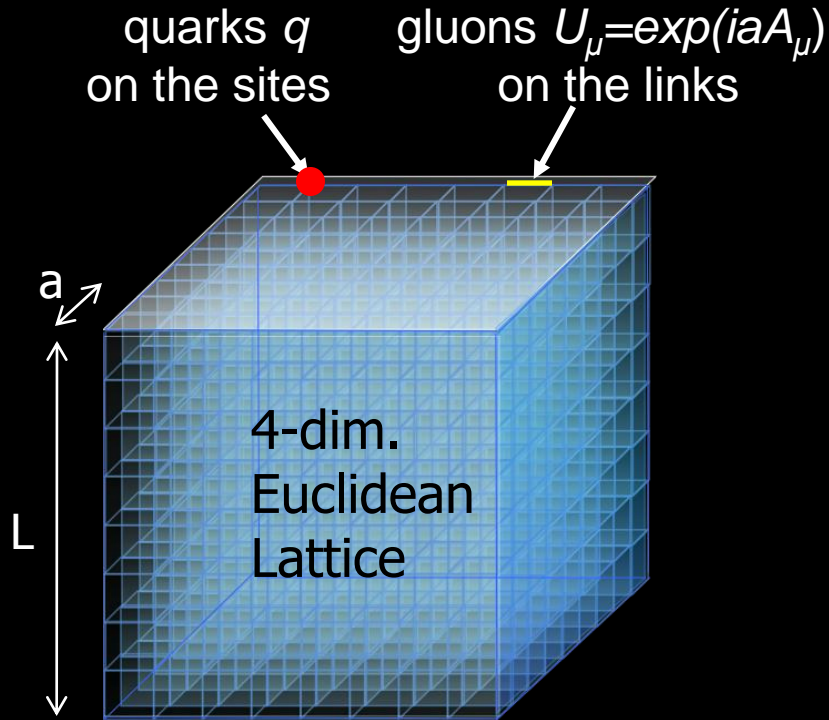


$$Z = \int [dU][dq d\bar{q}] \exp \left[- \int d\tau d^3x \mathcal{L}_E \right]$$

Monte Carlo method

Observable
= $O(g, m, a, L)$

$$\mathcal{L} = -\frac{1}{4}G_{\mu\nu}^a G_a^{\mu\nu} + \bar{q}\gamma^\mu (i\partial_\mu - g t^a A_\mu^a)q - m\bar{q}q$$



What can be done

- hadron properties & interactions
- hot plasma in equilibrium

What is difficult

- cold plasma
- phenomena far from equilibrium

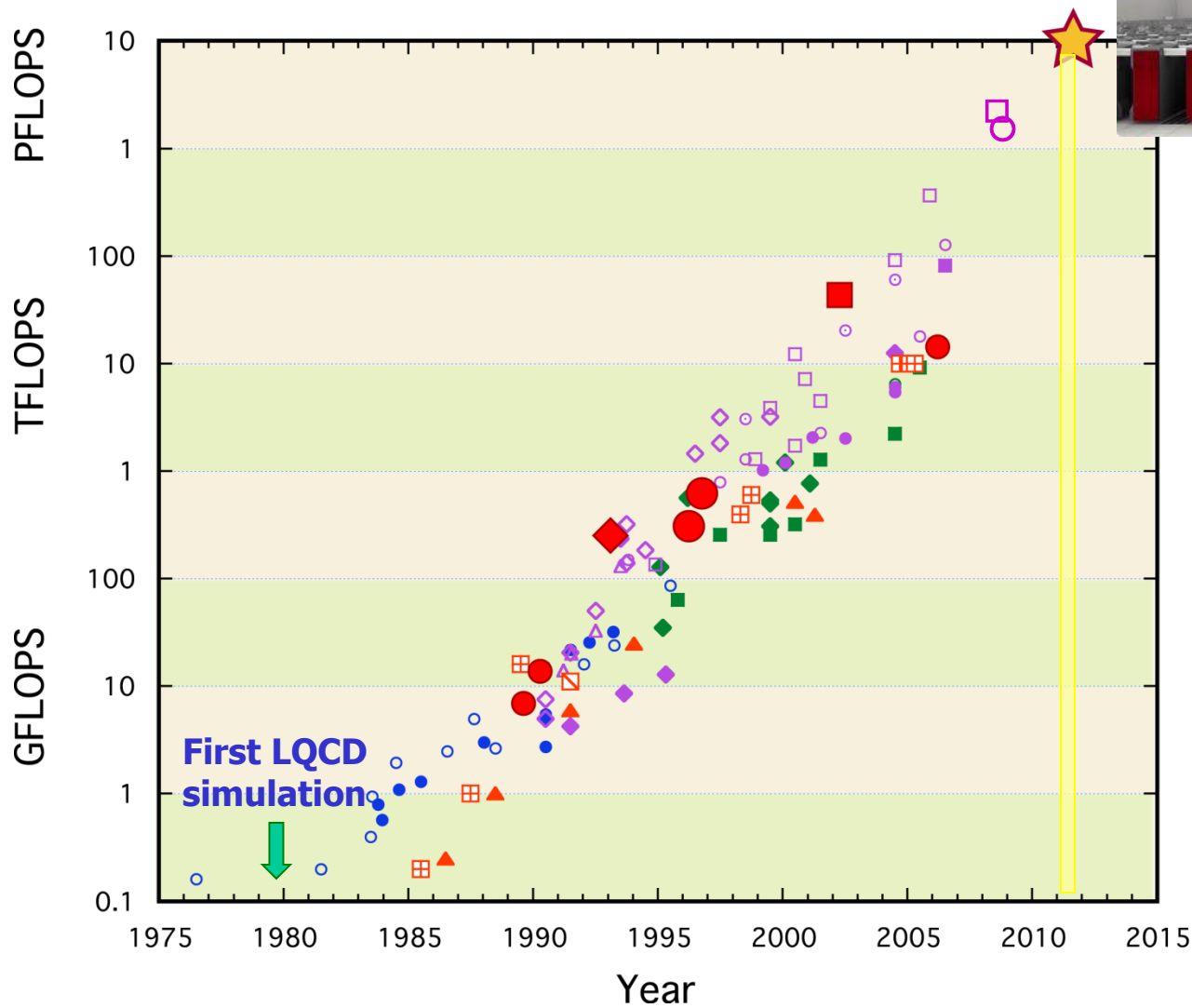
$$Z = \int [dU][dq d\bar{q}] \exp \left[- \int d\tau d^3x \mathcal{L}_E \right]$$

Monte Carlo method

Observable

$= O(g, m, a, L)$

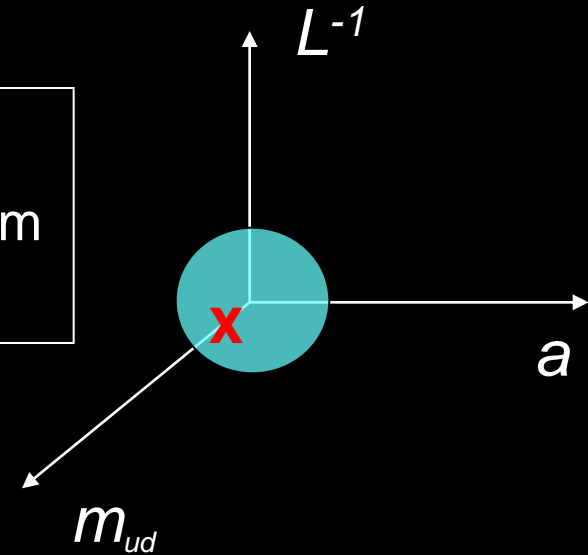
Supercomputer peak performance



**10PFlops K computer
(RIKEN)**

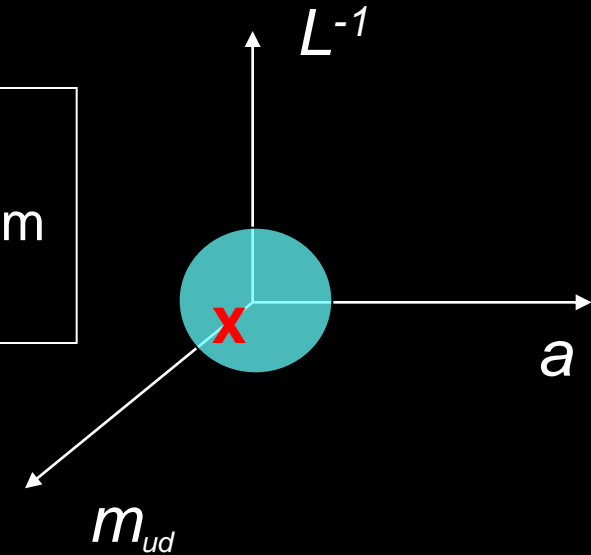
Three Important limits

$L^{-1} \rightarrow 0$ (thermodynamics limit) : finite size scaling
 $a \rightarrow 0$ (continuum limit) : asymptotic freedom
 $m \rightarrow 0$ (chiral limit) : chiral pert. theory



Three Important limits

$L^{-1} \rightarrow 0$ (thermodynamics limit) : finite size scaling
 $a \rightarrow 0$ (continuum limit) : asymptotic freedom
 $m \rightarrow 0$ (chiral limit) : chiral pert. theory



“Techniques”

Fermions:

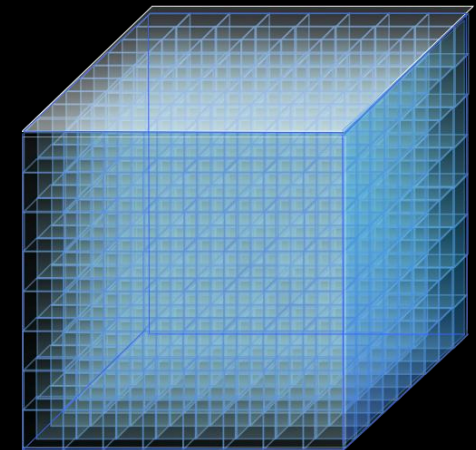
Staggered, Wilson, Domain-wall, Overlap
different ways of handling chiral symmetry

Improved actions:

stout, HEX, asktad, HISQ, clover, ...
different ways of reducing the discretization error

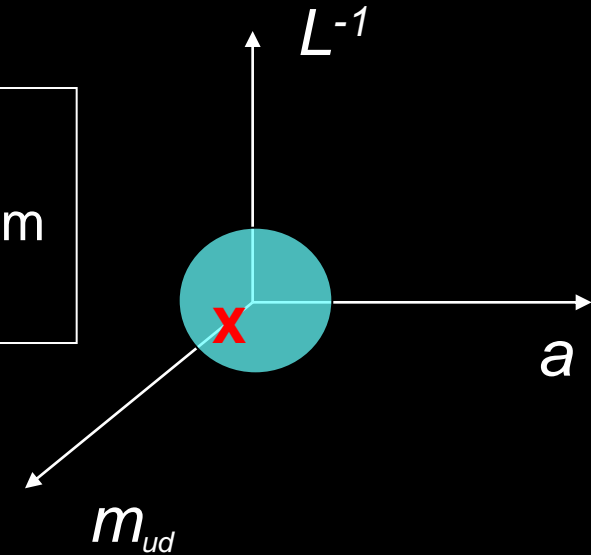
Advanced algorithms:

RHMC, DDHMC, LMA,
techniques to make the simulations fast and reliable



Three Important limits

$L^{-1} \rightarrow 0$ (thermodynamics limit) : finite size scaling
 $a \rightarrow 0$ (continuum limit) : asymptotic freedom
 $m \rightarrow 0$ (chiral limit) : chiral pert. theory



“Techniques”

Fermions:

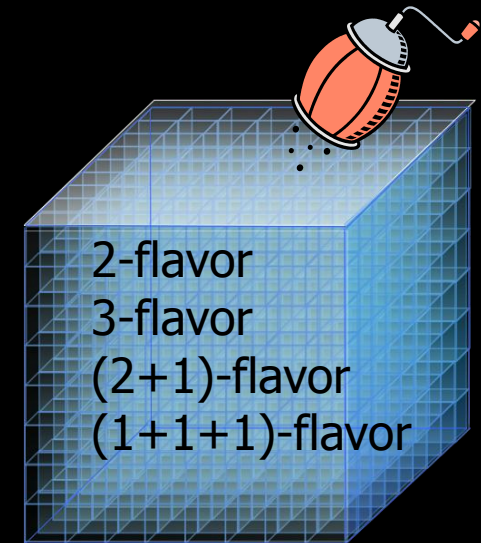
Staggered, Wilson, Domain-wall, Overlap
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Improved actions:

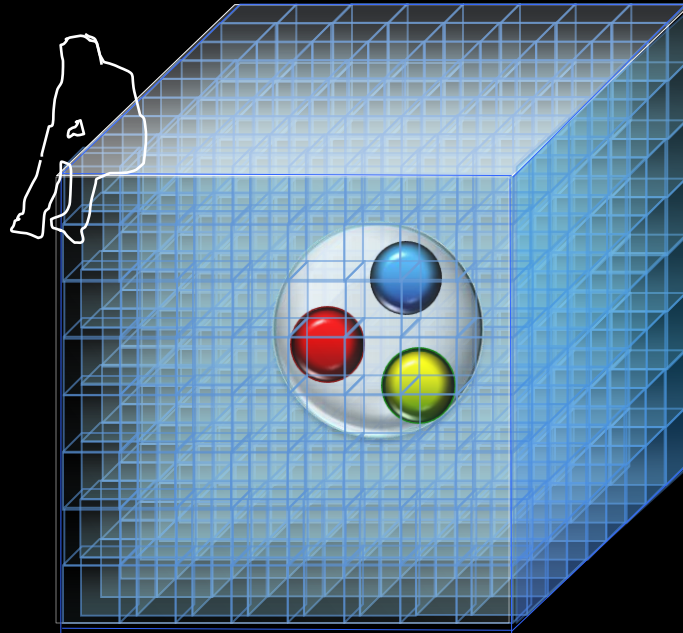
stout, HEX, asktad, HISQ, clover, ...
different ways of reducing the discretization error

Advanced algorithms:

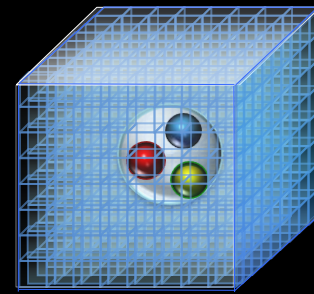
RHMC, DDHMC, LMA, ...
techniques to make the simulations fast and reliable



Precision Lattice QCD



Hadron masses @ 2009

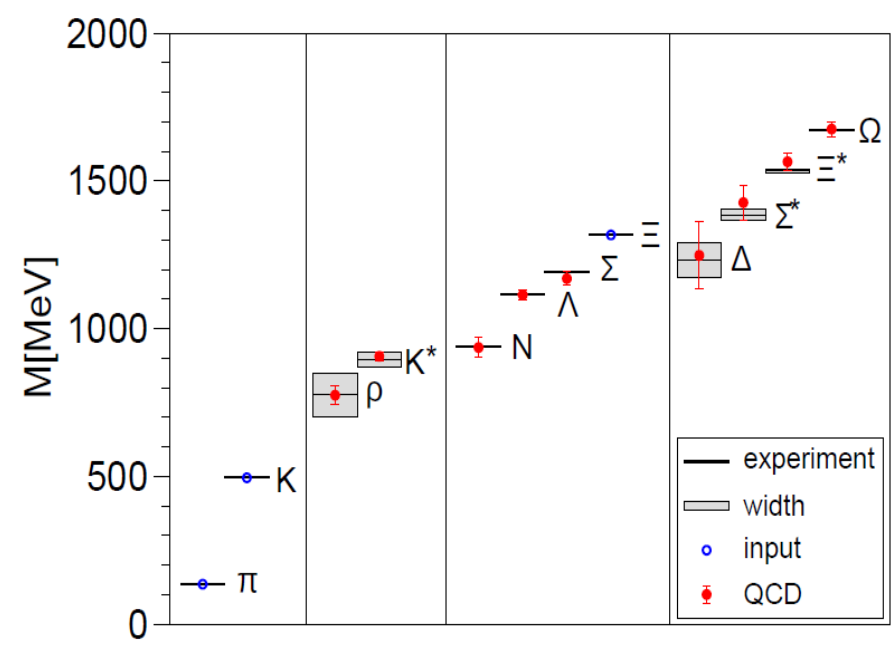
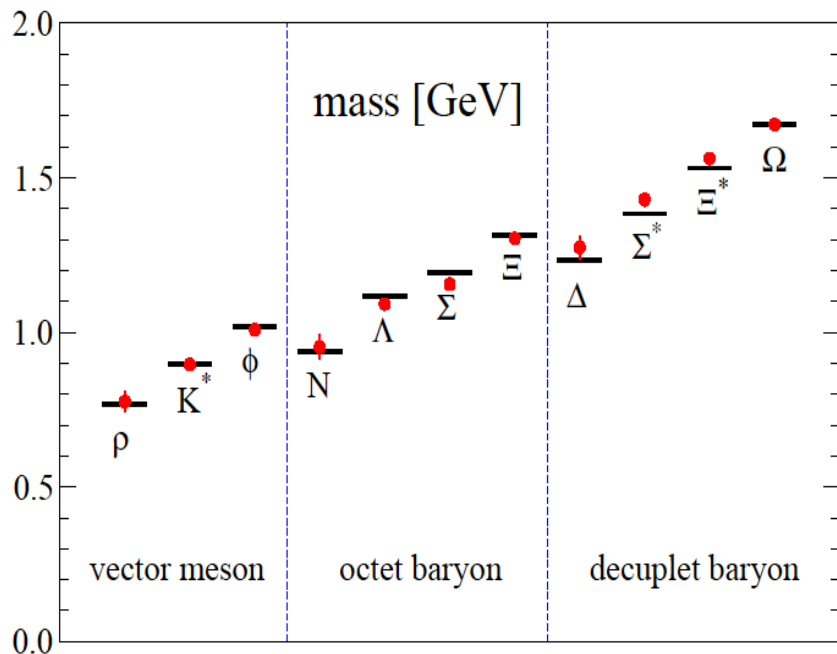


PACS-CS Collaboration,
Phys.Rev.D79(2009)034503

(2+1)-flavor, Wilson
 $L = 2.9$ fm, $a = 0.09$ fm
 $m_{\pi}(\text{min}) = 156$ MeV

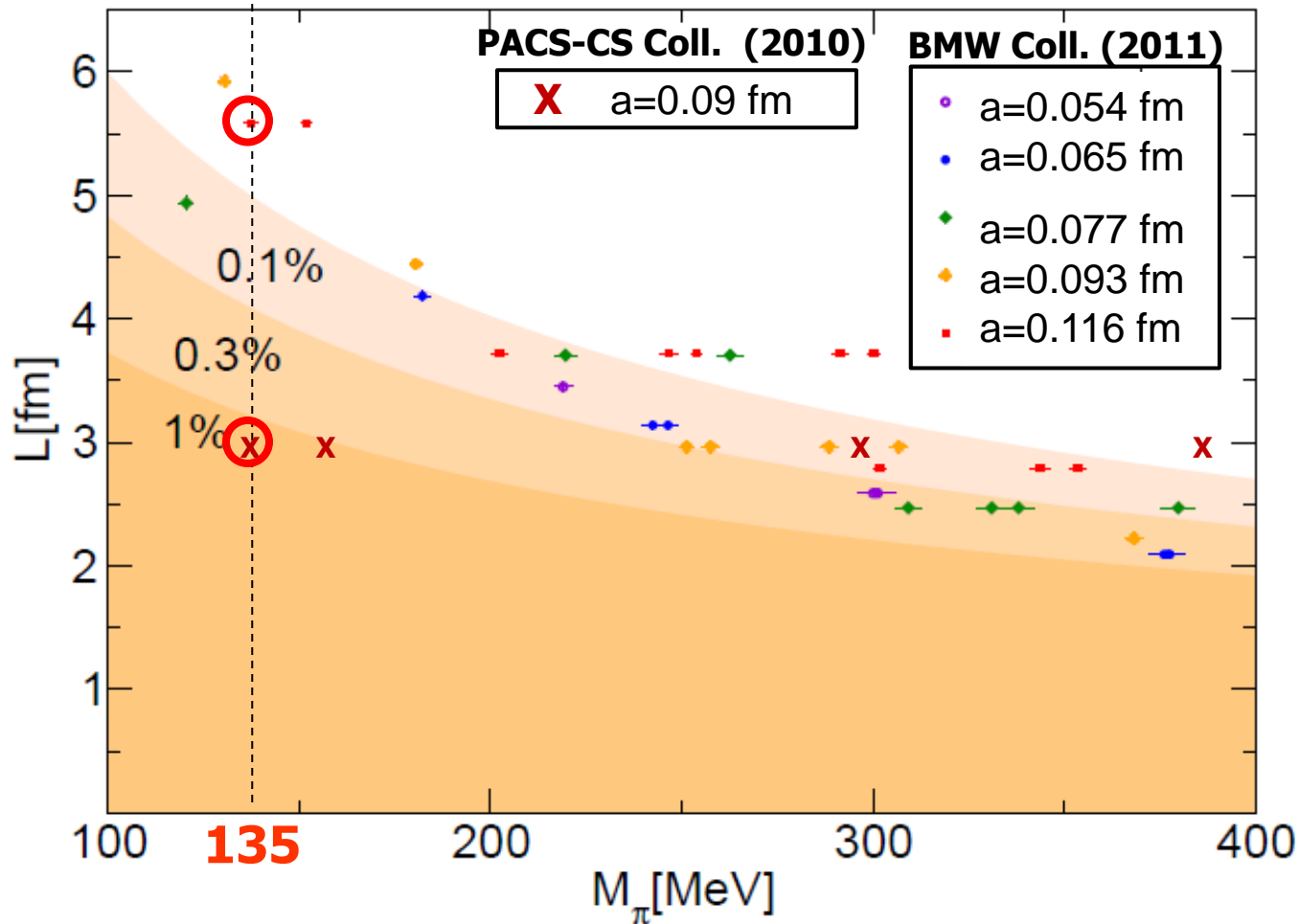
BMW Collaboration,
Science 322 (2008) 1224.

(2+1)-flavor, Wilson
 $L = (2.0 - 4.1)$ fm, $a = 0.065, 0.085, 0.125$ fm
 $m_{\pi}(\text{min}) = 190$ MeV

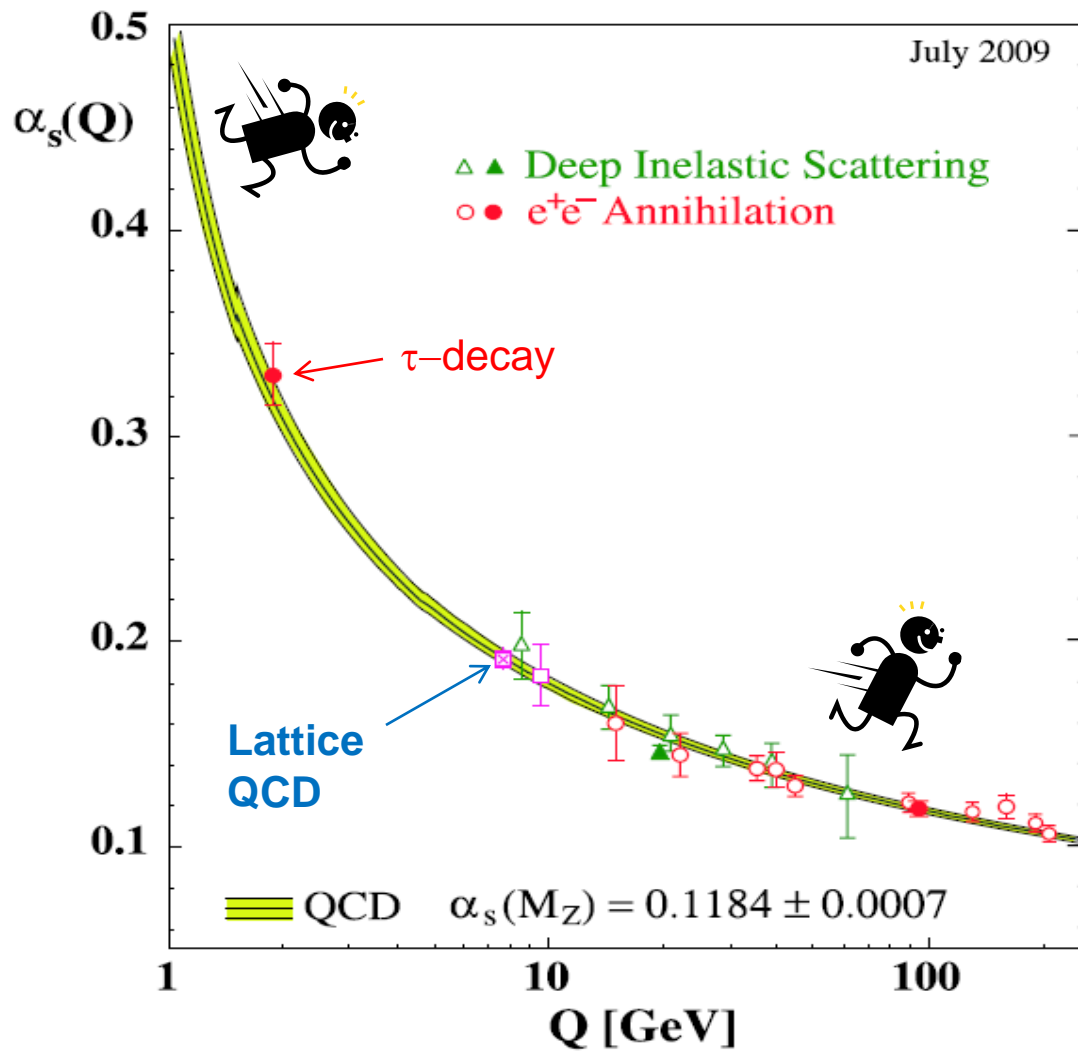


Hadron masses @ 2011

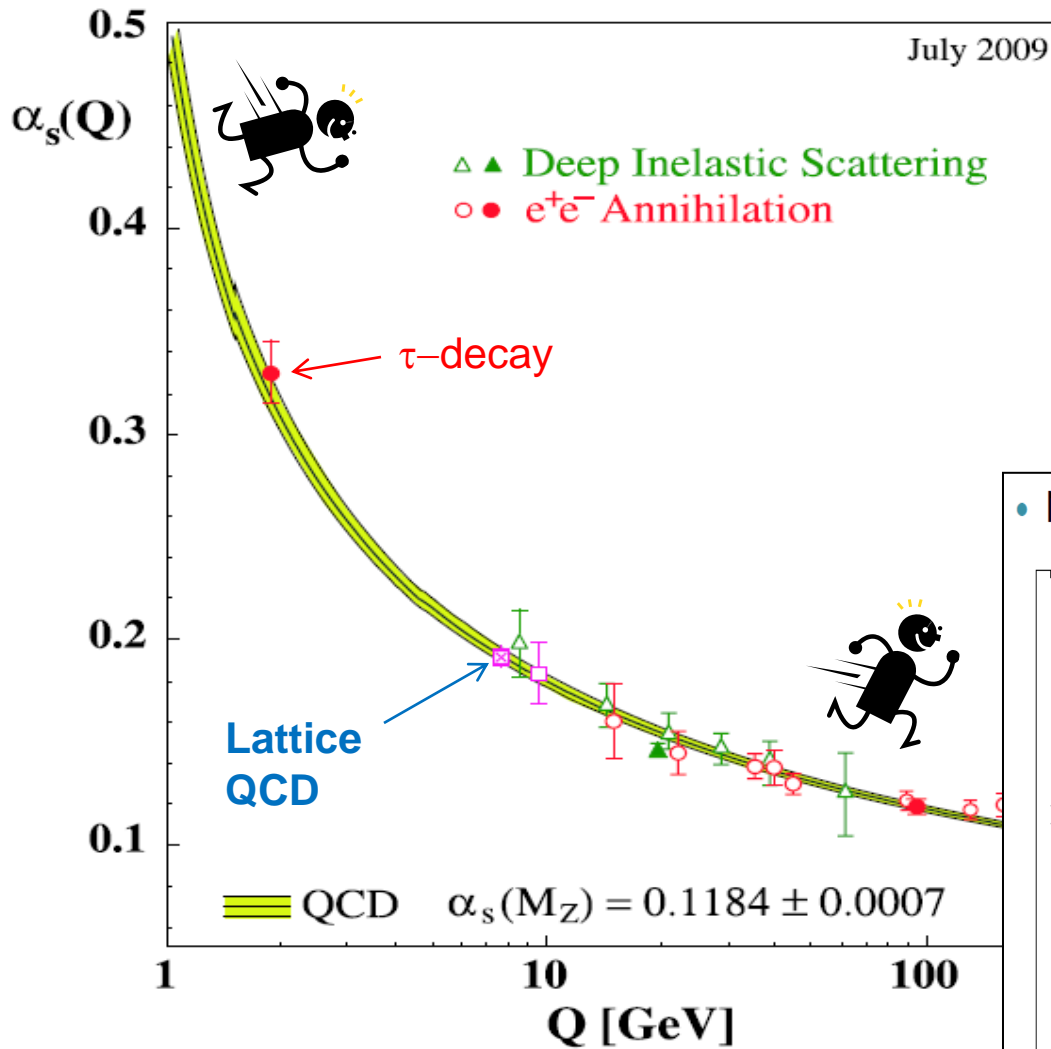
Physical point simulations in (2+1)-flavor QCD



QCD running coupling

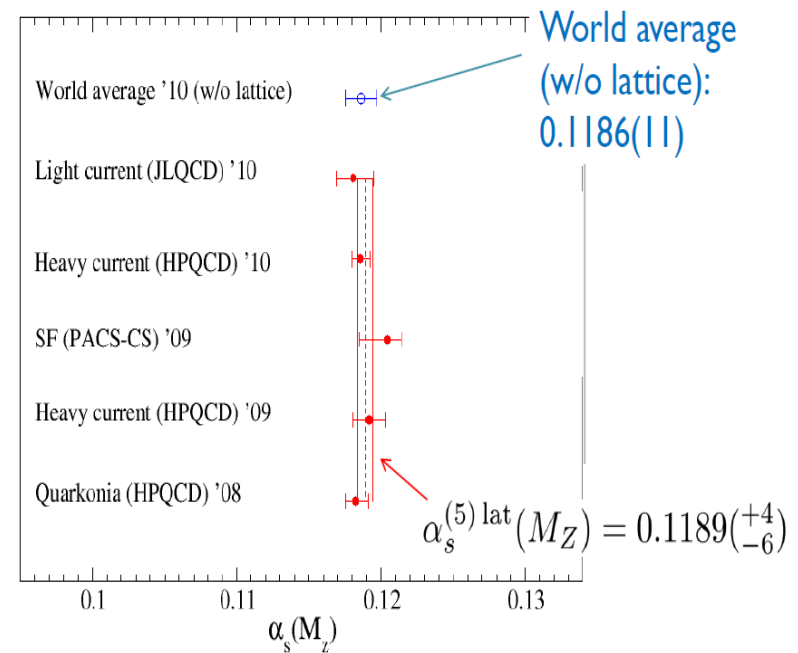


QCD running coupling



Shintani @Lattice2011

- $N_f=2+1$ on the lattice

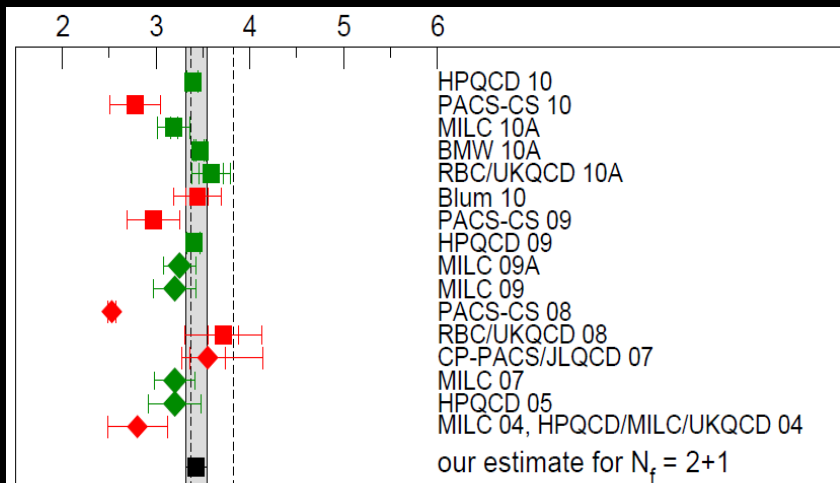


Light quark masses (MSbar, @2GeV)

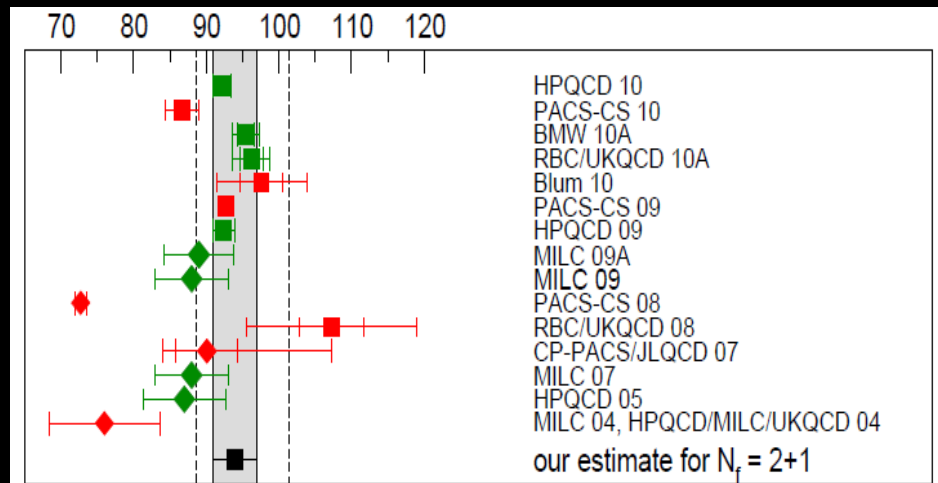
Summary by
FLAG working group
arXiv:1011.4408[hep-lat]



m_{ud} [MeV]



m_s [MeV]

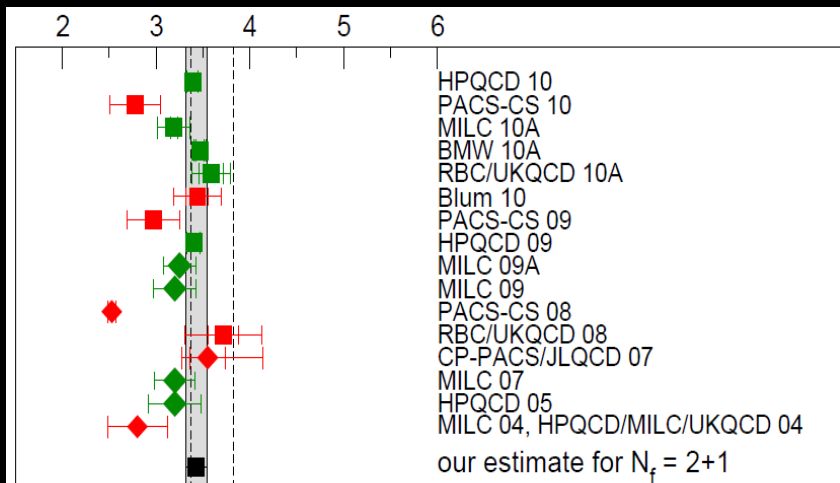


Light quark masses (MSbar, @2GeV)

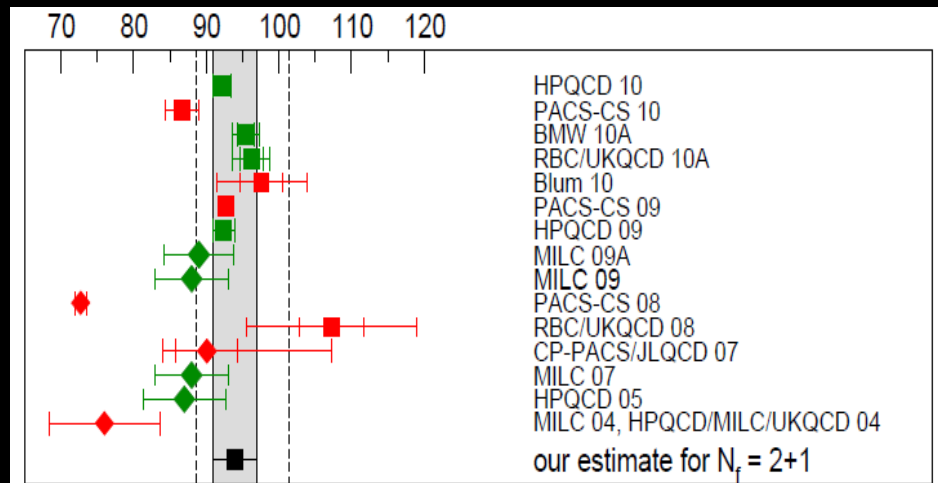
Summary by
FLAG working group
arXiv:1011.4408[hep-lat]



m_{ud} [MeV]



m_s [MeV]

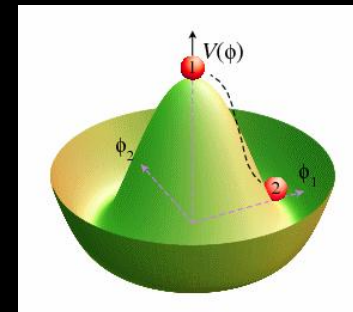
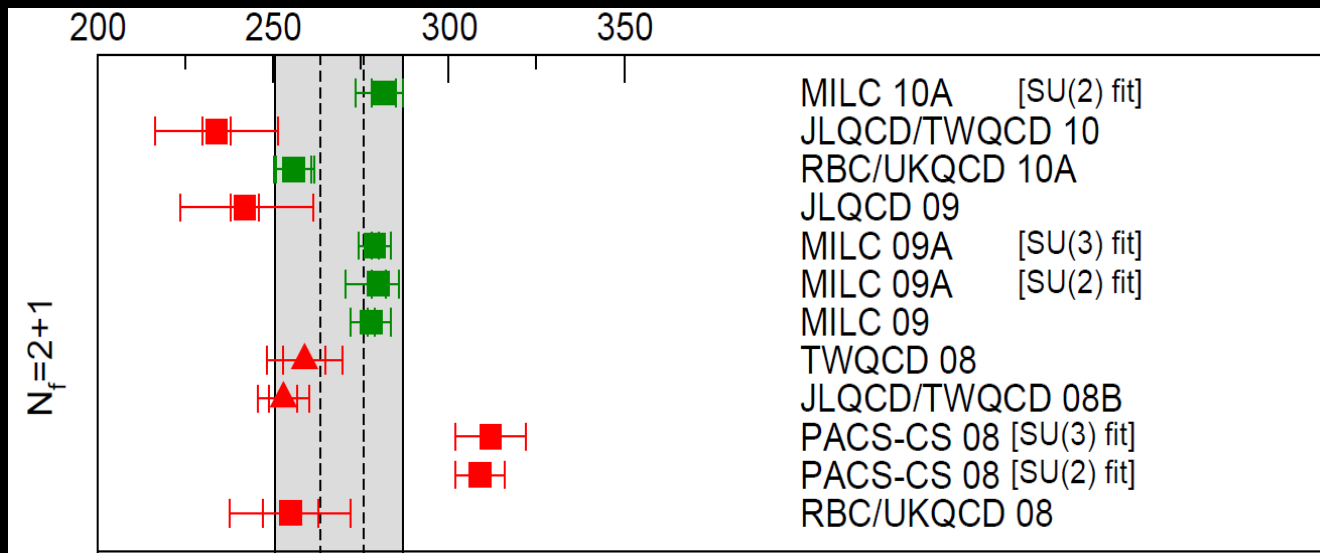


QCD(simulation)+QED(estimate)

N_f	m_u [MeV]	m_d [MeV]	m_{ud} [MeV]	m_s [MeV]	m_s/m_{ud}
2+1	2.19(15)	4.67(20)	3.42(11)	94(3)	27.4(4)

QCD+QED simulation has also been started
Blum et al., Phys. Rev. D82 (2010) 094508

Low energy constants



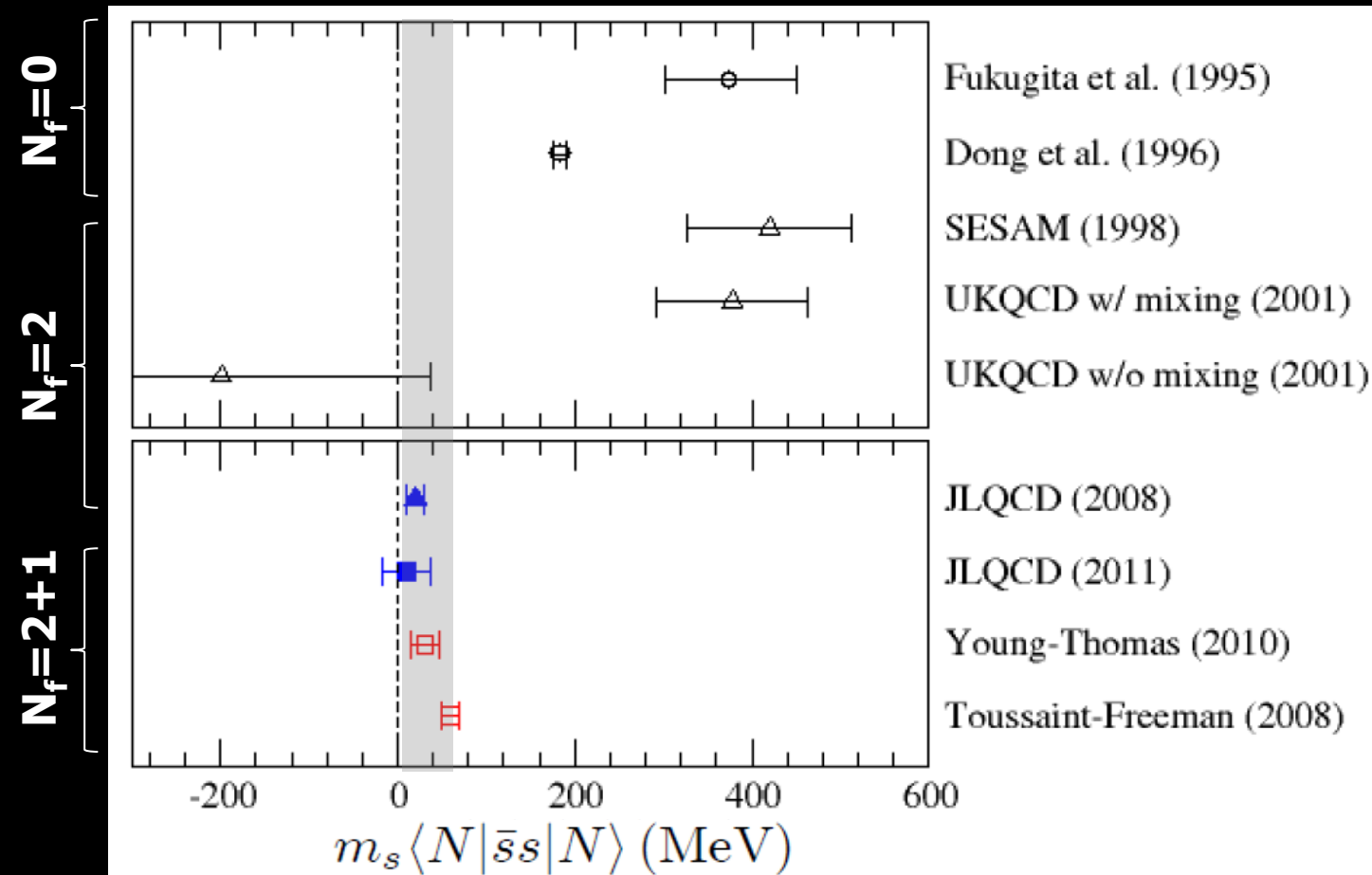
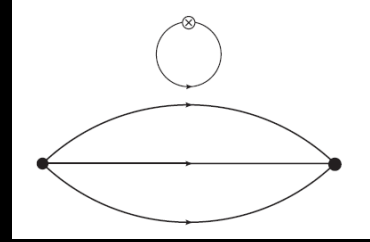
$$[-\langle \bar{q}q \rangle_{2\text{GeV}}]^{1/3} = (250 - 275) [\text{MeV}]$$

$L_{4-8}, f_+(0), f_K/f_\pi, B_K, \text{etc}$



More in arXiv:1011.4408 [hep-lat]
(FLAG working group)

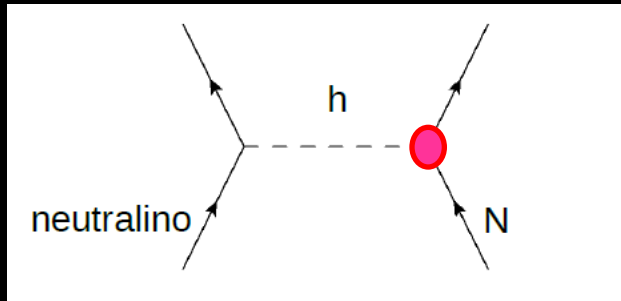
strangeness content of the proton



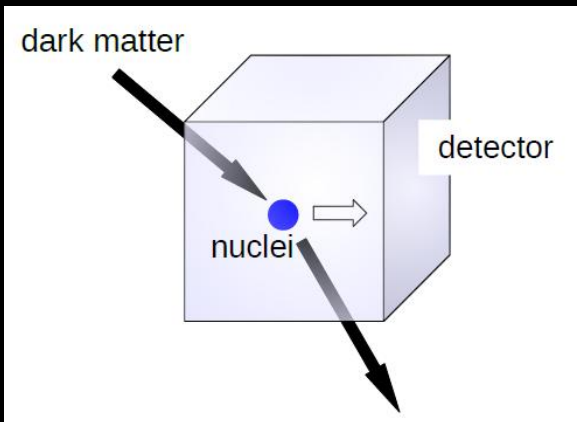
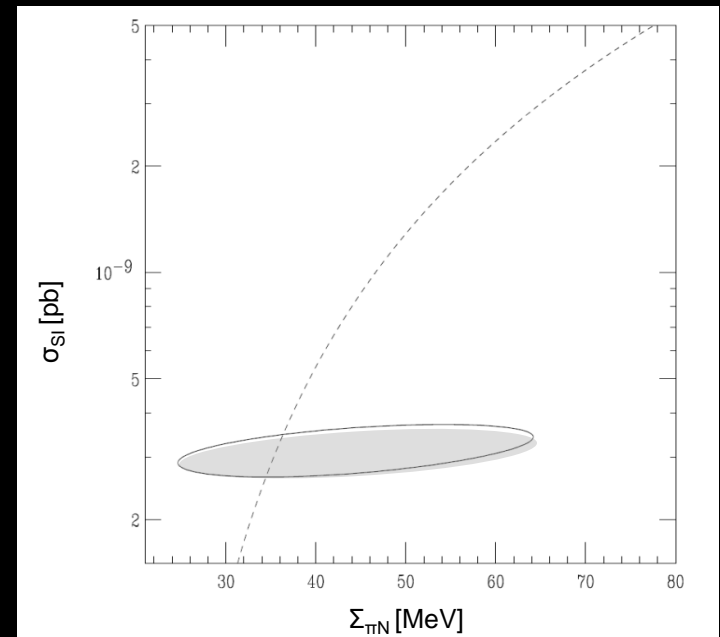
Takeda [JLQCD Coll.],
Phys.Rev. D83 (2011)
114506

$$m_s \langle N | \bar{s}s | N \rangle < 60 \text{ MeV} \Leftrightarrow y = \frac{2 \langle N | \bar{s}s | N \rangle}{\langle N | \bar{u}u + \bar{d}d | N \rangle} < 0.05$$

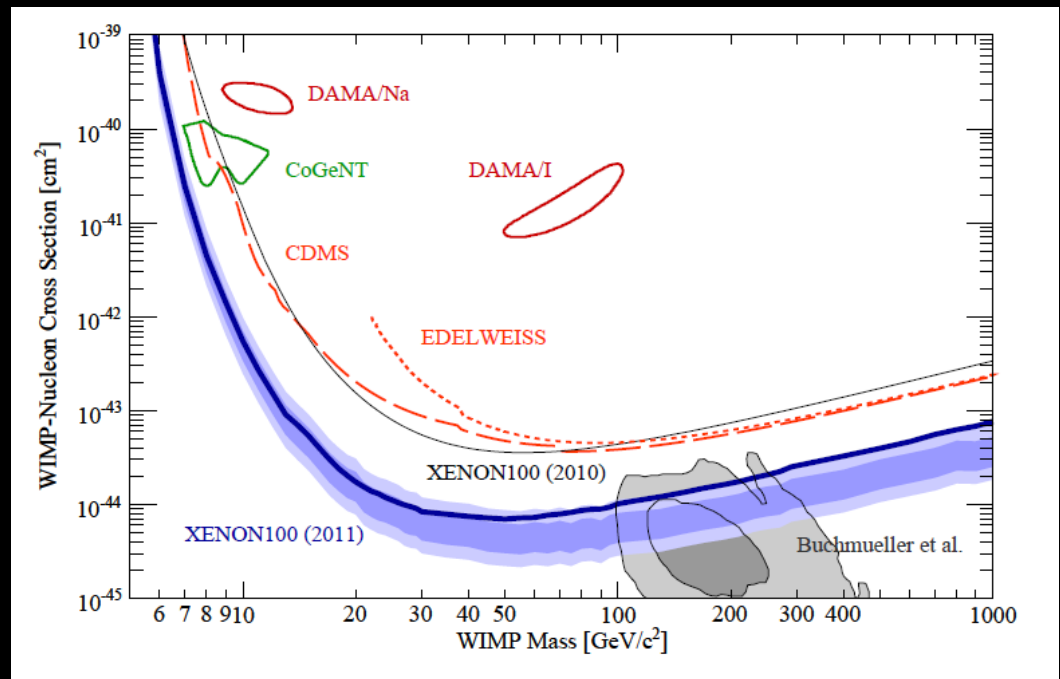
DM – Nucleon Interaction



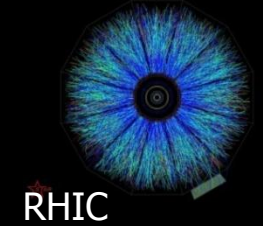
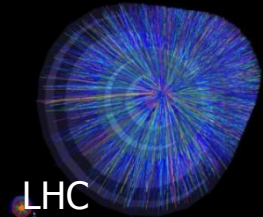
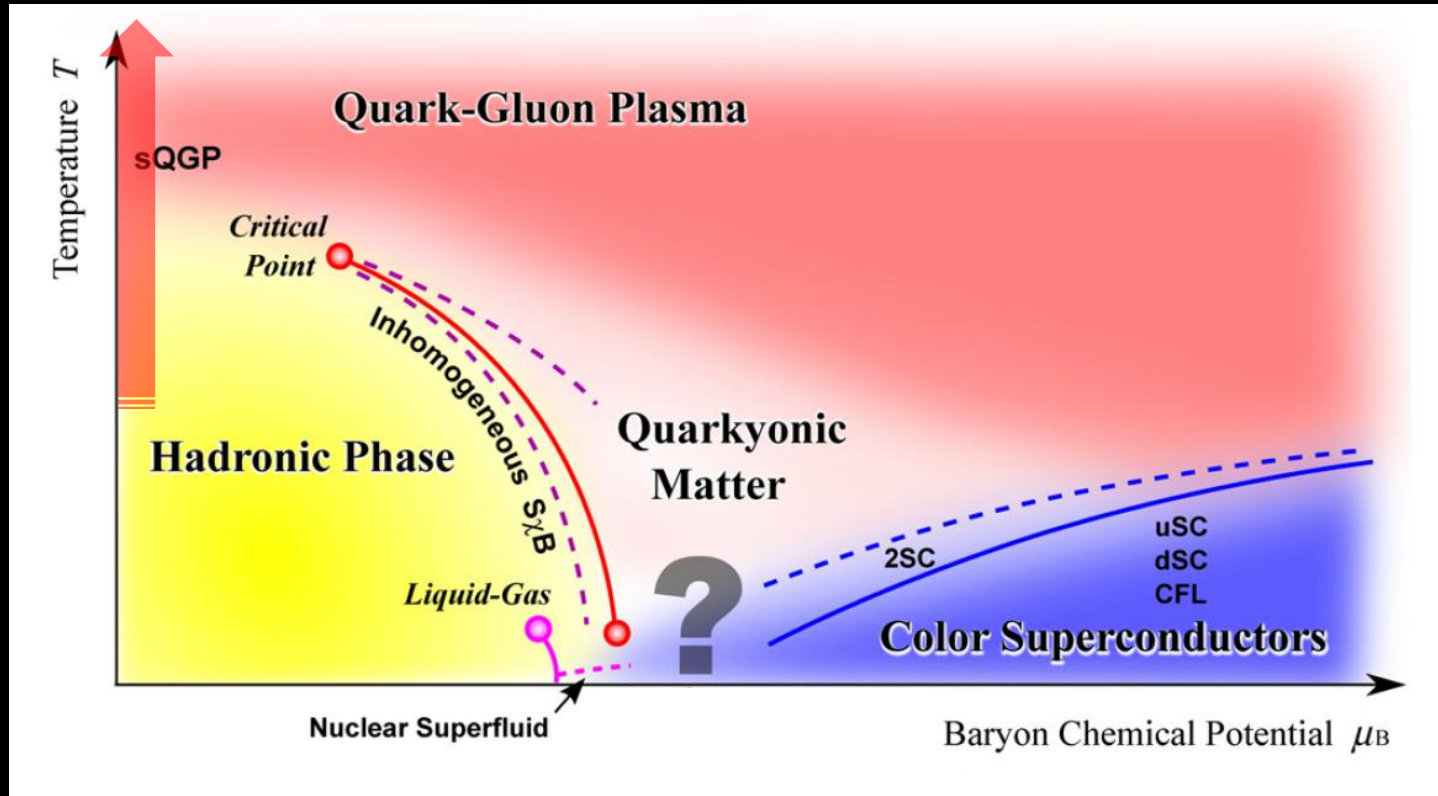
Giedt, Thomas, Young
Phys. Rev. Lett. 103 (2009) 201802



XENON100 Coll.:
Phys.Rev.Lett. 105 (2010) 131302
arXiv:1104.2549 [astro-ph.CO]
arXiv:1107.2155 [astro-ph.IM]

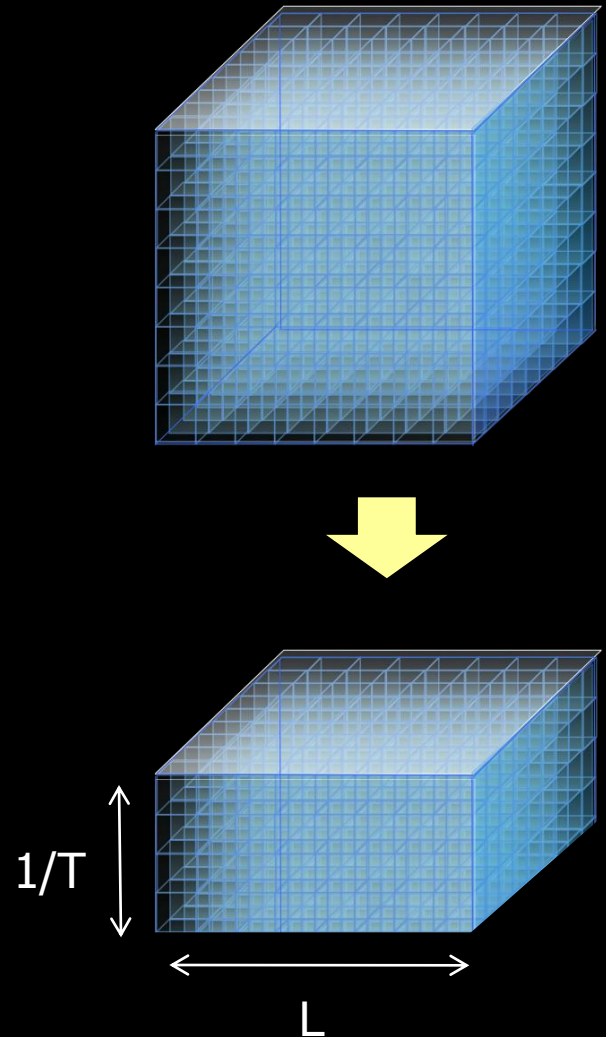
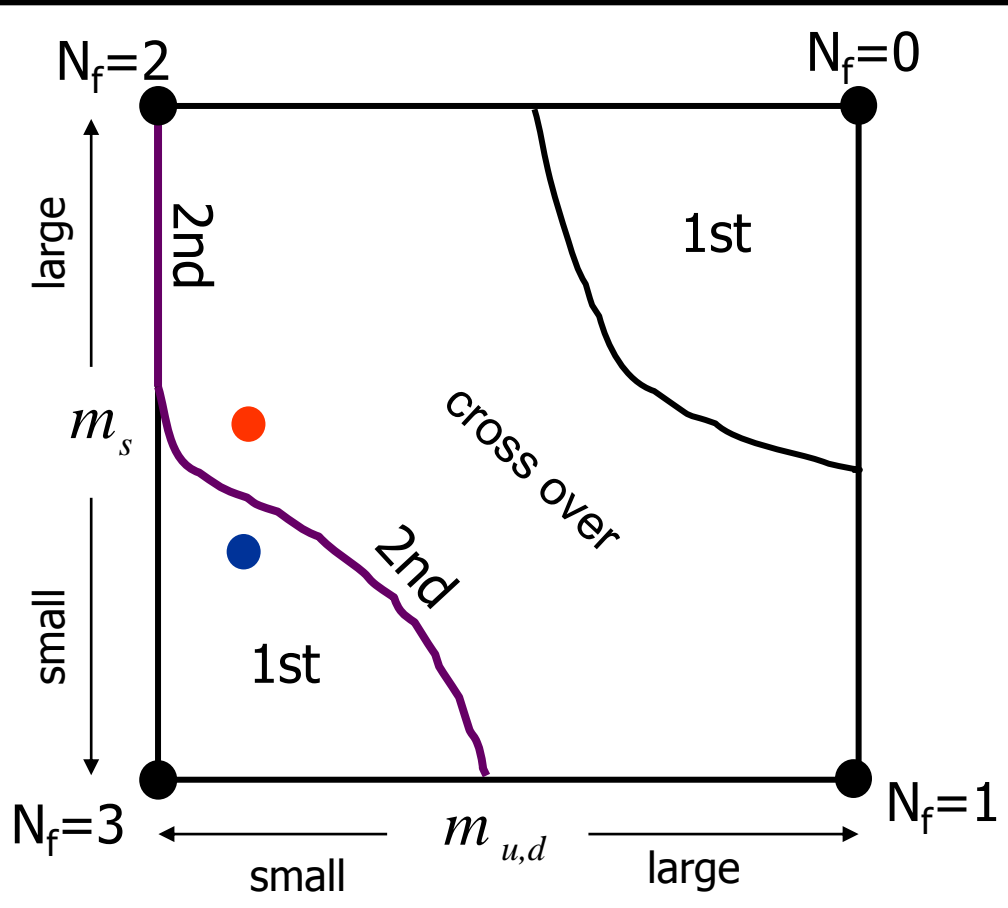


Thermal Lattice QCD



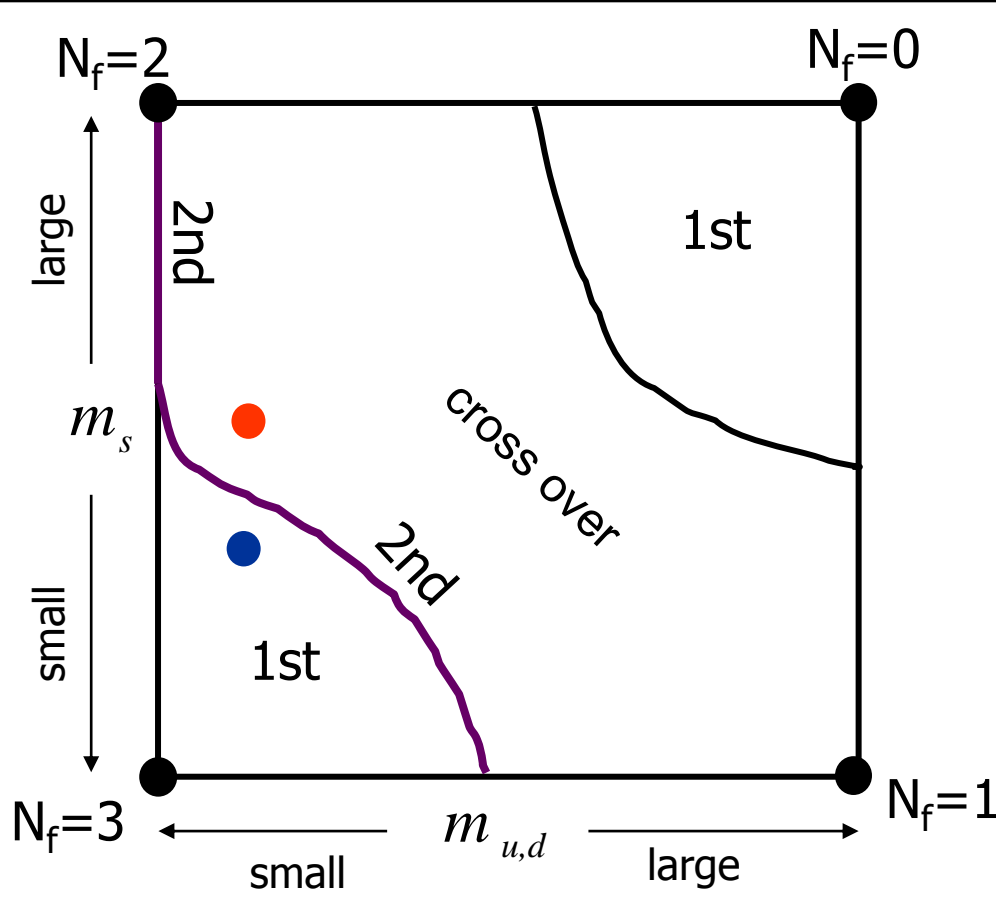
Thermal QCD transition at $\mu=0$

Columbia plot



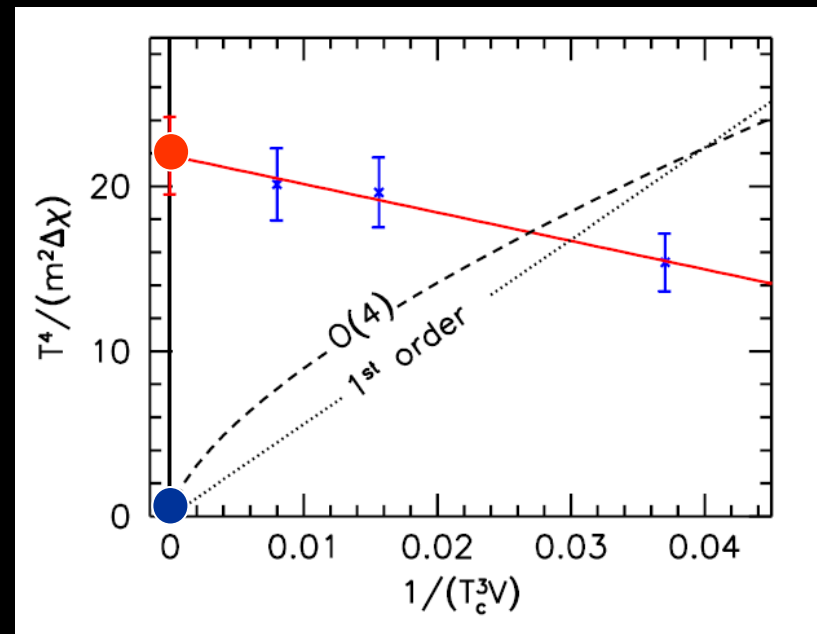
Thermal QCD transition at $\mu=0$

Columbia plot



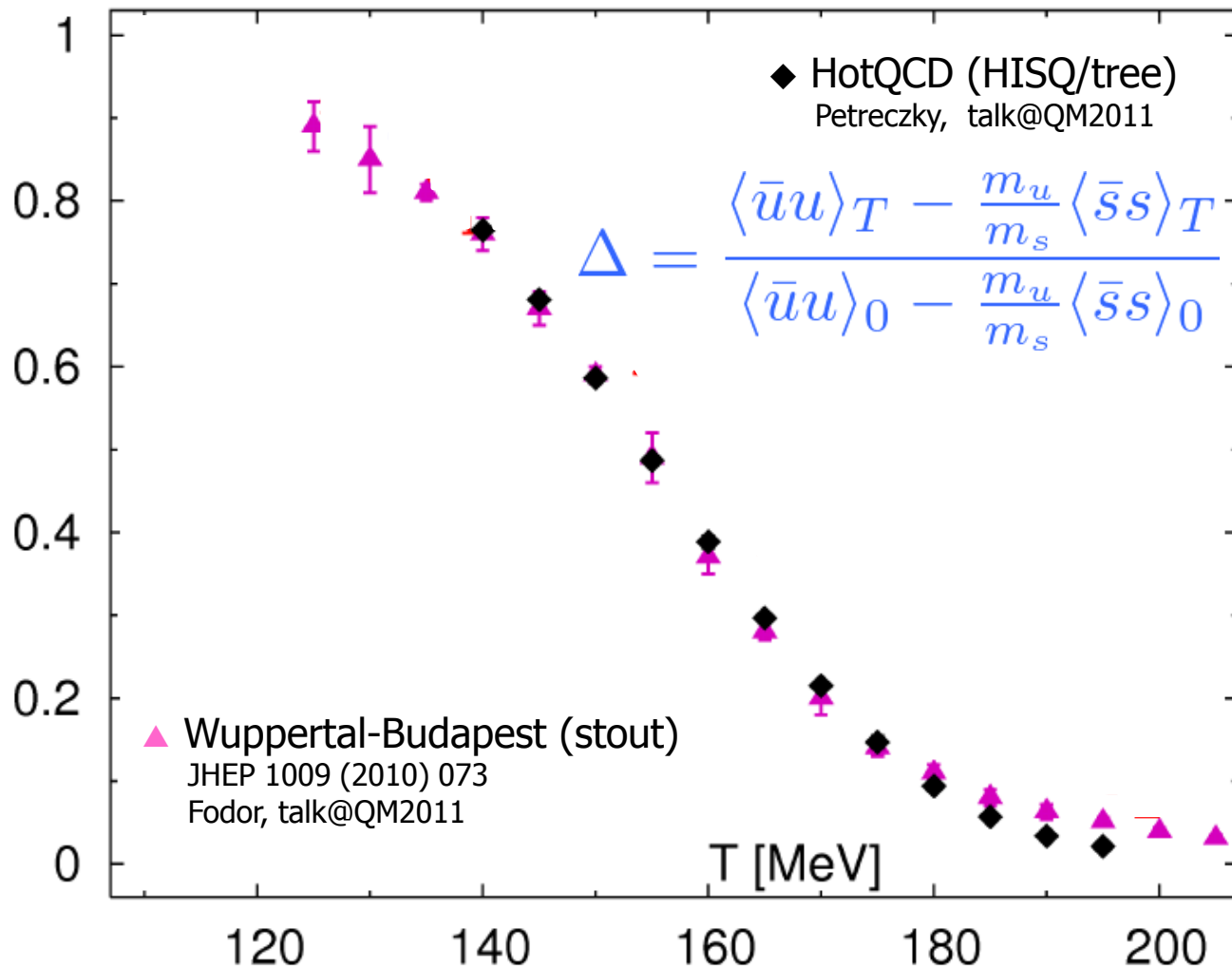
Finite size scaling

$$\chi_m = \frac{\partial^2 P}{\partial m_{ud}^2} \sim \begin{cases} V & \text{1st order} \\ V^{2/3} & \text{2nd order} \\ V^0 & \text{crossover} \end{cases}$$

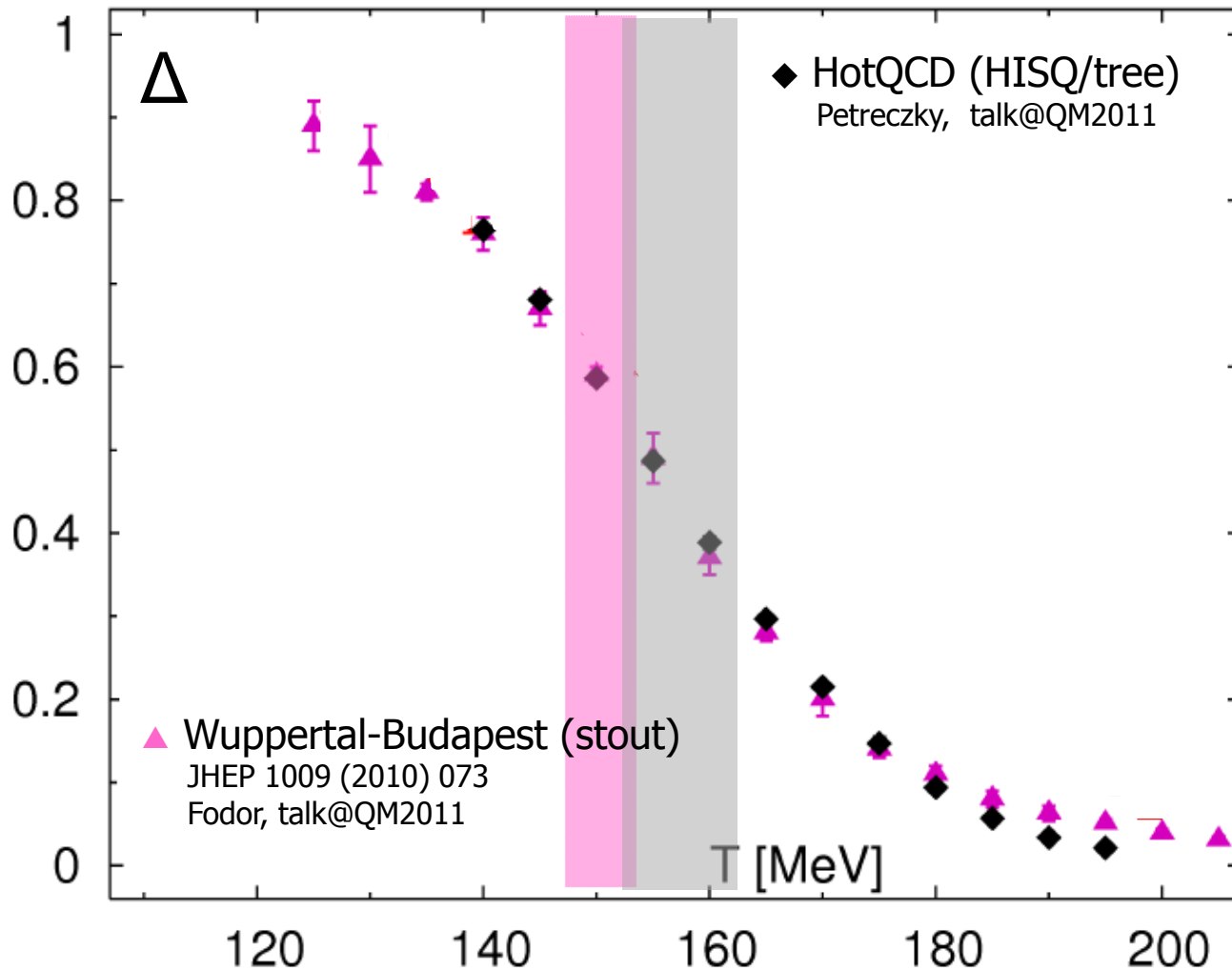


Budapest group, Nature 443 (2006) 675
Staggered, (2+1)-flavor, physical mass

Thermal chiral condensate at $m_\pi=135$ MeV



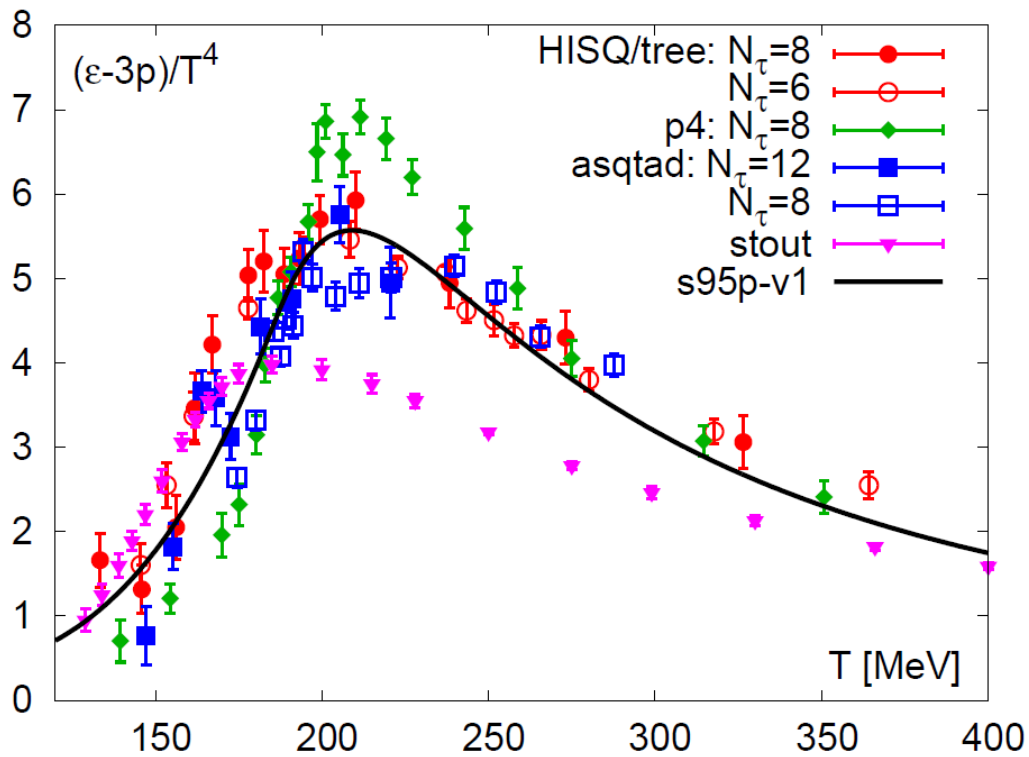
Pseudo-critical temperature T_{pc}



Chiral susceptibility peak (max. fluctuation) $\Rightarrow T_{pc}=150-160\text{MeV}$

Equation of state (EOS) : $p(T)$, $\epsilon(T)$

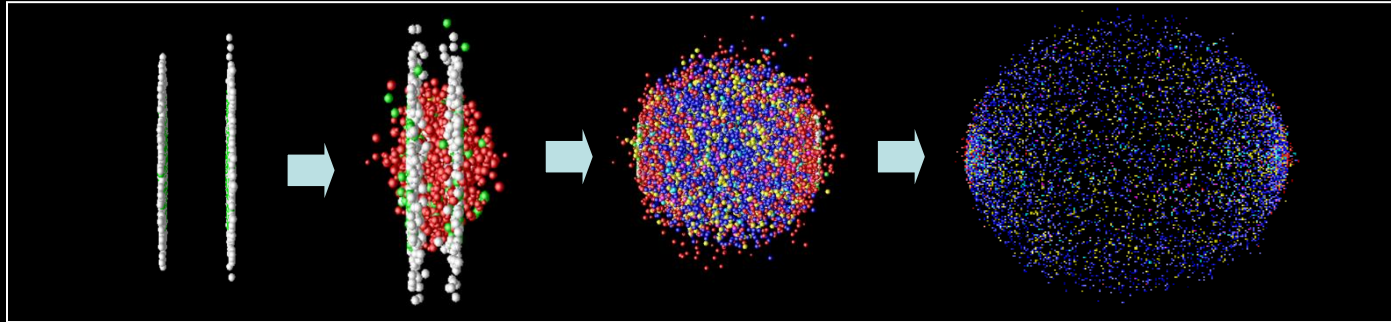
$$\frac{p(T)}{T^4} - \frac{P(T_0)}{T_0^4} = \int_{T_0}^T \frac{dT'}{T'} \frac{\epsilon(T') - 3p(T')}{T'^4}$$



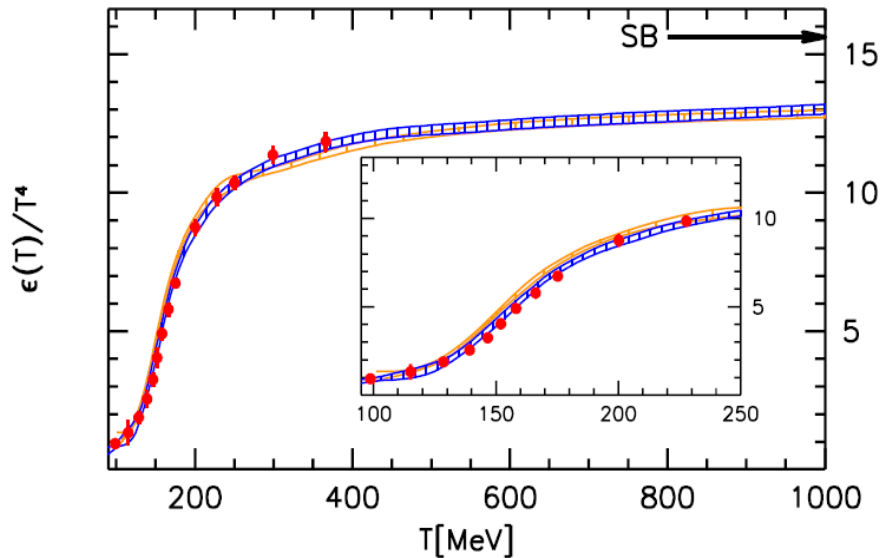
▼ Wuppertal-Budapest Coll.,
 JHEP 1011 (2010) 77
● ● ■ HotQCD Coll.,
 ArXiv: 1012.1257 [hep-lat]

- Discrepancy exists among different improved staggered actions
- Hopefully resolved in a $\rightarrow 0$ limit
- Need to be checked by Wilson, Domain-wall, Overlap

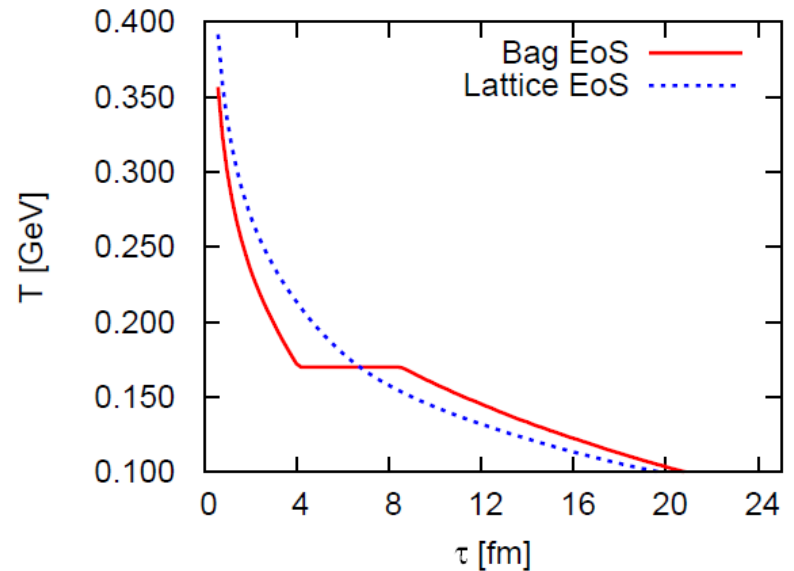
LQCD-EOS applied to RHIC



3D hydro.
 $\partial_\mu T^{\mu\nu}(p, \varepsilon)=0$

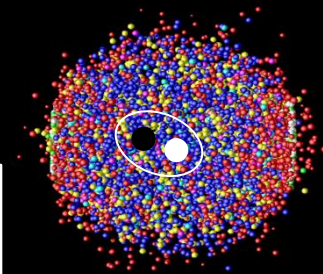


Wuppertal-Budapest's LQCD EOS
JHEP 1011 (2010) 77



Akamatsu, Hamagaki, Hirano, Hatsuda
arXiv:1107.36[nucl-th]

Heavy QQbar in QGP



Matsui & Satz,
PLB (1986)

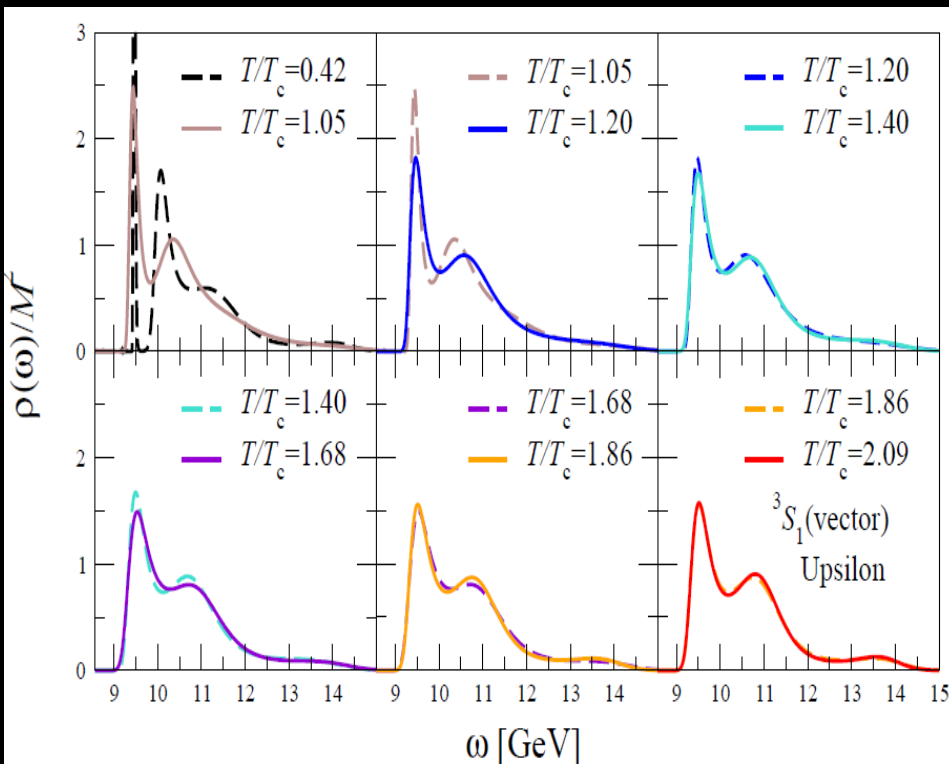
LQCD + Bayesian analysis \Rightarrow spectral function \Rightarrow dilepton rate

Asakawa, Nakahara, Hatsuda, Prog. Part. Nucl. Phys.46 (2001) 469

Υ spectral function on the lattice

$N_f=2$, $a_s=0.162$ fm, $\xi=6$, $L=1.94$ fm

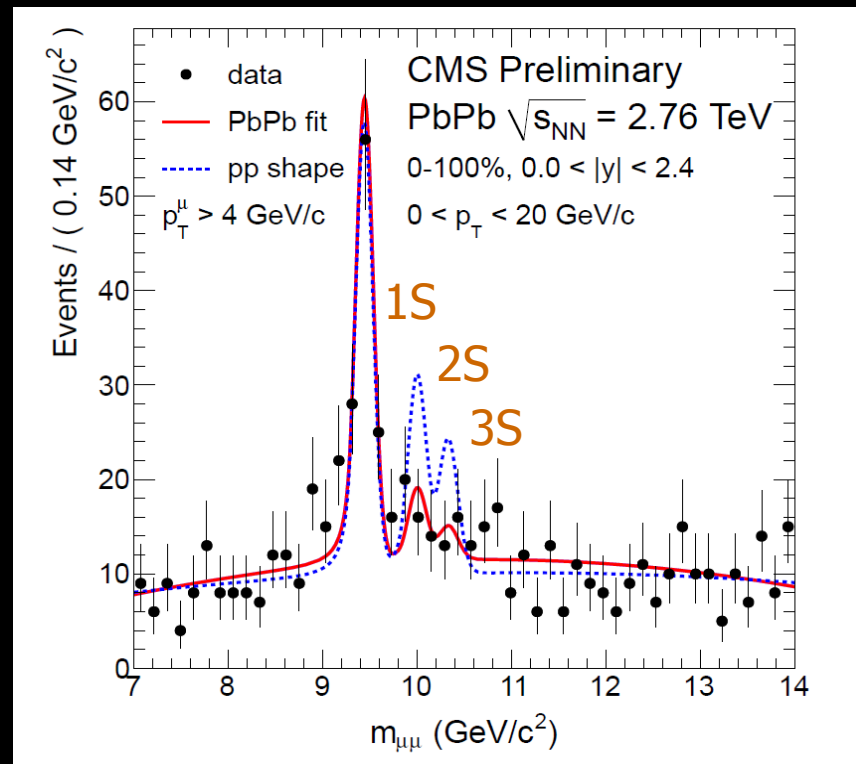
Aarts et al., arXiv:1109.4496 [hep-lat]



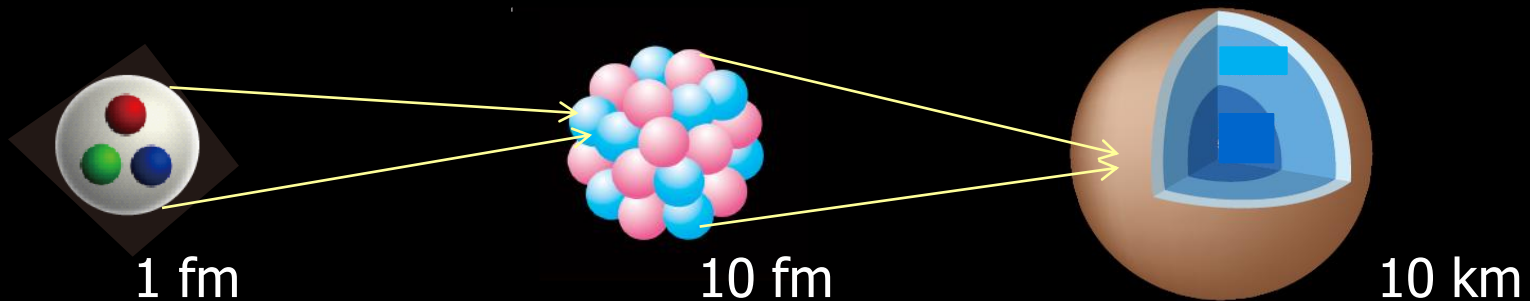
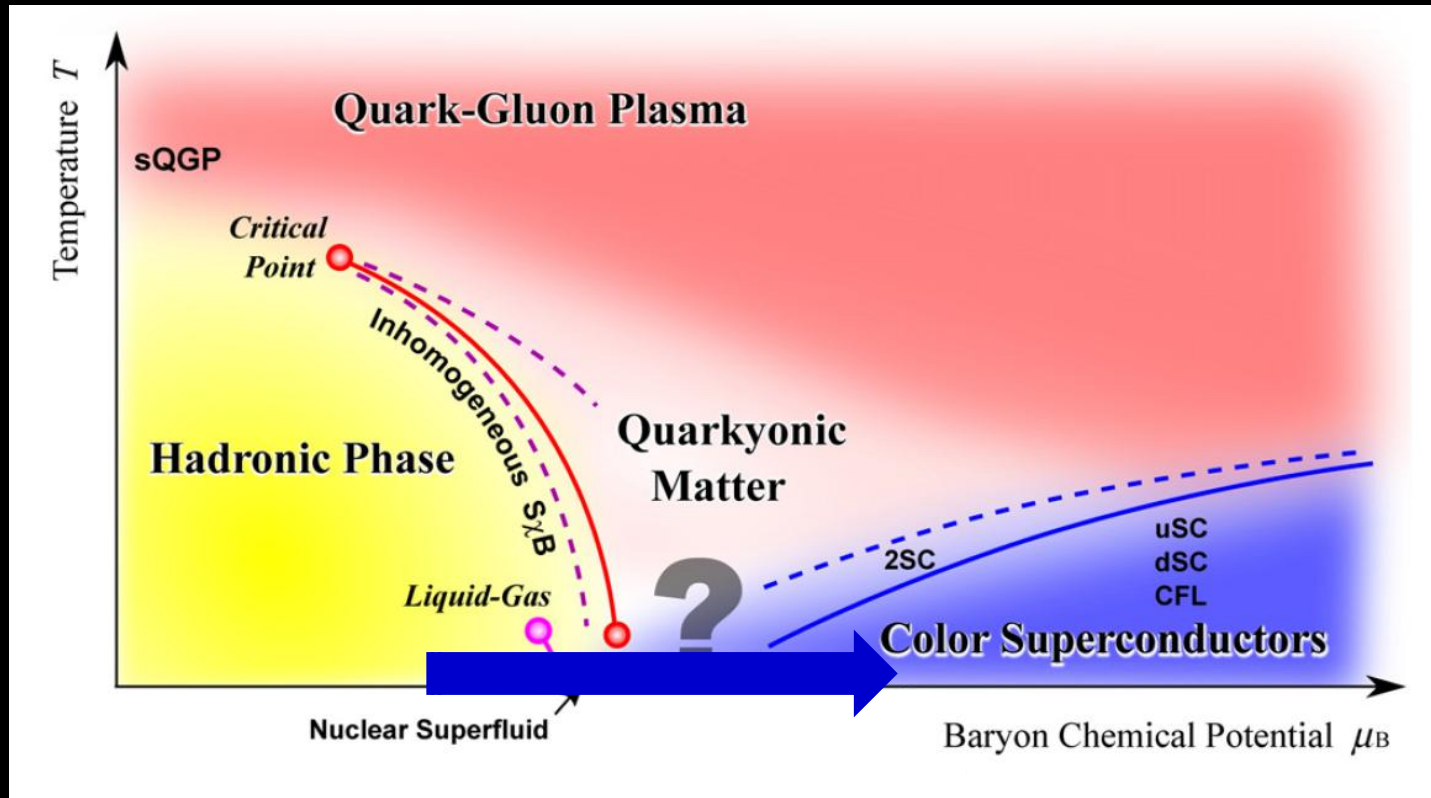
Υ in pp and PbPb collisions at LHC

CMS Coll., arXiv:1105.4894[nucl-ex]

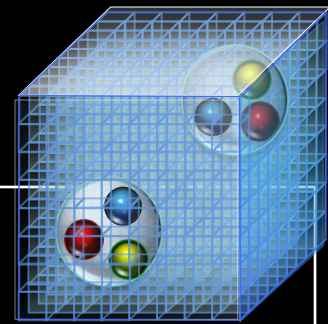
Wyslouch (for CMS), arXiv:1107.2895[nucl-ex]



Nuclear Lattice QCD



Nuclear Force from LQCD



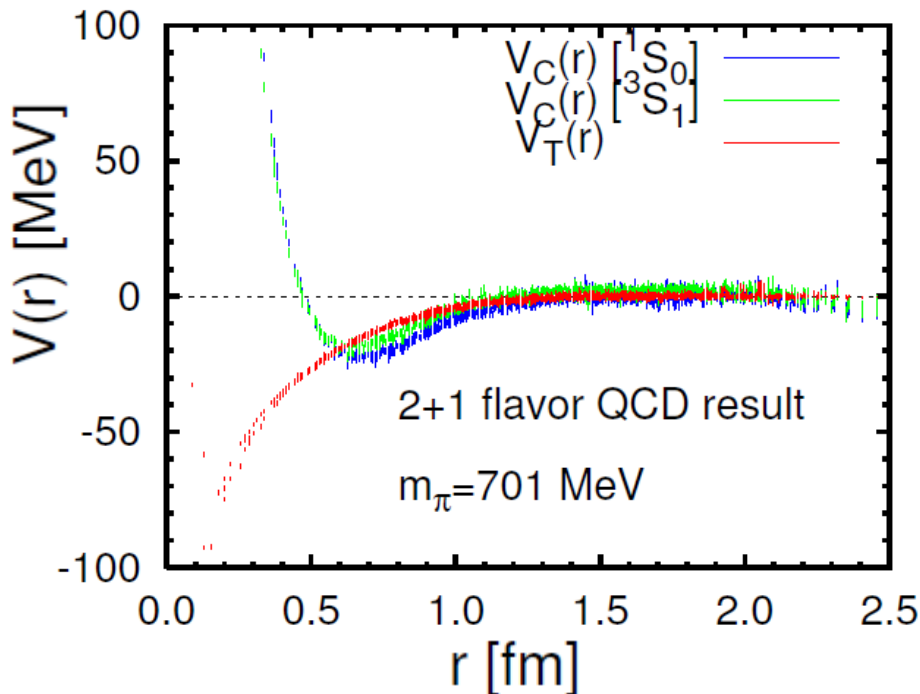
1. Low energy NN int. \Leftrightarrow NN potential

$$V(\vec{r}, \nabla) = V_C(r) + S_{12}V_T(r) + \vec{L} \cdot \vec{S} V_{LS}(r) + \{V_D(r), \nabla^2\} + \dots$$

2. NN potential from NN "wave function"

$$\phi(\vec{r}) = \langle 0 | N(\vec{x} + \vec{r}) N(\vec{x}) | 6q \rangle$$

Ishii, Aoki, Hatusda,
Phys.Rev.Lett. 99 (2007) 022001

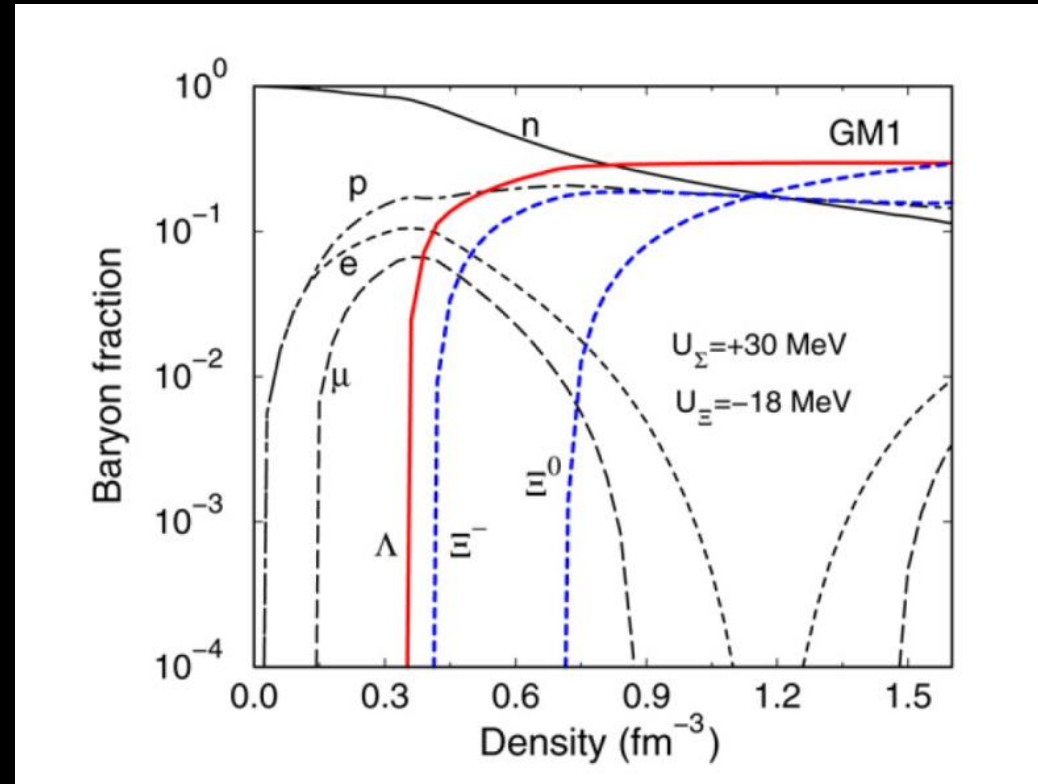
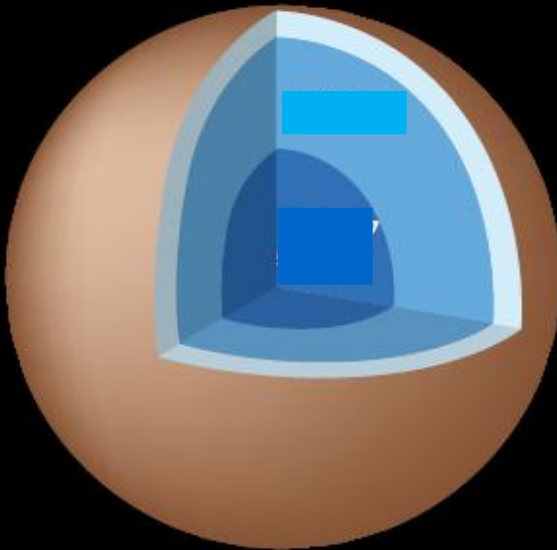


repulsive core
+ attractive well
+ tensor force
from LQCD

Inner core of neutron star

-- role of 2-body and 3-body forces --

Radius ~ 10 km
 Mass \sim solar mass
 Central density $\sim 10^{12}$ kg/cm³

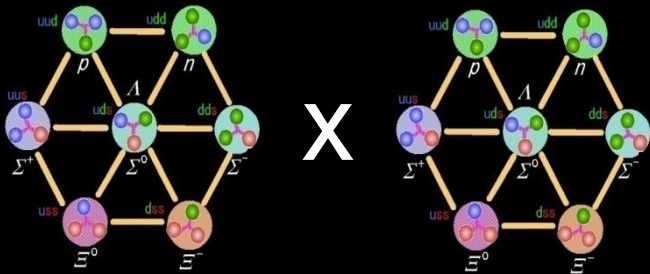


Schaffner-Bielich, Nucl. Phys.A 835, 279 (2010)

YN interaction \Leftrightarrow onset of hyperon mixture
 NNN (BBB) interaction \Leftrightarrow large max mass (e.g. $1.97(4) M_{\odot}$)

BB interactions in 3-flavor LQCD

1. Numerical experiments of YN & YY interactions
(not easily accessible in laboratory experiments)
2. Physical origin of the short range NN repulsion
3. Fate of H-dibaryon

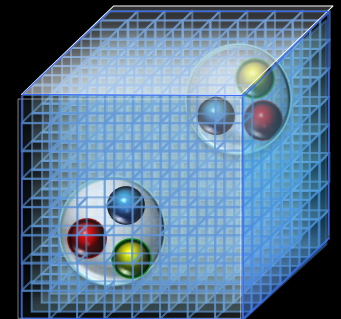
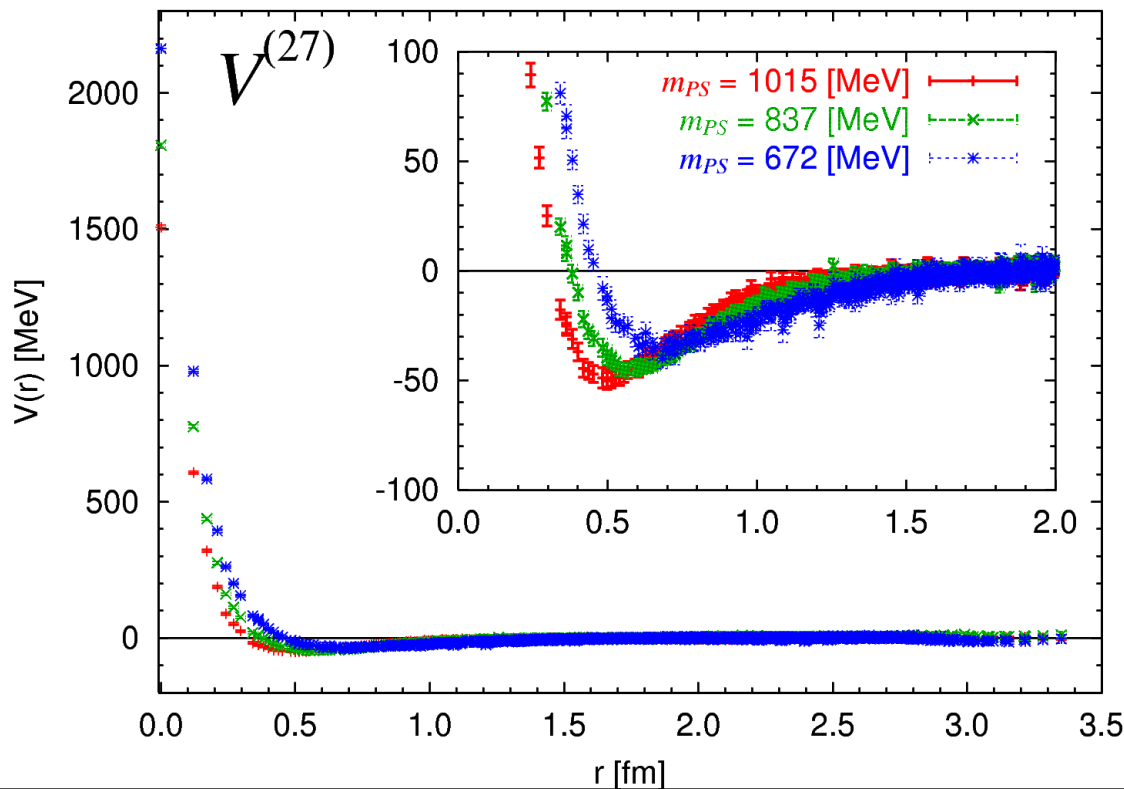


$$8 \times 8 = \underbrace{27 + 8s + 1}_{\text{Symmetric}} + \underbrace{10^* + 10 + 8a}_{\text{Anti-symmetric}}$$

Six independent potentials in the flavor-basis

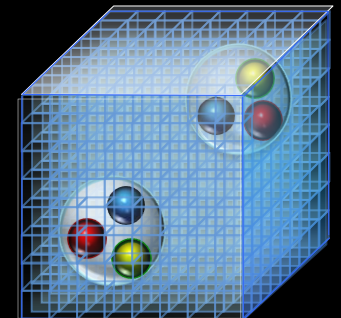
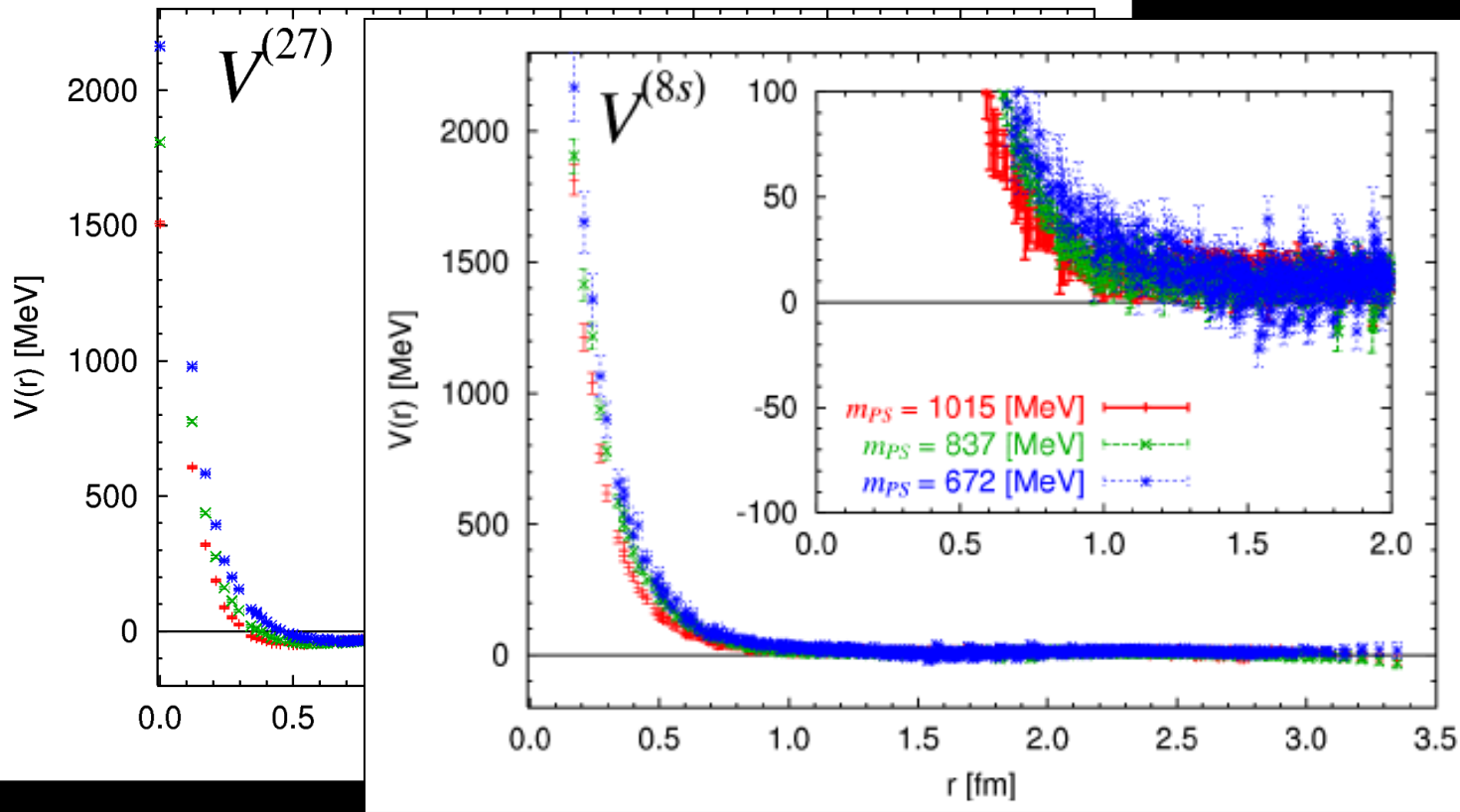
BB potentials in flavor-basis

Inoue et al. [HAL QCD Coll.]
Phys. Rev. Lett. 106 (2011) 162002



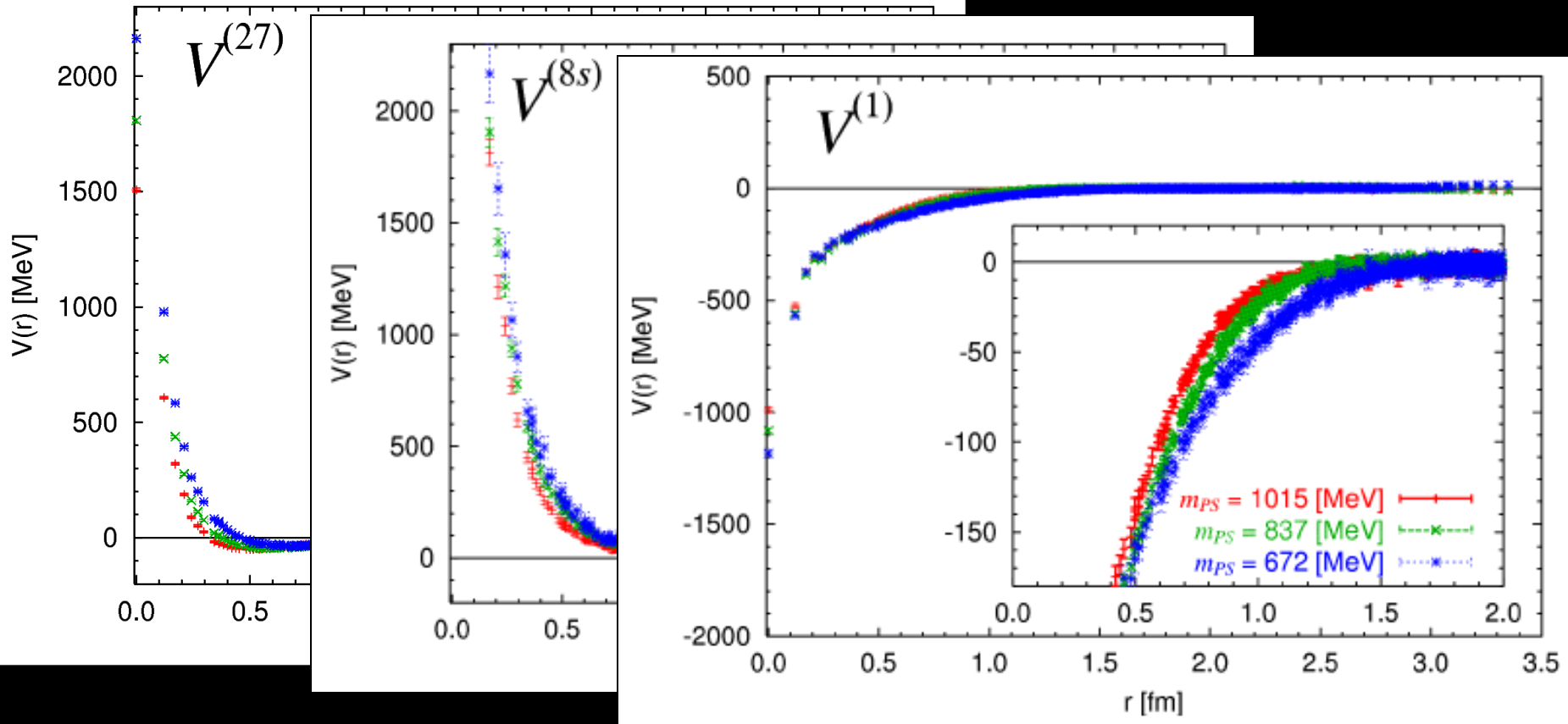
BB potentials in flavor-basis

Inoue et al. [HAL QCD Coll.]
Phys. Rev. Lett. 106 (2011) 162002



BB potentials in flavor-basis

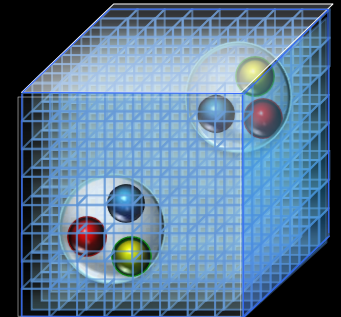
Inoue et al. [HAL QCD Coll.]
Phys. Rev. Lett. 106 (2011) 162002



Short range BB int. \Leftrightarrow Quark Pauli principle

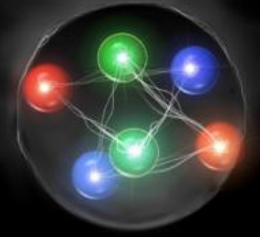
1 : allowed,
27 : partially blocked, 8_s : blocked

c.f. constituent quark model (Oka, Yazaki, Shimizu, ...)

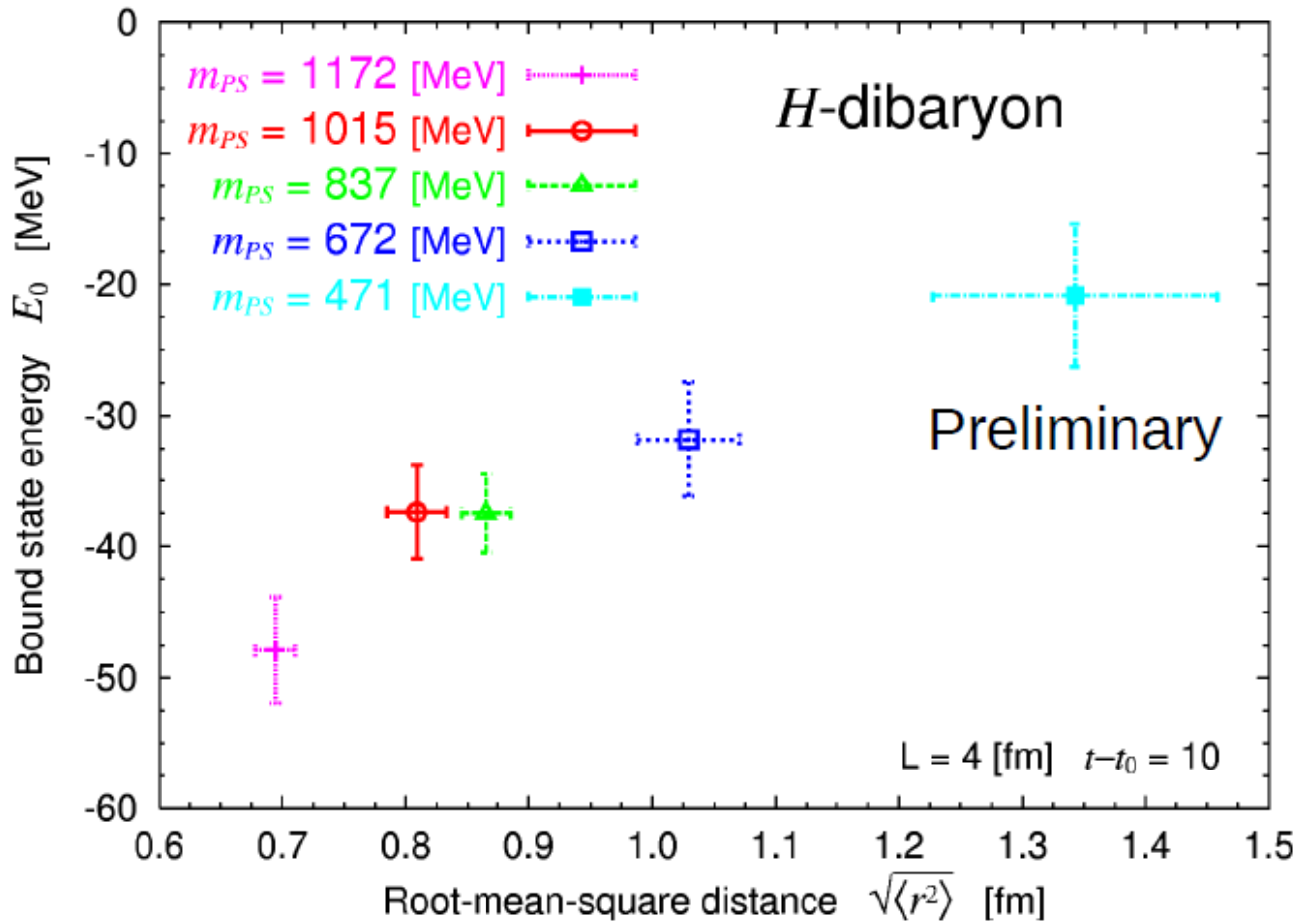


H-dibaryon from LQCD

-- binding energy vs. size --



Jaffe, Phys. Rev. Lett.
38 (1977) 195

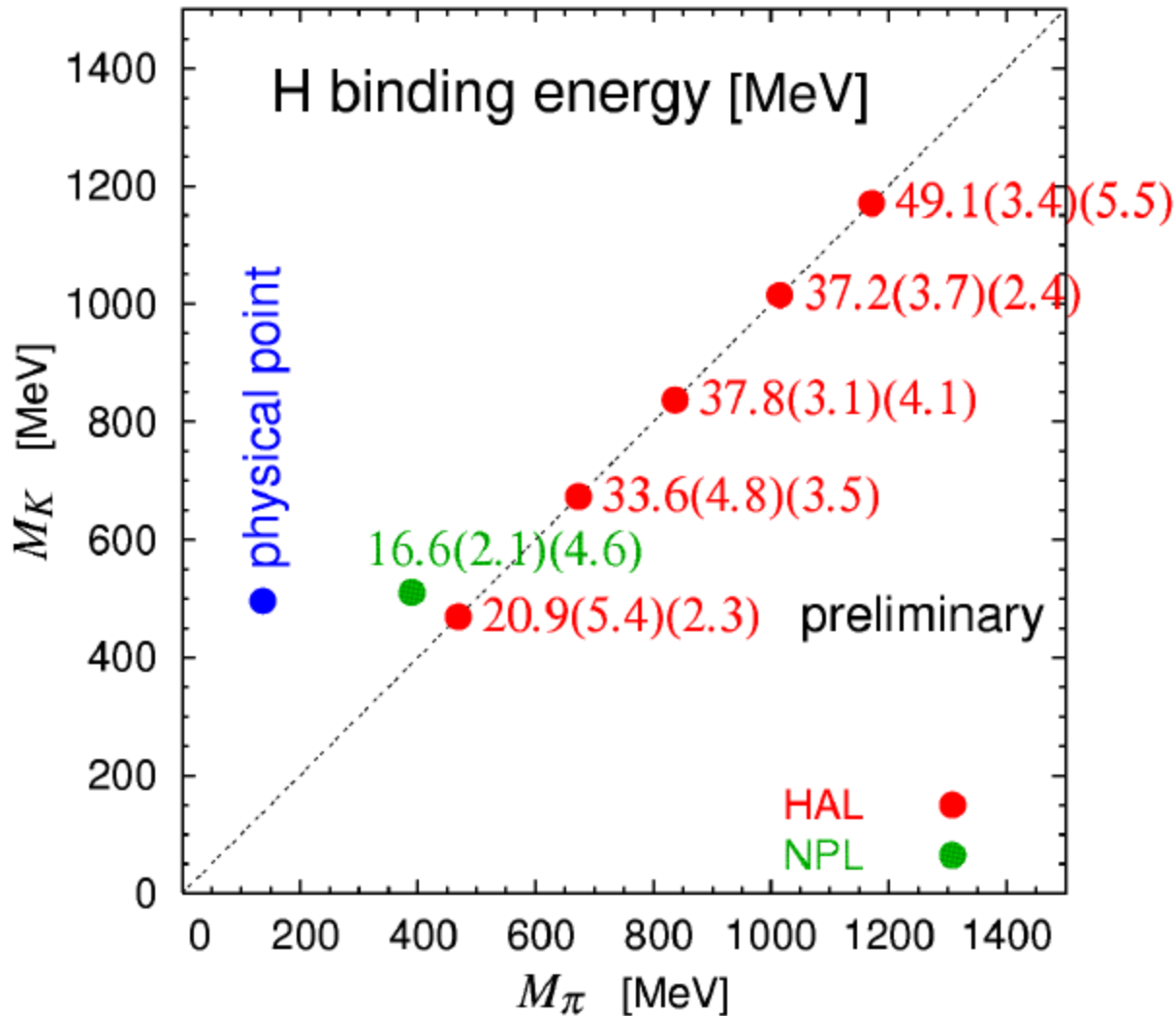


Inoue et al. [HAL QCD Coll.] Phys. Rev. Lett. 106 (2011) 162002

H dibaryon from LQCD



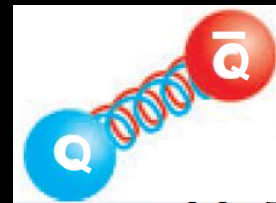
Jaffe, Phys. Rev. Lett.
38 (1977) 195



Inoue et al. [HAL QCD Coll.] Phys. Rev. Lett. 106 (2011) 162002

Beane et al. [NPLQCD Coll.] Phys. Rev. Lett. 106 (2011) 162001

“Constituent quark model” from LQCD ?



Quenched QCD: Coulomb gauge

$L = 3.3$ fm, $a = 0.104$ fm

(2+1)-flavor QCD: Coulomb gauge

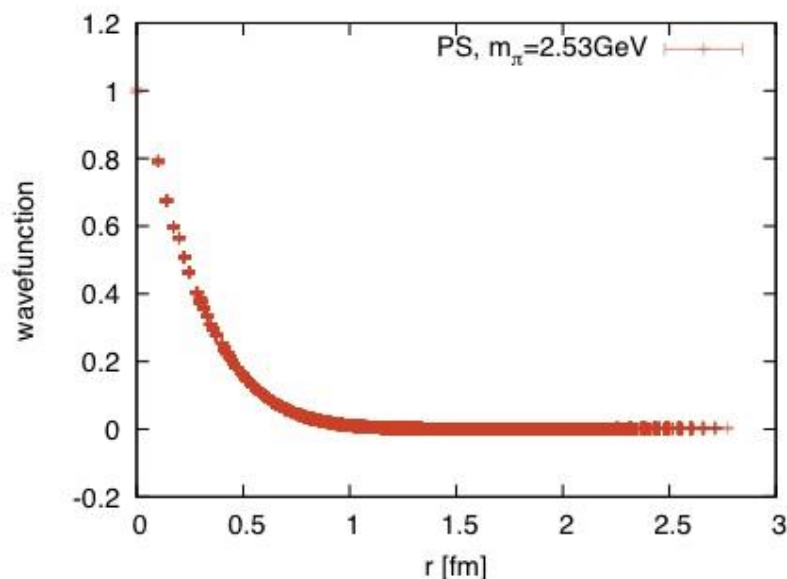
$L = 3$ fm, $a = 0.09$ fm

Ikeda & Iida, 1102.2097 [hep-lat]

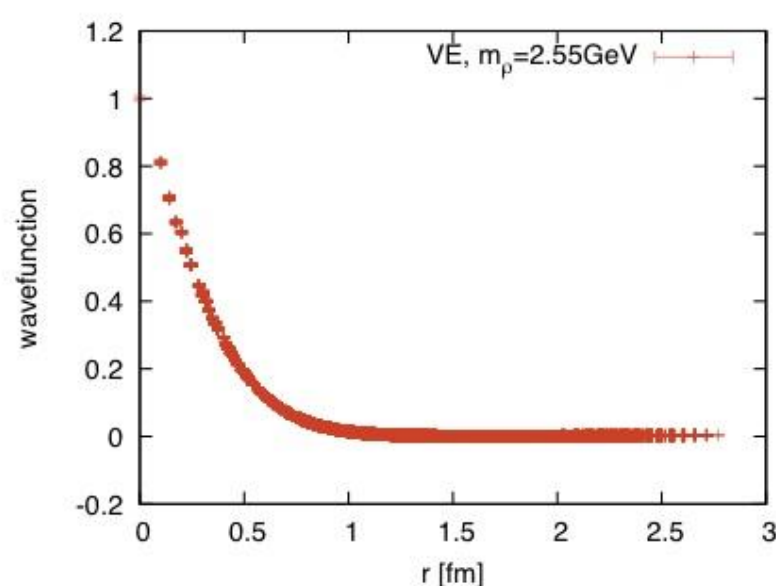
Kawanai & Sasaki, 1102.3246 [hep-lat]

Kawanai @ Lattice2011

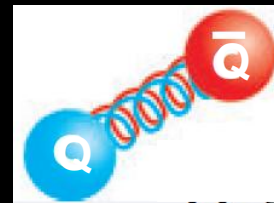
$J^P=0^-$ channel



$J^P=1^-$ channel



“Constituent quark model” from LQCD ?



Quenched QCD: Coulomb gauge

$L = 3.3$ fm, $a = 0.104$ fm

(2+1)-flavor QCD: Coulomb gauge

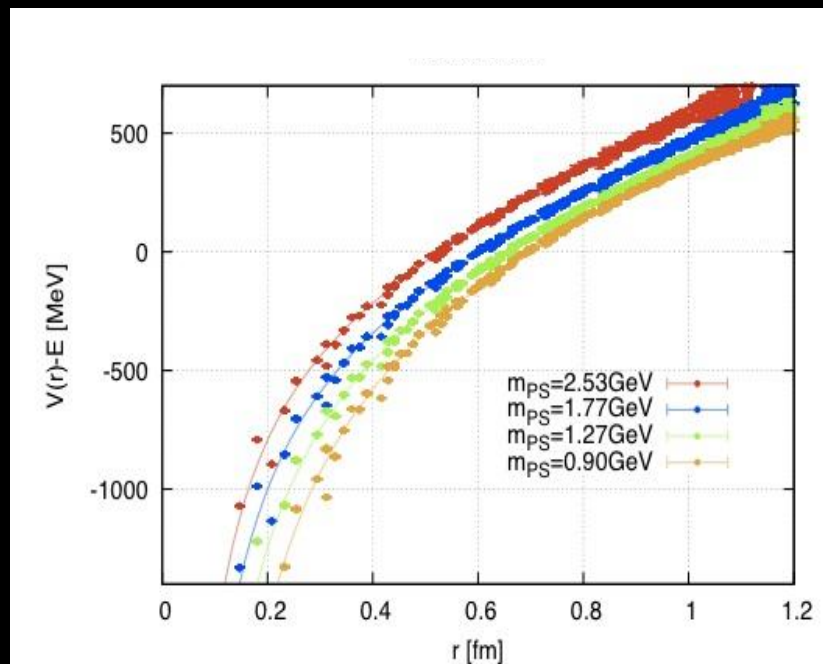
$L = 3$ fm, $a = 0.09$ fm

Ikeda & Iida, 1102.2097 [hep-lat]

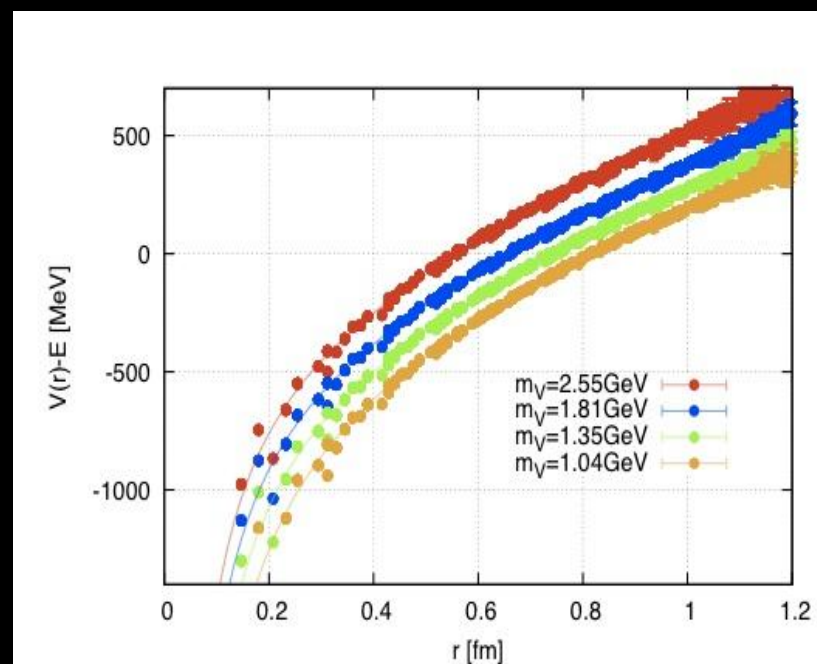
Kawanai & Sasaki, 1102.3246 [hep-lat]

Kawanai @ Lattice2011

$J^P=0^-$ channel

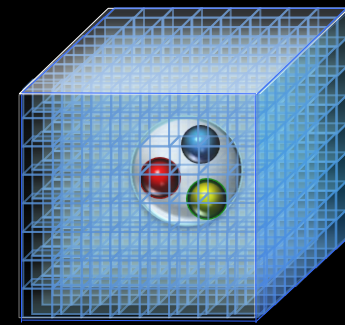


$J^P=1^-$ channel

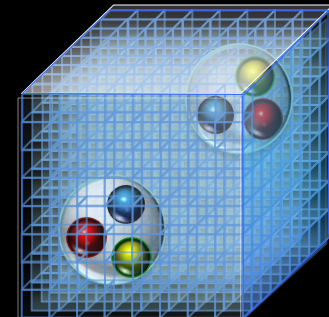


Coulomb+linear+spin-dep. potential between dynamical quarks

“Summary”



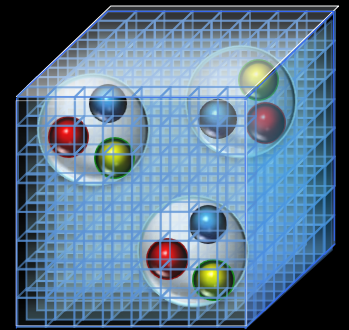
1. **LQCD provides precision computations**
(2+1)-flavor, $L=6\text{fm}$, $a=0.05\text{fm}$, $m_{\pi}=135\text{MeV}$
 α_s , $m_{u,d,s}$, low energy constants, ...
2. **LQCD provides inputs for PAN phenomenology**
 - dark matter (e.g. $s\bar{s}$ in the nucleon)
 - quark-gluon plasma (EOS, spectral function,...)
3. **LQCD provides qualitative pictures**
 - nucleon and hyperon forces
 - the constituent quark model



“Future” (10 Pflops era from 2012)

At next PANIC (2014), we would (like to) hear

1. Physical point simulations for many observables
no more chiral extrapolation
2. Simulations with “better” fermions
staggered, Wilson \rightarrow domain wall, overlap
3. Realistic BB, BBB forces & light nuclei
4. New ideas in $\mu/T > 1$?



Yamazaki et al. [PACS-CS Coll.],
Phys.Rev.D81(2010)111504

T. Doi [HAL QCD Coll.],
arXiv:1106.2276[hep-lat]

Backup slides

“Conclusion”


“The Scientist as Rebel”
by Freeman J. Dyson

It often happens that the understanding of **the mathematical nature of an equation** is impossible without a detailed understanding of its solutions. **The black hole** is a case in point.

One could say without exaggeration that **Einstein's equations of general relativity** were understood only at a very superficial level before the discovery of **the black hole**.

 → **String theory**

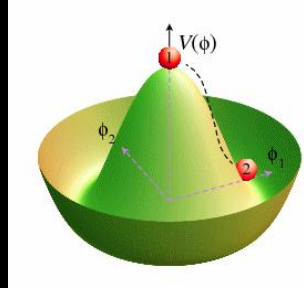
 → **D-brane**

 → **QCD**

 → **Nucleon? QGP?**

The progress of science requires the growth of understanding in both directions, downward from the whole to the parts and upward from the parts to the whole.

Chiral condensate



Gell-Mann-Oakes-Renner (GOR) formula (1968)

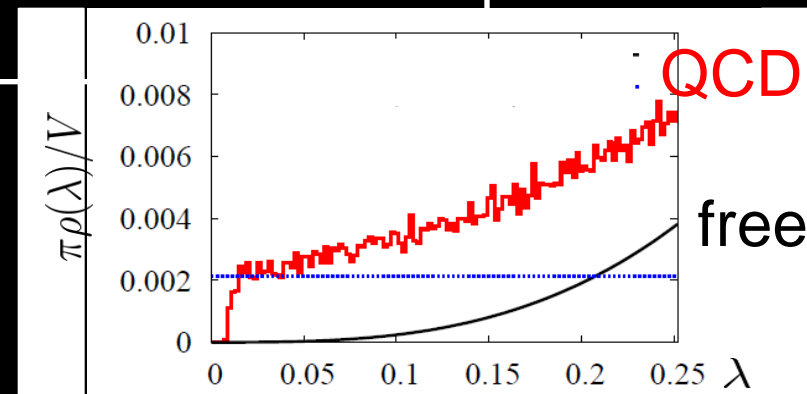
$$f_\pi^2 m_\pi^2 = -(m_u + m_d) \langle \bar{q}q \rangle_0 + O(m_{u,d}^2)$$

Di Vecchia-Veneziano formula (1980)

$$\chi_{\text{top}} = \frac{-\langle \bar{q}q \rangle_0}{1/m_u + 1/m_d + 1/m_s} + O(m_u^2)$$

Banks-Casher relation (1980)

$$\langle \bar{q}q \rangle = - \lim_{m \rightarrow 0} \int_0^\infty d\lambda \frac{2m}{\lambda^2 + m^2} \rho(\lambda) = -\pi \rho(0)$$



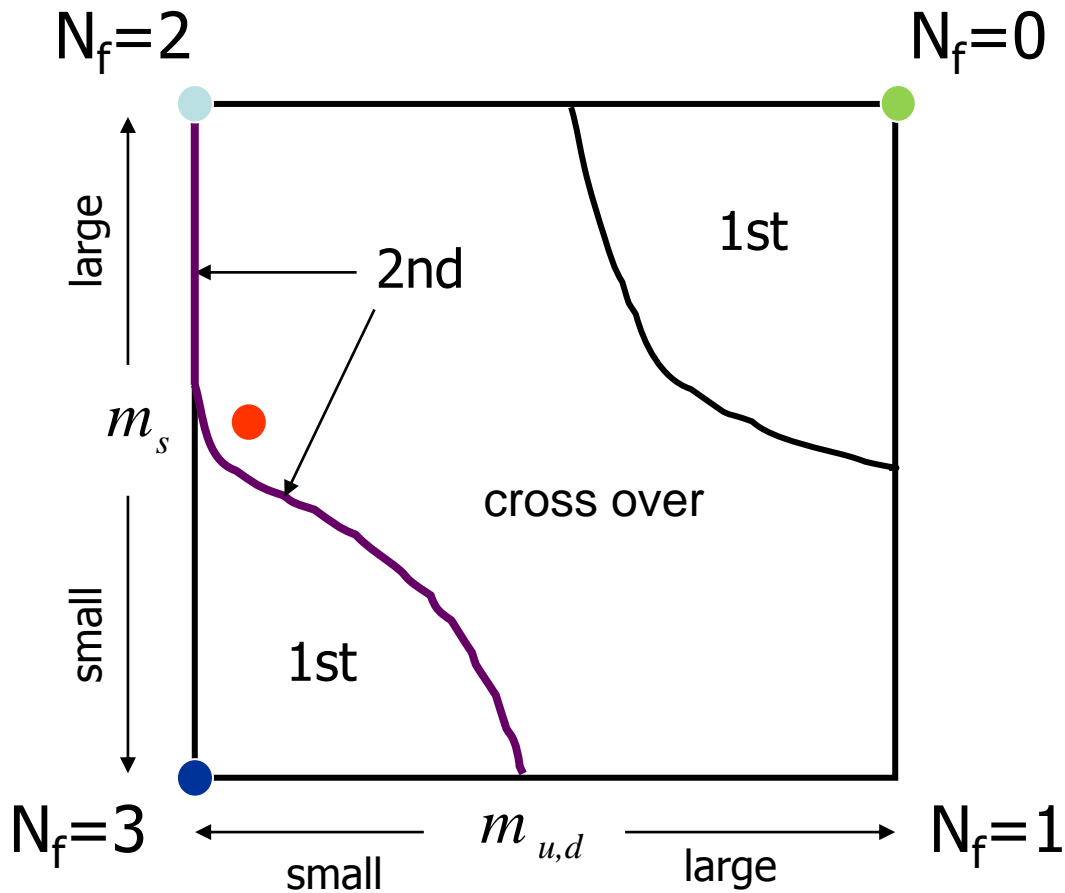
Order of the thermal QCD transition ($\mu=0$)

Svetitsky & Yaffe, NPB210 ('82)

Pisarski and Wilczek, PRD29 ('84)

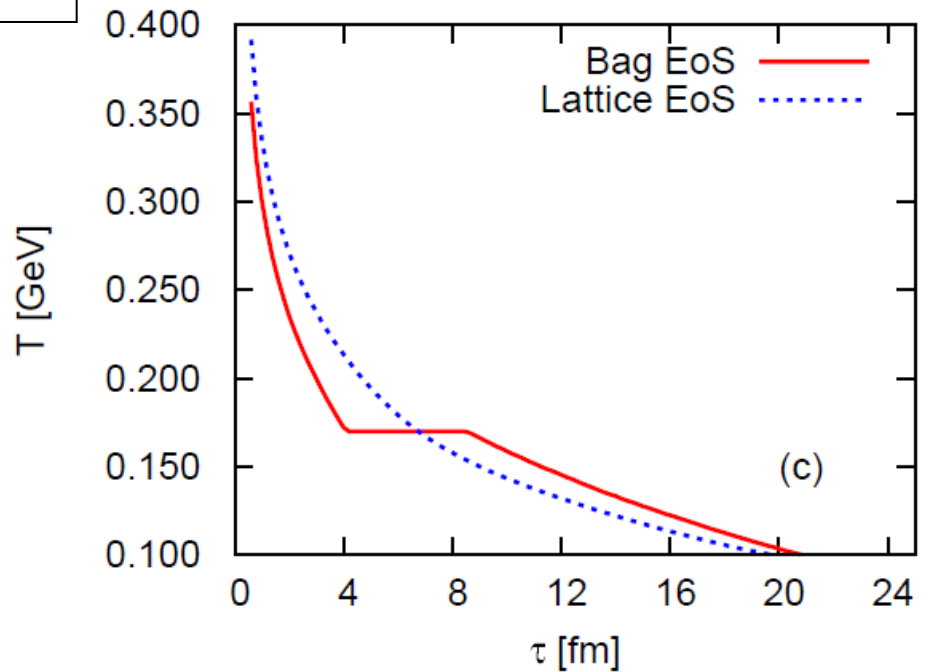
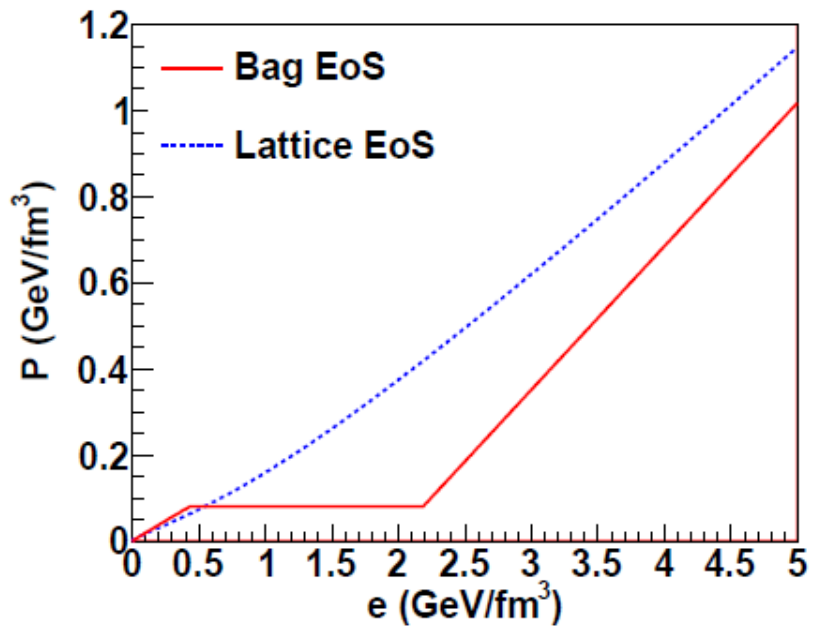
$$\mathcal{V} = \frac{a}{2}\phi^2 + \frac{b}{4}\phi^4 - h\phi$$

$$\mathcal{V} = \frac{a}{2}L^2 - \frac{c}{3}L^3 + \frac{b}{4}L^4 - hL$$



$$\mathcal{V} = \frac{a}{2}\sigma^2 - \frac{c}{3}\sigma^3 + \frac{b}{4}\sigma^4 - h\sigma$$

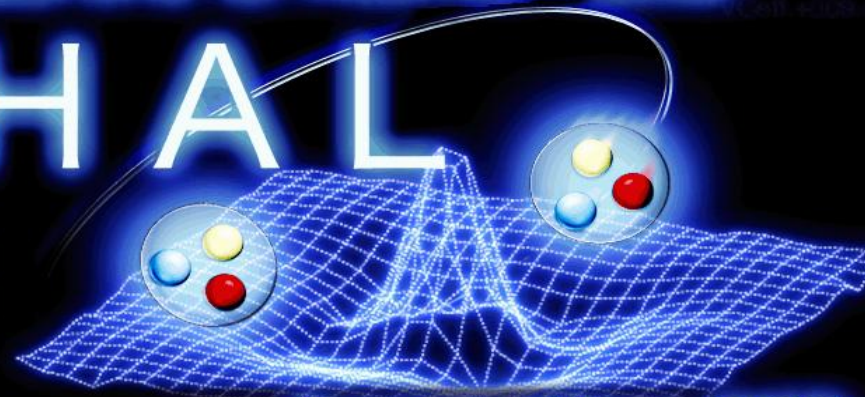
$$\mathcal{V} = -A\chi + \frac{a}{2}\chi^2 + \frac{b}{4}\chi^4$$



Akamatsu et al.,
arXiv:1107.36[nucl-th]

Hadrons to Atomic nuclei

HAL



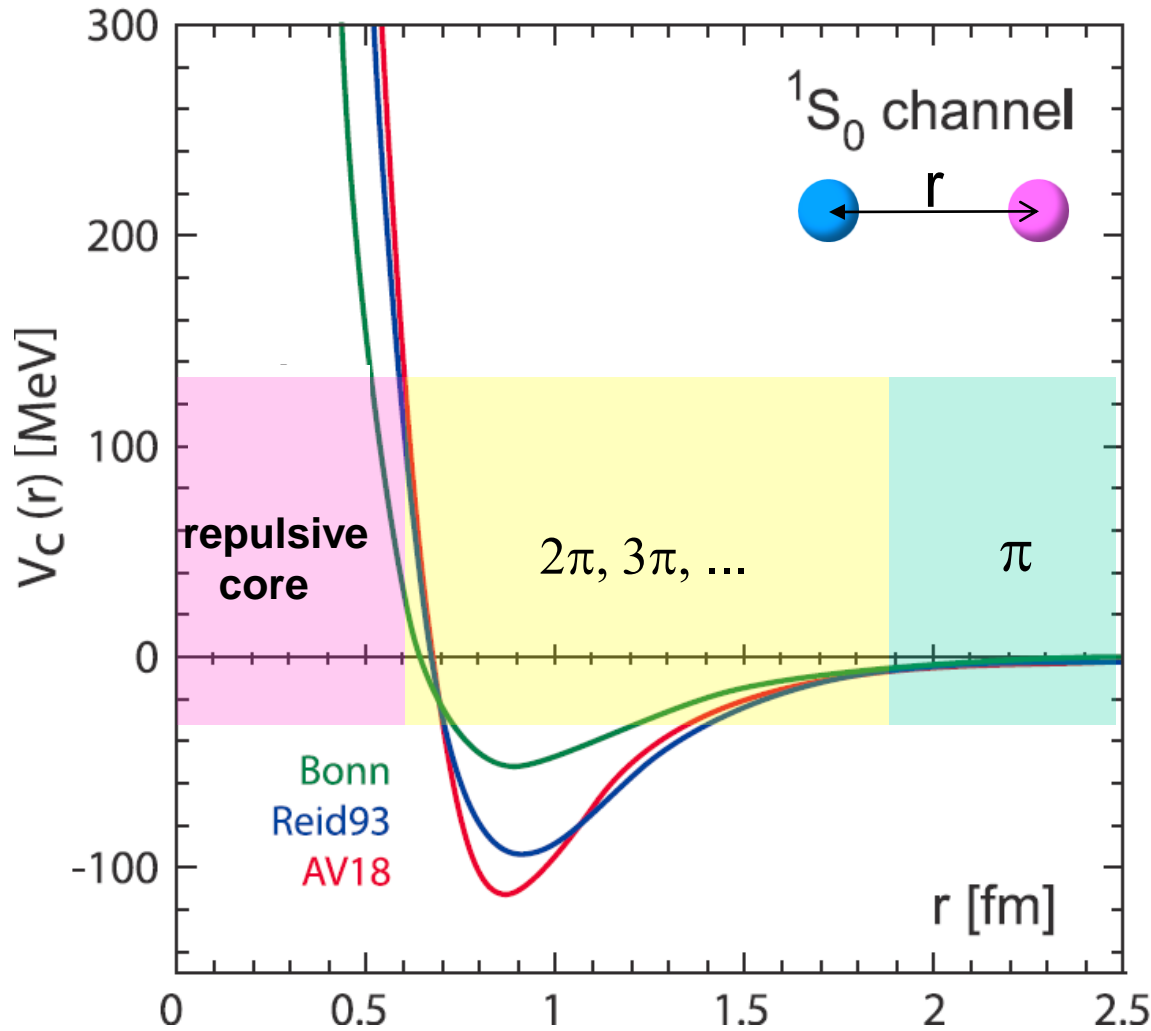
from Lattice QCD

Tohoku Univ.
Univ. Tsukuba
RIKEN
Nihon Univ.
Tokyo Inst. Tech.
CNS, Univ. Tokyo

H. Nemura
S. Aoki, N. Ishii, K. Sasaki
K. Murano, T. Hatsuda
T. Inoue
Y. Ikeda
T. Doi

Phenomenological NN potentials

(~40 parameters to fit 5000 phase shift data)



One-pion exchange
by Yukawa (1935)

Multi-pions
by Taketani et al. (1951)

Repulsive core
by Jastrow (1951)

Key channels in NN scattering ($^{2s+1}L_J$)

$$V(\vec{r}, \nabla) = V_C(r) + S_{12}V_T(r) + \vec{L} \cdot \vec{S} V_{LS}(r) + \{V_D(r), \nabla^2\} + \dots$$

LO

LO

NLO

NNLO

1S_0 Central force \longleftrightarrow nuclear BCS pairing

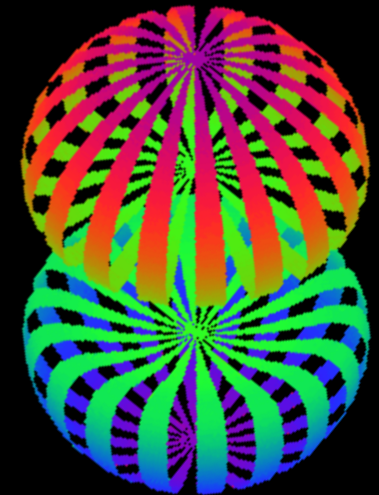
Bohr, Mottelson & Pines, Phys. Rev. 110 (1958)

3S_1 - 3D_1 Tensor force \longleftrightarrow deuteron binding

Pandharipande et al., Phys. Rev. C54 (1996)

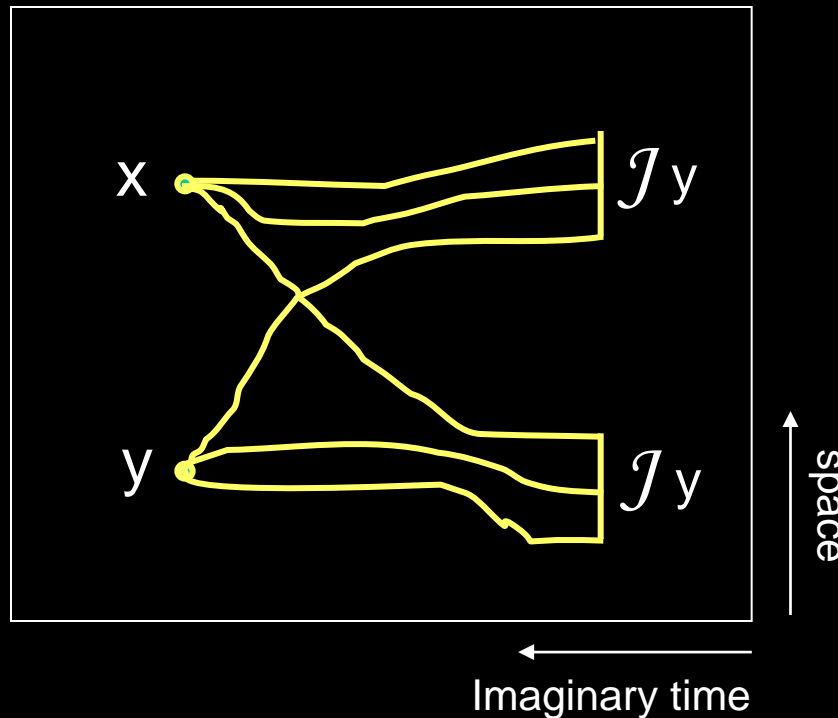
3P_2 - 3F_2 LS force \longleftrightarrow neutron superfluidity
in neutron stars

Tamagaki, Prog. Theor. Phys. 44 (1970)



Density profile
of the deuteron
with $S_z = \pm 1$

Equal-time NBS amplitude $\phi(\mathbf{r})$ in lattice QCD



+ all possible combinations

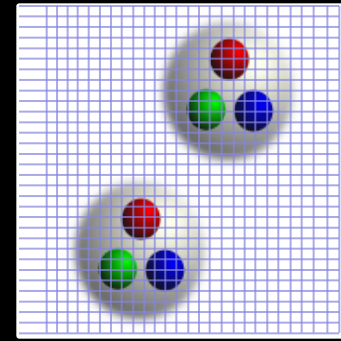
$$\begin{aligned}
 C_4(\mathbf{r}; t) &= \langle N_1(\mathbf{x}, t) N_2(\mathbf{y}, t) \mathcal{J}_1^\dagger(0) \mathcal{J}_2^\dagger(0) \rangle \\
 &= \sum_n \langle 0 | N_1(\mathbf{x}) N_2(\mathbf{y}) | n \rangle A_n e^{-E_n t} \longrightarrow \phi(\mathbf{r}) A_0 e^{-E_0 t}
 \end{aligned}$$

$\phi(r > R) \rightarrow$ phase shift : Luscher, Nucl. Phys. B354 (1991) 531

$\phi(r < R) \rightarrow$ potential : Ishii, Aoki & Hatsuda, PRL 99 (2007) 022001

Methods to extract NN interaction from LQCD

Luscher, Nucl. Phys. B354 (1991) 531



[1] Temporal correlation : $E_{NN}(L) \rightarrow$ NN phase shift

$$\frac{2Z_{00}(1, q)}{L\pi^{1/2}} = k \cot \delta_0(k)$$

- quenched QCD: CP-PACS Coll. (1995)
- full QCD: NPLQCD Coll. (2006-)

[2] Spatial correlation :

BS wave function \rightarrow NN potential \rightarrow observables

$$(E - H_0)\phi(\mathbf{r}) = \int U(\mathbf{r}, \mathbf{r}')\phi(\mathbf{r}')d\mathbf{r}'$$

← half off-shell T-matrix

- π - π system : CP-PACS Coll. (2005)
- NN system (quenched QCD) : Ishii, Aoki & T.H., PRL 99, 022001 (2007).
- NN, YN systems (full QCD): HAL QCD Coll. (2008-)

Systematic procedure to define the NN potential in lattice QCD

Full details, see
Aoki, Ishii & Hatsuda, 0909.5585 [hep-lat]

(i) **Choose your favorite operator:** e.g. $N(x) = \epsilon_{abc} q^a(x) q^b(x) q^c(x)$

observables do not depend on the choice
yet the local operator is useful

Nishijima, Haag, Zimmermann (1958)

(ii) **Measure the NBS amplitude:** $\phi(\vec{r}) = \langle 0 | N(\vec{x} + \vec{r}) N(\vec{x}) | 6q \rangle$

(iii) **Define the non-local potential:** $(E - H_0)\phi(\vec{r}) = \int U(r, \vec{r}') \phi(\vec{r}') d^3 r'$

(iv) **Velocity expansion :** $U(\vec{r}, \vec{r}') = V(\vec{r}, \nabla) \delta^3(\vec{r} - \vec{r}')$

$$V(\vec{r}, \nabla) = V_C(r) + S_{12} V_T(r) + \vec{L} \cdot \vec{S} V_{LS}(r) + \{V_D(r), \nabla^2\} + \dots$$

Okubo-Marshak (1958), Tamagaki-Watari (1967)

(v) **Calculate observables :** phase shifts, binding energies etc

Properties of lattice NN potential $U(r,r')$

$$U(\vec{r}, \vec{r}') = V(\vec{r}, \nabla) \delta^3(\vec{r} - \vec{r}')$$

[1] $U(r,r')$ is $N(x)$ -dependent

QM : $(\psi, V) \rightarrow$ observables

QFT : (asymptotic field, vertices) \rightarrow observables

$(N(x), U(r,r')) \rightarrow$ observables

[2] $U(r,r')$ is E -independent

non-locality can be determined order by order

[3] $U(r,r')$ has minor volume dependence

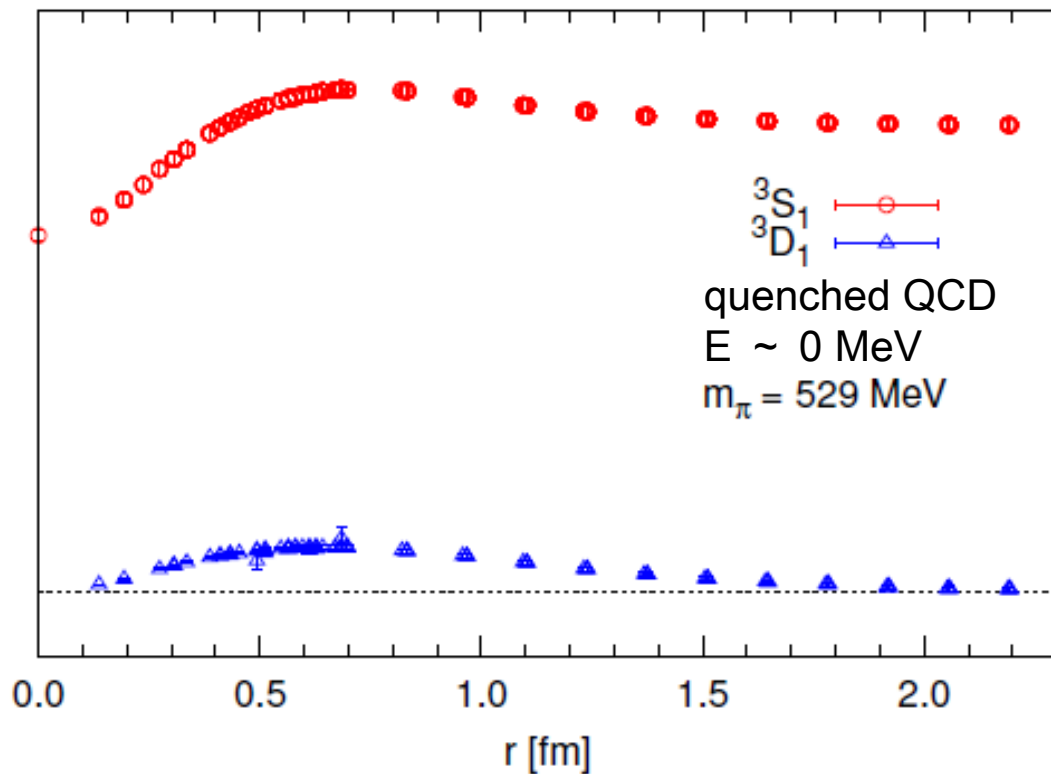
Wave function is sensitive to the volume

Potential is insensitive to the volume

remember the deuteron !

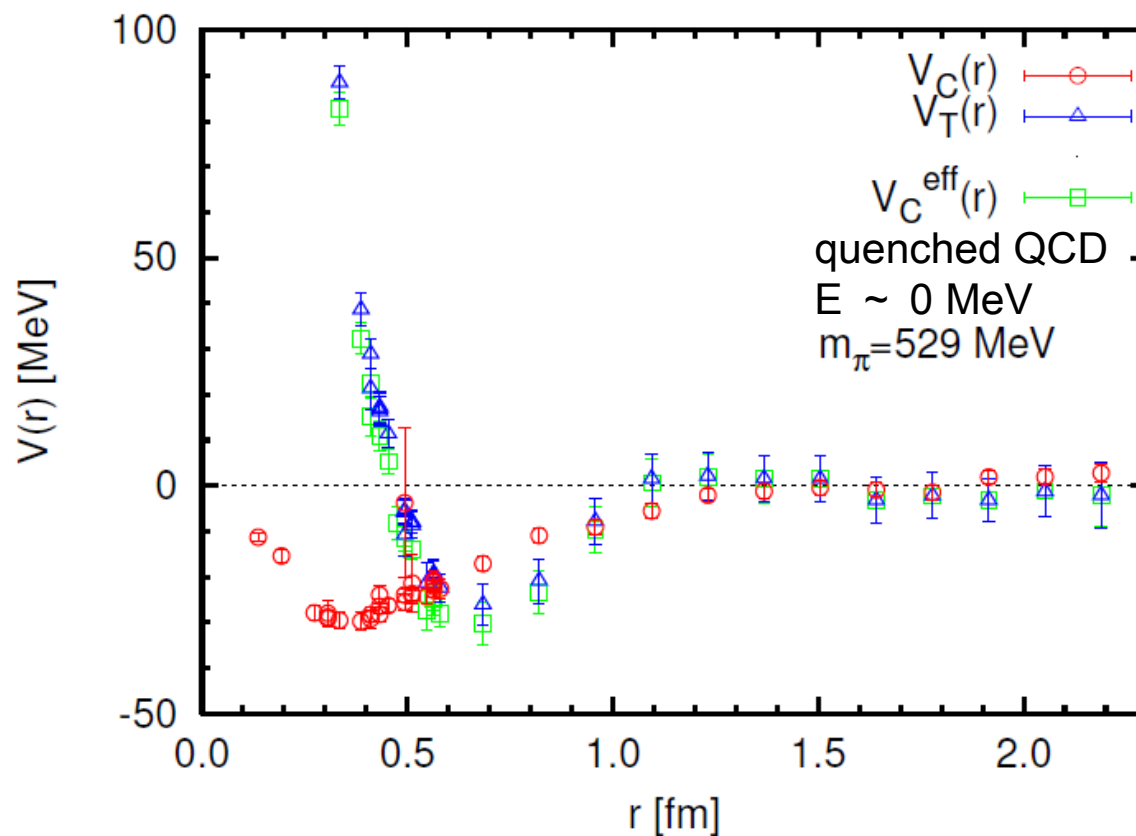
Central & tensor potentials : $V_C(r)$ & $V_T(r)$

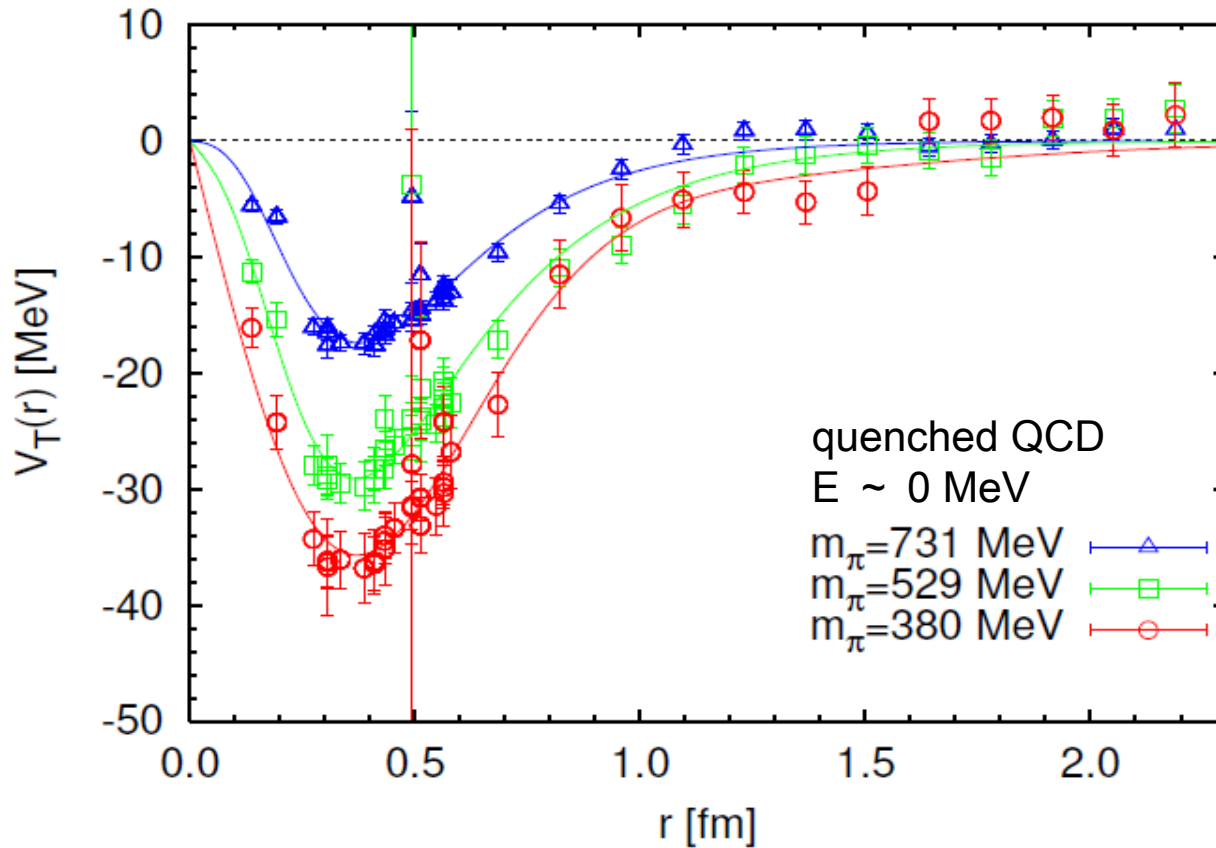
Aoki, Ishii & Hatsuda,
0909.5585 [hep-lat]



Central & tensor potentials : $V_C(r)$ & $V_T(r)$

Aoki, Ishii & Hatsuda,
0909.5585 [hep-lat]



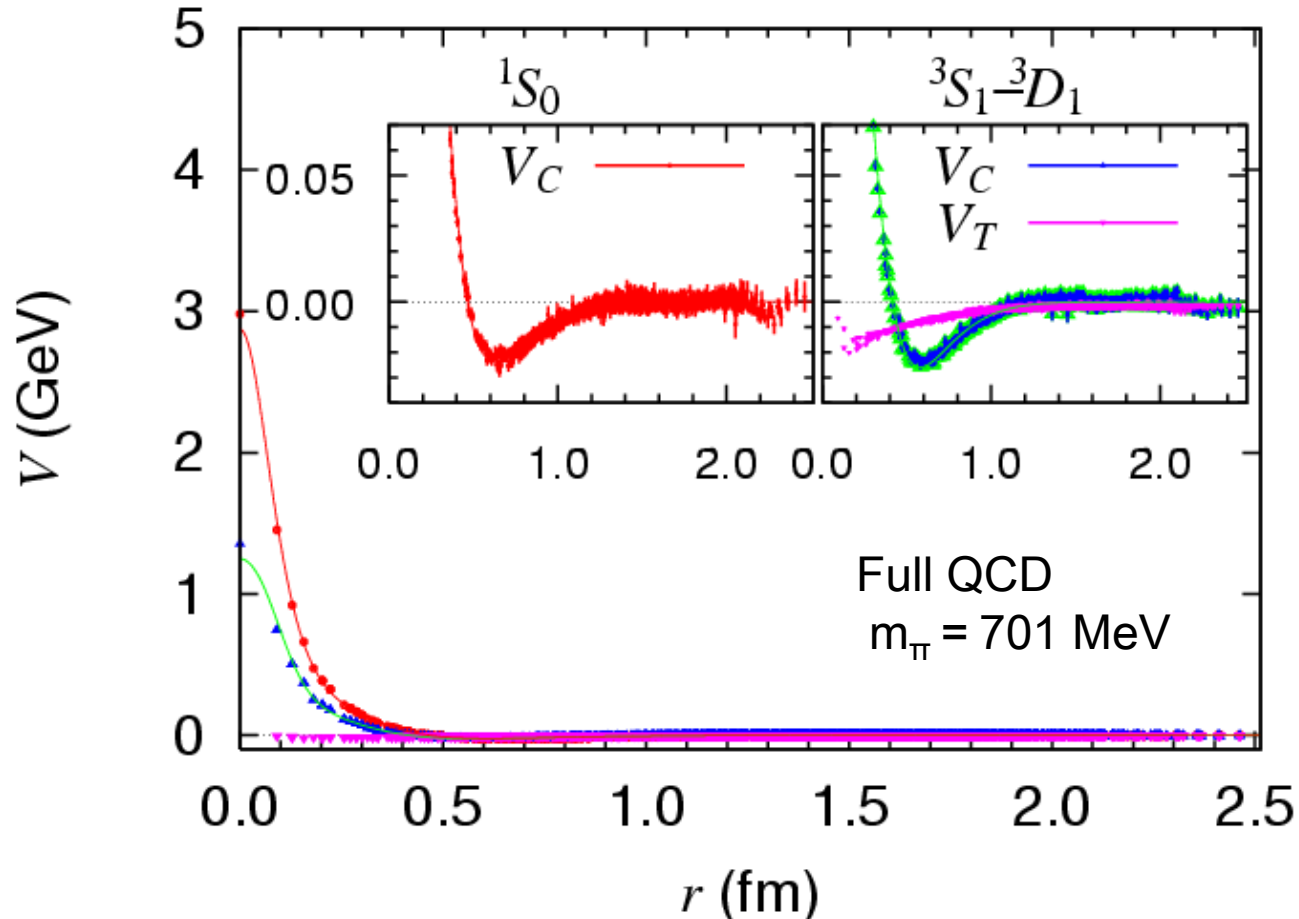


fit function

$$V_T(r) = b_1(1 - e^{-b_2 r^2})^2 \left(1 + \frac{3}{m_\rho r} + \frac{3}{(m_\rho r)^2} \right) \frac{e^{-m_\rho r}}{r} + b_3(1 - e^{-b_4 r^2})^2 \left(1 + \frac{3}{m_\pi r} + \frac{3}{(m_\pi r)^2} \right) \frac{e^{-m_\pi r}}{r},$$

- Rapid quark-mass dependence of $V_T(r)$
- Evidence of the one-pion-exchange

ΛN interaction



- Repulsive core + attractive well
- Weak tensor force
- Overall attraction

irreducible BB source operator

$$\overline{BB}^{(27)} = +\sqrt{\frac{27}{40}} \overline{\Lambda\Lambda} - \sqrt{\frac{1}{40}} \overline{\Sigma\Sigma} + \sqrt{\frac{12}{40}} \overline{N\Xi} \quad \text{or} \quad +\sqrt{\frac{1}{2}} \overline{p\bar{n}} + \sqrt{\frac{1}{2}} \overline{\bar{n}p}$$

$$\overline{BB}^{(8_s)} = -\sqrt{\frac{1}{5}} \overline{\Lambda\Lambda} - \sqrt{\frac{3}{5}} \overline{\Sigma\Sigma} + \sqrt{\frac{1}{5}} \overline{N\Xi}$$

$$\overline{BB}^{(1)} = -\sqrt{\frac{1}{8}} \overline{\Lambda\Lambda} + \sqrt{\frac{3}{8}} \overline{\Sigma\Sigma} + \sqrt{\frac{4}{8}} \overline{N\Xi} \quad \text{with}$$

$$\overline{\Sigma\Sigma} = +\sqrt{\frac{1}{3}} \overline{\Sigma^+\Sigma^-} - \sqrt{\frac{1}{3}} \overline{\Sigma^0\Sigma^0} + \sqrt{\frac{1}{3}} \overline{\Sigma^-\Sigma^+}$$

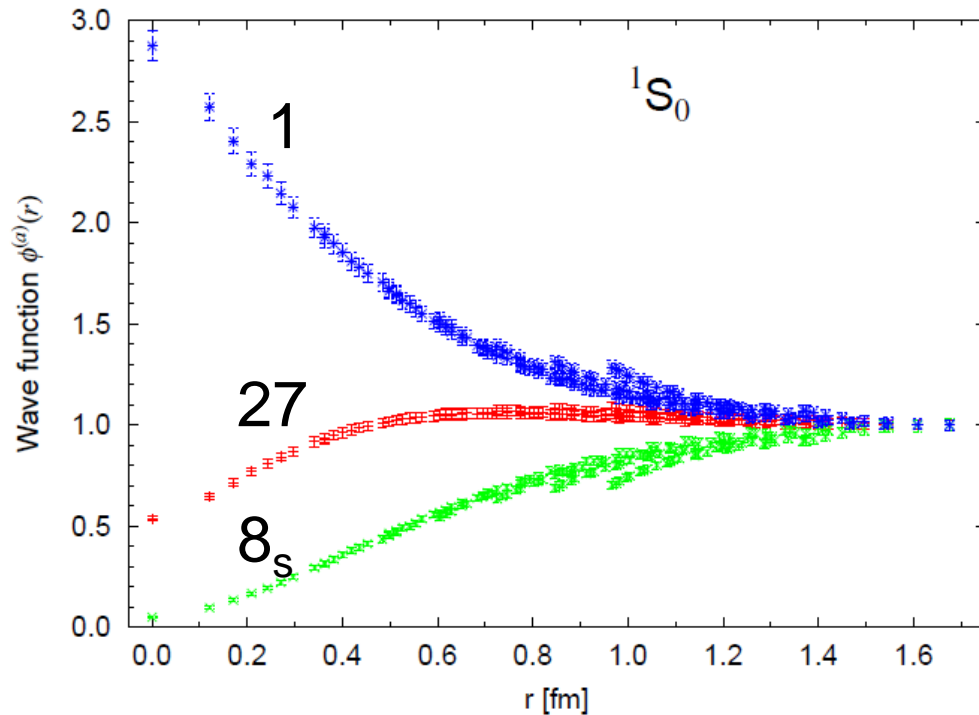
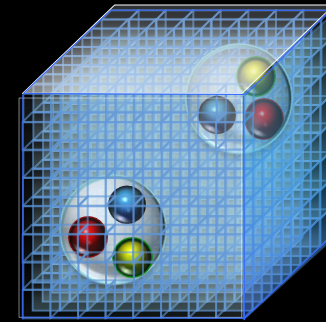
$$\overline{BB}^{(10^*)} = +\sqrt{\frac{1}{2}} \overline{p\bar{n}} - \sqrt{\frac{1}{2}} \overline{\bar{n}p}$$

$$\overline{N\Xi} = +\sqrt{\frac{1}{4}} \overline{p\Xi^-} + \sqrt{\frac{1}{4}} \overline{\Xi^-p} - \sqrt{\frac{1}{4}} \overline{\bar{n}\Xi^0} - \sqrt{\frac{1}{4}} \overline{\Xi^0\bar{n}}$$

$$\overline{BB}^{(10)} = +\sqrt{\frac{1}{2}} \overline{p\overline{\Sigma^+}} - \sqrt{\frac{1}{2}} \overline{\overline{\Sigma^+}p}$$

$$\overline{BB}^{(8_a)} = +\sqrt{\frac{1}{4}} \overline{p\overline{\Xi^-}} - \sqrt{\frac{1}{4}} \overline{\overline{\Xi^-}p} - \sqrt{\frac{1}{4}} \overline{\bar{n}\overline{\Xi^0}} + \sqrt{\frac{1}{4}} \overline{\overline{\Xi^0}\bar{n}}$$

BB wave functions in flavor-basis



Iwasaki + clover
(CP-PACS/JLQCD config.)
 $L=1.9$ fm, $a=0.12$ fm, $16^3 \times 32$
 $m_\pi=835$ MeV, $m_B=1752$ MeV

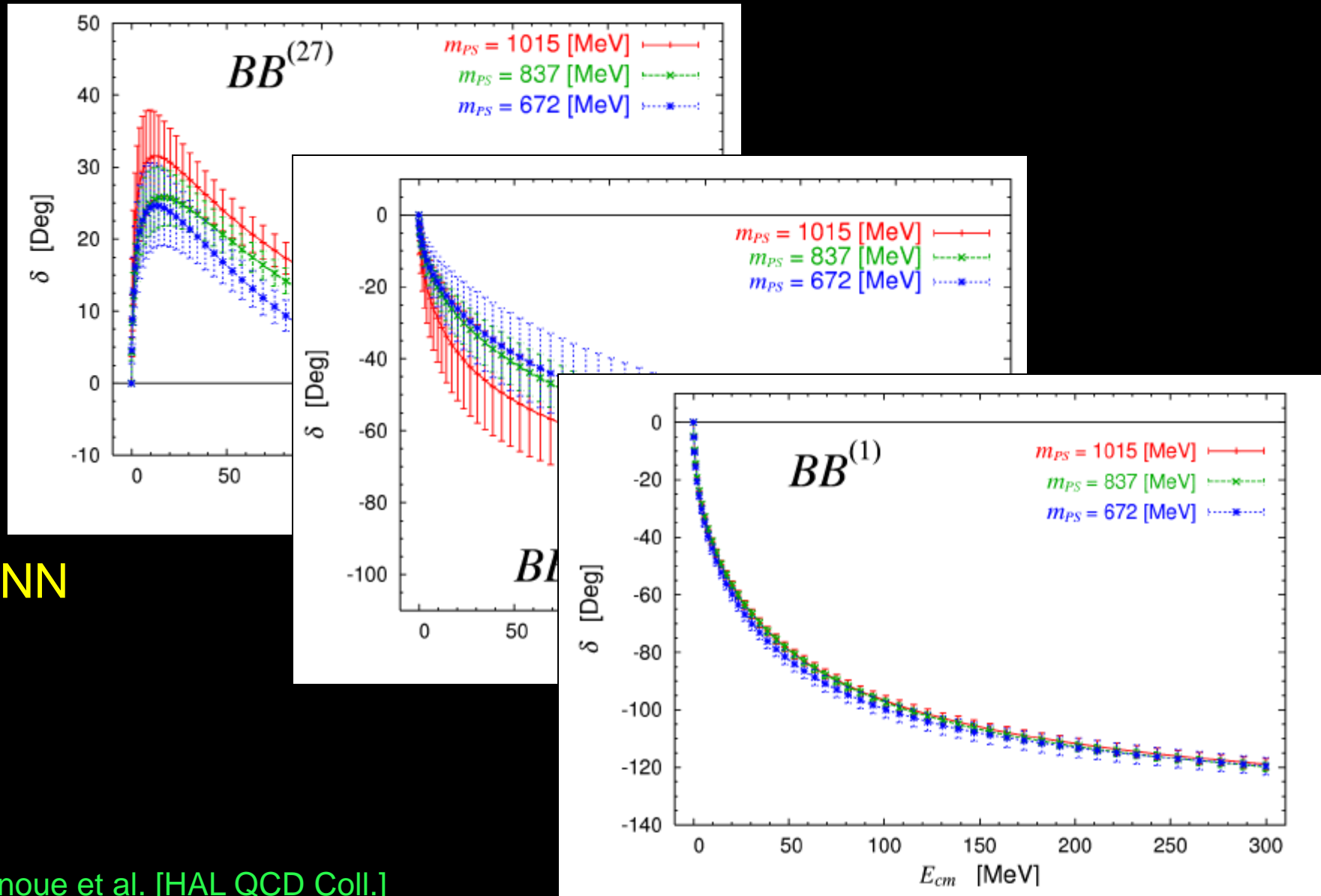
Inoue et al. (HAL QCD Coll.)
Prog. Theor. Phys. 124 (2010) 591

Origin of the short range BB int. \Leftrightarrow Quark Pauli principle !

1 : allowed, 27 : partially blocked, 8_s : blocked

c.f. constituent quark model (Oka, Yazaki, Shimizu, ...)

BB phase shifts in flavor-basis (1S_0 channel)



NN