

# Global analysis of fragmentation functions: AKK08 & ABKK12

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# Global Analyses of Unpolarized FFs

light charged (l.c.h.)  $\pi^\pm, K^\pm, p/\bar{p}$ ,  
strange neutral (s.n.h.)  $K_S^0, \Lambda/\bar{\Lambda}$

## Albino-Kramer-Kniehl (2005/6)

- $e^+e^-$ : OPAL etc.,  $\sqrt{s} = M_Z$
- OPAL  $u, d, s$  tagged
- $h^+, h^-$  not distinguished
- $x > 0.1$

## Hirai-Kumano-Nagai-Sudoh (2007)

- $e^+e^-$ : OPAL, TASSO, etc.
- OPAL  $u, d, s$  tagged
- FF exp. errors
- $h^+, h^-$  separately
- $x > .05$

## de Florian-Sassot-Stratmann (2007)

- $e^+e^-$ : OPAL, TASSO, etc.
- OPAL  $u, d, s$  tagged
- $pp$ : BRAHMS, STAR, PHENIX
- $ep$ : HERMES ( $\pi^\pm, K^\pm$ )
- $h^+, h^-$  separately
- $x > .05$

# Summary of AKK08 fit [NPB803(2008)42]

## Experiment

- l.c.h. ( $\pi^\pm, K^\pm, p/\bar{p}$ ) + s.n.h. ( $K_S^0, \Lambda/\bar{\Lambda}$ )
- $e^+e^-$ ,  $x > .05$ :
- $\sqrt{s} = M_Z$  (OPAL, ALEPH...)
- $l, c, b$  tag (OPAL:  $u, d, s, c, b$ )
- $+ 12\text{GeV} < \sqrt{s} < M_Z$  (TASSO...)
- $pp$ : BRAHMS, STAR, PHENIX ( $p_T > 2$  GeV)
- $p\bar{p}$ : CDF, 2005 ( $K_S^0, \Lambda/\bar{\Lambda}$ )
- norm. errors  $\in$  covariance matrix
- exclude  $ep$ : HERMES ( $\pi^\pm, K^\pm$ ):
  - $Q \lesssim 2$  GeV — fixed order?
  - $ep(Q) \sim e^+e^- (\sqrt{s})$
  - compare low  $Q$  with low  $\sqrt{s}$
  - different *systematic* theory error

## Theory

- FF basis:
  - $\pm$  sum ( $h^\pm$ ) ( $e^+e^-$  ∴ accurate)
  - $\pm$  diff ( $\Delta_c h^\pm$ ) ( $pp$  ∴ less so)
- non pert. constraints:
  - Strong:  $\text{FF}(\sqrt{2}\text{GeV}) = Nx^\alpha(1-x)^\beta$ ,  
 $\times(1 + \gamma(1-x)^\delta)$  for  $h^\pm$
  - Weak: SU(2) isospin  $u \leftrightarrow d \in \pi^\pm$
  - Weaker:  $D_{\bar{q}}^{(\Delta_c)h^\pm} = -D_q^{(\Delta_c)h^\pm}$  (QCD)
- hadron mass:
  - fitted ( $e^+e^-$ )  
 $\rightarrow$  subtract low  $x, \sqrt{s}$  effects
  - baryons accurate (+1%)
  - mesons OK
  - large  $x$  res. ( $e^+e^-$ ) — improves fit

# General calculation

Factorization theorem I

$$IS \rightarrow \sum_j \text{parton } j + X \rightarrow h + X$$

$$\frac{d\sigma^h}{dx}(x, E) = \sum_j \int_x^1 \frac{dz}{z} \frac{d\sigma^j}{d(x/z)} \left( \frac{x}{z}, a_s(E), \frac{E}{M_f} \right) D_j^h(z, M_f)$$

$$(E = \sqrt{s}, p_T)$$

Neglect higher twist  $O(1/E)$

**Scale:**  $\mu^2, M_f^2 = kE^2$

$$(k = 1, 1/4, 4)$$

**Scheme:**  $\overline{\text{MS}}$

Factorization theorem II

DGLAP evolution

$$\frac{d}{d \ln M_f} D_i^h(x, M_f) = \sum_j \int_x^1 \frac{dz}{z} P_{ij} \left( \frac{x}{z}, a_s(M_f) \right) D_j^h(z, M_f)$$

$D_i^h(x, M_f)$  universal

$D_q^h(x, M_f < m_q^t \sim m_q)$ :  
Extrinsic = 0,

Intrinsic not in evolution

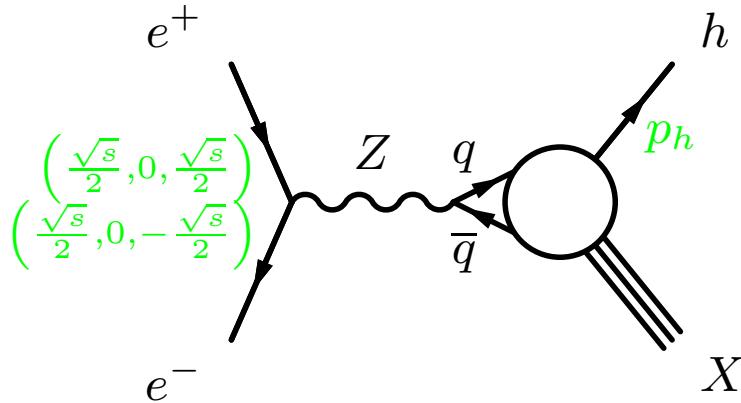
Calculations to NLO

$$\text{CTEQ6.5S0 PDFs} \rightarrow \Lambda_5^{\overline{\text{MS}}} = 226 \text{ MeV}$$

# Processes

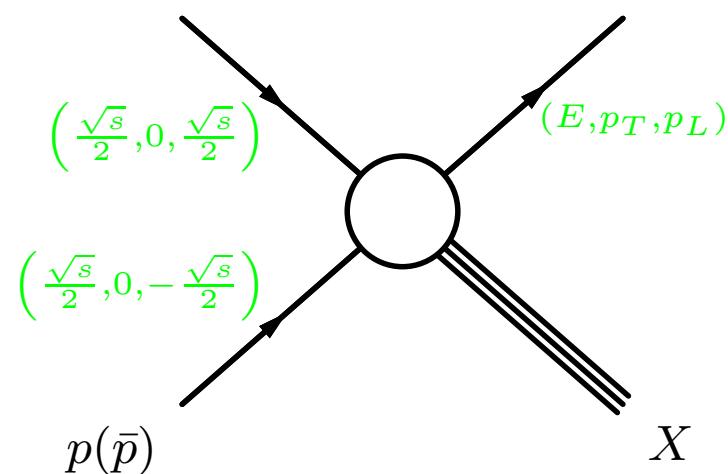
$$e^+ e^- \rightarrow h + X$$

$$\frac{d\sigma^h}{dx} \left( x = \frac{2|\mathbf{p}_h|}{\sqrt{s}}, \sqrt{s} \right) =$$



$$pp(\bar{p}) \rightarrow h + X$$

$$\frac{d^2\sigma^h}{dp_T dy} \left( x = \frac{2p_T}{\sqrt{s}} \cosh y, y = \frac{1}{2} \ln \frac{E + p_L}{E - p_L}, \sqrt{s} \right) =$$



ALEPH, DELPHI, TASSO, OPAL, ...  
 $\rightarrow$  precise  $q \rightarrow h^\pm$  FFs (no  $s - d$ )

- CDF (2005) ( $p\bar{p} \rightarrow K_S^0, \Lambda/\bar{\Lambda}$ ,  $\sqrt{s} = 630$  GeV)  
 $\rightarrow$  FF shape + relative (not overall) norms
- BRAHMS, PHENIX, STAR (2006,2007) @ RHIC  
 $(pp, \sqrt{s} = 200$  GeV)  
 $\rightarrow$  remaining FFs (gluon,  $q \rightarrow \Delta_c h^\pm$ )

# $pp(\bar{p})$ in Mellin Space

$N$ - (Mellin) space ( $f(N) = \int_0^1 dx x^{N-1} f(x)$ ):

$x$ -convolutions  $\longrightarrow$  products in  $N$ -space

$N$ -space calculations fast and accurate

$pp(\bar{p})$  XS from convolution with FFs:

$$F\left(x = \frac{2p_T}{\sqrt{s}} \cosh y\right) = \int_x^1 \frac{dz}{z} \widehat{F}\left(\frac{x}{z}\right) z D(z)$$
$$\therefore F(N) = \widehat{F}(N) D(N+1)$$

$\widehat{F}(N)$  analytic calculation impossible

solution:  $\widehat{F}$  as Chebyshev expansion

$$\widehat{F}(z) / \ln(1 - x/z) = \sum_k c_k T_k(X)$$

$-1 < X < 1 \leftarrow$  linear map  $\rightarrow x < z < 1$

$\longrightarrow$  semi-analytic  $\widehat{F}(N)$

$\longrightarrow F(x)$  calc. is quick, to few parts per mil

# Systematic Errors

$$P(\{f_i^e\}, \{f_i^t\}) \propto \exp \left[ -\frac{1}{2} \chi^2 = \sum_i \left( \frac{f_i^t - f_i^e}{\sigma_i} \right)^2 \right]$$

$pp(\bar{p})$ , low  $\sqrt{s}$   $e^+e^-$ : norm. uncertainty(s)  
= systematic error(s)

systematic error,  $K$ th source:

- $\Delta_K f_i^e = \lambda_K \sigma_i^K$

- $P(\lambda_K) \propto \exp [-\frac{1}{2} \lambda_K^2]$   
 $(\rightarrow \text{expect } |\lambda_K| \lesssim 1)$

$$\chi^2 = \sum_i \left( \frac{f_i^t - f_i^e - \sum_K \lambda_K \sigma_i^K}{\sigma_i} \right)^2 + \sum_K \lambda_K^2$$

Minimize  $\chi^2$  w.r.t.  $\lambda_K$  (Max prob):

- $\lambda_K = \sum_i \frac{f_i^t - f_i^e}{\sigma_i^2} \left( \sigma_i^K - \sum_{j k L} \sigma_i^L \sigma_j^L (\mathcal{C}^{-1})_{jk} \sigma_k^K \right)$

- $\chi^2 = \sum_{ij} (f_i^t - f_i^e) (\mathcal{C}^{-1})_{ij} (f_j^t - f_j^e)$

( Covariance  $C_{ij} = \sigma_i^2 \delta_{ij} + \sum_K \sigma_i^K \sigma_j^K$  )

Unknown systematic effects:

Cancel for data sets from large no. of different exps.(?)

# Systematic Errors — $\pi^\pm$

Collaboration	$\sqrt{s}$ (GeV)	# data	Norm. (%)	$\chi^2_{\text{DF}}$	$\lambda_K$	Shift (%)
TASSO	12	5	20	0.50	0.21	4.3
TASSO	14	10	8.5	0.92	-1.26	-10.7
TASSO	22	1	6.3	0.01	-0.08	-0.5
TASSO	30	4	20	0.57	0.69	13.8
TASSO	34	10	6	1.07	0.62	3.7
TASSO	44	7	6	1.99	0.66	3.9
ALEPH	91.2	22	3	0.61	-0.55	-1.6
BRAHMS, $y \in [2.9, 3]$ $y \in [3.25, 3.35]$	200	8	11,7,8(13), 2,1(3)	0.96	-1.76, -1.12, -1.22, -0.32, -0.13	-19.4, -7.9, -9.7, -0.6, -0.1
PHENIX ( $\pi^0$ ), $ \eta  < 0.35$		7		2.68	-2.01, -1.28, -1.80, -0.37, -0.32	-22.2, -9.0, -14.4, -0.7, -0.3
STAR ( $\pi^0$ ), $\eta = 3.3$	200	4	16	0.70	-0.70	-11.3
STAR ( $\pi^0$ ), $\eta = 3.8$	200	2	16	0.57	-0.31	-5.0
STAR, $ y  < 0.5$	200	10	11.7	0.49	-0.34	-3.9

# Systematic Errors — $p/\bar{p}$

Collaboration	$\sqrt{s}$ (GeV)	# data	Norm. (%)	$\chi^2_{\text{DF}}$	$\lambda_K$	Shift (%)
TASSO	12	3	20	0.49	-0.31	-6.2
TASSO	14	9	8.5	2.33	-0.58	-4.9
TASSO	22	9	6.3	1.36	-0.98	-6.2
TASSO	30	3	20	0.52	-0.82	-16.5
JADE	34	2	14	6.04	-0.82	-16.5
TASSO	34	7	6	1.12	-1.63	-9.8
ALEPH	91.2	18	3	0.62	-1.77	-5.3
BRAHMS, $y \in [2.9, 3]$ $y \in [3.25, 3.35]$	200	7	11,7,8(13), 2,1(3)	2.38	-2.66, -1.69, -1.65, -0.48, -0.13	-29.3, -11.8, -13.2, -1.0, -0.1
		5		5.21	-3.64, -2.32, -2.47, -0.66, -0.26	-40.1, -16.2, -19.7, -1.3, -0.3
STAR, $ y  < 0.5$	200	8	11.7	3.08	-1.89	-22.1

# Hadron Mass Effects

Particle	Fitted mass (MeV)	True mass (MeV)
$\pi^\pm$	156.9	139.6
$K^\pm$	340.7	493.7
$p/\bar{p}$	949.2	938.3
$K_S^0$	363.2	497.6
$\Lambda/\bar{\Lambda}$	1127.0	1115.7

$\pi^\pm$ : large excess

→ decays from much heavier  $\rho(770) \rightarrow \pi^+ + \pi^-$

$K$ s: large undershoot

→ complicated decay channels, e.g.  $K$  resonance  $\rightarrow \pi + K$

baryons ( $p/\bar{p}$ ,  $\Lambda/\bar{\Lambda}$ ):  $\simeq 1\%$  excess

→ decays from slightly heavier resonances

good environment to study partonic fragmentation

# Hadron Mass Effects

factorization theorem:  $x$  = available c.m. light cone momenta ( $p^0 + |\mathbf{p}|$ ) fraction

$$e^+ e^- \rightarrow h(p_h) + X$$

$$p_h = \left( + = \frac{x\sqrt{s}}{\sqrt{2}}, - = \frac{m_h^2}{\sqrt{2}x\sqrt{s}}, (1, 2) = \mathbf{0} \right)$$

$\rightarrow x$  from  $x_E, x_p$

$$\text{e.g. } x_p = x \left( 1 - \frac{m_h^2}{sx^2} \right)$$

$\rightarrow$  low  $x, \sqrt{s}$  effect

$$pp(\bar{p}) \rightarrow h(p) + X$$

Partonic c.m. frame: parton  $\mathbf{l} \parallel$  hadron  $\mathbf{p}$

$$\rightarrow x = \frac{p^+}{l^+} = \frac{\sqrt{|\mathbf{p}|^2 + m_h^2} + |\mathbf{p}|}{2|\mathbf{l}|}$$

$$\rightarrow \frac{d\mathbf{p}^3}{E} = \frac{|\mathbf{p}|^2}{\sqrt{|\mathbf{p}|^2 + m_h^2}} d|\mathbf{p}| d\Omega$$

$$\rightarrow \frac{d\mathbf{l}^3}{|\mathbf{l}|} = |\mathbf{l}| d|\mathbf{l}| d\Omega$$

mass correction factor

$$E \frac{d^3}{d\mathbf{p}^3} \sigma^{h_1 h_2 \rightarrow h(p) + X}(p, \sqrt{s}) =$$

$$\sum_{i_1 i_2} \int dx_1 dx_2 f_{i_1}^{h_1}(x_1, M_f^2) f_{i_2}^{h_2}(x_2, M_f^2)$$

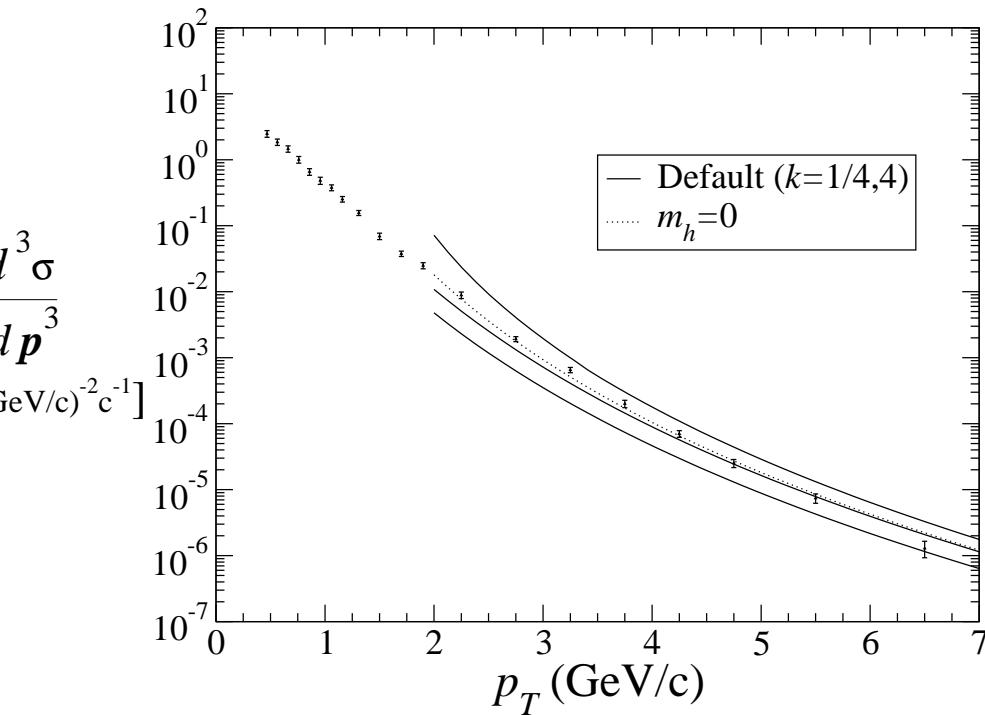
$$\times \sum_i \int dx D_i^h(x, M_f^2)$$

$$\times \frac{1}{R^2(x_1, x_2, y, m_h^2/p_T^2)} \frac{1}{x^2} |\mathbf{l}| \frac{d^3}{d\mathbf{l}^3} \sigma^{i_1 i_2 \rightarrow i(l) + X}(l, x_1 x_2 \sqrt{s})$$

# Hadron Mass Effects

(use true mass)

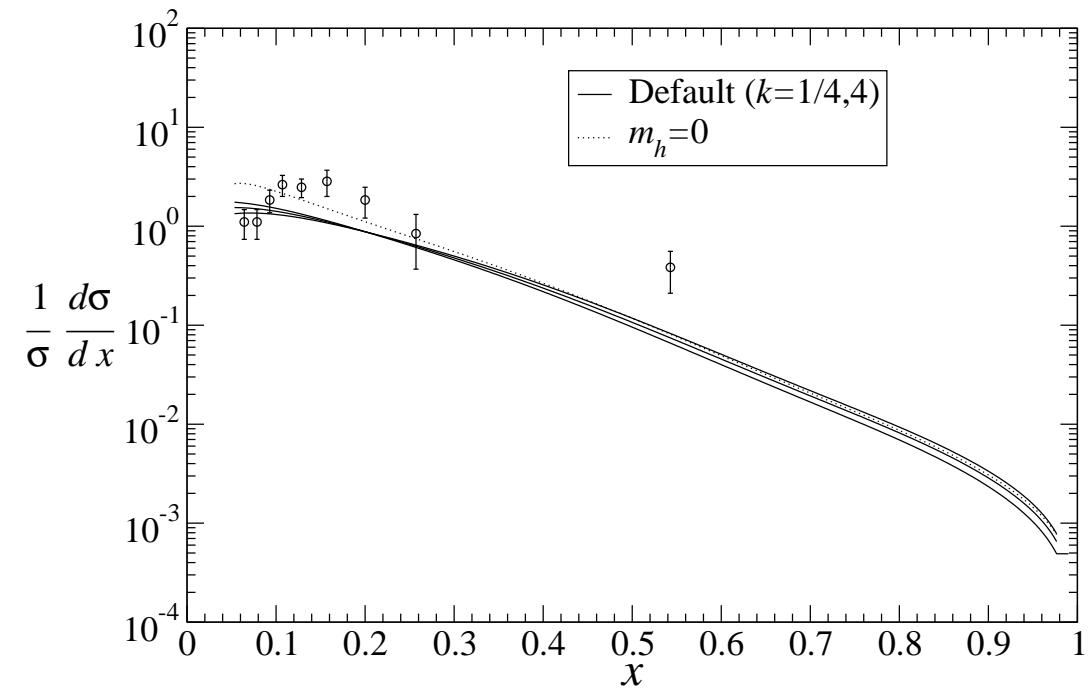
$pp \rightarrow p/\bar{p} + X$ , STAR (-0.5 <  $y$  < 0.5),  $\sqrt{s}=200$  GeV



similar results at high rapidity

(fit mass)

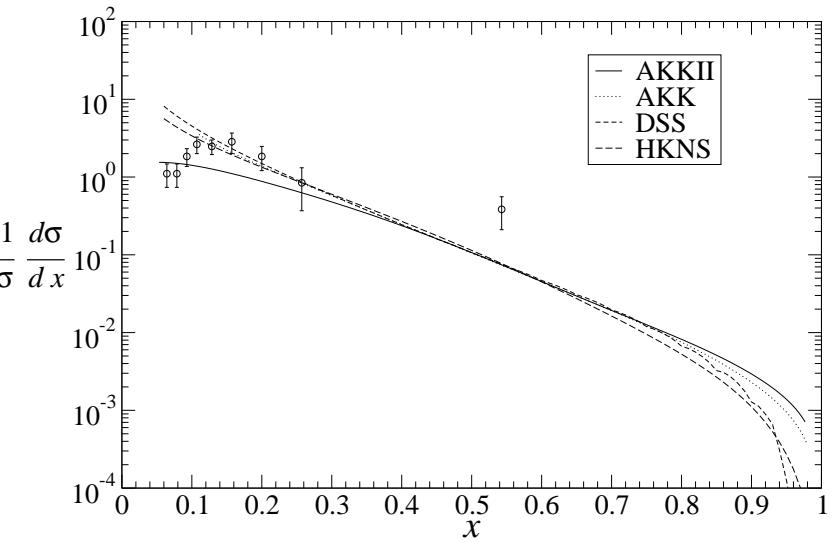
$e^+ e^- \rightarrow p/\bar{p} + X$ , TASSO,  $\sqrt{s}=14$  GeV



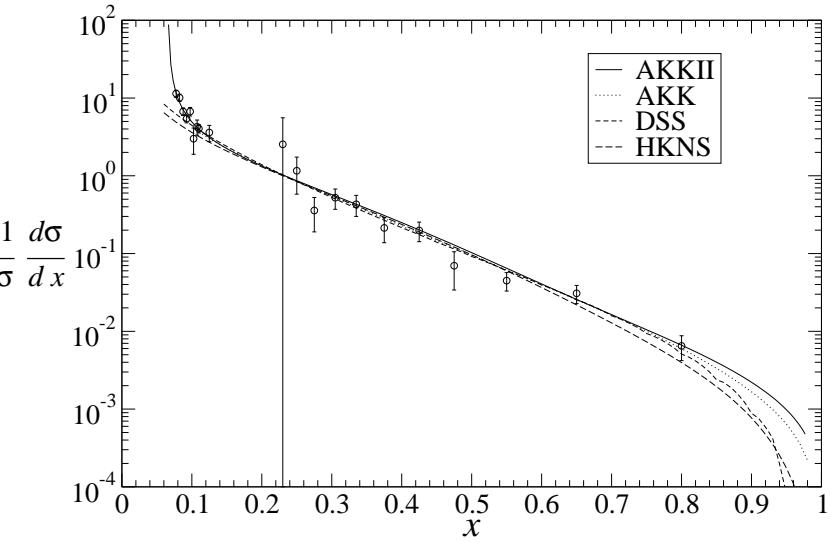
$\pi^\pm$ : mass effects negligible

# Hadron Mass Effects

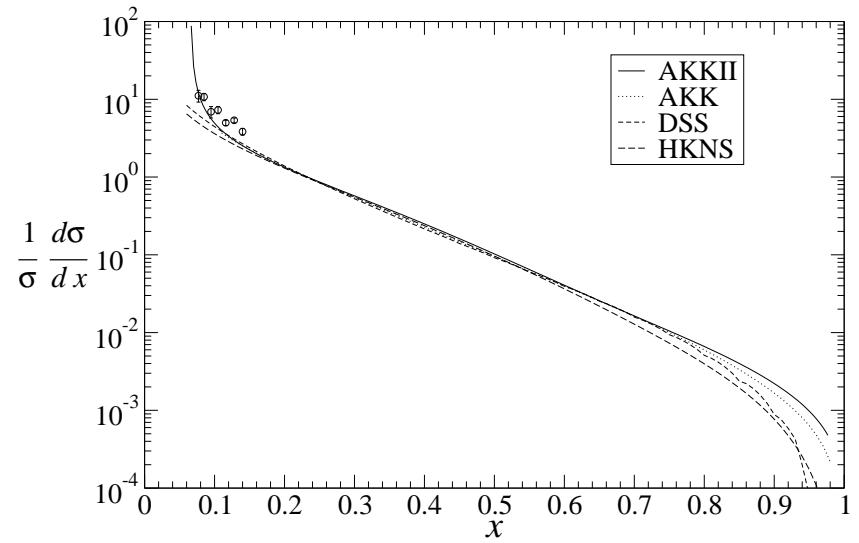
$e^+ e^- \rightarrow p/\bar{p} + X$ , TASSO,  $\sqrt{s}=14$  GeV



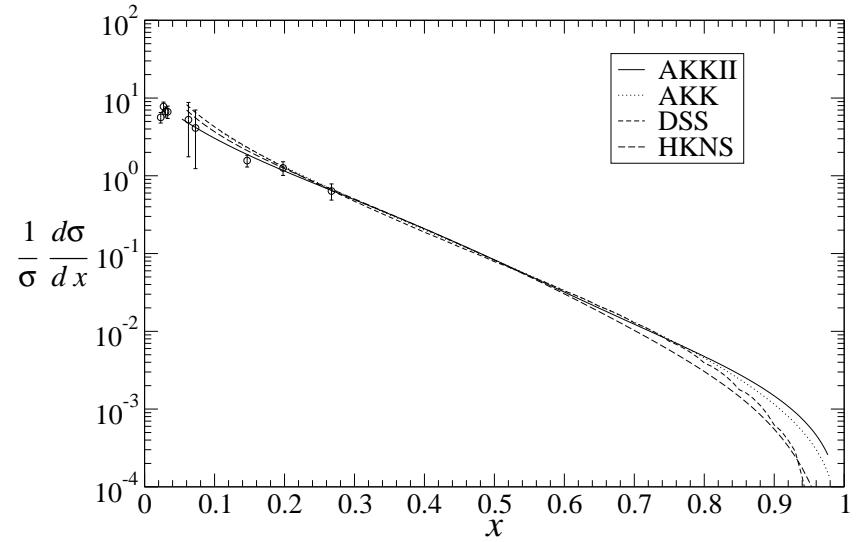
$e^+ e^- \rightarrow p/\bar{p} + X$ , TPC,  $\sqrt{s}=29$  GeV



$e^+ e^- \rightarrow p/\bar{p} + X$ , HRS,  $\sqrt{s}=29$  GeV



$e^+ e^- \rightarrow p/\bar{p} + X$ , TOPAZ,  $\sqrt{s}=58$  GeV



# Large $x$ Resummation — $e^+e^-$

$e^+e^-$  quark coefficient function  $C_q$

fixed order approach fails at large  $N$ :

$$C_q(a_s, N) \simeq 1 + A a_s \ln N + B a_s^2 \ln^2 N + \dots$$

general solution: resum  $(a_s \ln N)^n$  into single function  $f(a_s \ln N)$

(also  $a_s(a_s \ln N)^n$  into  $a_s g(a_s \ln N)$  etc.):

$$C_q(a_s, N) = C_q^r(a_s, N) \times \left( \sum_n a_s^n C_q^{\text{FO}(n)}(N) \right)$$

NLO solution: resum  $(a_s \ln N)^n$  ( $r = 0$ ),  $a_s(a_s \ln N)^n$  ( $r = 1$ ) via Cacciari et al. result

$$C_q(a_s, N) = C_q^{r=0,1}(a_s, N) \times \left( 1 + a_s(C_q^{(1)} - C_q^{r=0,1(1)}) \right)$$

also resum DGLAP evolution in  $e^+e^-$  and  $pp(\bar{p})$  [AKK PRL100(2008)192002]

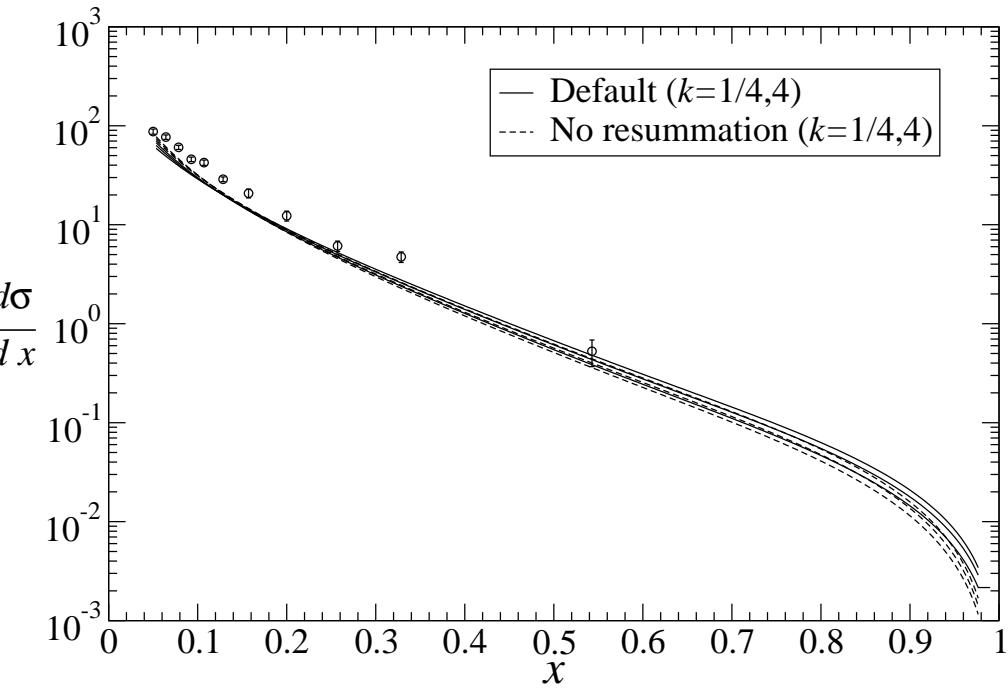
$pp(\bar{p})$ , finite  $\delta$ (rapidity): no resummation (de Florian+Vogelsang, PRD71:114004, 2005)

effects:

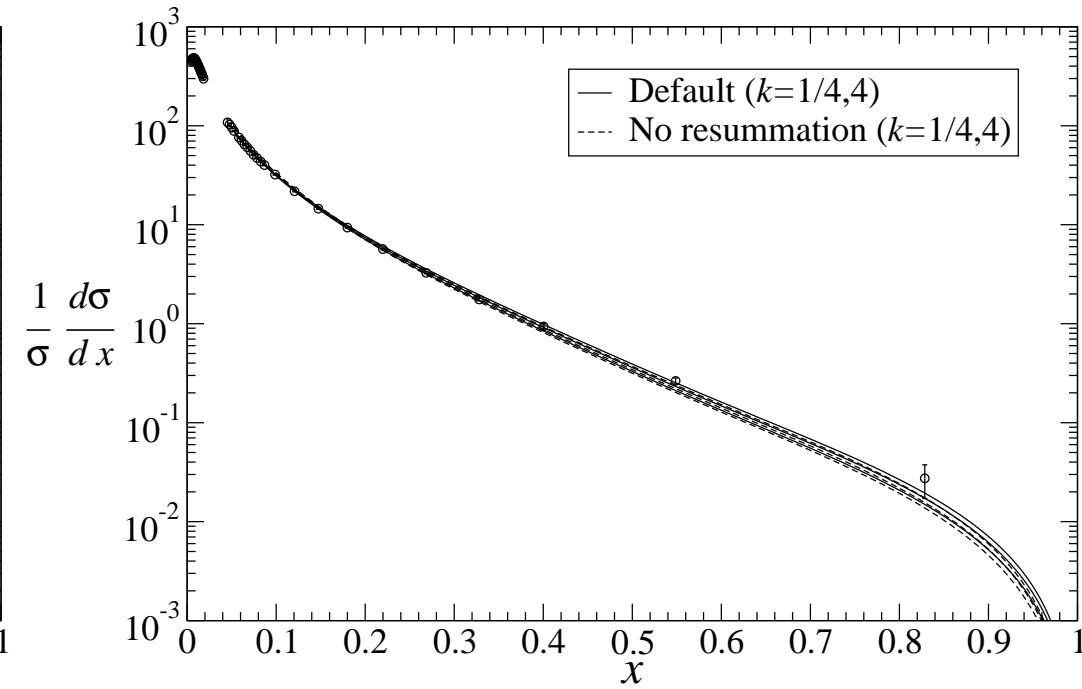
- enhances XS at large  $x$  (more so at small  $\sqrt{s}$ )
- reduces theoretical error(?)

# Large $x$ Resummation — $e^+e^-$

$e^+e^- \rightarrow \pi^\pm + X$ , TASSO,  $\sqrt{s}=14$  GeV



$e^+e^- \rightarrow \pi^\pm + X$ , OPAL,  $\sqrt{s}=91$  GeV



$H$	$\chi^2$	
	Main fit	Unres. fit
$\pi^\pm$	519.7	520.8
$K^\pm$	417.3	488.2
$p/\bar{p}$	525.1	538.0
$K_S^0$	318.6	318.8
$\Lambda/\bar{\Lambda}$	272.3	326.0

# Partonic contributions — $pp$

$$F^h(pp \rightarrow h + X) = \sum_i \widehat{F}^i(pp \rightarrow i + X) \otimes D_i^h(i \rightarrow h + X)$$

( $\widehat{F}$  contains PDFs  $f_j$ ,  $f_k$ , and  $jk \rightarrow i + X$  subprocesses)

types" of fragmenting parton  $i$ :  $u - \bar{u}$ ,  $d - \bar{d}$  (valence) and sea ( $g$ ,  $2\bar{u}$ ,  $2\bar{d}$ ,  $2s = 2\bar{s}$ , ...)

→ partonic decomposition of  $F^h$ :

$$F^{h^\pm} = \underbrace{\widehat{F}^{u_v}}_{\widehat{F}^{u_v} = \widehat{F}^u - \widehat{F}^{\bar{u}} > 0} D_u^{h^\pm} + \underbrace{\widehat{F}^{d_v}}_{\widehat{F}^{d_v} = \widehat{F}^d - \widehat{F}^{\bar{d}} > 0} D_d^{h^\pm} + \sum_{i=g, q_{\text{sea}}} \widehat{F}^i D_i^{h^\pm}$$

$$\widehat{F}^{u_v} > \widehat{F}^{d_v}:$$

$$D_u^{\pi^+} = D_d^{\pi^-} (\gg D_u^{\pi^-} = D_d^{\pi^+}) \rightarrow (pp \rightarrow u_v \rightarrow \pi^\pm) > (pp \rightarrow d_v \rightarrow \pi^\pm) \rightarrow pp \rightarrow \pi^+ > pp \rightarrow \pi^-$$

$$D_u^{p/\bar{p}} > D_d^{p/\bar{p}} \rightarrow (pp \rightarrow u_v \rightarrow p/\bar{p}) >> (pp \rightarrow d_v \rightarrow p/\bar{p}), D_{u,d}^p \gg D_{u,d}^{\bar{p}} \rightarrow pp \rightarrow p \gg pp \rightarrow \bar{p}$$

$$D_{s,\bar{s}}^{K^\pm} > D_u^{K^\pm} \gg D_d^{K^\pm} \rightarrow (pp \rightarrow \text{sea} \rightarrow K^\pm) \gg (pp \rightarrow u_v \rightarrow K^\pm) \gg (pp \rightarrow d_v \rightarrow K^\pm)$$

$$D_{u,d}^{\pi^\pm, p/\bar{p}} \gg D_{\text{sea}}^{\pi^\pm, p/\bar{p}} \text{ but } \widehat{F}^{u_v, d_v} \ll \widehat{F}^{\text{sea}} \text{ — sea or valence?}$$

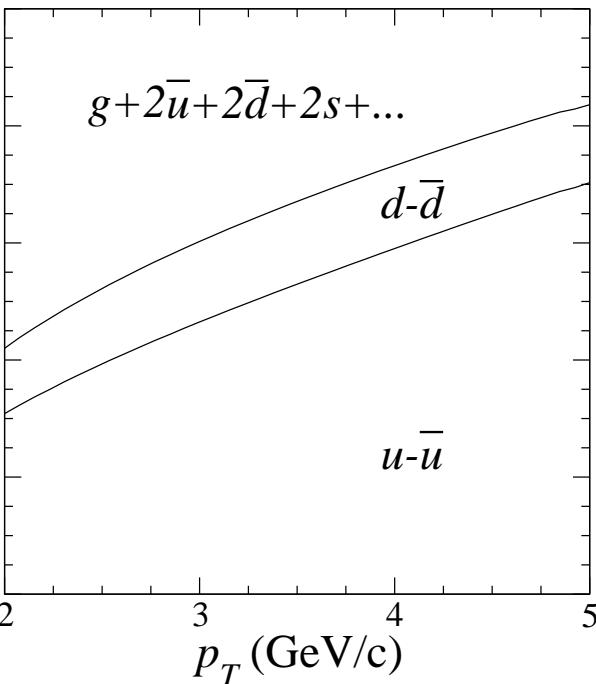
more physical, gives same qualitative results, bad for graphics:

$$F_{pp} = \underbrace{(F_{p\bar{p}} - F_{\bar{p}p})}_{\propto f_{u_v}^2} \Big|_{u_v} + \underbrace{(F_{p\bar{p}} - F_{\bar{p}p})}_{\propto f_{d_v}^2} \Big|_{d_v} + (F_{pp} + F_{\bar{p}\bar{p}} - F_{p\bar{p}})$$

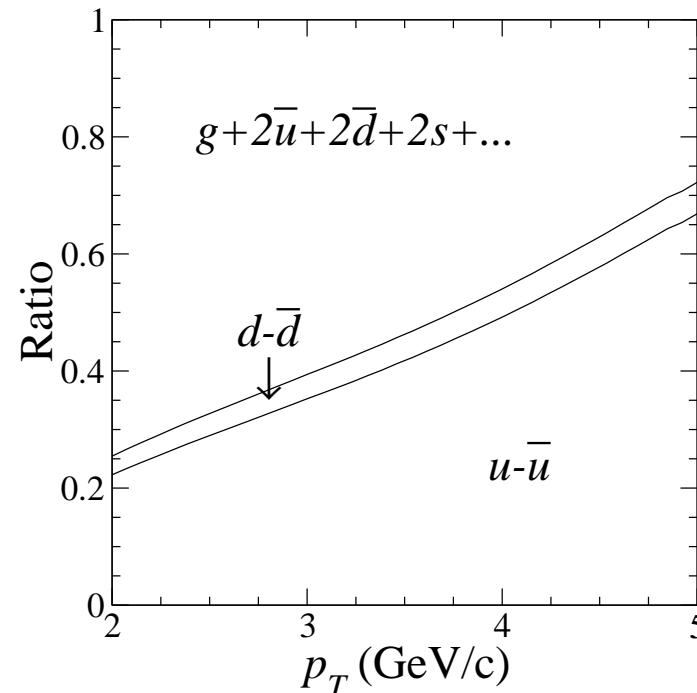
# Partonic contributions — $pp$

charge-sign unidentified

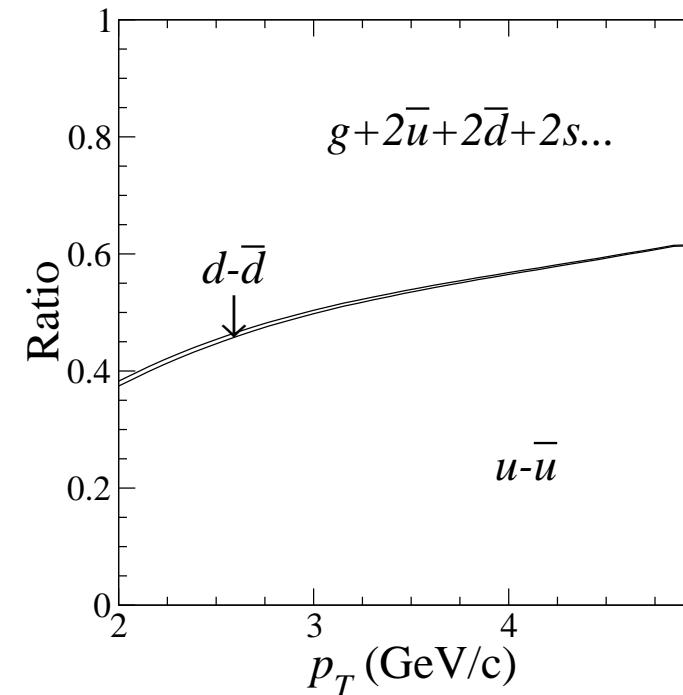
$\nu \rightarrow \pi^\pm + X \text{ (} 2.9 < y < 3 \text{), } M_f = p_T, \sqrt{s} = 200 \text{ GeV}$



$pp \rightarrow p/\bar{p} + X \text{ (} 2.9 < y < 3 \text{), } M_f = p_T, \sqrt{s} = 200 \text{ GeV}$



$pp \rightarrow K^\pm + X \text{ (} 2.9 < y < 3 \text{), } M_f = p_T, \sqrt{s} = 200 \text{ GeV}$



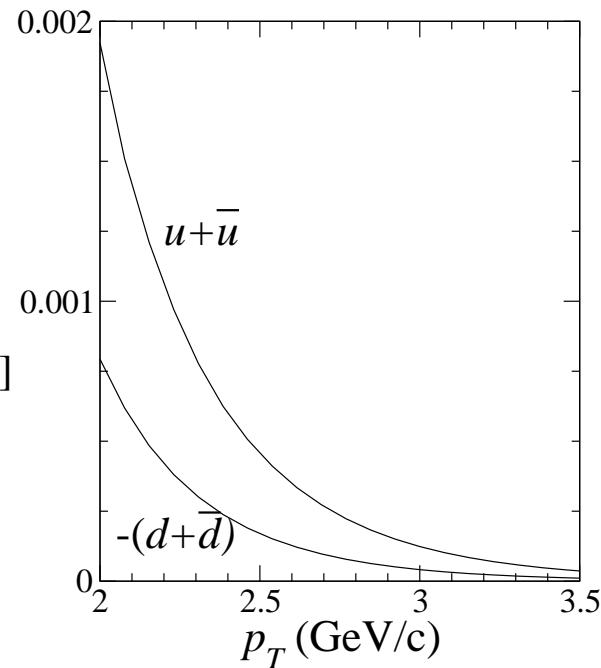
less valence at lower rapidity  
(physically,  $pp$  sea dominates)

# Partonic contributions — $pp$

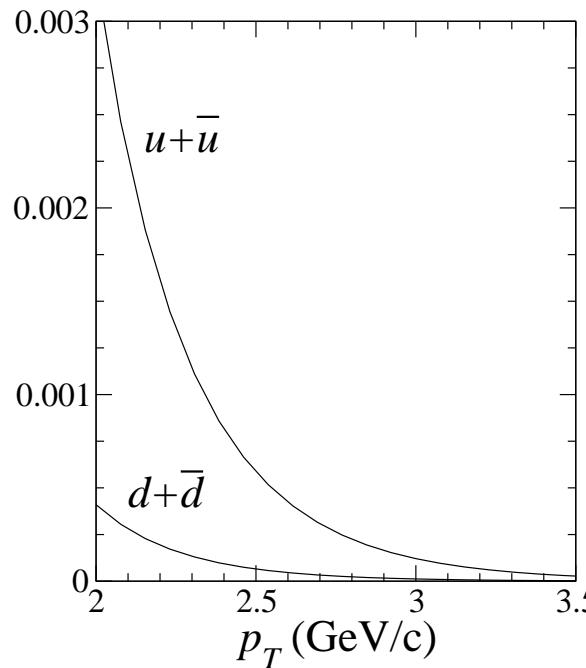
charge-sign assymetry

$$F^{\Delta_c h^\pm} = \hat{F}^{u_v} D_u^{\Delta_c h^\pm} + \hat{F}^{d_v} D_d^{\Delta_c h^\pm}$$

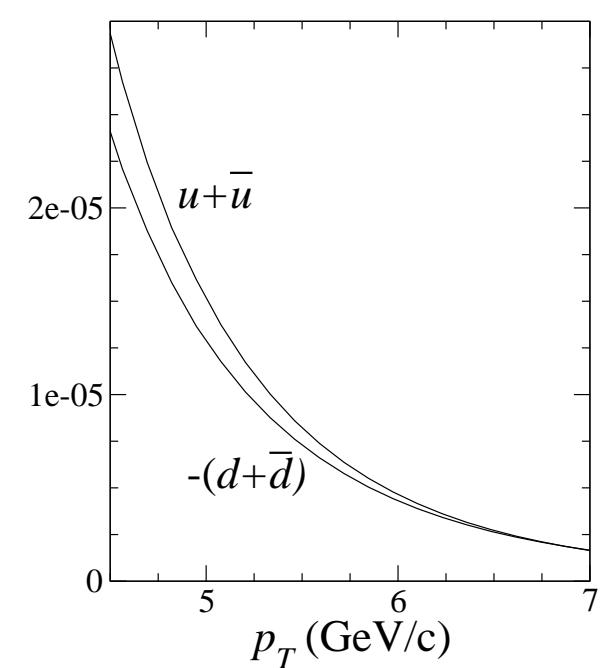
$pp \rightarrow \Delta_c \pi^\pm + X$  ( $2.9 < y < 3$ ),  $\sqrt{s}=200$  GeV



$pp \rightarrow \Delta_c p/\bar{p} + X$  ( $2.9 < y < 3$ ),  $\sqrt{s}=200$  GeV



$pp \rightarrow \Delta_c p/\bar{p} + X$  ( $-0.5 < y < 0.5$ ),  $\sqrt{s}=200$  GeV

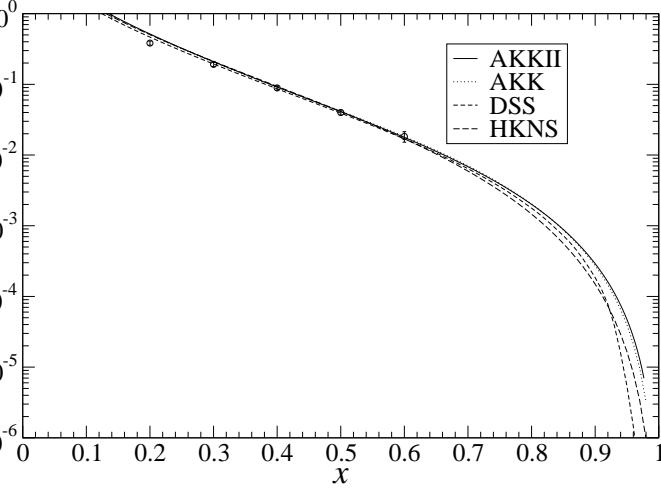


negative  $d, \bar{d} \rightarrow \Delta_c p/\bar{p}$  unexpected

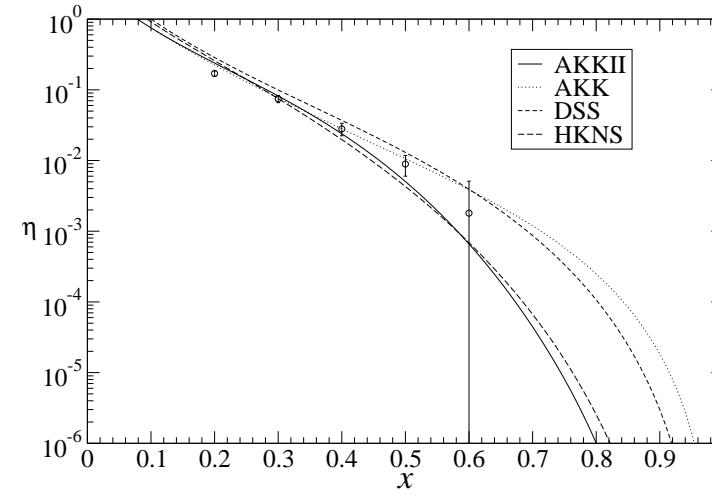
# Partonic contributions — $e^+e^-$

OPAL tagging probabilities are obvious physical candidates

$e^+e^- \rightarrow u\bar{u} \rightarrow \pi^\pm + X$ , OPAL,  $\sqrt{s}=91$  GeV

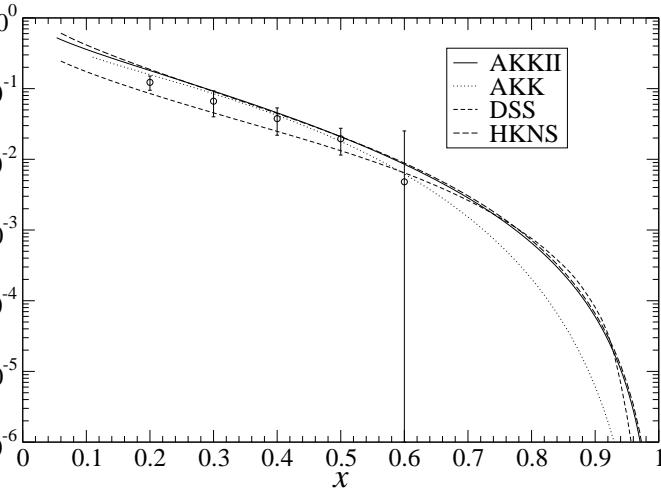


$e^+e^- \rightarrow s\bar{s} \rightarrow \pi^\pm + X$ , OPAL,  $\sqrt{s}=91$  GeV

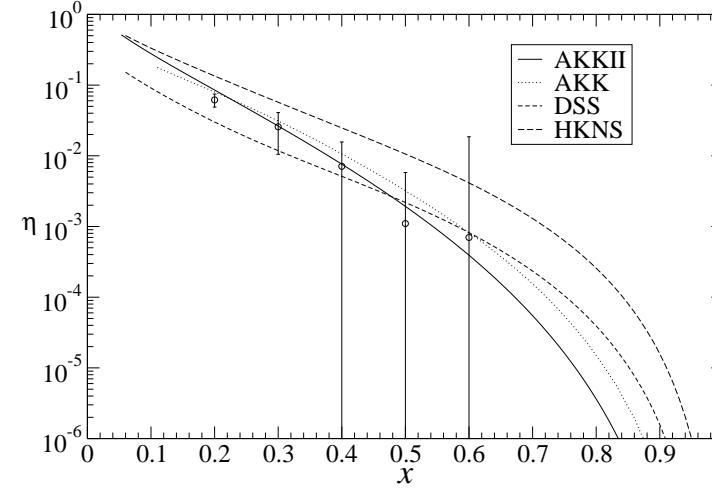


unfavoured quarks:  
large  $x$  uncertain

$e^+e^- \rightarrow u\bar{u} \rightarrow K^\pm + X$ , OPAL,  $\sqrt{s}=91$  GeV

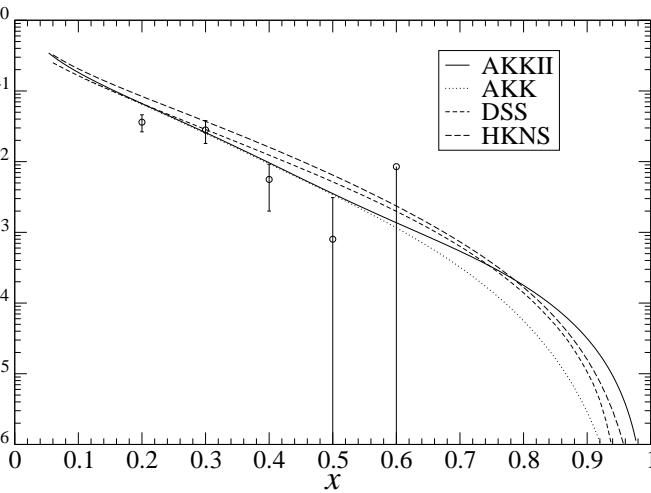


$e^+e^- \rightarrow d\bar{d} \rightarrow K^\pm + X$ , OPAL,  $\sqrt{s}=91$  GeV

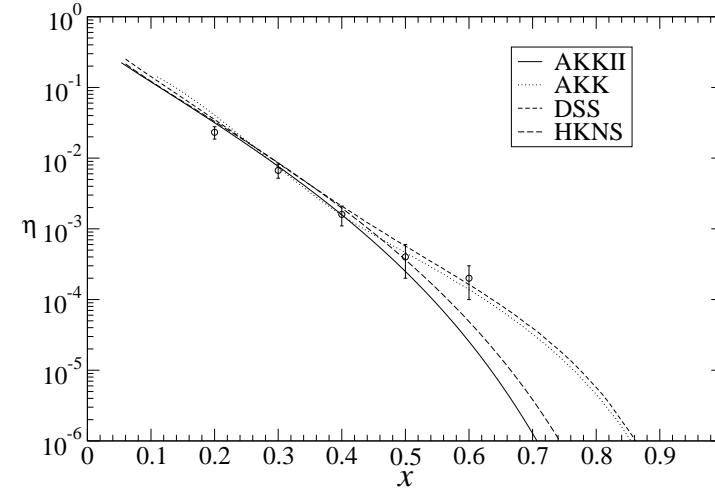


# Partonic contributions — $e^+e^-$

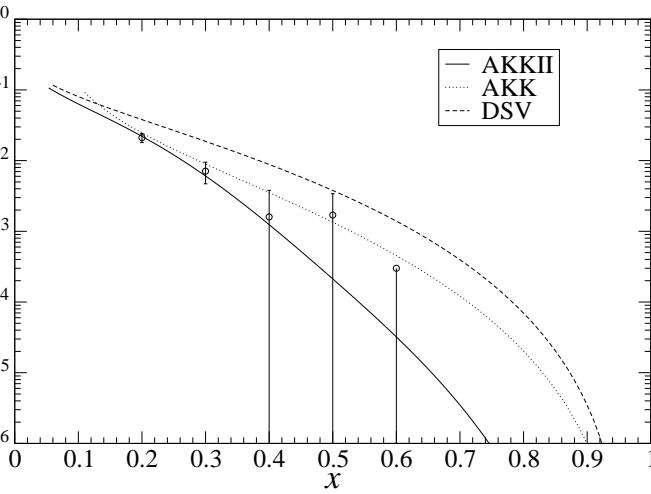
$e^+e^- \rightarrow d\bar{d} \rightarrow p/\bar{p} + X$ , OPAL,  $\sqrt{s}=91$  GeV



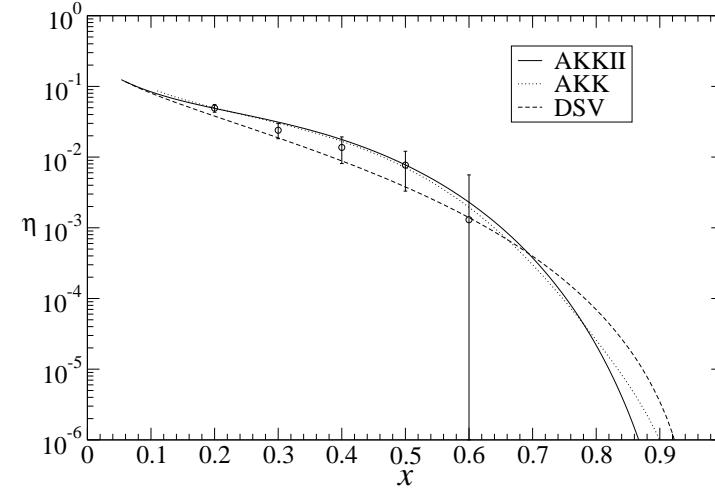
$e^+e^- \rightarrow b\bar{b} \rightarrow p/\bar{p} + X$ , OPAL,  $\sqrt{s}=91$  GeV



$e^+e^- \rightarrow u\bar{u} \rightarrow \Lambda/\bar{\Lambda} + X$ , OPAL,  $\sqrt{s}=91$  GeV



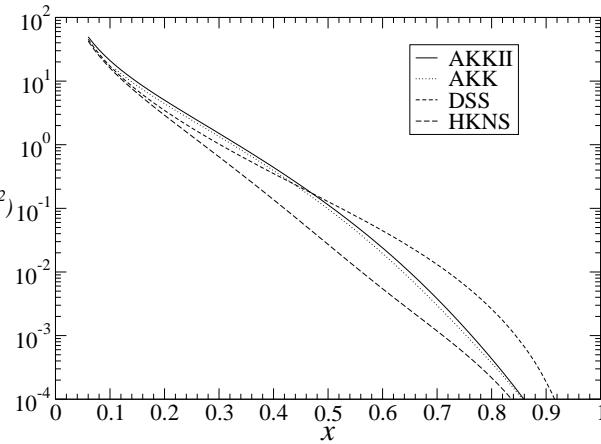
$e^+e^- \rightarrow s\bar{s} \rightarrow \Lambda/\bar{\Lambda} + X$ , OPAL,  $\sqrt{s}=91$  GeV



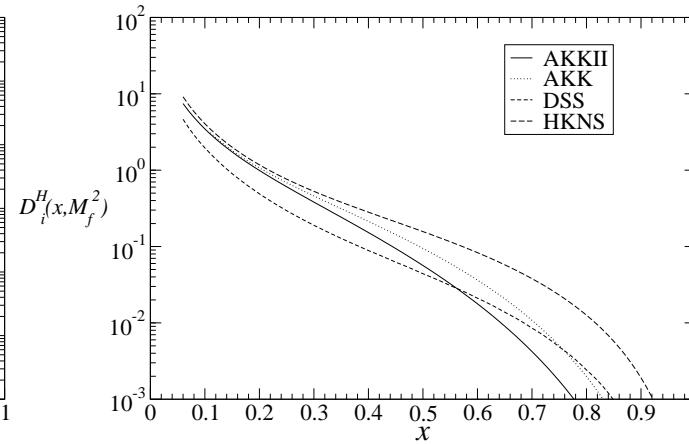
$\Lambda/\bar{\Lambda}$ :  $s$  more favoured than  $u, d$

# Gluon FFs

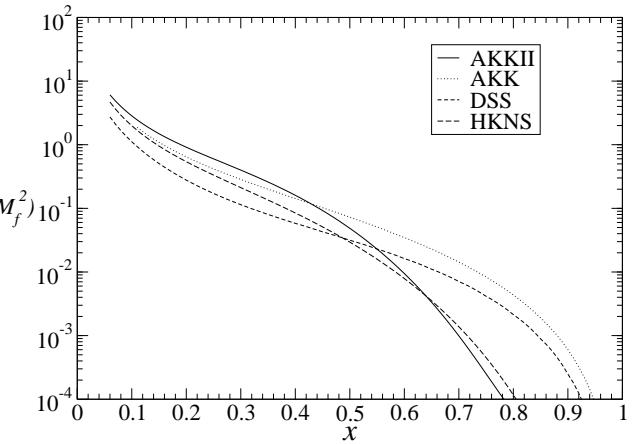
$H = \pi^\pm, i=g, M_f = 91.2 \text{ GeV}$



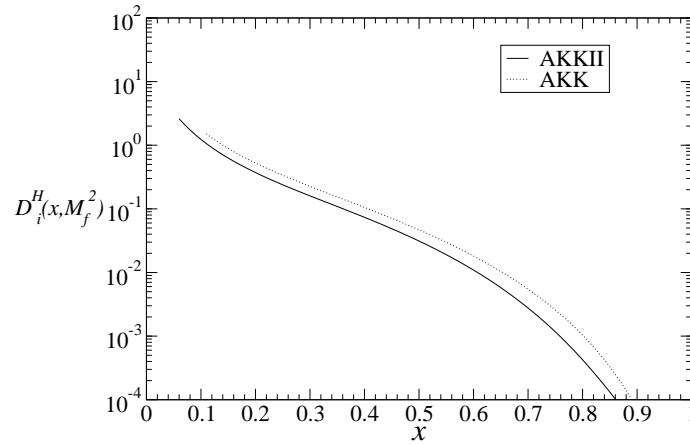
$H = K^\pm, i=g, M_f = 91.2 \text{ GeV}$



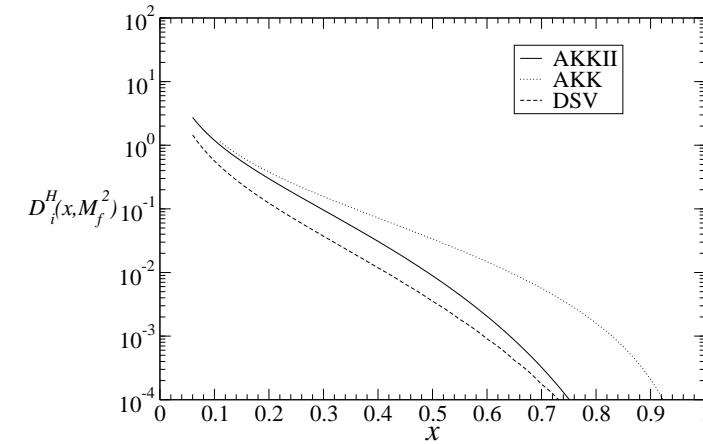
$H = p\bar{p}, i=g, M_f = 91.2 \text{ GeV}$



$H = K_S^0, i=g, M_f = 91.2 \text{ GeV}$

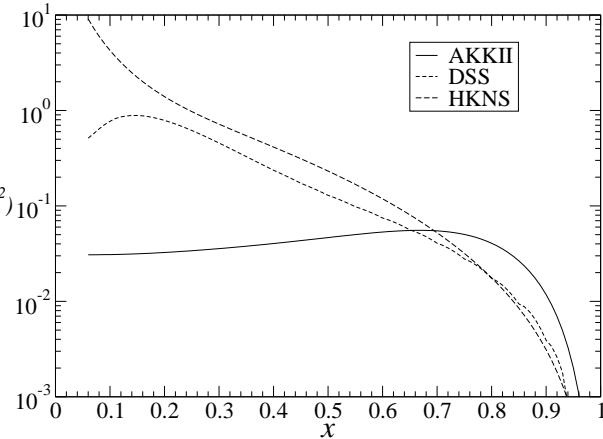


$H = \Lambda/\bar{\Lambda}, i=g, M_f = 91.2 \text{ GeV}$

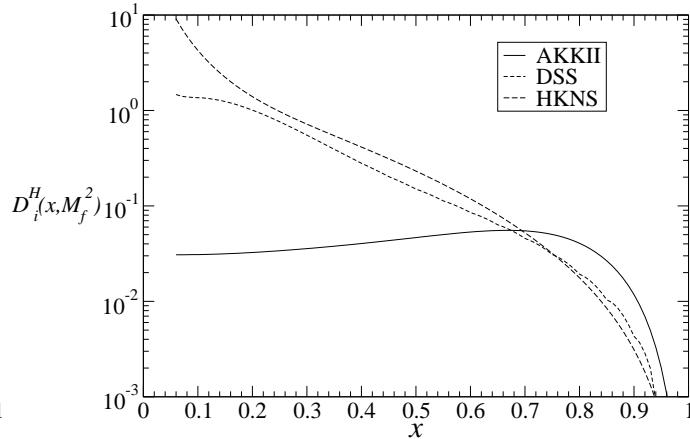


# Charge-sign assymetry FFs

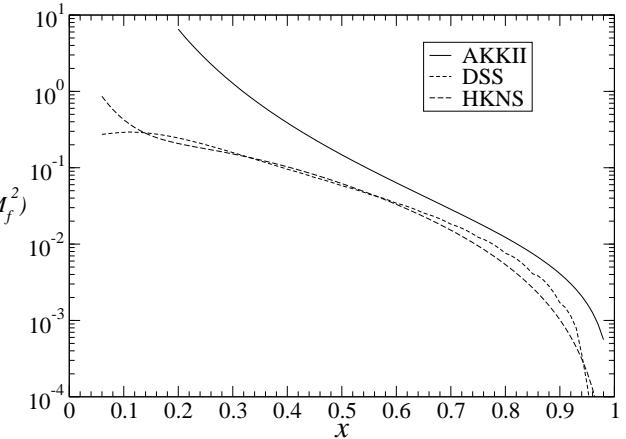
$H = \Delta_c \pi^\pm, i=u, M_f=91.2 \text{ GeV}$



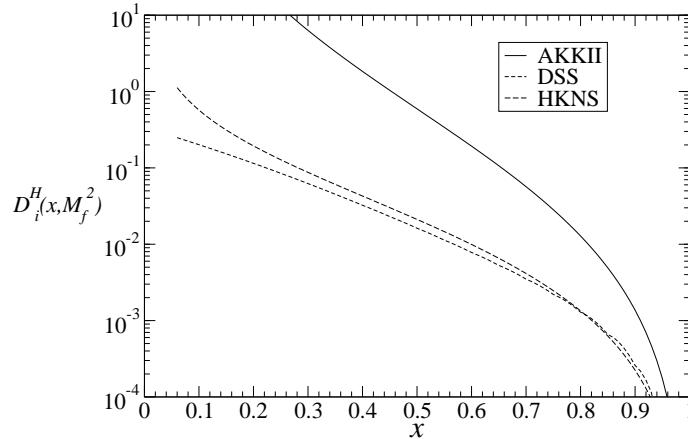
$H = \Delta_c \pi^\pm, i=\bar{d}, M_f=91.2 \text{ GeV}$



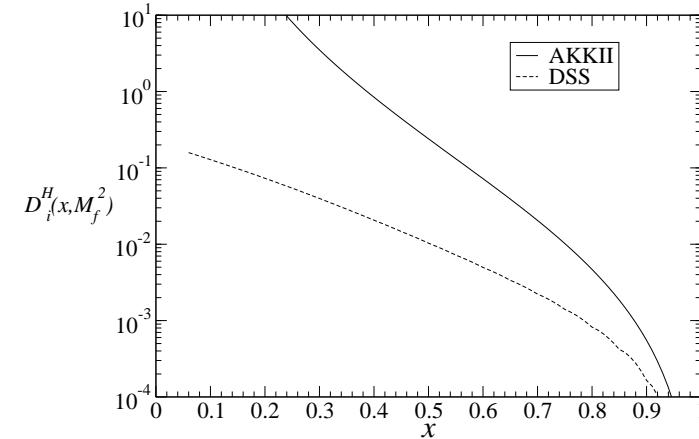
$H = \Delta_c K^\pm, i=u, M_f=91.2 \text{ GeV}$



$H = \Delta_c p/\bar{p}, i=u, M_f=91.2 \text{ GeV}$



$H = \Delta_c p/\bar{p}, i=d, M_f=91.2 \text{ GeV}$



BRAHMS → higher FF

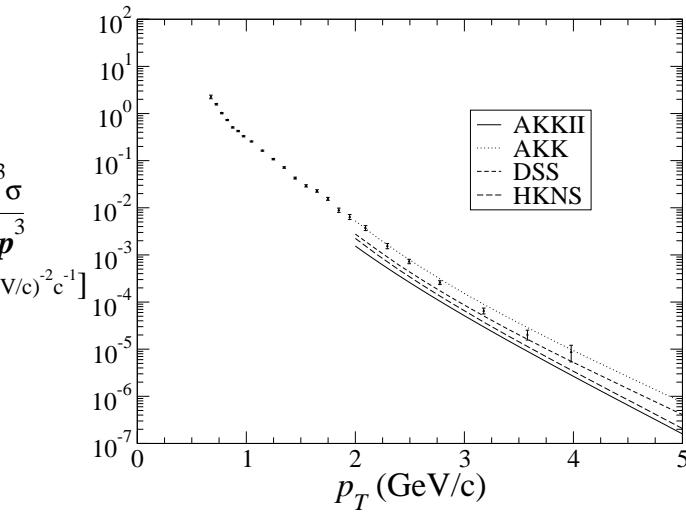
(HKNS: negative)

HKNS: no  $\Delta_c h^\pm$  data

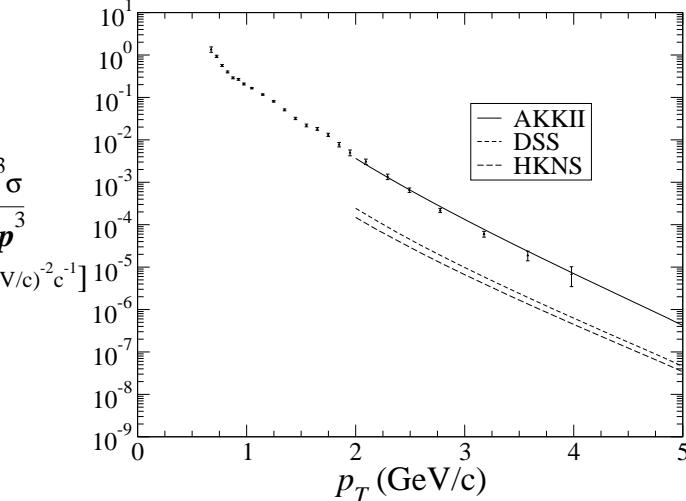
HKNS ∼ DSS → ∼ non perturbative assumptions

# $pp \rightarrow$ baryon issues

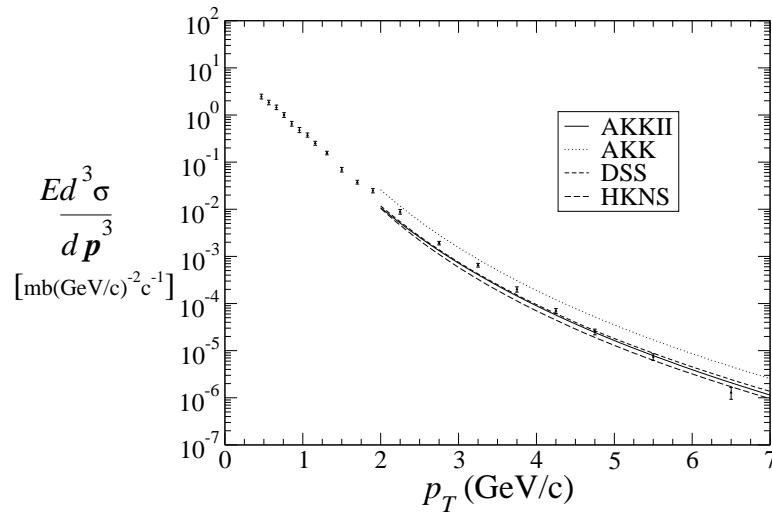
$pp \rightarrow p/\bar{p} + X$ , BRAHMS ( $2.9 < y < 3$ ),  $\sqrt{s}=200$  GeV



$pp \rightarrow \Delta_c p/\bar{p} + X$ , BRAHMS ( $2.9 < y < 3$ ),  $\sqrt{s}=200$  GeV

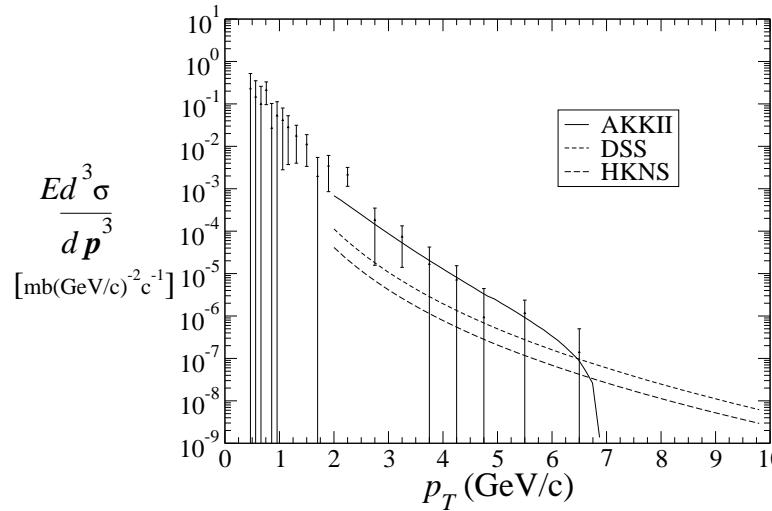


$pp \rightarrow p/\bar{p} + X$ , STAR ( $-0.5 < y < 0.5$ ),  $\sqrt{s}=200$  GeV



new fit:  
slightly lower  $p/\bar{p}$

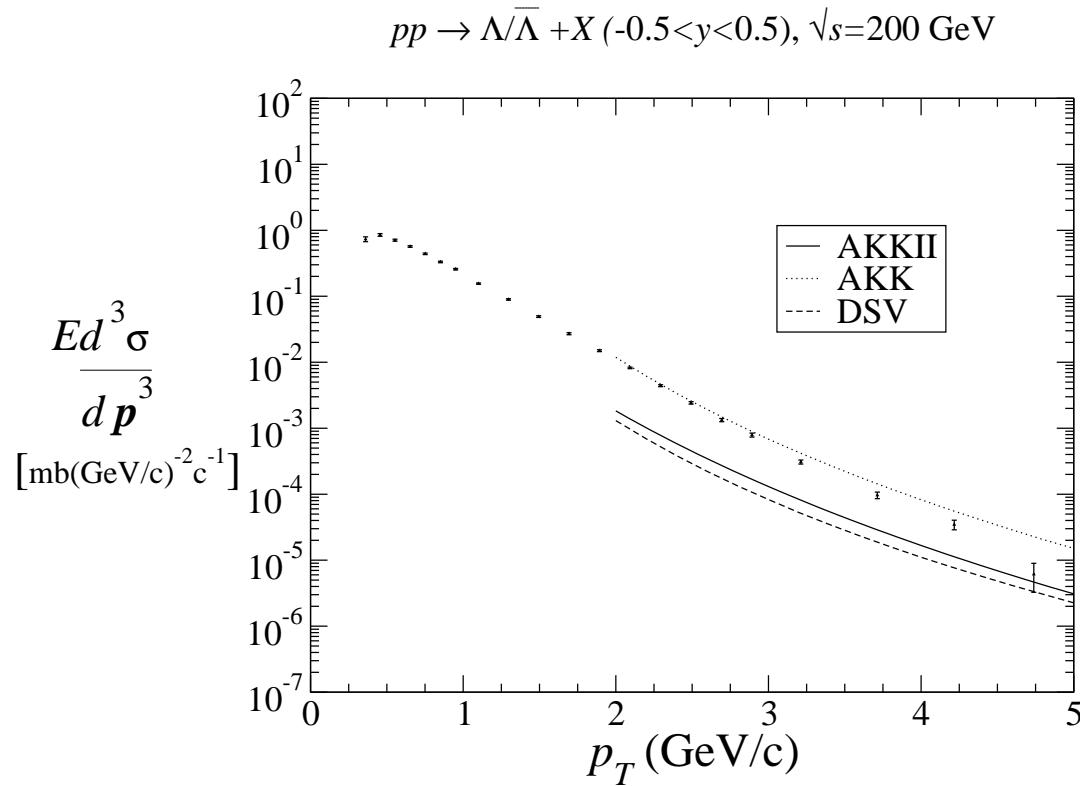
$pp \rightarrow \Delta_c p/\bar{p} + X$ , STAR ( $-0.5 < y < 0.5$ ),  $\sqrt{s}=200$  GeV



much higher  
 $\Delta_c p/\bar{p}$

# $pp \rightarrow$ baryon issues

previous AKK gluon FF:  $\Lambda/\bar{\Lambda}$  fixed to 1/3  $p/\bar{p}$   
→ good agreement with STAR

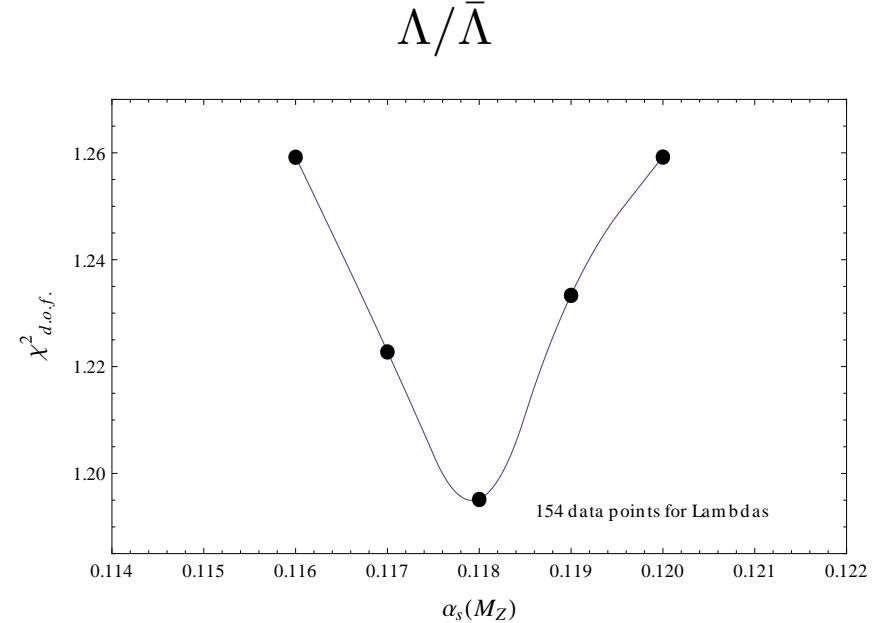
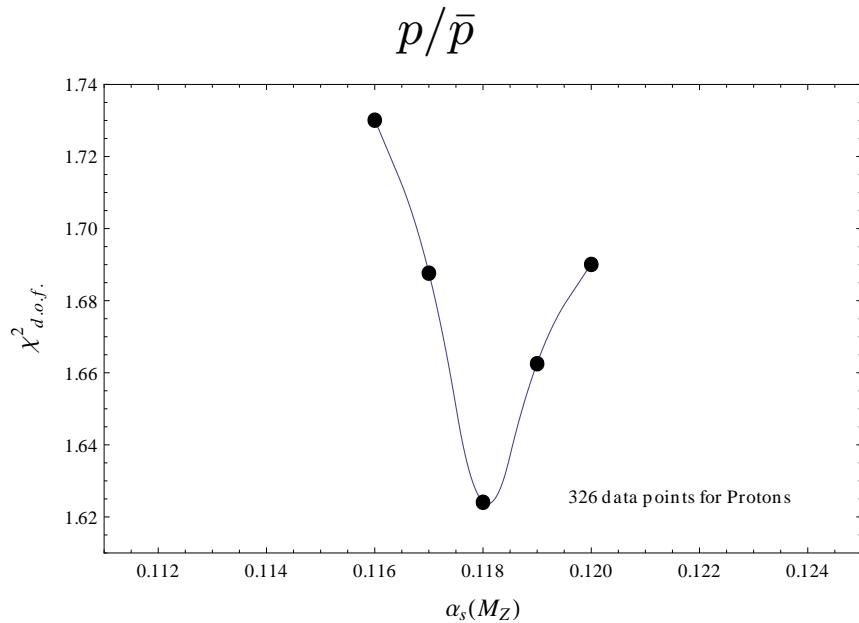
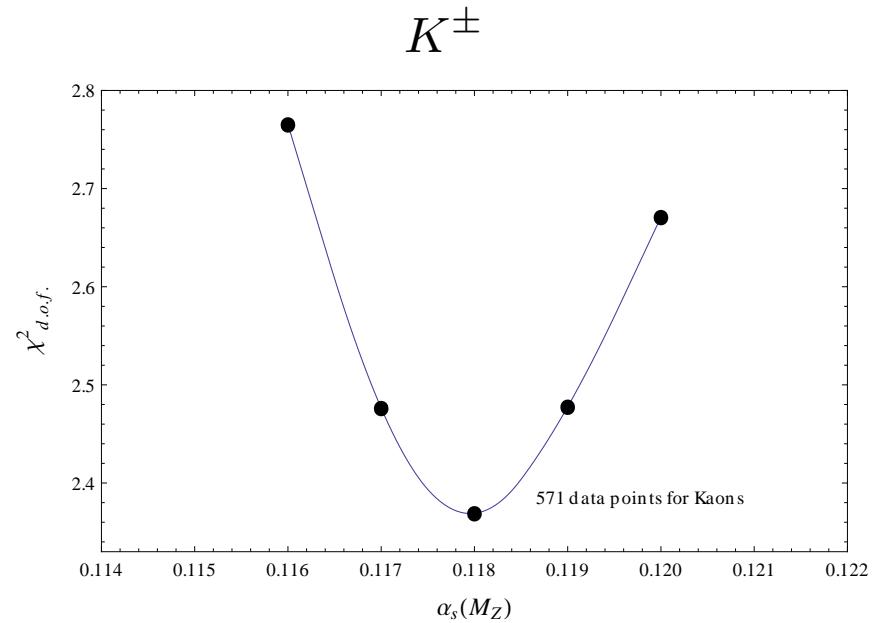
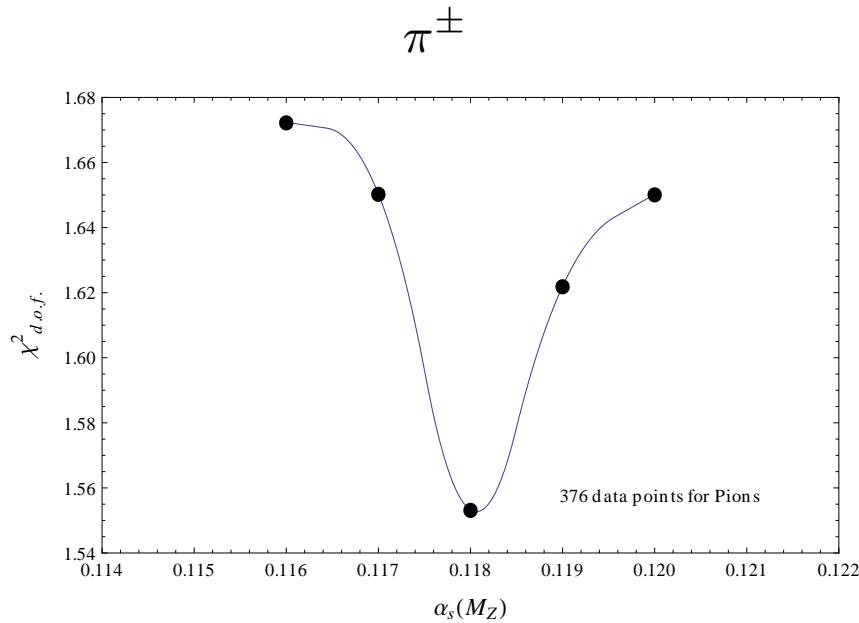


full error analysis required to determine (in)consistency between  $e^+e^-$  and  $pp$

# ABKK12 Update [in preparation]

- AKK08 uses CTEQ6.5S0 PDFs w/  $\Lambda_{\overline{\text{MS}}}^{(5)} = 226 \text{ MeV}$  for incoming  $p/\bar{p}$   
 $\leadsto \alpha_s^{(5)}(m_Z)$  not determined by fit
- Determine  $\alpha_s^{(5)}(m_Z)$  from scaling violations in FFs
- Exclude  $pp, p\bar{p}$  data from fit, but retain all  $e^+e^-$  data
- Include precise Belle data on  $\pi^\pm, K^\pm, p/\bar{p}$  inclusive production from below the  $\Upsilon(4S)$  resonance [[hep-ex/0406017](#)]
- Adopt theoretical input from AKK08
- Individual fits to  $\pi^\pm, K^\pm, p/\bar{p}, \Lambda/\bar{\Lambda}$  all yield  $\alpha_s^{(5)}(m_Z) = 0.118$

# ABKK12 Update (cont.)



# Summary

- FFs for  $\pi^\pm, K^\pm, p/\bar{p}, K_S^0$  and  $\Lambda/\bar{\Lambda}$ ,
- FFs for  $\Delta_c \pi^\pm, \Delta_c K^\pm, \Delta_c p/\bar{p}$ , separate fits
- data from  $e^+e^-$ ,  $pp(\bar{p})$
- hadron mass effects, fitted in  $e^+e^-$
- slight overshoot for pions: decays from heavier particles?
- mass deficiency for kaons: complicated decay channels?
- mass good for baryons,  
better stage for partonic fragmentation?
- investigate baryon problems further  
namely  $\Lambda/\bar{\Lambda}$  undershoot, -ve valence  $d \rightarrow \Delta_c p/\bar{p}$  @ STAR
- $e^+e^-$  large  $x$  resummation generally improves fit
- norm. error as systematic errors in covariance matrix,  $\lambda_K$  fitted

# Ongoing and future work

- $\alpha_s$  determination
- full error analysis for FFs  
→ (in)consistencies with other FFs + data
- additional theory input
- future high  $Q^2$   $ep$  from HERA — hadron species and charge identification,  
measurement of charge-sign asymmetry?