

研究計画B03:

冷却原子を用いた中性子過剰な低密度核物質の 状態方程式

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理化学研究所(理論) : 中務孝(連携)

☆冷却原子 \leftrightarrow 原子核

岡山大学(理論) : 水島健(連携)

☆p波相互作用

APCTP(理論) : 渡辺元太郎(協力)

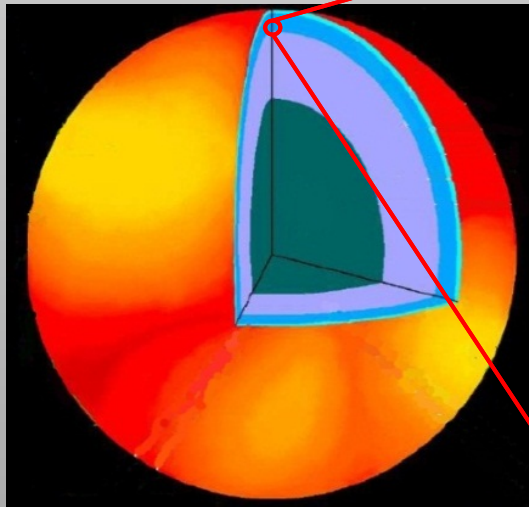
☆冷却原子 \leftrightarrow 原子核

Our mission

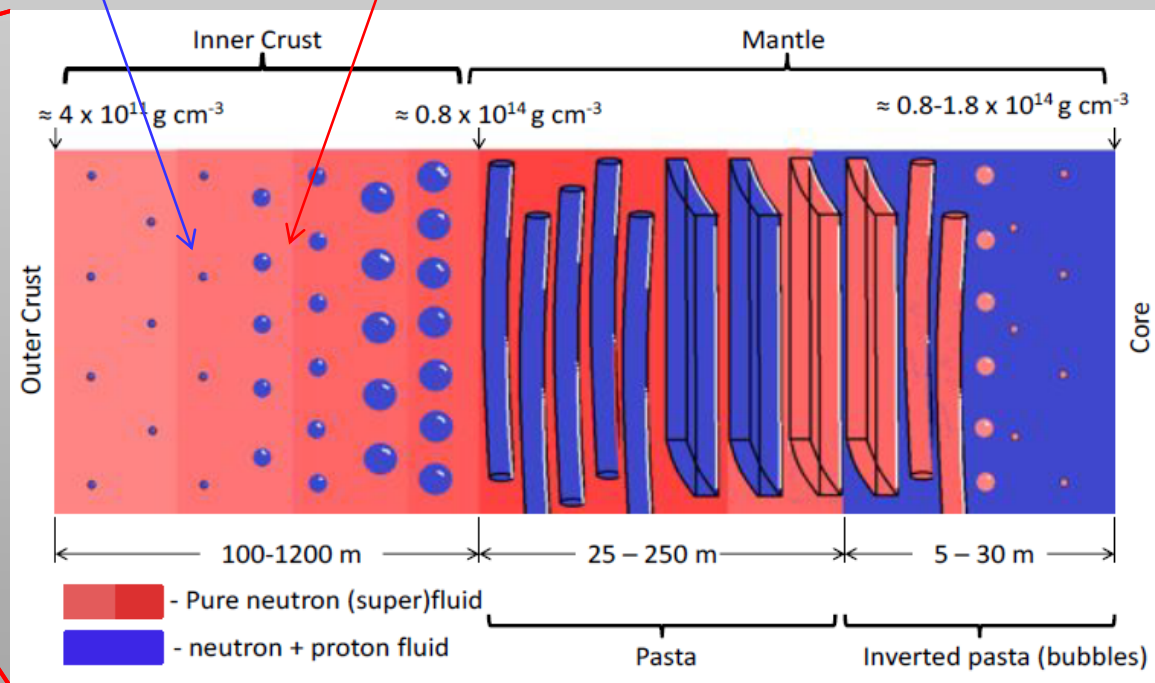
Simulation of neutron matter using cold Fermi gas

bcc crystal lattice of nuclei

Dripped neutrons from nuclei



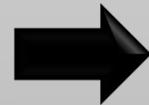
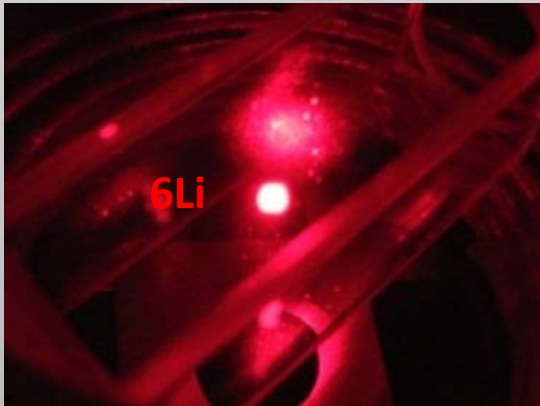
Neutron star



inner crust region [Newton, *et al.*, arXiv:1112.2018]

Cold atom system

Laser cooling
(magneto-optical trap)

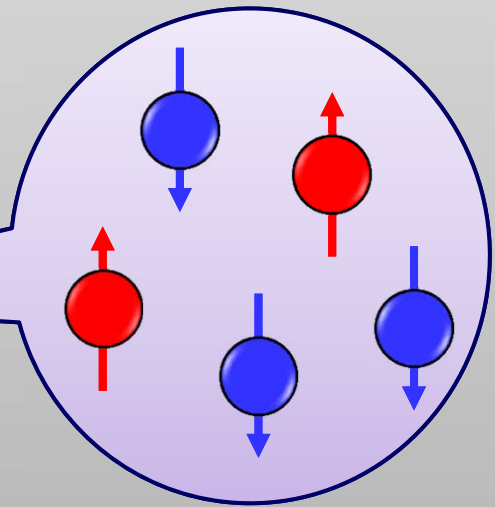


Optical dipole trap

$N \sim 10^4$

$T \sim 100\text{nK}$

Two-component Fermi gas



- Approximately **homogeneous** system
- **Dilute** two-components Fermi systems
- **Tunable interactions** by Feshbach resonances
- **Universal** many-body system
- **Precise measurements**

Cold atom vs. Neutron matter

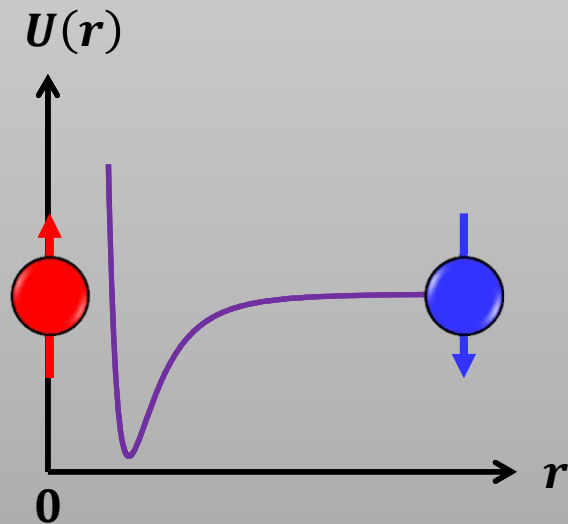
	Cold Fermi atom	Neutron matter
Interaction	s-wave	s-wave
Temperature : T	$\sim 10^{-7}$ K	$\sim 10^9$ K
Fermi temperature : T_F	$\sim 10^{-6}$ K	$\sim 10^{10}$ K
Interparticle distance : $\sim k_F^{-1}$	100nm	6~3fm
Scattering length : a	$-\infty \sim \infty$ (Feshbach resonance)	-18.5 ± 0.3 fm
Effective range : r_e	4.7nm	2.75 ± 0.11 fm
↓	↓	↓
Temperature : T/T_F	10~0.05	0.1~0
Interaction : $-1/k_F a$	$-\infty$ (BEC limit) ~ $+\infty$ (BCS limit)	0.28~0.04
Effective range : $k_F r_e$	0.05	0.53~3.3
Phase transition : T_C/T_F	~0.2	~0.1
Superfluid gap : Δ/E_F	~0.6	~0.2
Lattice potential	Optical lattice, Ion crystal	Nuclei crystal

Universal many-body Fermi system

Why can we simulate neutron matter using cold atoms ?

Because both systems belong to the **universal** system !

Basic concept



The potential shape depends on
details of the particle

Parameterization of
scattering



Ultralow temperature
Diluteness
Two-body elastic collision

Two-body s-wave phase shifts
No details of the particle

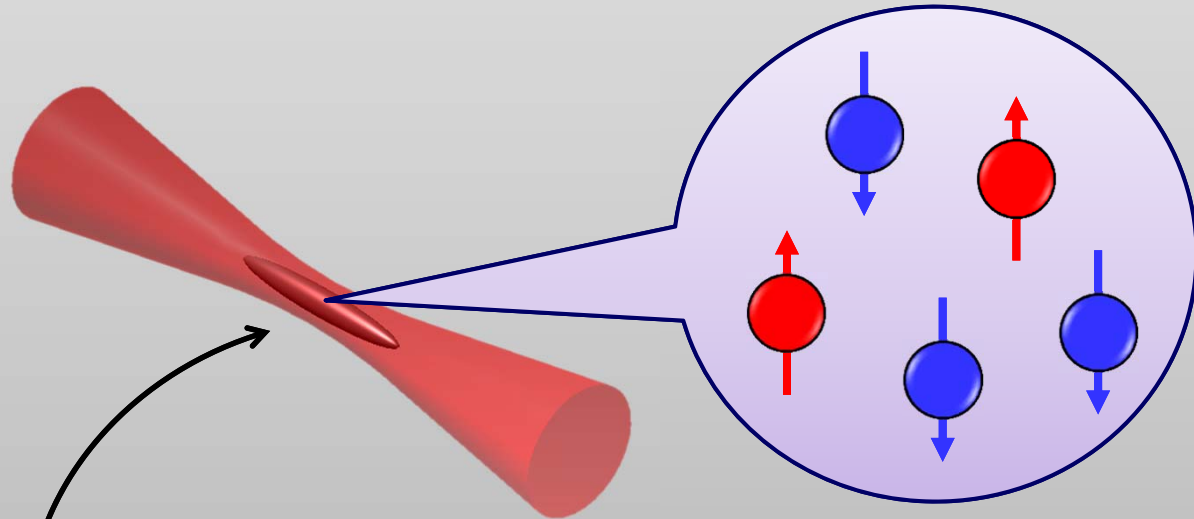
$$k \cot \delta_k = -\frac{1}{a} + \frac{1}{2} r_e k^2 + \mathcal{O}(k^4)$$

Scattering length

Effective range

Simulation of neutron matter using cold atoms

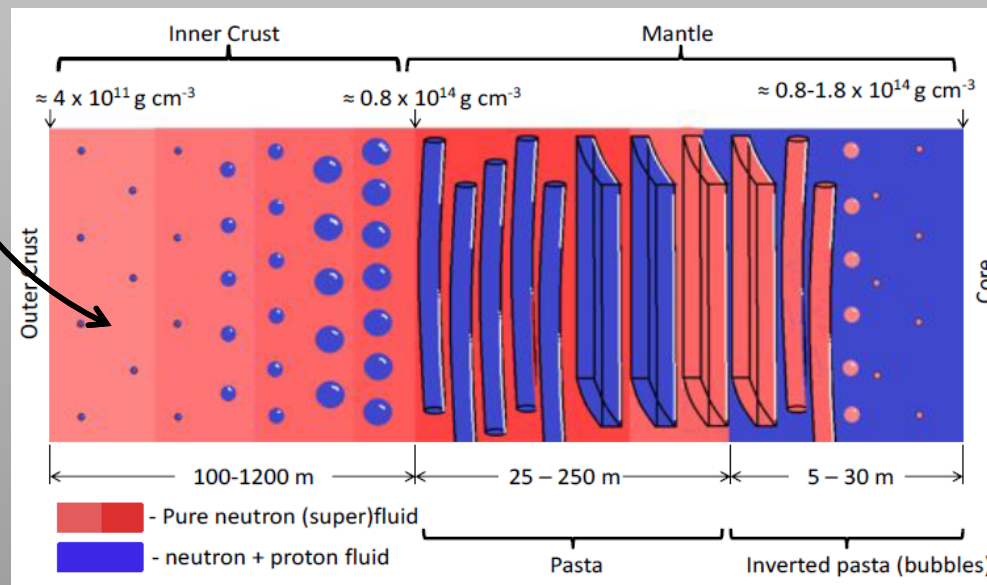
Cold atom



Internal energy :

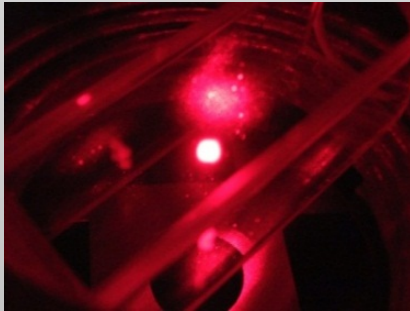
$$\frac{E}{E_F} \left(\frac{T}{T_F}, -\frac{1}{k_F a}, k_F r_e \right)$$

Inner crust



Our goal

Cold atoms

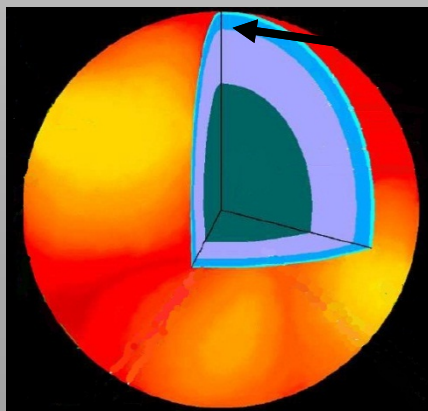


Measurements



- Universal equation of state
- Internal energy
- Specific heat
- Critical temperature
- s,p-wave superfluid
- Superfluid gap
- Superfluid density
- ⋮
- Benchmark for theories

Inner crust of neutron stars



EOS

Cooling curve

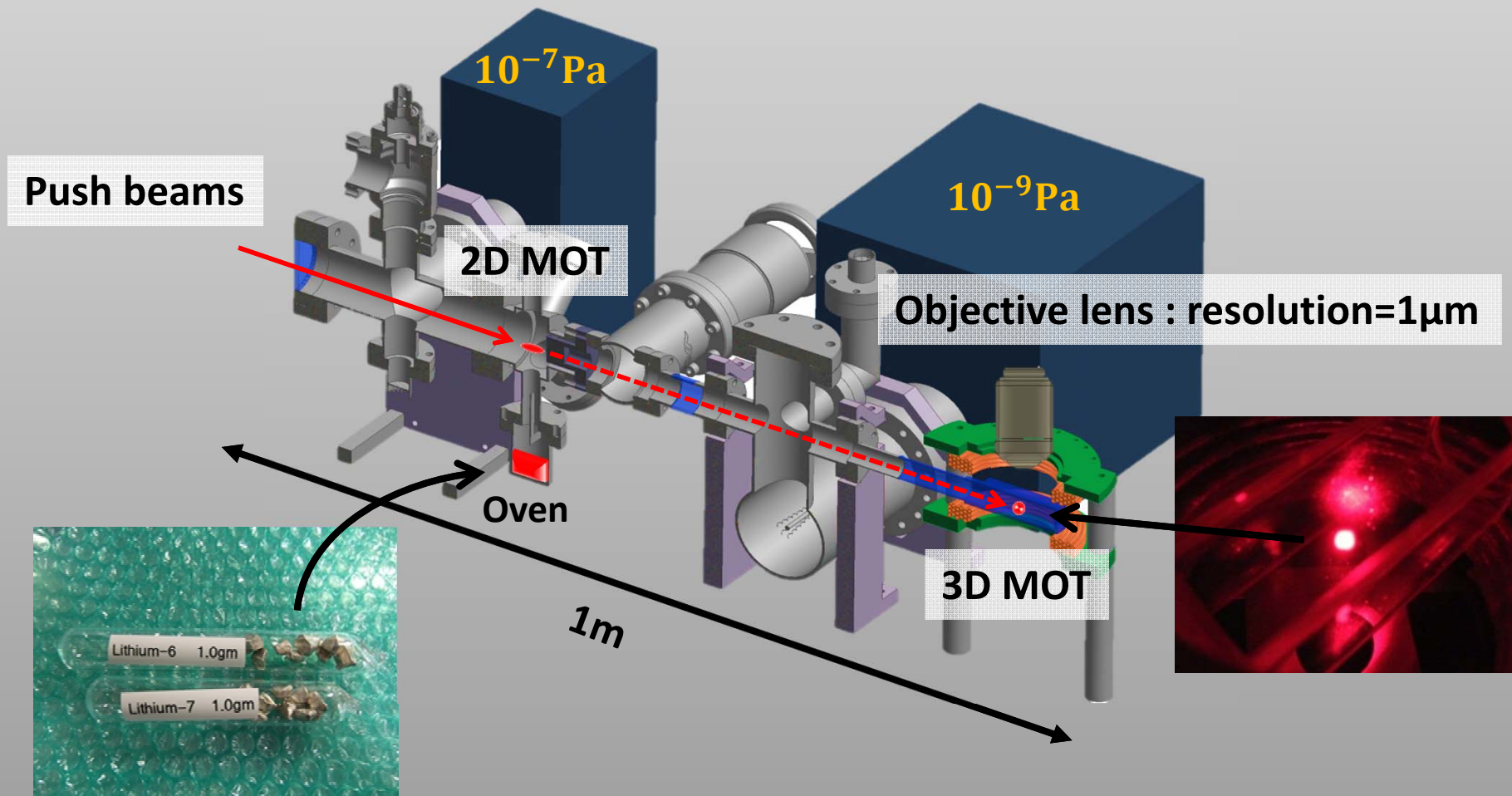
- ◆ Correction of the effective range
- ◆ Lattice of neutron-rich nuclei
- ◆ Protons

Theories



Cold Atom Lab at University of Tokyo

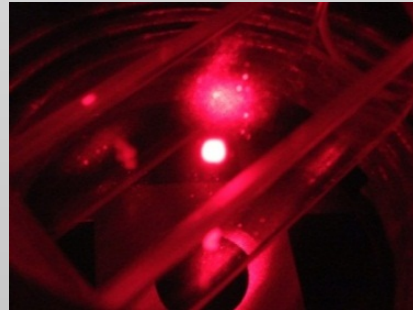
s-wave interaction, ${}^6\text{Li}$ - ${}^7\text{Li}$ (Fermi-Bose) mixture



Present status

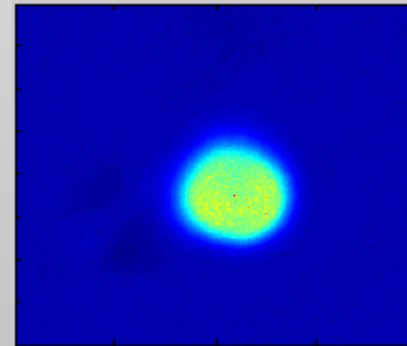


2011/April



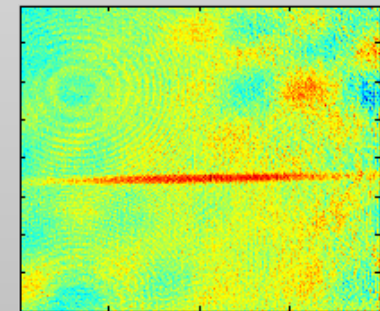
Simultaneous MOT
of 6Li and 7Li

$$N_{6\text{Li}} \sim N_{7\text{Li}} \sim 10^8$$
$$T \sim 1\text{mK}$$



Compression-MOT

$$N_{6\text{Li}} \sim N_{7\text{Li}} \sim 10^8$$
$$T \sim 200\mu\text{K}$$



Optical dipole trap

$$N_{6\text{Li}} \sim 10^6$$
$$T \sim 200\mu\text{K}$$

Last week!

Next

preparation of two-component
Fermi gas

Simultaneous optical dipole
trap of 6Li and 7Li

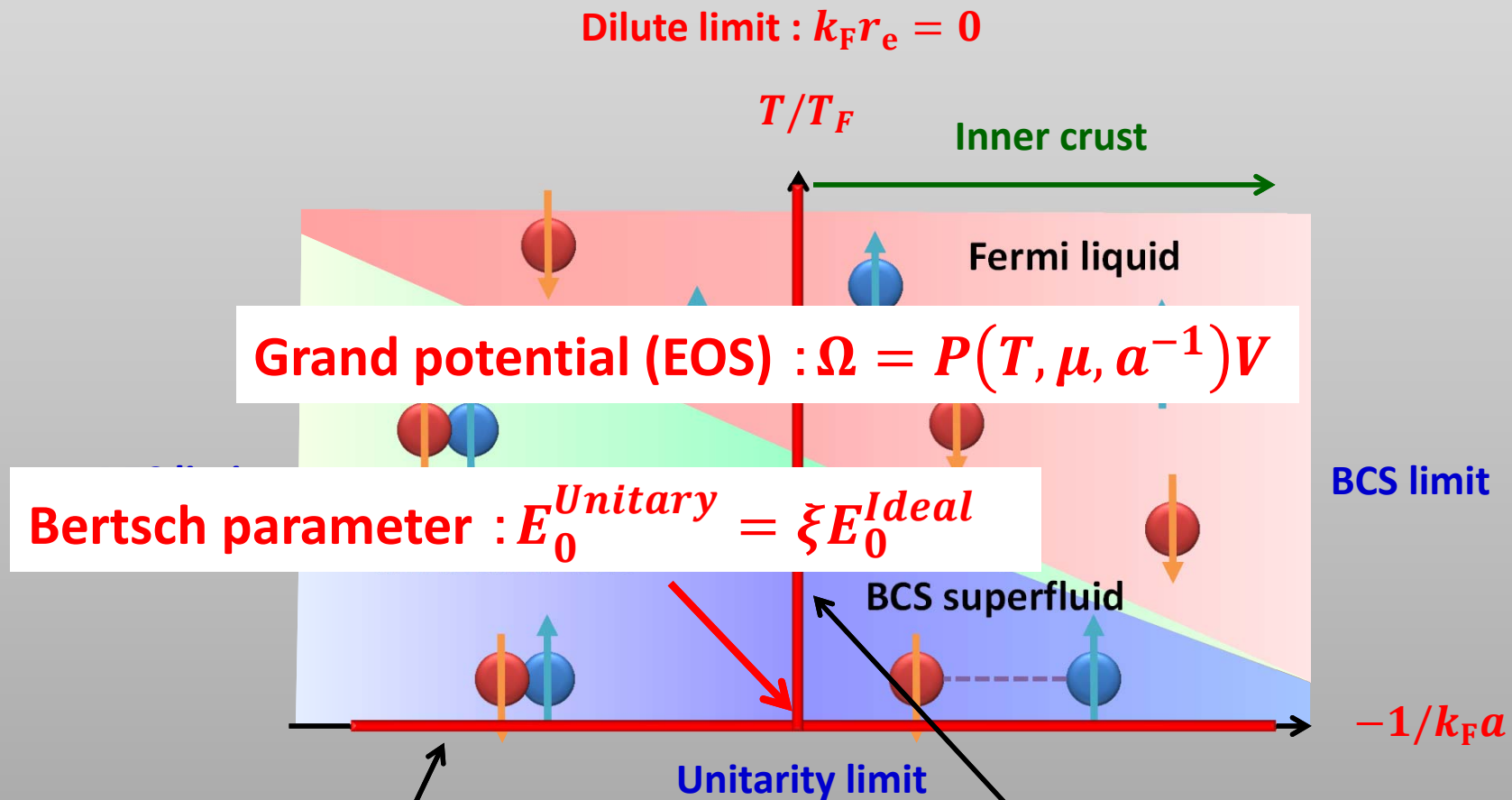
High-resolution imaging system



Evaporative cooling

→ **Degenerate Fermi gas**

BCS-BEC crossover in 2-component Fermi systems



[N. Navon, *et al.*, Science **328**, 729 (2010)]

[M. Horikoshi, *et al.*, Science **327**, 442 (2010)]

[S. Nascimbène, *et al.*, Nature **463**, 1057 (2010)]

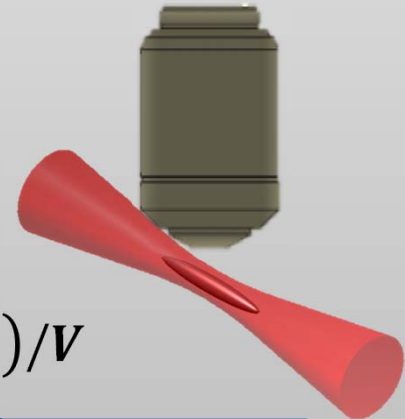
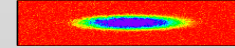
[M. Ku, *et al.*, Science **335**, 563 (2012)]

Our method to determine the EOS

EOS for dilute Fermi gases : $P(T, \mu, \alpha^{-1})$



But, cold atoms are trapped in a harmonic trap
The gas has **inhomogeneous density distribution**



Position dependent EOS : $P(\mathbf{r}) = \Omega(T, \mu(\mathbf{r}), \alpha(B)^{-1})/V$

- **Local pressure $P(\mathbf{r})$** : density distribution is pressure distribution
⇒ **Resolution of imaging determine precision of the EOS**

[Tin-Lun Ho. Nature Physics 6, 131 (2010)]

- **Temperature T** : mixing 7Li for thermometry into 6Li

[S. Nascimbène, *et al.*, Nature 463, 1057 (2010)]

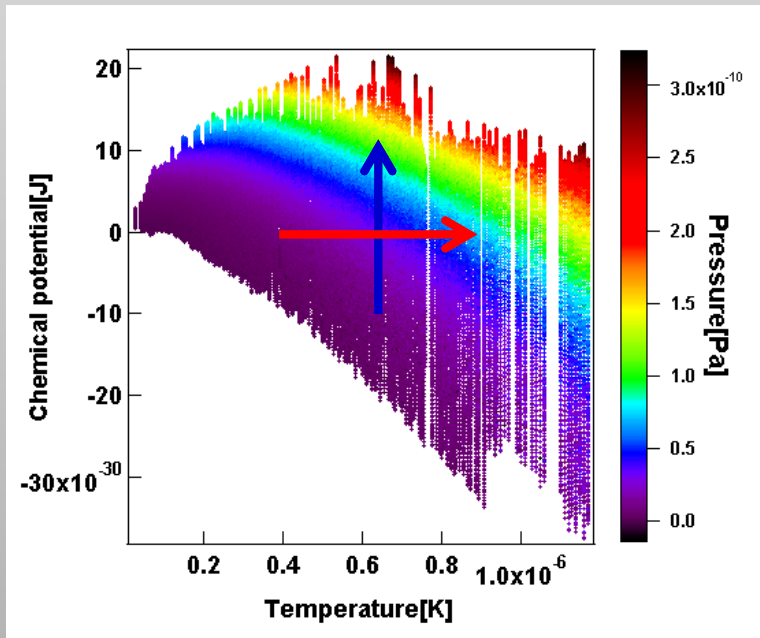
- **Local chemical potential $\mu(\mathbf{r})$** : LDA $\Rightarrow \mu(\mathbf{r}) = \mu(0) - U_{\text{trap}}(\mathbf{r})$

Thermodynamic relation for harmonically trapped system :

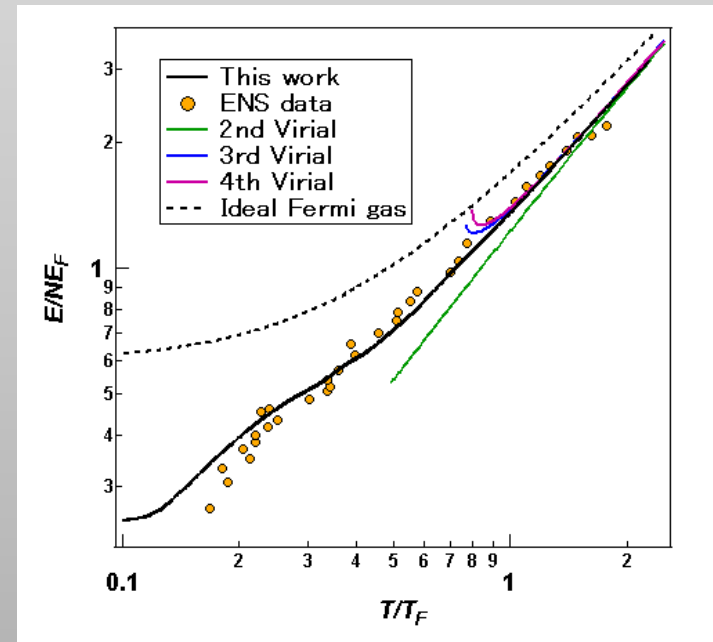
$$E_{\text{rel}} = TS + \mu(0)N - \frac{5}{3}E_{\text{pot}}$$

EOS of unitary Fermi gases

EOS via ENS's route : $P(\mu, T, a^{-1}=0)$



Internal energy



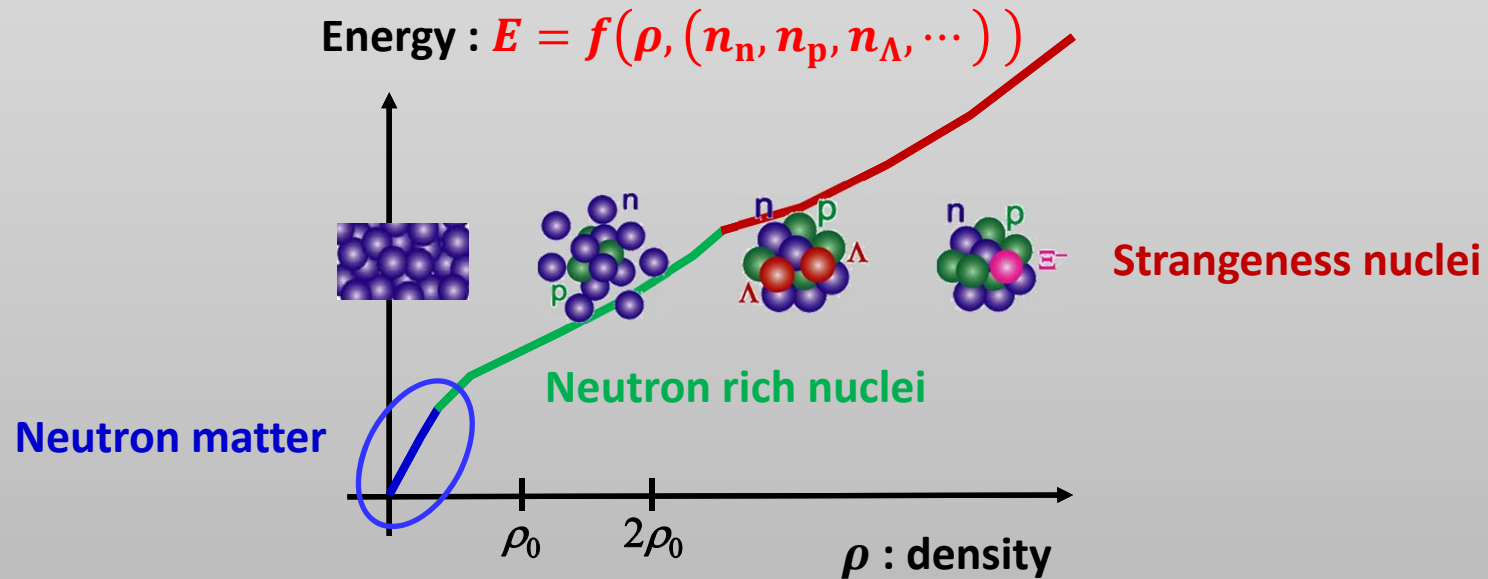
Particle density : $n = (dP/d\mu)_{T,a}$

Entropy density : $s = (dP/dT)_{\mu,a}$

Internal energy density : $\varepsilon = Ts + \mu n - P$

Analyzed data correspond to
[M. Horikoshi, *et al.*,
Science **327**, 442 (2010)]

Determination of EOS for neutron matter



Cold Fermi atoms :

$$\frac{E}{E_F} \left(\frac{T}{T_F}, -\frac{1}{k_F a} \right)$$

Correction

Neutron matter :

$$\frac{E}{E_F} \left(\frac{T}{T_F}, -\frac{1}{k_F a}, k_F r_e \right)$$

Neutron stars :

$$E = f(\rho, (n_n))$$

$$\frac{\partial E}{\partial(k_F r_e)} = \zeta \quad (\zeta = 0.127(4) @ |\infty| = 0, T = 0)$$

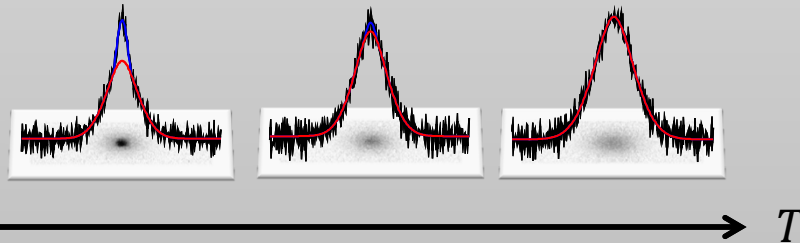
[Forbes, *et al.*, arXiv:1205.4815]

[Werner and Castin, Phys. Rev. A 86, 013626 (2012)]

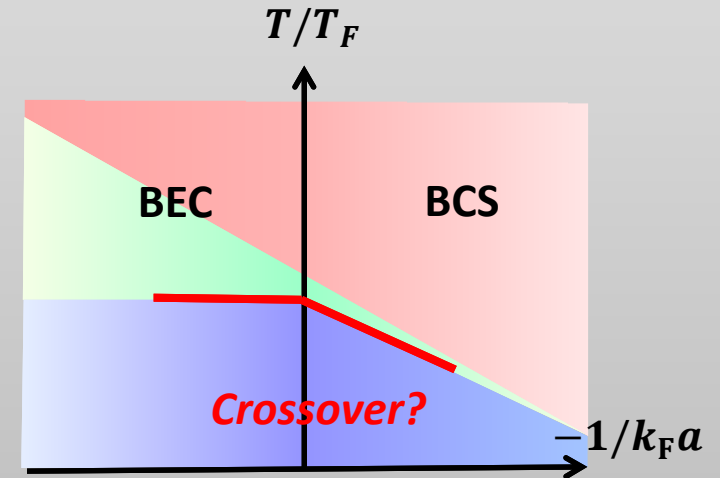
Superfluid phase transition

Emergence of cooper-pair with zero center-of-mass momentum

[M.Greiner *et al.*, *Nature* **426**, 537]

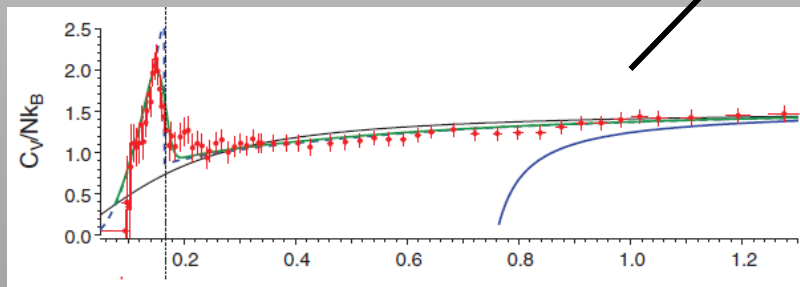


7Li thermometry



Lambda transition of heat capacity
Imaging resolution is 1.4 μ m

[Mark J. H. Ku, *et al.*, **335**, 563 (2012)]



T/T_F

Challenging issue

Universality class over the crossover

Critical exponent for correlation length :

$$\xi \propto \left| \frac{T-T_C}{T_C} \right|^{-\nu} \quad (\nu = 0.67 \text{ for 3D-XY model})$$

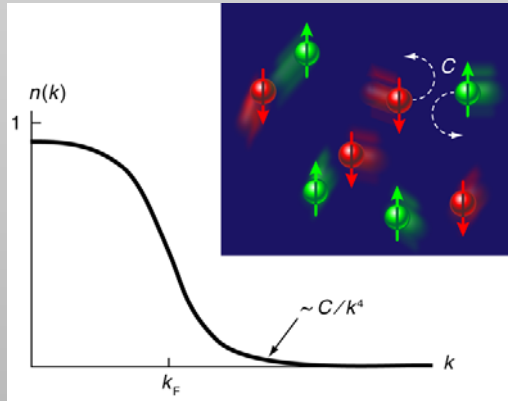
Critical exponent for heat capacity :

$$C_V \propto \left| \frac{T-T_C}{T_C} \right|^{-\alpha} \quad (\alpha = -0.013 \text{ for 3D-XY model})$$

$$\alpha = 2 - 3\nu$$

Universal many-body function, Tan's contact

In 2005, Shina Tan suggested that, the tail of momentum distribution
 [Shina Tan, cond-mat/0508320] for the universal 2-component Fermi system become



Physics **3**, 48 (2010)

Tan's Contact

$$n_{\uparrow\text{or}\downarrow}(k_F, |a|^{-1}, \Lambda_T^{-1} < k < R_{\text{vdw}}^{-1}) = \frac{C(x, \theta)}{k^4}$$

Fourier transform

$$C \equiv 4\pi k_F N h(x, \theta)$$

The origin is two-particle density matrix at short range :

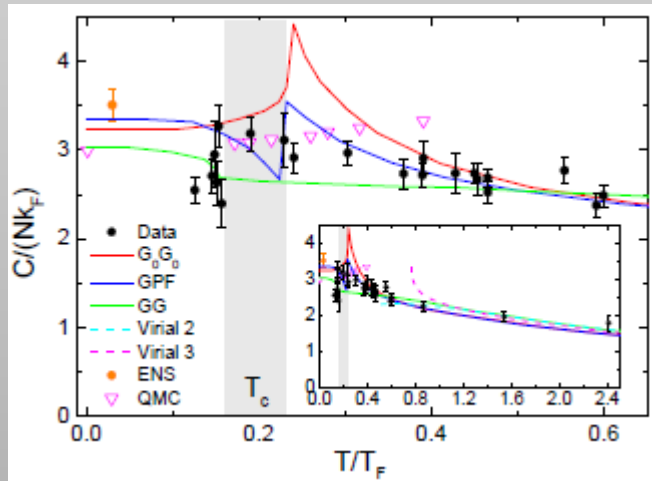
$$|\phi_{\text{pair}}(r)|^2 = \langle \psi_{\uparrow}^{\dagger}(\mathbf{r}_1) \psi_{\downarrow}^{\dagger}(\mathbf{r}_2) \psi_{\downarrow}(\mathbf{r}_2) \psi_{\uparrow}(\mathbf{r}_1) \rangle$$

$$\xrightarrow{a, k_F^{-1} \gg r \equiv |r_1 - r_2| \gtrsim R_{\text{vdw}}} 4\pi k_F N \cdot h(x, \theta) \cdot \left| \frac{\phi(r)}{4\pi} \right|^2, \left(\begin{array}{l} x \equiv -\frac{1}{k_F a}, \theta \equiv \frac{T}{T_F} \\ \phi(r) = \frac{1}{r} - \frac{1}{a} \end{array} \right)$$

Universal many-body function

Universal many-body function, Tan's contact

Recent measurement from JILA
at unitarity limit ($|a| = \infty$)

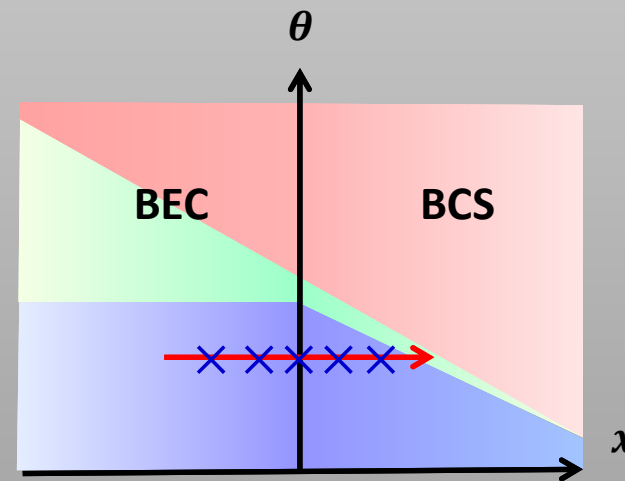


[Yoav Sagi, et al., arXiv:1208.2067]

Peak at the critical temperature?
Contact has some critical exponent?

Some Tan's relation

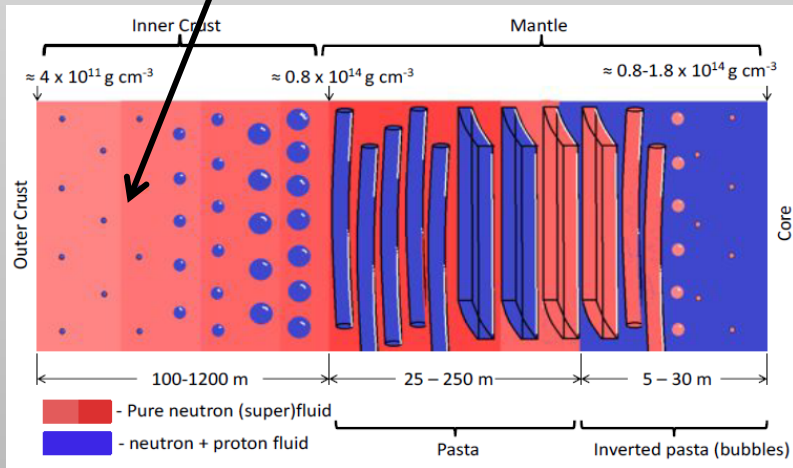
- Adiabatic relation : $\left. \frac{\partial E}{\partial x} \right|_{\theta} = 2\varepsilon_F N h(x, \theta) > 0$
- P - E relation : $PV = \frac{2}{3} (E - N\varepsilon_F x h(x, \theta))$



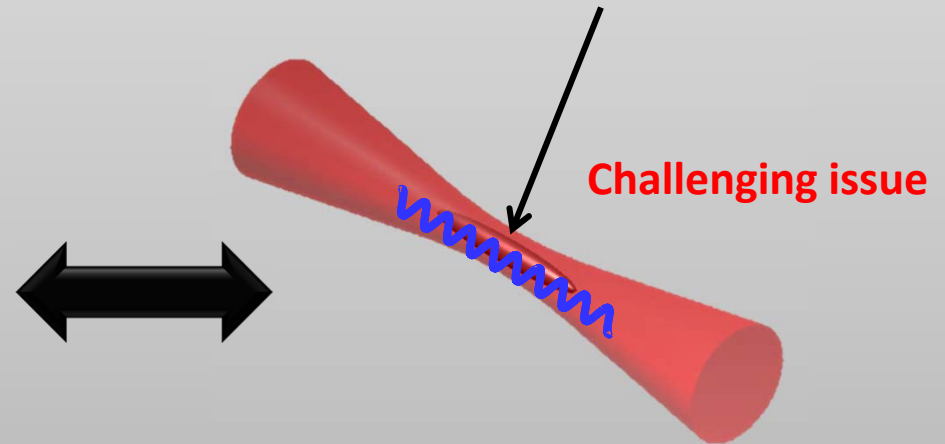
Verification of experimental result

Periodic potential

Dripped neutrons are moving in periodic potential by nuclei crystal



Tunable periodic potential by standing wave of laser beams

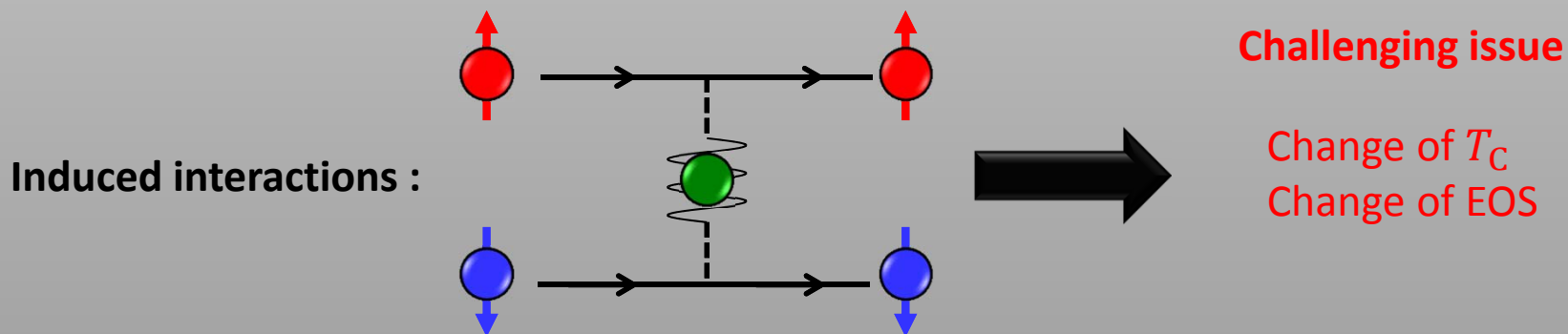


[G. Watanabe, *et al.*, Phys. Rev. Lett. **107**, 270404 (2011)]

	Cold atom	Neutron matter
Fermion	6Li, 40K	neutrons
Interparticle distance : $\sim k_F^{-1}$	100nm	6~3fm
Periodic potential : d	500nm	15~50fm ?
Ratio : $k_F d$	5	5~10
Lattice potential	Optical lattice	Nuclei crystal

Induced interaction : effects of Protons

	Cold atom	Neutron rich matter
Mixture	6Li \uparrow -6Li \downarrow -7Li 6Li \uparrow -6Li \downarrow -6Li \rightarrow	n \uparrow -n \downarrow -p \uparrow
Type	Femi-Bose Fermi-Fermi	Fermi-Fermi
Scattering length : a_{nn}	$ \infty $ (unitarity limit)	-18.5 ± 0.3 fm
Scattering length : a_{np}	2nm (Femi-Bose) -174 nm($\uparrow\rightarrow$), -851 nm($\downarrow\rightarrow$)	5.423fm (triplet) -23.749 fm (singlet)
Interparticle distance : $\sim k_F^{-1}$	100nm	6~3fm



[H. Heiselberg, et al., Phys. Rev. Lett. **85**, 2418 (2000)]

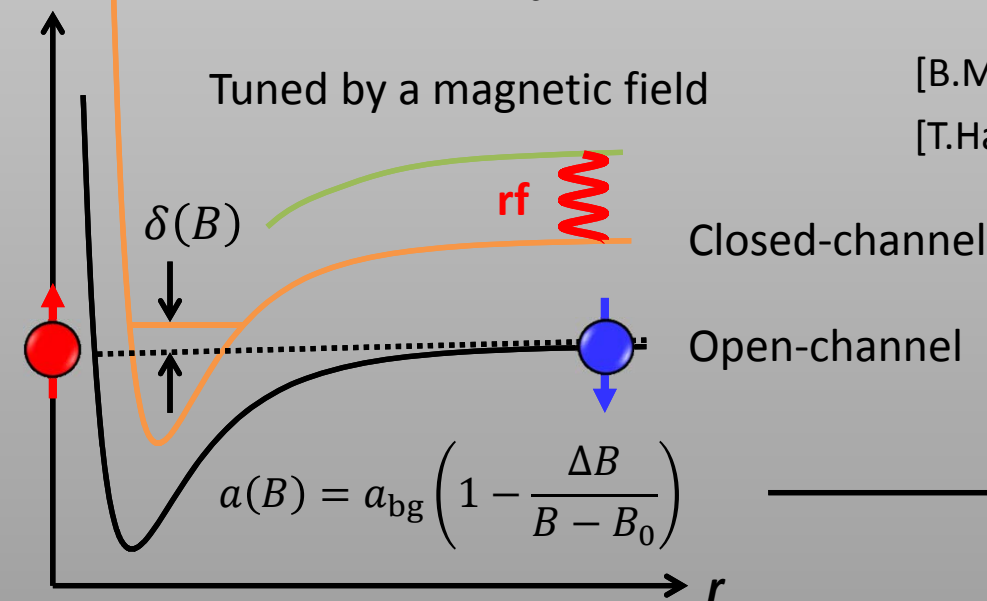
Tuning of effective range

	Cold atom	Neutron matter
Fermion	6Li	neutrons
Effective range: $k_F r_e$	0.05	0.53~3.3

Feshbach resonance

Challenging issue

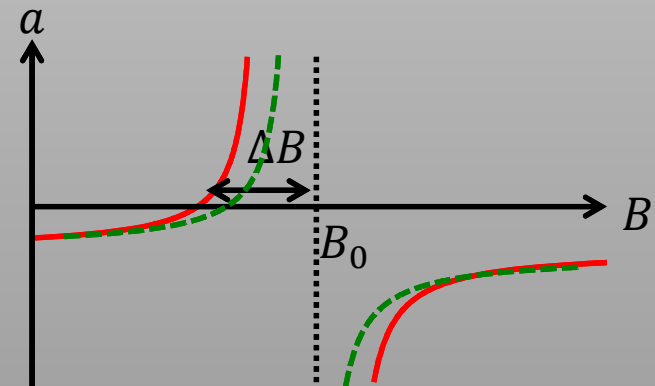
Effective range: $r_e \propto \Delta B^{-1}$



Manipulation of Feshbach resonances by rf field or dc-electric field

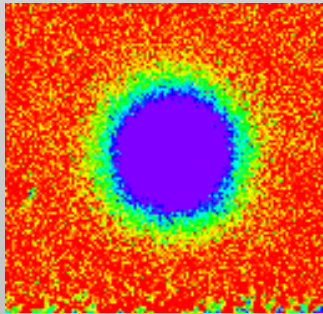
[B.Marcelis, et al., Phys. Rev. Lett. **100**, 153201 (2008)]

[T.Hanna, et al., New J. Phys. **12** 083031 (2010)]

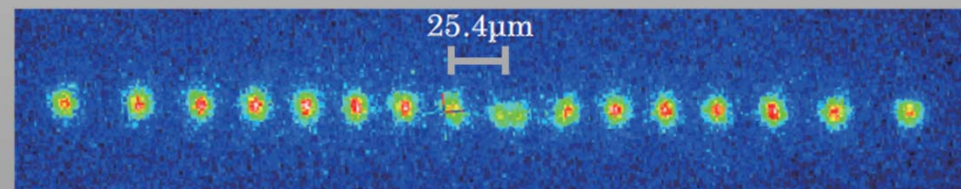
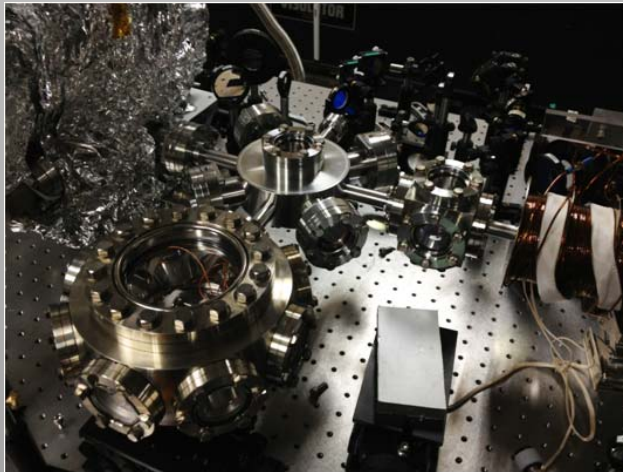
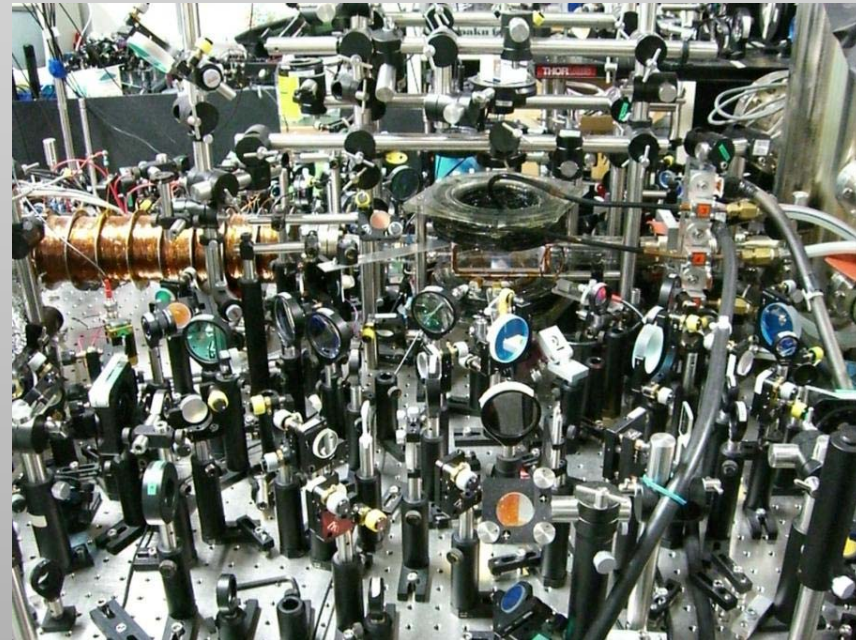


Cold Atom Lab at University of Electro-Communications

p-wave interaction, ${}^6\text{Li}$ - ${}^{40}\text{Ca}^+$ (Fermi-Ion) mixture



Interacting ${}^6\text{Li}$ by p-wave collisions

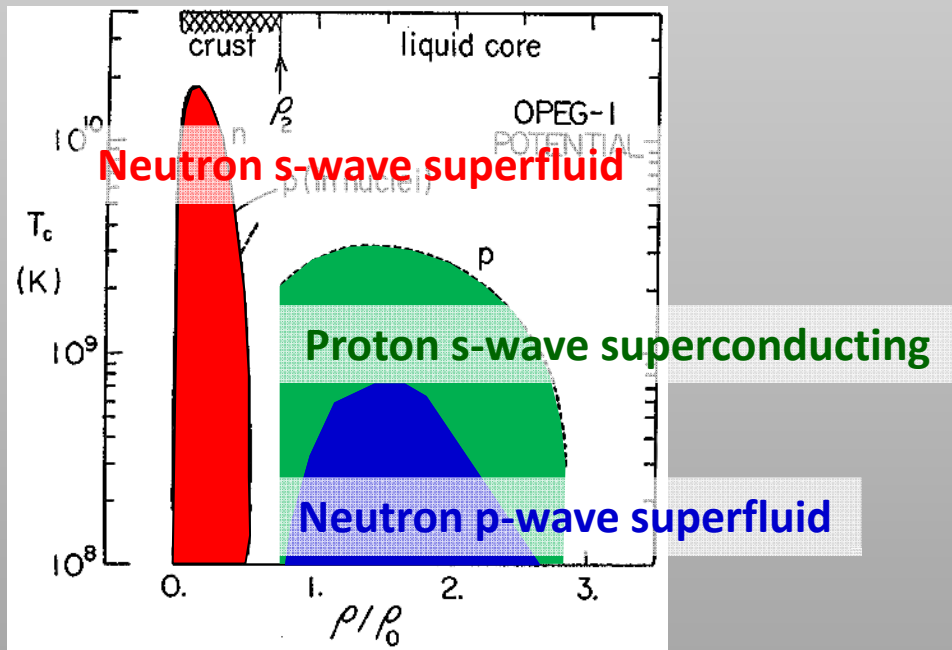


Crystalized cold ions

p-wave superfluidity in neutron stars

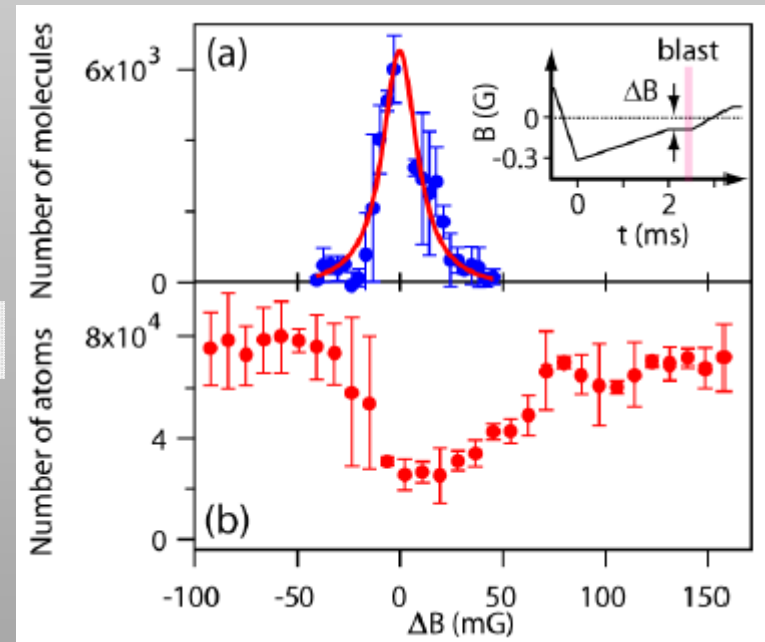
Missions : EOS of many-body Fermi system with **p-wave interactions**
Realization of **p-wave superfluid**

Superfluid phase diagram in neutron stars



[T. Takatsuka and R. Tamagaki, Progress of Theoretical Physics Supplement 112, 27 (1993)]

Cold 6Li atom with p-wave Feshbach resonance



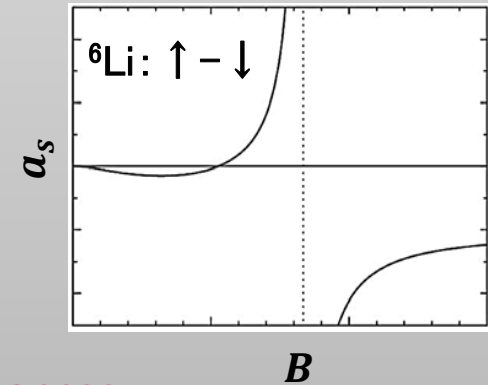
[Y.Inada, *et al.*, Phys. Rev. Lett. **101**, 100401 (2008)]

What is p-wave scattering parameters ?

By knowing parameters of s-wave Feshbach resonances, studies of many-body physics using s-wave interactions have progressed dramatically

s-wave interaction :

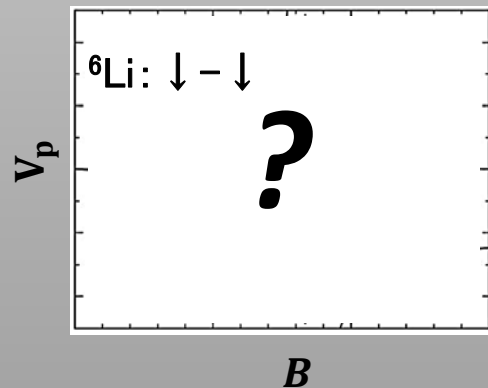
$$f_s(k, B) = \frac{1}{-\frac{1}{a_s(B)} + \frac{1}{2}r_e k^2 - ik} \quad a_s(B) = a_{bg} \left(1 - \frac{\Delta B}{B - B_{res}} \right)$$



Today, no one knows parameters of p-wave Feshbach resonances

p-wave interaction :

$$f_p(k, B) = \frac{1}{-\frac{1}{V_p(B)} + \frac{1}{2}r_e k^2 - ik^3} \quad V_p(B) = V_{bg} \left(1 - \frac{\Delta B}{B - B_{res}} \right)$$

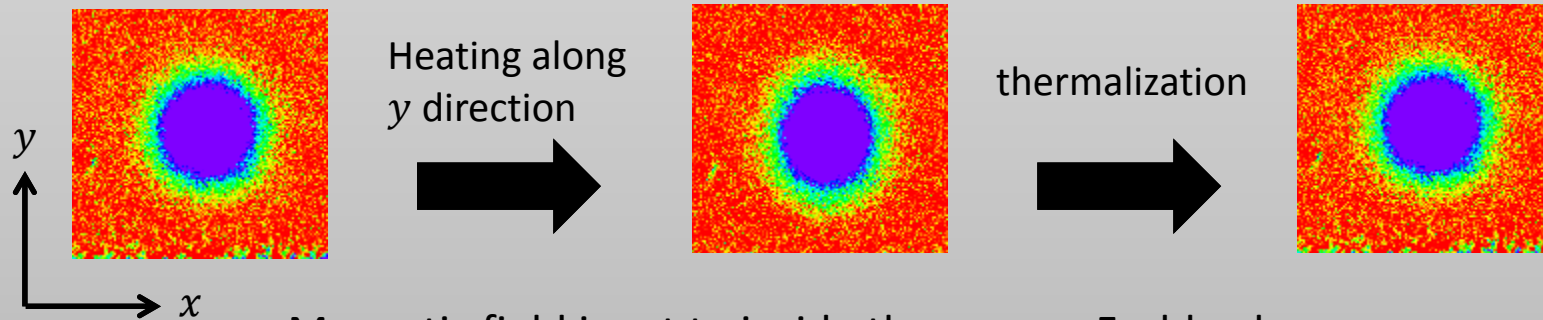


First of all, we must determine them

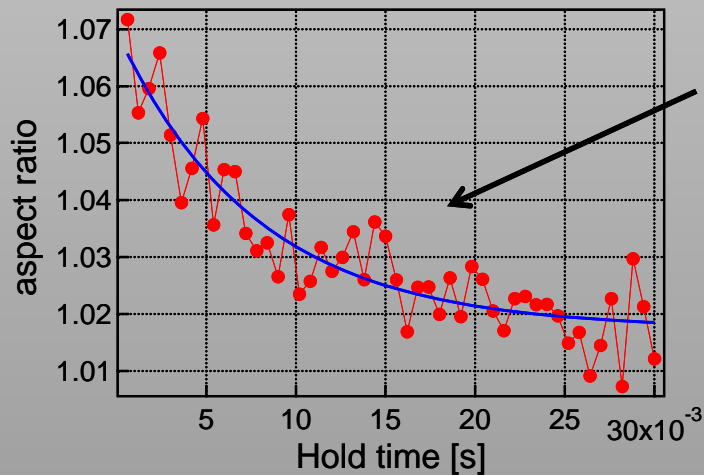
Determination of p-wave scattering parameters

Method : measurement of thermalization rate

Momentum(temperature) space



Magnetic field is set to inside the p-wave Feshbach resonance



Exponential decay with a thermalization rate of τ

Collisional cross section averaged over the trap : $\bar{\sigma}(T, B) = \frac{2.7\tau}{\bar{n}\bar{v}}$

Elastic collision cross section : $\sigma = 4\pi|f_p(k)|$

Preliminary data

Elastic collision cross section : $\sigma = 4\pi |f_p(k, V_p(B).r_e)|^2$

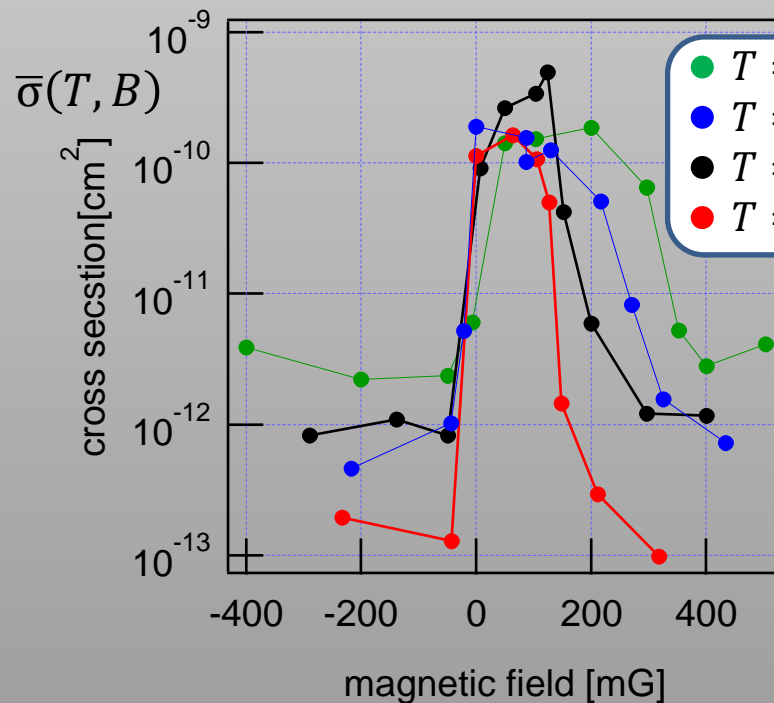
Thermal averaging under the assumption of Maxwell-Boltzman distribution



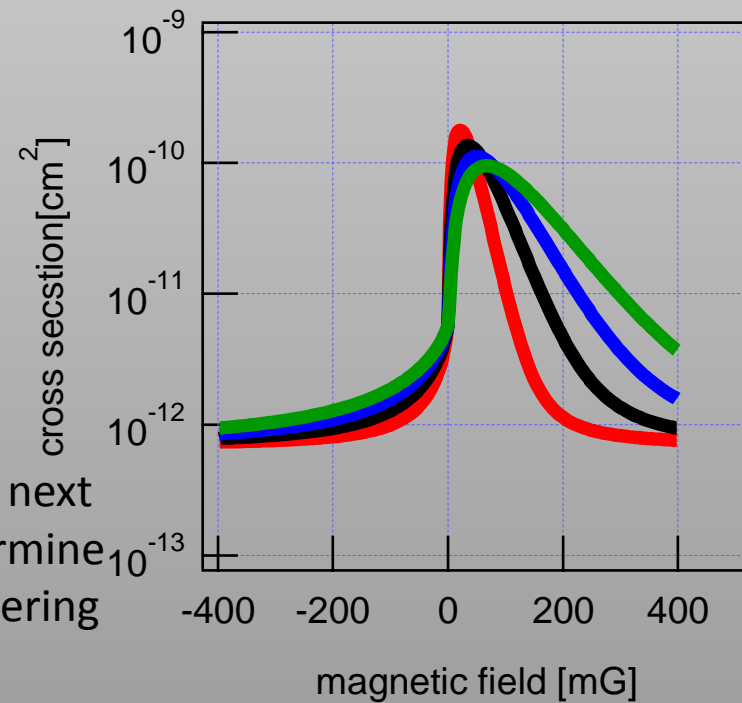
$$\langle \sigma(T, V_p(B).r_e) \rangle = (k_B T)^{-2} \int_0^\infty 4\pi |f_p(E, V_p(B).r_e)|^2 E e^{-E/k_B T} dE$$

Experimental data

$\longleftrightarrow \bar{\sigma}(T, B)$

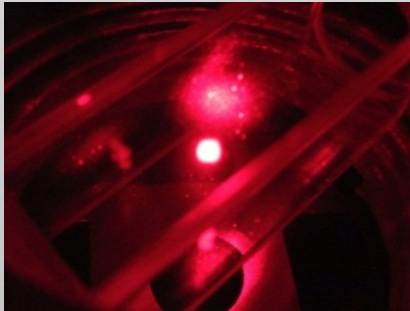


Fitting is the next step to determine p-wave scattering parameters



Summary

Cold atoms

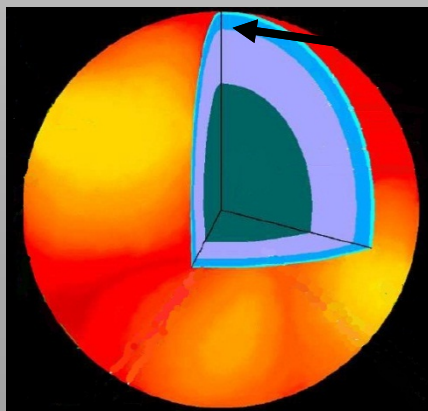


Measurements



- Universal equation of state
- Internal energy
- Specific heat
- Critical temperature
- s,p-wave superfluid
- Superfluid gap
- Superfluid density
- ⋮
- Benchmark for theories

Inner crust of neutron stars



EOS

Cooling curve



- ◆ Correction of the effective range
- ◆ Lattice of neutron-rich nuclei
- ◆ Protons

Theories

